



From quarks and gluons to hadrons and multiquarks

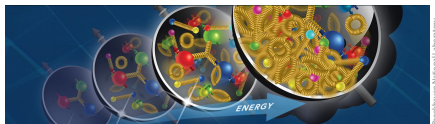
Gernot Eichmann

LIP Lisboa

LIP Seminar
June 4, 2020

QCD and nuclear physics

- Quark-gluon structure of hadrons



- From quarks and gluons to nuclei

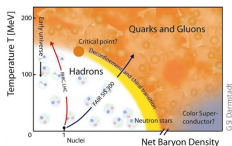


- Mass generation and confinement

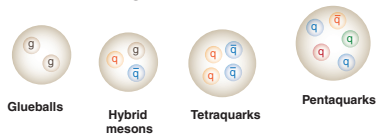
Higgs

QCD

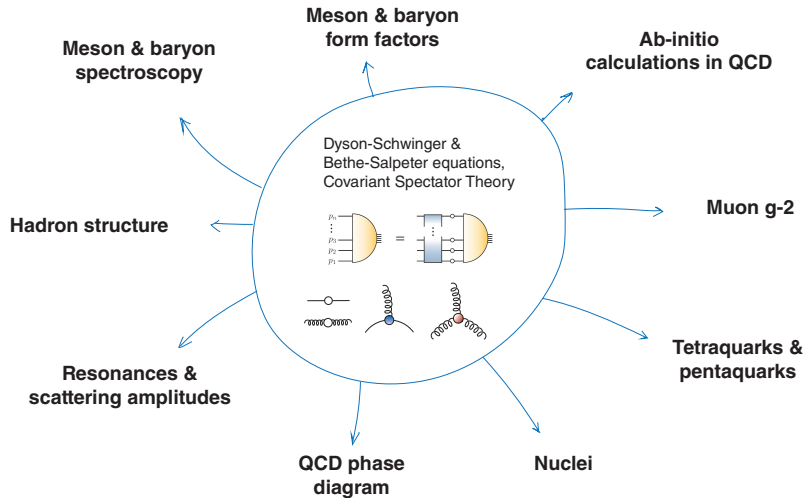
- The phases of QCD



- Understanding exotic hadrons

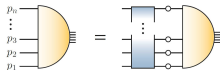


Our research



Methods

- Hadronic **bound-state equations** (BSEs, Faddeev eqs, ...)



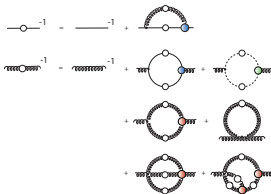
“QFT analogue of Schrödinger eq.”

- hadron masses & “wave functions”
- **spectroscopy calculations**

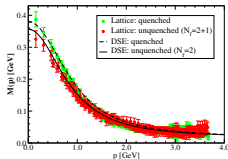
- Similar: **CST** (Covariant spectator theory)

Stadler & Gross, PRL 78 (1997), PRC 78 (2008),
Leitão, Stadler, Peña, Biernat, PLB 764 (2017)

- Ingredients: **QCD’s n-point functions**, Satisfy Dyson-Schwinger equations (**DSEs**): QCD’s quantum eqs. of motion

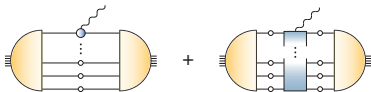


...



→ running **quark mass**

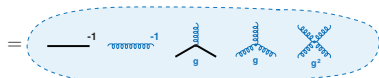
- Structure calculations: form factors, PDFs, GPDs, TMDs, two-photon processes, ...



Everything runs

QCD's classical action:

$$S = \int d^4x \left[\bar{\psi} (\not{\partial} + ig\mathcal{A} + m) \psi + \frac{1}{4} F_{\mu\nu}^a F_a^{\mu\nu} \right]$$

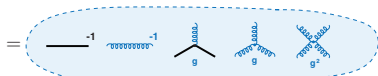


Tree-level propagators & vertices

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QCD's classical action:

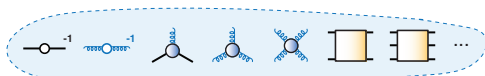
$$S = \int d^4x [\bar{\psi} (\not{\partial} + ig\mathbf{A} + m) \psi + \frac{1}{4} F_{\mu\nu}^a F_a^{\mu\nu}]$$



Tree-level propagators & vertices

Quantum “effective” action:

$$\int \mathcal{D}[\psi, \bar{\psi}, A] e^{-S} = e^{-\Gamma}$$

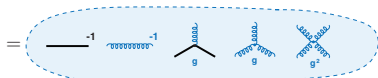


Dressed propagators & vertices

Everything runs

QCD's classical action:

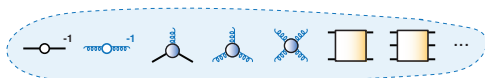
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Tree-level propagators & vertices

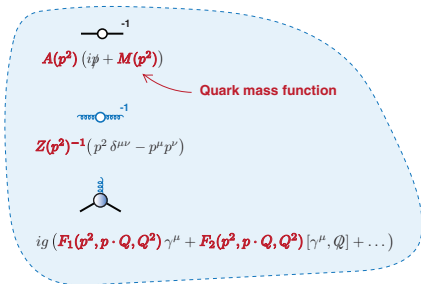
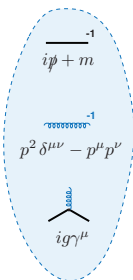
Quantum "effective" action:

$$\int \mathcal{D}[\psi, \bar{\psi}, A] e^{-S} = e^{-\Gamma}$$



Dressed propagators & vertices

Quantum also means **structure**:



- **everything runs** with momentum
- **dressing functions** encode everything we can know about quarks & gluons
- encode **full QFT**

Unwrapping the quark DSE

$$\text{---}\bigcirc\text{---}^{-1} = \text{---}^{-1} - \text{---}\bigcirc\text{---}^{-1} - \text{---}\bigcirc\text{---}^{-1} - \text{---}\bigcirc\text{---}^{-1} - \text{---}\bigcirc\text{---}^{-1} - \text{---}\bigcirc\text{---}^{-1} - \dots$$

- DSE generates every diagram in **perturbation theory**:

$$S^{-1} = S_0^{-1} - \Sigma \Rightarrow S = S_0 + S_0 \Sigma S$$

$$= S_0 + S_0 \Sigma S_0 + S_0 \Sigma S_0 \Sigma S_0 + \dots$$

$$\text{---}\bigcirc\text{---} = \text{---} + \text{---}\bigcirc\text{---} + \text{---}\bigcirc\text{---} + \text{---}\bigcirc\text{---} + \dots$$

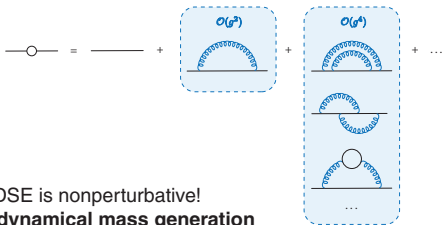
Unwrapping the quark DSE

$$\text{---}\text{---}\text{---}^{-1} = \text{---}\text{---}\text{---}^{-1} - \text{---}\text{---}\text{---}\text{---}\text{---}\text{---}$$

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- But DSE is nonperturbative!
E.g. **dynamical mass generation**

$$S(p)^{-1} = A(p^2) (i\not{p} + M(p^2)) \Rightarrow M(p^2) \sim \text{Tr} S(p)^{-1}$$

Massless quarks:

$$\text{---}\text{---}\text{---}^{-1} \quad \text{---}\text{---}\text{---}$$

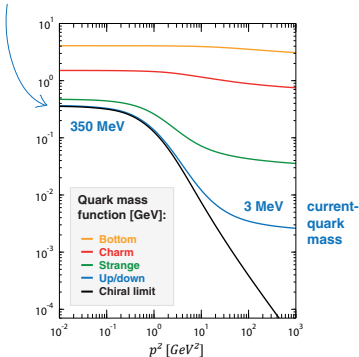
$$i\not{p} \quad ig\gamma^\mu$$

Every perturbative diagram
contains odd # gamma matrices
 \Rightarrow mass function always zero!

Unwrapping the quark DSE

$$\text{---} \circ \text{---}^{-1} = \text{---} \text{---}^{-1} - \text{---} \circ \text{---} \text{---}^{-1}$$

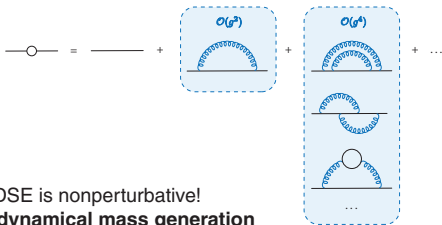
“constituent-quark mass”:
nonperturbative effect



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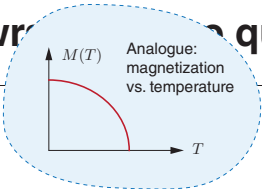
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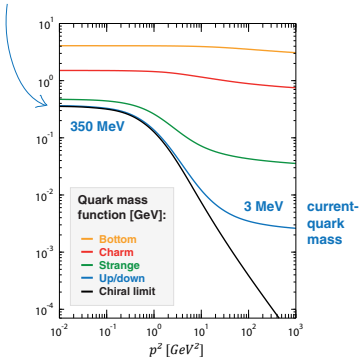
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Every perturbative diagram
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Unwrapping the quark DSE



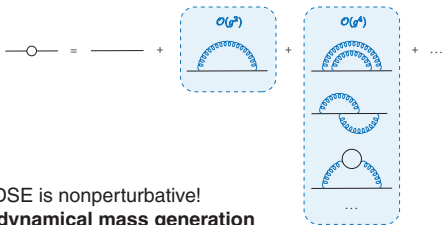
“constituent-quark mass”: nonperturbative effect



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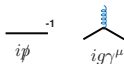
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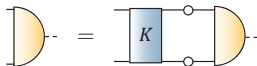
Massless quarks:



Every perturbative diagram contains odd # gamma matrices \Rightarrow mass function always zero!

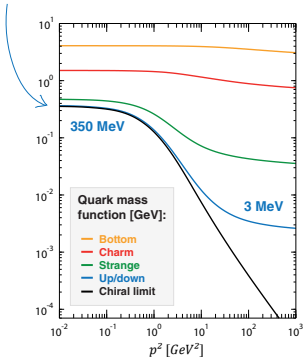
Massless pion

- Bethe-Salpeter equation for **mesons**:



Mass generation for hadrons,
but pion is **Goldstone boson**

“constituent-quark mass”:
nonperturbative effect

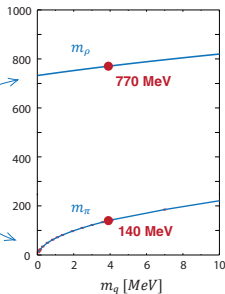


current-quark mass

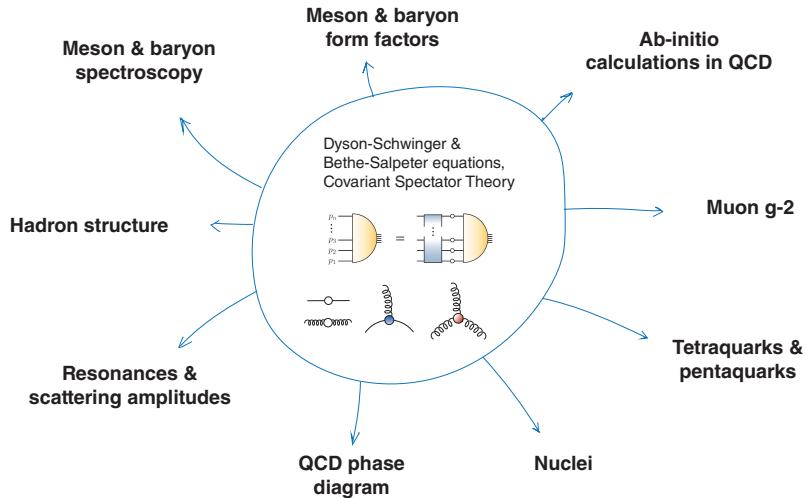
Effects are connected:
Dynamical breaking
of chiral symmetry

$$m_\pi^2 \sim m_q$$

Meson masses [MeV]

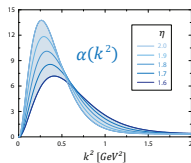
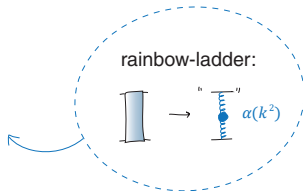
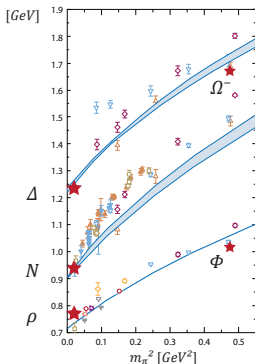
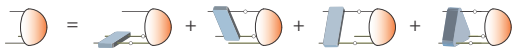


Our research



Baryon spectroscopy

- Solution of nucleon's covariant **3-body Faddeev equation**
 GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)

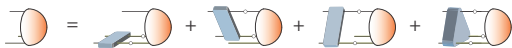


Maris, Tandy, PRC 60 (1999),
 Qin et al., PRC 84 (2011)

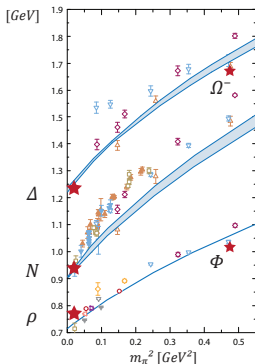
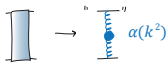
Baryon spectroscopy

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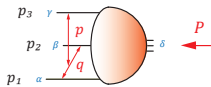
GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)



rainbow-ladder:



Keep full structure of baryon's Faddeev amplitude:



$$\Psi_{\alpha\beta\gamma\delta}(p, q, P) = \sum_i f_i(p^2, q^2, p \cdot q, p \cdot P, q \cdot P) \tau_i(p, q, P)_{\alpha\beta\gamma\delta}$$

Lorentz-invariant dressing functions

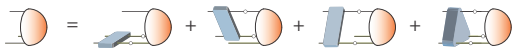
Dirac-Lorentz tensors:
64 for J = 1/2,
128 for J = 3/2

Review: GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, Prog. Part. Nucl. Phys. 91 (2016), 1606.09602

Baryon spectroscopy

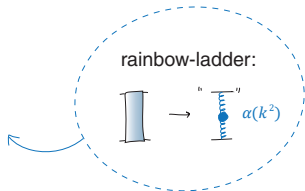
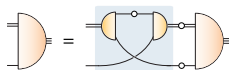
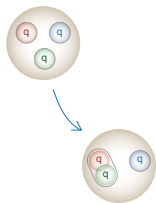
- Solution of nucleon's covariant **3-body Faddeev equation**

GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)



- Simpler: **quark-diquark Faddeev equation**

Oettel et al., PRC 58 (1998), Cloet et al., FBS 46 (2009), Chen et al., PRD 97 (2018)



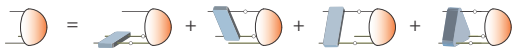
Ping-pong exchange
between quark and diquark

Calculate all ingredients in rainbow-ladder:
can compare **quark-diquark** and **3-body** directly
GE, Krassnigg, Schwinzerl, Alkofer, Ann. Phys. 323 (2008)

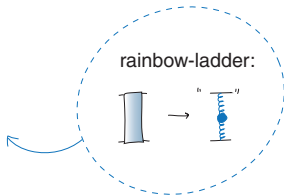
Baryon spectroscopy

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GE, Alkofer, Nicmorus, Krassnigg, PRL 104 (2010)

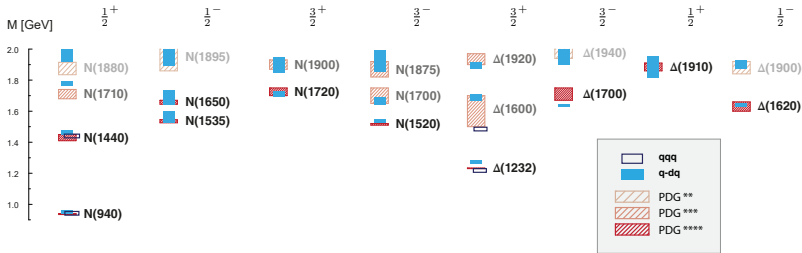


rainbow-ladder:



- Similar results in **quark-diquark** description

GE, Fischer, Sanchis-Alepuz, PRD 94 (2016)

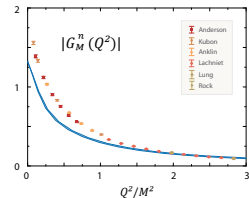
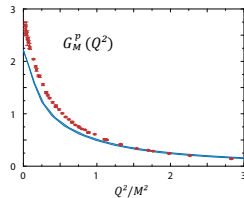
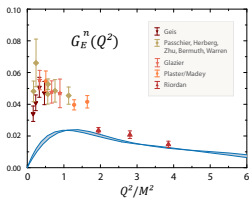
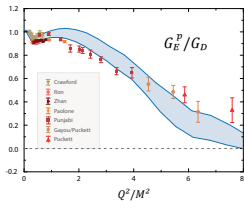


Baryon form factors

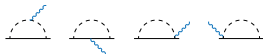
Nucleon em. form factors from three-quark equation

GE, PRD 84 (2011)

$$J^\mu = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4}$$



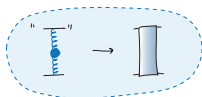
- “Quark core without pion cloud”



- similar: $N \rightarrow \Delta \gamma$ transition, axial & pseudoscalar FFs, octet & decuplet em. FFs

Review: GE, Sanchis-Alepuz, Williams, Fischer, Alkofer, PPNP 91 (2016), 1606.09602

Ab-initio calculations



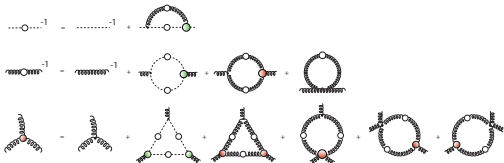
- **Three-gluon vertex** with all tensors:
zero crossing at low momenta

GE, Williams, Alkofer, Vujanovic, PRD 89 (2014)



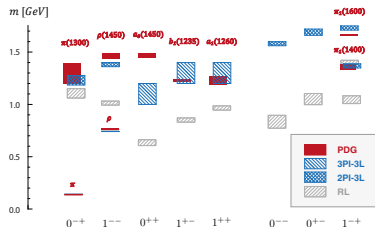
- **Mass generation** in Yang-Mills sector

GE, Pawłowski, Silva, in preparation

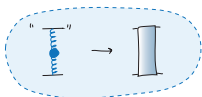


- Beyond rainbow-ladder calculations
improve **light-meson spectrum**

Williams, Fischer, Heupel, PRD 93 (2016)



Ab-initio calculations



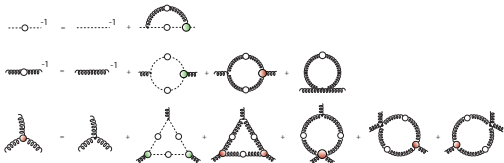
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GE, Williams, Alkofer, Vujanovic, PRD 89 (2014)



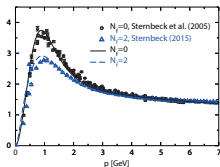
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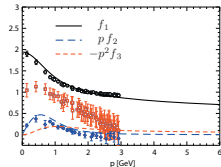


- Beyond rainbow-ladder calculations so far (mostly) only for **mesons**
- Need more interplay between **DSEs, FRG, lattice**
Huber, 2003.13703 [hep-ph]
Cyrol, Mitter, Pawłowski, Strodthoff, PRD 97 (2018)
Oliveira, Silva, Skullerud, Sternbeck, PRD 99 (2019)

Gluon propagator



Quark-gluon vertex



GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, Prog. Part. Nucl. Phys. 91 (2016)

Muon g-2

- Muon anomalous magnetic moment:**
total SM prediction deviates from exp. by $\sim 4\sigma$,
theory uncertainty dominated by **QCD**



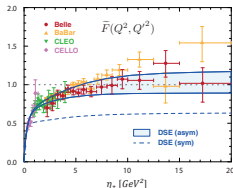
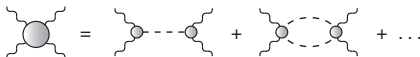
Hadronic vacuum polarization



Hadronic light-by-light scattering

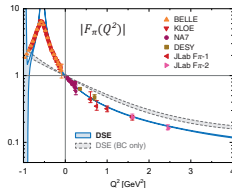
- Dispersion theory:** LbL amplitude through intermediate onshell mesons

Colangelo, Hoferichter, Kubis, Procura, Stoffer, PLB 738 (2014)



Pion transition form factor

Pion em. form factor



Knowledge of FFs in spacelike region helps constrain theory uncertainty

GE, Fischer, Weil, Williams, PLB 774 (2017), PRD 96 (2017), PLB 797 (2019), PRD 101 (2020)

Jegerlehner, Nyffeler, Phys. Rept. 477 (2009), Jegerlehner 1705.00263

$a_\mu [10^{-10}]$

QED: 11 658 471.9 (0.0)

EW: 15.4 (0.1)

Hadronic:

• VP (LO+HO) 680.1 (3.4)

• **LbL 10.3 (2.9) ?**

SM: 11 659 177.6 (4.4)

Exp: 11 659 208.9 (6.3)

Diff: 31.3 (7.7)

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Jegerlehner, Nyffeler,
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$$a_\mu [10^{-10}]$$



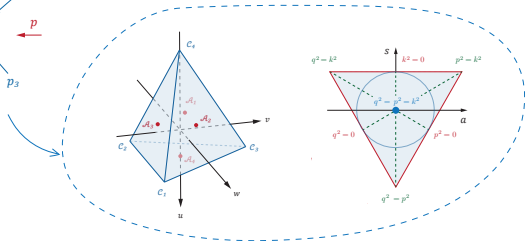
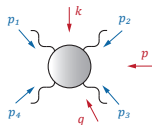
Hadronic vacuum polarization



Hadronic light-by-light scattering

- Structure of LbL amplitude from **permutation group S4**

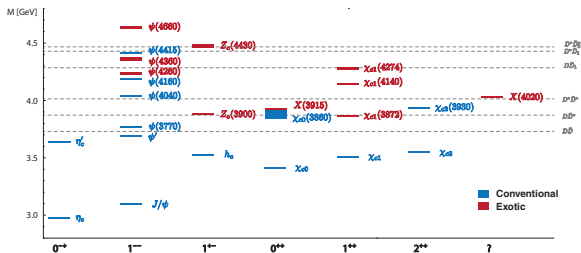
GE, Fischer, Heupel, PRD 92 (2015)



136 independent tensors ...

QED:	11 658 471.9	(0.0)
EW:	15.4	(0.1)
Hadronic:		
• VP (LO+HO)	680.1	(3.4)
• LbL	10.3	(2.9) ?
SM:	11 659 177.6	(4.4)
Exp:	11 659 208.9	(6.3)
Diff:	31.3	(7.7)

Exotic mesons



- Several tetraquark candidates in **charmonium spectrum**: X(3872), X(3915), Zc(3900), ...
- Z states cannot be $c\bar{c}$ since they carry charge
- Oldest tetraquark candidates: **light scalar mesons**

Reviews:

Chen, Chen, Liu, Zhu,
Phys. Rept. 639 (2016), 1601.02092

Lebed, Mitchell, Swanson
PPNP 93 (2017), 1610.04528

Esposito, Pilloni, Polosa,
Phys. Rept 668 (2017), 1611.07920

Guo, Hanhart, Meißner et al.,
Rev. Mod. Phys. 90 (2018), 1705.00141

Ali, Lange, Stone,
PPNP 97 (2017), 1706.00610

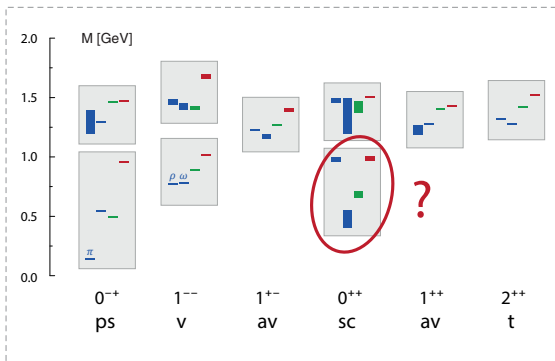
Olsen, Skwarnicki, Zieminska,
Rev. Mod. Phys. 90 (2019), 1708.04012

Liu, Chen, Chen, Liu, Zhu,
PPNP 107 (2019), 1903.11976

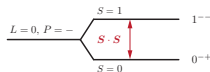
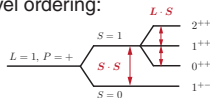
Brambilla, Eidelman, Hanhart et al.,
1907.07583

Light scalar mesons

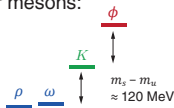
Light meson spectrum (PDG):
grouped with J^{PC} and flavor content



- Nonrelativistic level ordering:



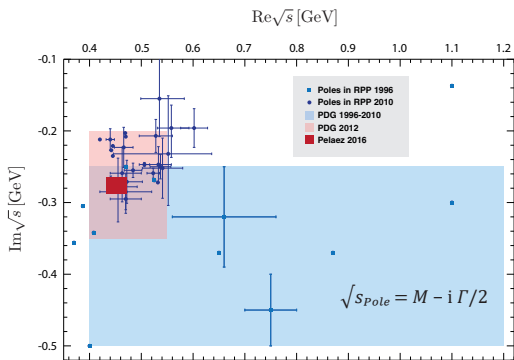
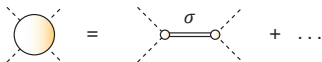
- Vector mesons:



- Pseudoscalar mesons?
spontaneous chiral symmetry breaking & axial anomaly
- Scalar mesons..?

Light scalar mesons

$\sigma/f_0(500)$ is a **resonance** in $\pi\pi$ scattering:



Pelaez, Phys. Rept. 658 (2016) 1

- **PDG 2010:** “ $f_0(600)$ ”

$$\sqrt{s} \sim (400 \dots 1200) - i(250 \dots 500) \text{ MeV}$$

- **Dispersive analyses:**

$$\sqrt{s} \sim 450(20) - i 275(10) \text{ MeV}$$

Caprini, Colangelo, Leutwyler 2006

Garcia-Martin, Kaminski, Pelaez, Ruiz de Elvira 2011

Moussallam 2011

Masjuan, Ruiz de Elvira, Sanz-Cillero 2014

Pelaez 2016

- **PDG 2012:** “ $f_0(500)$ ”

$$\sqrt{s} \sim (400 \dots 550) - i(200 \dots 350) \text{ MeV}$$

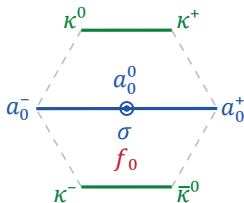
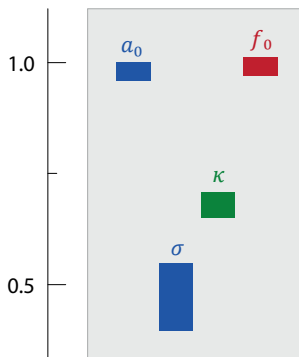
- Pole locations from **lattice QCD**

Briceno, Dudek, Edwards, Wilson, PRL 118 (2017),

PRD 97 (2018)

Light scalar mesons

Light scalar (0^{++}) mesons don't fit into the conventional meson spectrum:



f_0 (980 MeV) $s\bar{s}$
 κ (680 MeV) $u\bar{s}, d\bar{s}$
 a_0 (980 MeV) } $u\bar{u}, d\bar{d}, u\bar{d}$
 σ (500 MeV) }

- Why are a_0, f_0 mass-degenerate?
- Why are their **decay widths** so different?

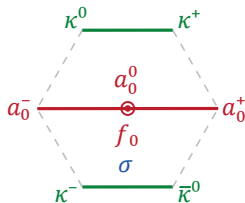
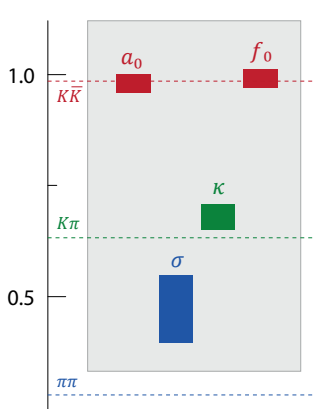
$$\Gamma(\sigma, \kappa) \approx 550 \text{ MeV}$$

$$\Gamma(a_0, f_0) \approx 50\text{--}100 \text{ MeV}$$

- Why are they so **light**?
 Scalar mesons ~ **p-waves**, should have masses similar to axialvector & tensor mesons ~ 1.3 GeV

Light scalar mesons

What if they were **tetraquarks** (diquark-antidiquark)? Jaffe 1977, Close, Tornqvist 2002, Maiani, Polosa, Riquer 2004



f_0 (980 MeV) } $us\bar{u}s, \dots$
 a_0 (980 MeV) }
 κ (800 MeV) } $us\bar{u}d, \dots$
 σ (500 MeV) } $ud\bar{u}d$

- Explains **mass ordering & decay widths**:
 f_0 and a_0 couple to $K\bar{K}$, large widths for σ, κ

- Alternative: **meson molecules?**
Weinstein, Isgur 1982, 1990; Close, Isgur, Kumano 1993

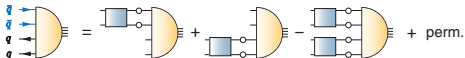
- Non- $q\bar{q}$ nature** of σ supported by dispersive analyses, unitarized ChPT, large N_c , extended linear σ model, quark models
Pelaez, Phys. Rept. 658 (2016)



Tetraquarks

- Light scalar mesons (σ , κ , a_0 , f_0) as **four-quark states**:

GE, Fischer, Heupel, PLB 753 (2016)



$$\Gamma(p, q, k, P) = \sum_i f_i(p^2, q^2, k^2, \{\omega_j\}, \{\eta_j\}) \tau_i(p, q, k, P) \otimes \text{Color} \otimes \text{Flavor}$$

9 Lorentz invariants:

$$p^2, \quad q^2, \quad k^2, \quad P^2 = -M^2$$

$$\omega_1 = q \cdot k \quad \eta_1 = p \cdot P$$

$$\omega_2 = p \cdot k \quad \eta_2 = q \cdot P$$

$$\omega_3 = p \cdot q \quad \eta_3 = k \cdot P$$

256 Dirac-Lorentz tensors

2 Color tensors:

$$3 \otimes \bar{3}, \quad 6 \otimes \bar{6} \text{ or}$$

$$1 \otimes 1, \quad 8 \otimes 8$$

(Fierz-equivalent)

- Group momentum variables into multiplets of **permutation group S4**:
can switch off groups of variables without destroying symmetries

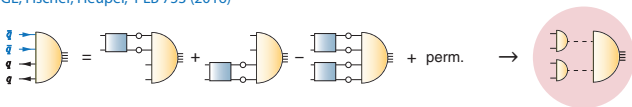
GE, Fischer, Heupel, PRD 92 (2015)

$$f_i(\mathcal{S}_0, \nabla, \blacklozenge, \circ)$$

Tetraquarks

- Light scalar mesons (σ , κ , a_0 , f_0) as **four-quark states**:

GE, Fischer, Heupel, PLB 753 (2016)



- BSE dynamically generates **meson poles** in BS amplitude:

$$f_i(\mathcal{S}_0, \nabla, \triangle, \circ) \rightarrow 1500 \text{ MeV}$$

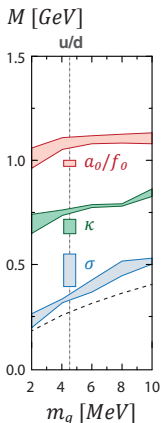
$$f_i(\mathcal{S}_0, \nabla, \triangle, \circ) \rightarrow 1500 \text{ MeV}$$

$$f_i(\mathcal{S}_0, \nabla, \triangle, \circ) \rightarrow 1200 \text{ MeV}$$

$$f_i(\mathcal{S}_0, \nabla, \triangle, \circ) \rightarrow \mathbf{350 \text{ MeV !}}$$



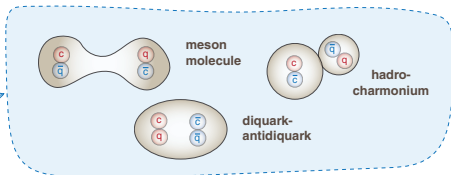
- “Light scalar mesons” look like **meson molecules**, diquark-antidiquark components almost negligible
- Lightness is inherited from pseudoscalar Goldstone bosons!



Tetraquarks

- Heavy-light **four-quark states**:
what is their internal decomposition?

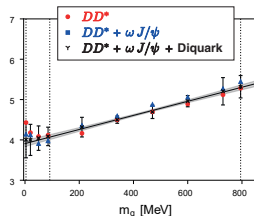
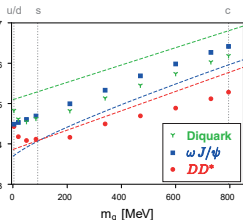
$cq\bar{q}\bar{c}$



- **Four-quark BSE**: all mix together

$M_{cq\bar{q}\bar{c}}$ [GeV]
in $X(3872)$
channel

--- $m_g + m_A$
--- $m_\omega + m_{J/\psi}$
--- $m_D + m_{D^*}$

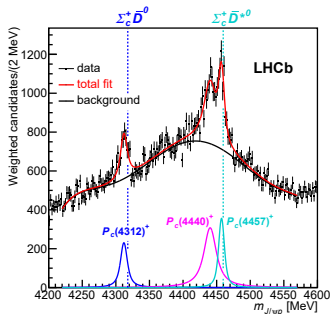


Wallbott, GE, Fischer,
PRD 100 (2019),
2003.12407 [hep-ph]

$cq\bar{q}\bar{c} \rightarrow$ strong meson-meson
component: DD^* for
 $X(3872)$, $Z_c(3900)$:

$cc\bar{q}\bar{q} \rightarrow$
diquarks also play role

Pentaquarks?



Aaij et al., PRL 112 (2019)

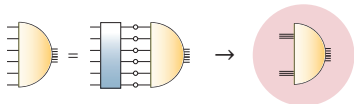


Madalena Lourenço

<https://www.youtube.com/user/projetoMEFT>

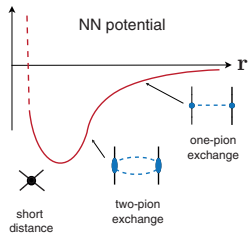
Nucleons in nuclei

Transition from quarks & gluons to **light nuclei**:

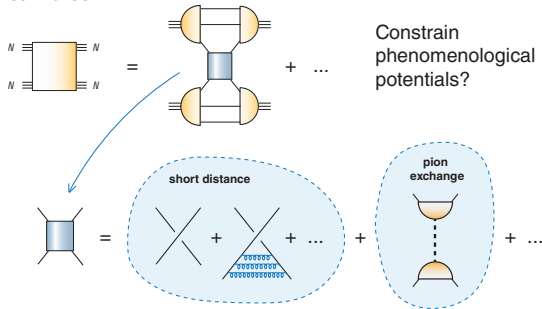


- Relativistic structure of **deuteron**
- Exotic dibaryons and hypernuclei
- **Short-range correlations**
- **EMC effect:** overlapping nucleons in nuclei?

Microscopic origins of **short-range nuclear force**?



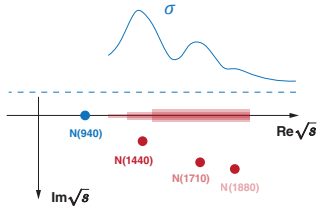
Weise, Nucl. Phys. A805 (2008)



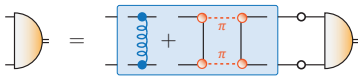
Constrain phenomenological potentials?

Resonances

- Most hadrons are **resonances** and decay
 \Leftrightarrow poles in complex momentum plane



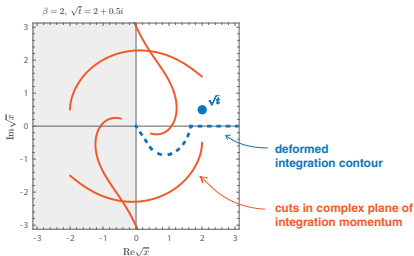
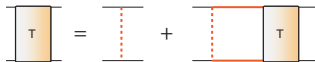
- BSE kernel must be aware of decay channels



real poles produce cuts \Rightarrow threshold,
 ρ meson becomes resonance

Williams, PLB 798 (2019)

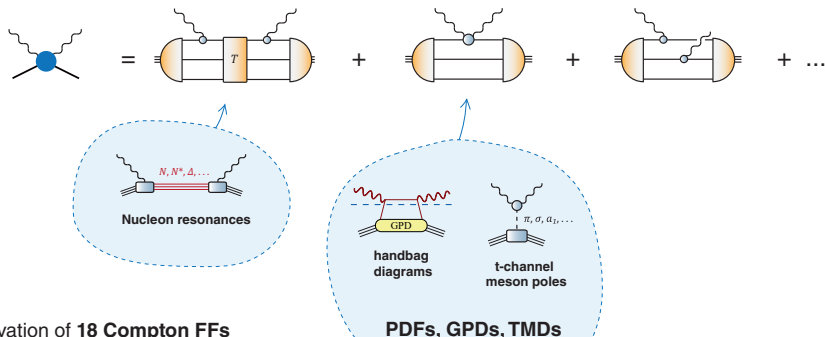
- Solution of **4d scattering equation**,
 contour deformations to go beyond thresholds
[GE, Duarte, Peña, Stadler, PRD 100 \(2019\)](#)



Could be extended to coupled channels...

Hadron structure

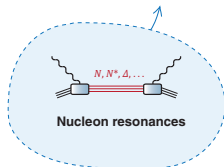
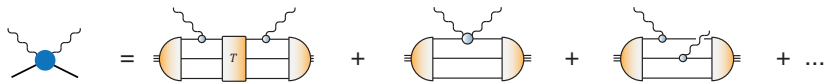
- Systematic derivation how external currents couple to hadrons in DSE/BSE approach, e.g. **nucleon Compton scattering** [GE, Fischer, PRD 85 \(2012\), PRD 87 \(2013\)](#)



- Derivation of **18 Compton FFs** free of kinematic constraints \rightarrow polarizabilities, DVCS, RCS, ... [GE, Ramalho, PRD 98 \(2018\)](#)

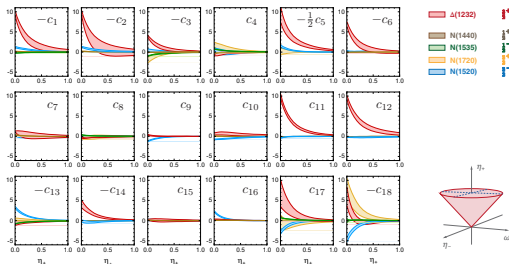
Hadron structure

- Systematic derivation how external currents couple to hadrons in DSE/BSE approach, e.g. **nucleon Compton scattering** [GE, Fischer, PRD 85 \(2012\)](#), [PRD 87 \(2013\)](#)



- Derivation of **18 Compton FFs** free of kinematic constraints \rightarrow polarizabilities, DVCS, RCS, ... [GE, Ramalho, PRD 98 \(2018\)](#)

Nucleon resonance contributions:



Thank you!

INT Program INT-20-2c

Accessing and Understanding the QCD Spectra

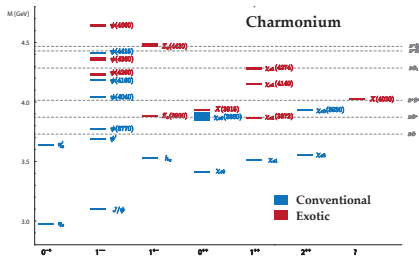
August 17 - September 4, 2020

Organizers:

Raul Briceño (Old Dominion & Jefferson Lab)

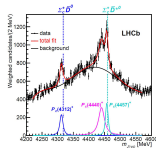
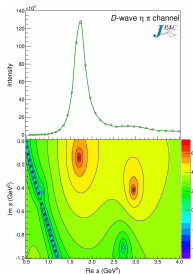
Gernot Eichmann (LIP & IST Lisboa)

Alessandro Pilloni (ECT* Trento)



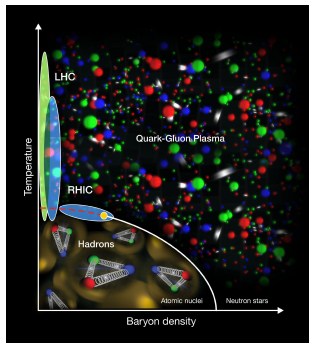
$$|n \dots |n\rangle_{\text{QCD}} = c_0 \text{ (diagram)} + c_1 \text{ (diagram)} + c_2 \text{ (diagram)} + c_3 \text{ (diagram)} + c_4 \text{ (diagram)} + \dots$$

... perhaps there is a hierarchy [e.g. $c_4 \gg c_3 \gg c_2 \gg c_1 \gg c_0$]

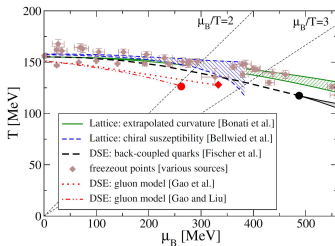


Backup slides

QCD phase diagram



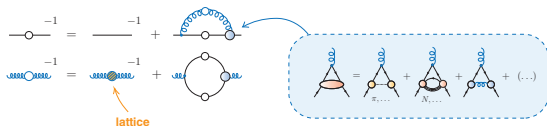
Search for **critical endpoint (CEP)** from DSEs & lattice:



Fischer, Prog. Part. Nucl. Phys. 105 (2019)

Location of CEP sensitive to **baryons**?

GE, Fischer, Welzbacher, PRD 93 (2016)



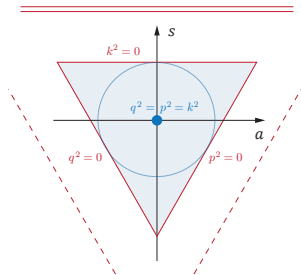
Structure of the amplitude

- **Singlet:** symmetric variable, carries overall scale:

$$\mathcal{S}_0 = \frac{1}{4} (p^2 + q^2 + k^2)$$

- **Doublet:** $\mathcal{D}_0 = \frac{1}{4\mathcal{S}_0} \begin{bmatrix} \sqrt{3}(q^2 - p^2) \\ p^2 + q^2 - 2k^2 \end{bmatrix}$

Mandelstam triangle,
outside: **meson and diquark poles!**

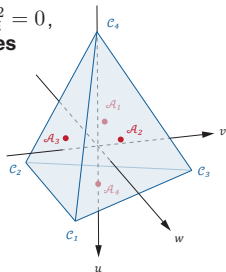


Lorentz invariants can be grouped into **multiplets of the permutation group S4**:

GE, Fischer, Heupel, PRD 92 (2015)

- **Triplet:** $\mathcal{T}_0 = \frac{1}{4\mathcal{S}_0} \begin{bmatrix} 2(\omega_1 + \omega_2 + \omega_3) \\ \sqrt{2}(\omega_1 + \omega_2 - 2\omega_3) \\ \sqrt{6}(\omega_2 - \omega_1) \end{bmatrix}$

tetrahedron bounded by $p_i^2 = 0$,
outside: **quark singularities**

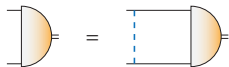


- **Second triplet:** 3dim. sphere

$$\mathcal{T}_1 = \frac{1}{4\mathcal{S}_0} \begin{bmatrix} 2(\eta_1 + \eta_2 + \eta_3) \\ \sqrt{2}(\eta_1 + \eta_2 - 2\eta_3) \\ \sqrt{6}(\eta_2 - \eta_1) \end{bmatrix}$$

Bound states & resonances

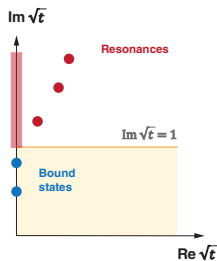
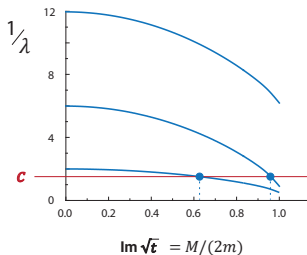
- Homogeneous BSE:



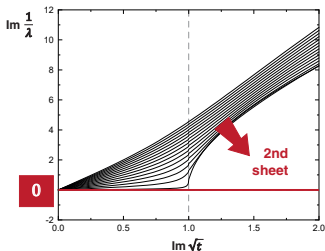
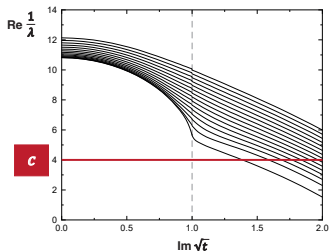
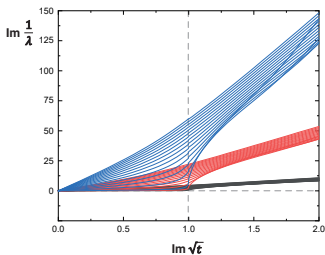
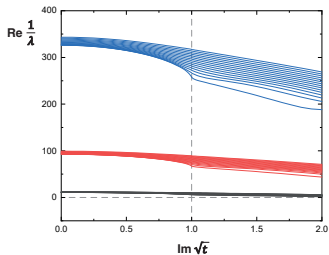
$$\Rightarrow \psi(t) = \mathbf{c} K G_o(t) \psi(t)$$

$$\psi(X, Z, t) = \mathbf{c} \int dx \int dz K(X, x, Z, z, t) G_o(x, z, t) \psi(x, z, t)$$

$$\Rightarrow \frac{1}{\lambda(t)} \stackrel{!}{=} \mathbf{c}$$



BSE Eigenvalues

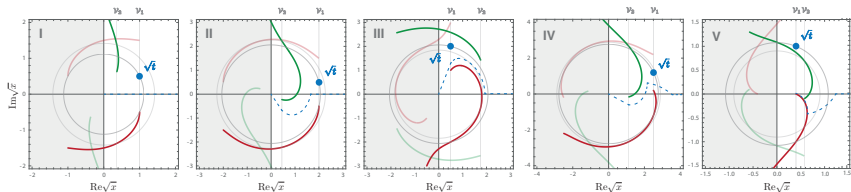


$$\frac{1}{\lambda(t)} = c + 0 \cdot i$$

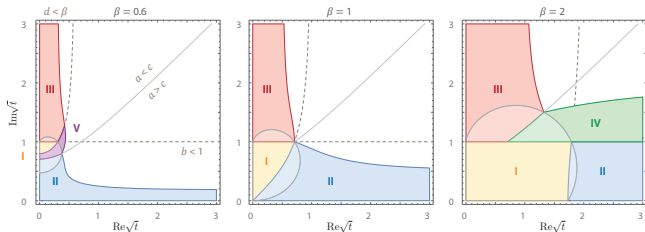
still valid for
complex poles:
 can detect
 resonances from
homogeneous BSE

Contour deformation

For onshell scattering amplitude more complicated:



Can still cover **parts** of complex t plane:



Benchmarks

• Binding energies

$$c = 1, \beta = 0.5$$

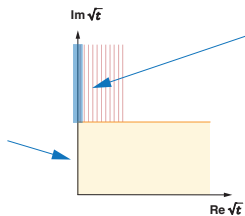
$\text{Im}\sqrt{t}$	π/λ_0 this work	π/λ_0 [1, 2]	π/λ_0 [3]
0.999	1.18(3)	1.211	1.216
0.995	1.43(1)	1.440	1.440
0.99	1.623	1.624	1.623
0.95	2.498	2.498	2.498
0.90	3.251	3.251	3.251
0.80	4.416	4.416	4.416
0.75	4.901	4.901	4.901
0.6	6.094	6.096	6.094
0.4	7.205	7.206	7.204
0.2	7.849	7.850	7.849
0	8.061	8.062	8.061

[1, 2] Kusaka, Simpson, Williams, PRD 56 (1997)
Karmanov, Carbonell, EPJ A 28 (2006)

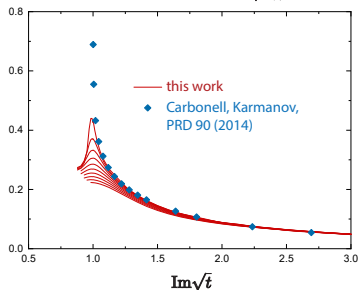
[3] Frederico, Salmè, Viviani, PRD 89 (2014)

• Phase shifts

$$f_i(t) = \frac{1}{2i\tau(t)} \left[e^{2i\delta_i(t)} - 1 \right]$$

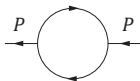


$\text{Re } \delta_0(t)/\pi$ $c = 1.2/\pi, \beta = 0.5$



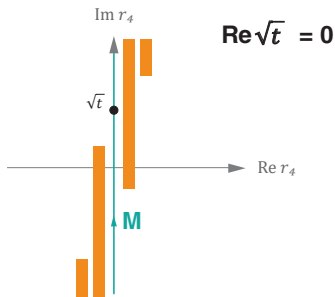
Two poles

Consider two-point function (current correlator, self energy, vacuum polarization, ...)



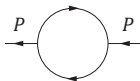
$$\int d^4k \frac{1}{k_+^2 + m^2} \frac{1}{k_-^2 + m^2}$$

$$\int_{-\infty(i-\epsilon)}^{\infty(i-\epsilon)} dr_4 \dots$$



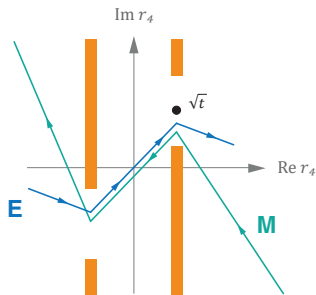
Two poles

Consider two-point function (current correlator, self energy, vacuum polarization, ...)



$$\int d^4k \frac{1}{k_+^2 + m^2} \frac{1}{k_-^2 + m^2}$$

$$\int_{-\infty(i-\epsilon)}^{\infty(i-\epsilon)} dr_4 \dots$$



So:

E = M

$\int d^3\mathbf{k} \int_{-\infty}^{\infty} dk_4$... close contours analytically, pick up **residues**

$\int_{-\infty}^{\infty} dk_4 \int d^3\mathbf{k}$... avoid cuts by numerical **contour deformation**

Suggestions for better wording:

"We need XY ~~in Minkowski space~~"

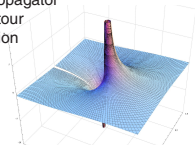
"We calculate XY ~~directly in Minkowski space~~"

... *in the full kinematical domain*

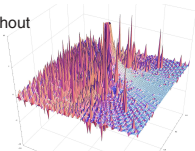
... *above threshold*

... *using residue calculus*

Quark propagator
with contour
deformation



... and without



Complex eigenvalues?

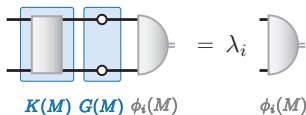
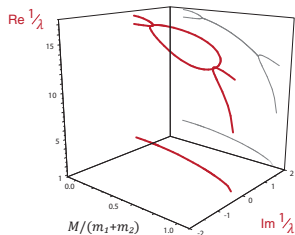
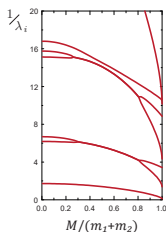
Excited states: some EVs are complex conjugate?

Typical for **unequal-mass** systems, already in Wick-Cutkosky model

Wick 1954, Cutkosky 1954

Connection with “**anomalous**” states?

Ahlig, Alkofer, Ann. Phys. 275 (1999)



If $G = G^\dagger$ and $G > 0$:

Cholesky decomposition $G = L^\dagger L$

$$K L^\dagger L \phi_i = \lambda_i \phi_i$$

$$(L K L^\dagger) (L \phi_i) = \lambda_i (L \phi_i)$$

\Rightarrow Hermitian problem with same EVs!

K and G are Hermitian (even for unequal masses!) but KG is not

Complex eigenvalues?

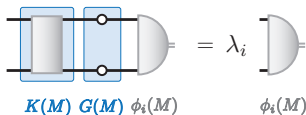
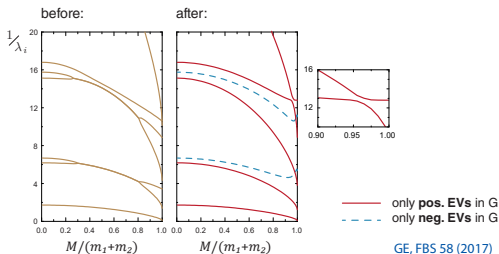
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\Rightarrow Hermitian problem with same EVs!

\Rightarrow all EVs strictly **real**

\Rightarrow level repulsion

\Rightarrow “anomalous states” removed?

Complex eigenvalues?

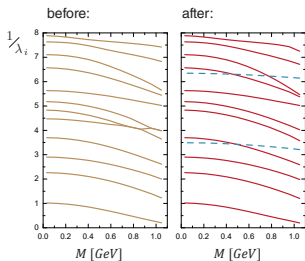
Excited states: some EVs are complex conjugate?

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Wick 1954, Cutkosky 1954

Connection with **“anomalous” states?**

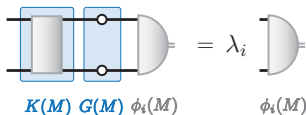
Ahlig, Alkofer, Ann. Phys. 275 (1999)



Eigenvalue spectrum for pion channel

GE, FBS 58 (2017)

— only pos. EVs in G
 - - - only neg. EVs in G



If $G = G^\dagger$ and $G > 0$:

Cholesky decomposition $G = L^\dagger L$

$$K L^\dagger L \phi_i = \lambda_i \phi_i$$

\Rightarrow Hermitian problem with same EVs!

$$(L K L^\dagger) (L \phi_i) = \lambda_i (L \phi_i)$$

K and G are Hermitian (even for unequal masses!) but KG is not

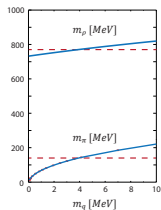
\Rightarrow all EVs strictly **real**

\Rightarrow level repulsion

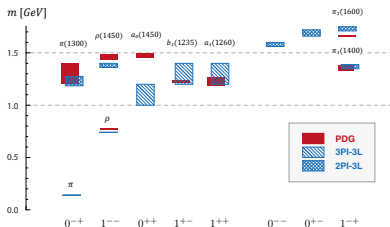
\Rightarrow “anomalous states” removed?

Mesons

- Pion is **Goldstone boson**: $m_\pi^2 \sim m_q$



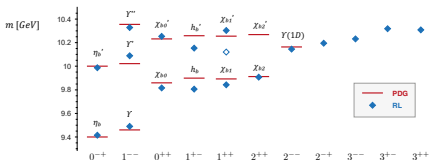
- Light meson spectrum** beyond rainbow-ladder



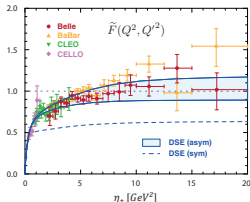
Williams, Fischer, Heupel, PRD 93 (2016)
 GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016)

- Bottomonium spectrum**

Fischer, Kubrak, Williams, EPJ A 51 (2015)



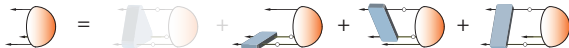
- Pion transition form factor**



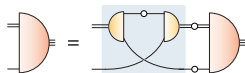
GE, Fischer, Weil, Williams, PLB 774 (2017)

Bound-state equations

Bethe-Salpeter equation for baryons: [GE, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 \(2016\), 1606.09602](#)



Quark-diquark approximation:

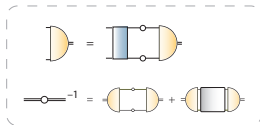


Rainbow-ladder:

Diagram illustrating the rainbow-ladder approximation. It shows a ghost loop (represented by a circle with a line) equal to a ghost loop (represented by a circle with a line) plus a ghost loop with a quark line (represented by a circle with a line and a quark line).

Diagram illustrating the rainbow-ladder approximation. It shows a ghost loop (represented by a circle with a line) equal to a ghost loop with a quark line (represented by a circle with a line and a quark line).

Maris, Tandy, PRC 60 (1999),
Qin et al., PRC 84 (2011)

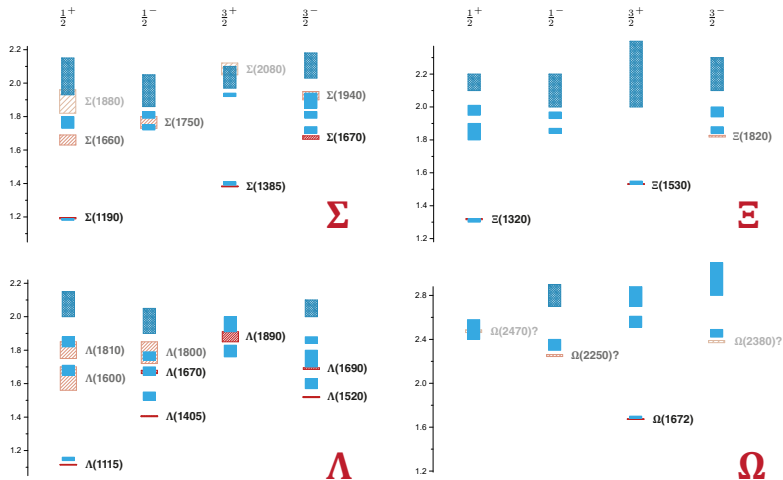


DSE / BSE / Faddeev landscape



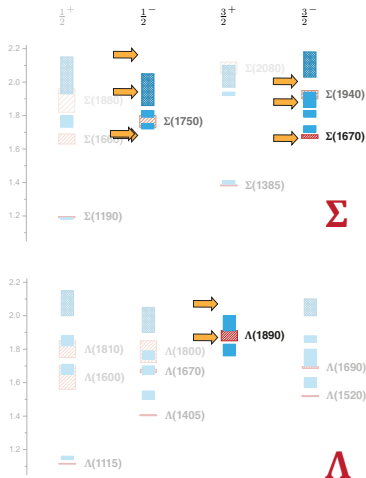
		NJL / contact	q-dq model	DSE (RL)		DSE (bRL)
u/d	N, Δ masses	✓	✓	✓	✓	✓
	N, Δ em. FFs	✓	✓	✓	✓	
	$N \rightarrow \Delta\gamma$	✓	✓	✓	✓	
	N^*, Δ^* masses (+)	✓	✓	✓	✓	
	$N \rightarrow N^*\gamma$	✓	✓			
	N^*, Δ^* masses (-)	✓	✓	✓	✓	
	$N \rightarrow N^*\gamma$					
s	ground states	✓	✓	✓	✓	✓... before 2015 ✓... after 2015
	excited states	✓	✓	✓	✓	
	em. FFs & TFFs				✓	
c, b	ground states	✓	✓		✓	
	excited states		✓		✓	
		Cloet, Thomas, Roberts, Bashir, Segovia, Chen, Wilson, Lu, ...	Oettel, Alkofer, Roberts, Cloet, Segovia, Chen, El-Bennich, ...	GE, Alkofer, Nicmoros, Sanchis-Alepuz, Fischer	GE, Sanchis-Alepuz, Fischer, Alkofer, Qin, Roberts	Sanchis-Alepuz, Williams, Fischer

Strange baryons



GE, Fischer, FBS 60 (2019), Fischer, GE, PoS Hadron 2017

Strange baryons



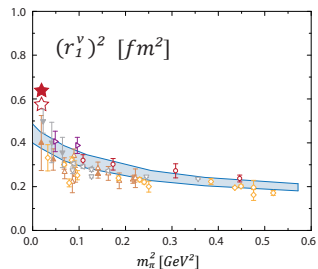
New states from Bonn-Gatchina
[Sarantsev et al., 1907.13387 \[nucl-ex\]](#)

GE, Fischer, FBS 60 (2019), Fischer, GE, PoS Hadron 2017

Nucleon em. form factors

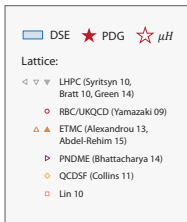
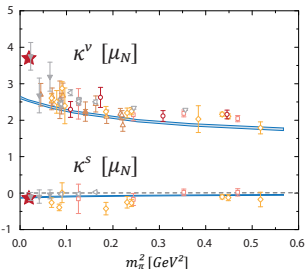
Nucleon charge radii:

isovector (p-n) Dirac (F1) radius

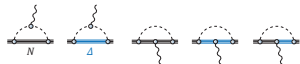


Nucleon magnetic moments:

isovector (p-n), isoscalar (p+n)



- **Pion-cloud effects** missing (\Rightarrow divergence!), agreement with lattice at larger quark masses.



- **But:** pion-cloud cancels in $\kappa^s \Leftrightarrow$ quark core

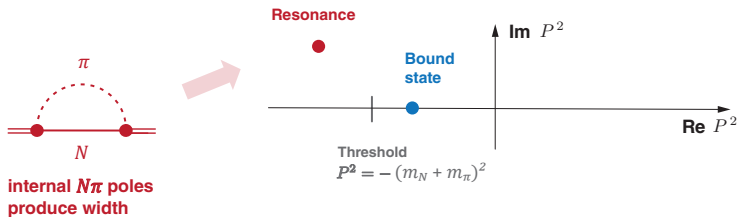
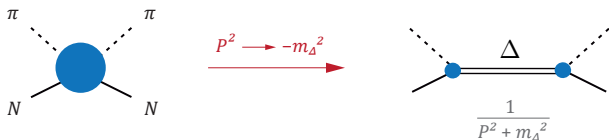
Exp: $\kappa^s = -0.12$

Calc: $\kappa^s = -0.12(1)$



GE, PRD 84 (2011)

Resonances?

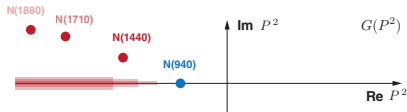
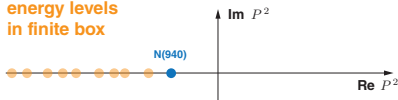


Resonances?

Lattice QCD:

$$\langle \dots \rangle = \int \mathcal{D}[\psi, \bar{\psi}, A] e^{-S[\psi, \bar{\psi}, A]} (\dots)$$

energy levels
in finite box



- **Finite volume:**
bound states & scattering states



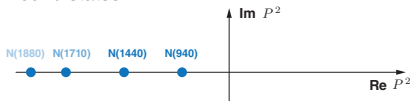
vary volume,
Luescher method

- **Infinite volume:**
Bound states, resonances,
branch cuts

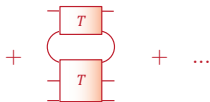
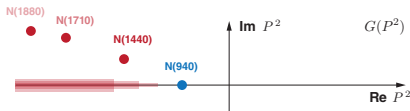
Resonances?

In terms of quarks and gluons?

Bound states:



Resonances by meson-baryon interactions:



Both **bound states** and **resonances** must be generated from quark-gluon structure!



Analogue for $\rho \rightarrow \pi\pi$:

[Williams, 1804.11161 \[hep-ph\]](#),
[Miramontes, Sanchis-Alepuz, 1906.06227 \[hep-ph\]](#)