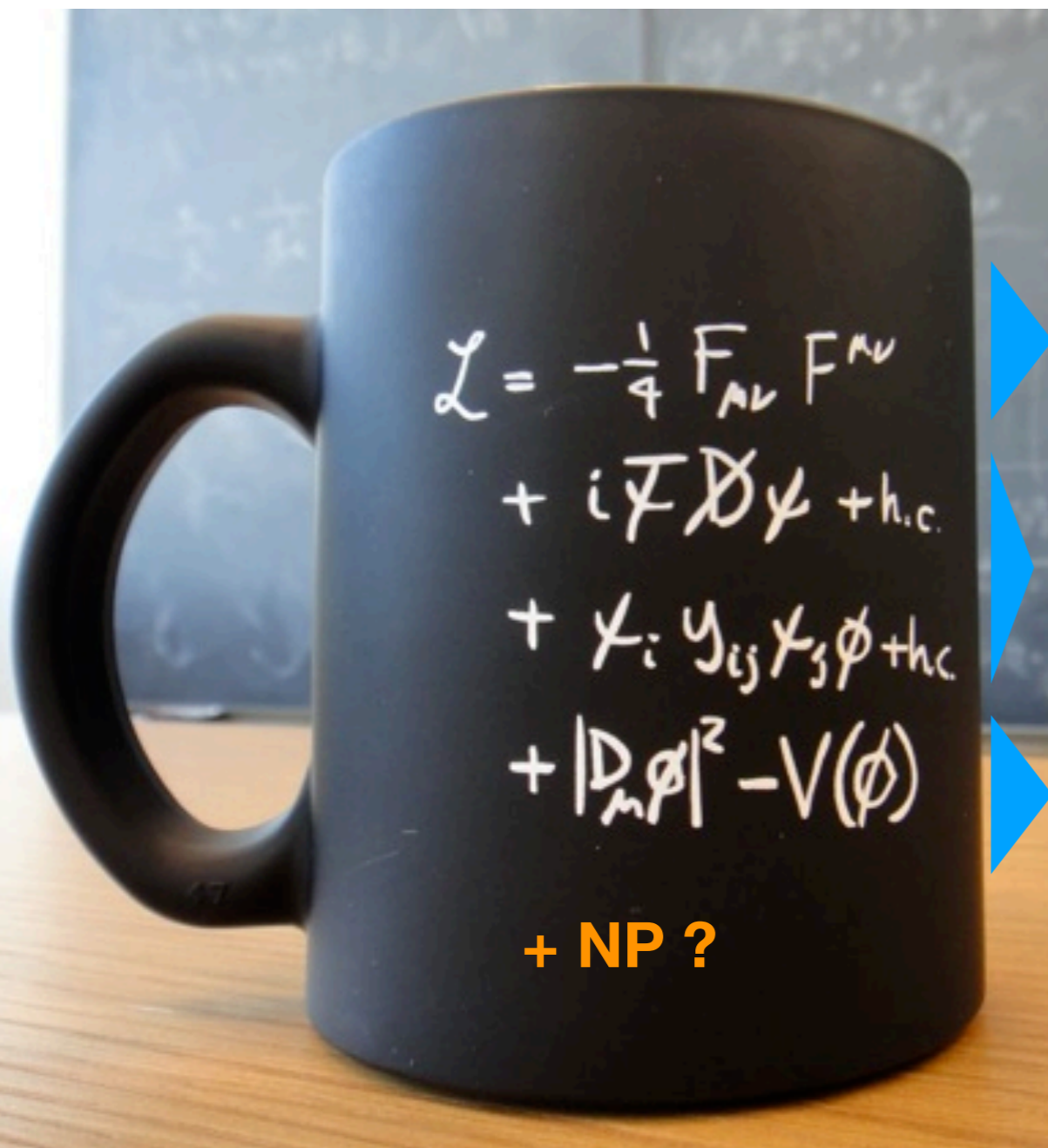


# the standard model (of particle physics)

## The SM Lagrangian



**Gauge sector**  
(spin 1)

**Flavor sector**  
Fermion (spin 1/2)  
dynamics & mass

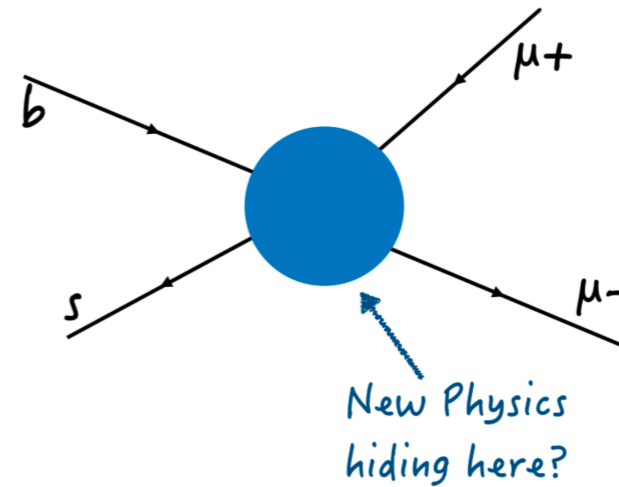
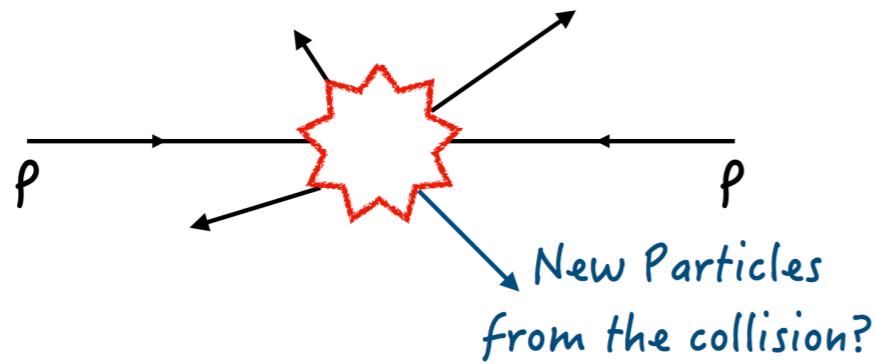
**Higgs sector**  
(spin 0)

describes everything  
experimentally confirmed  
before 2012  
Yukawa coupling w/ scalar  
(new interaction type)  
Scalar self-interaction  
and with gauge bosons

$$\begin{aligned} \mathcal{L}_{SM} = & -\frac{1}{2}\partial_\nu g_\mu^a \partial_\nu g_\mu^a - g_s f^{abc} \partial_\mu g_\nu^a g_\mu^b g_\nu^c - \frac{1}{4}g_s^2 f^{abc} f^{ade} g_\mu^b g_\nu^c g_\mu^d g_\nu^e - \partial_\nu W_\mu^+ \partial_\nu W_\mu^- - \\ & M^2 W_\mu^+ W_\mu^- - \frac{1}{2}\partial_\nu Z_\mu^0 \partial_\nu Z_\mu^0 - \frac{1}{2c_w^2} M^2 Z_\mu^0 Z_\mu^0 - \frac{1}{2}\partial_\mu A_\nu \partial_\mu A_\nu - igc_w (\partial_\nu Z_\mu^0 (W_\mu^+ W_\nu^- - \\ & W_\nu^+ W_\mu^-) - Z_\nu^0 (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + Z_\nu^0 (W_\nu^+ \partial_\mu W_\mu^- - W_\mu^- \partial_\nu W_\nu^+)) - \\ & ig_s (\partial_\mu A_\nu (W_\mu^+ W_\nu^- - W_\nu^+ W_\mu^-) - A_\nu (W_\mu^+ \partial_\nu W_\mu^- - W_\mu^- \partial_\nu W_\mu^+) + A_\nu (W_\nu^+ \partial_\mu W_\mu^- - \\ & W_\mu^- \partial_\nu W_\nu^+) + \frac{1}{2}g^2 W_\mu^+ W_\nu^- W_\mu^+ W_\nu^- + g^2 c_w^2 (Z_\mu^0 W_\mu^+ Z_\nu^0 W_\nu^- - \\ & (A_\mu W_\mu^+ A_\nu W_\nu^- - A_\mu A_\nu W_\mu^+ W_\nu^-) + g^2 s_w c_w (A_\mu Z_\nu^0 (W_\mu^+ W_\nu^- - \\ & W_\nu^+ W_\mu^-) - \frac{1}{2}\partial_\mu H \partial_\mu H - 2M^2 \alpha_h H^2 - \partial_\mu \phi^+ \partial_\mu \phi^- - \frac{1}{2}\partial_\mu \phi^0 \partial_\mu \phi^0 - \\ & + \frac{2M}{g} H + \frac{1}{2}(H^2 + \phi^0 \phi^0 + 2\phi^+ \phi^-) + \frac{2M^4}{g^2} \alpha_h - \\ & g\alpha_h M (H^3 + H\phi^0 \phi^0 + 2H\phi^+ \phi^-) - \\ & + 4(\phi^+ \phi^-)^2 + 4(\phi^0)^2 \phi^+ \phi^- + 4H^2 \phi^+ \phi^- + 2(\phi^0)^2 H^2) - \\ & gMW_\mu^+ W_\mu^- H - \frac{1}{2}g\frac{M}{c_w} Z_\mu^0 Z_\mu^0 H - \\ & \partial_\mu \phi^+ \partial_\mu \phi^- - \phi^+ \partial_\mu \phi^0 - W_\mu^- (\phi^0 \partial_\mu \phi^+ - \phi^+ \partial_\mu \phi^0) + \\ & H) + W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H) + \frac{1}{2}g\frac{1}{c_w} (Z_\mu^0 (H \partial_\mu \phi^0 - \phi^0 \partial_\mu H) + \\ & + W_\mu^+ (H \partial_\mu \phi^- - \phi^- \partial_\mu H) + W_\mu^- (H \partial_\mu \phi^+ - \phi^+ \partial_\mu H)) + \\ & Z_\mu^0 (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) - ig_s A_\mu (\psi_i^\dagger \gamma^\mu \psi_j + \psi_j^\dagger \gamma^\mu \psi_i) + \\ & + 2\phi^+ \phi^-) - \frac{1}{8}g^2 \frac{s_w}{c_w^2} Z_\mu^0 Z_\mu^0 (H^2 + \phi^0 \phi^0) + \frac{1}{2}(2s_w - 1) \phi^+ \phi^-) - \\ & \phi^+) - \frac{1}{2}ig^2 \frac{s_w}{c_w} (W_\mu^+ \phi^- - W_\mu^- \phi^+) + \frac{1}{2}g^2 s_w A_\mu \phi^0 (W_\mu^+ \phi^- + \\ & A_\mu H) (W_\mu^+ \phi^- - W_\mu^- \phi^+) - g^2 \frac{s_w}{c_w} (2c_w^2 - 1) Z_\mu^0 A_\mu \phi^+ \phi^- - \\ & \psi_i^\dagger (\bar{\sigma}^\mu \gamma^\mu q_j^\sigma) g_\mu^a - \bar{e}^\lambda (\gamma^\mu \partial + m_e^\lambda) e^\lambda - \bar{\nu}^\lambda (\gamma^\mu \partial + m_\nu^\lambda) \nu^\lambda - \bar{u}_j^\lambda (\gamma^\mu \partial + \\ & \lambda) d_j^\lambda + ig_s s_w (\bar{u}_j^\lambda \gamma^\mu q_j^\sigma) g_\mu^a + \bar{e}^\lambda (\gamma^\mu \partial + m_e^\lambda) e^\lambda + (d_j^\lambda \gamma^\mu (\frac{1}{2}s_w^2 - 1 - \gamma^5) d_j^\lambda) + \\ & + \frac{ig}{2} W_\mu^+ (\bar{u}_j^\lambda \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (d_j^\lambda \gamma^\mu (\frac{1}{2}s_w^2 - 1 - \gamma^5) d_j^\lambda) + \\ & U^{lep \dagger} \gamma^\mu (1 + \gamma^5) \nu^\lambda) + (\bar{d}_j^\kappa C_{\kappa\lambda}^\dagger \gamma^\mu (1 + \gamma^5) u_j^\lambda) + \\ & (\bar{\nu}^\lambda U_{\lambda\kappa}^{lep} \gamma^\mu (1 - \gamma^5) e^\kappa) + m_\nu^\lambda (\bar{\nu}^\lambda U_{\lambda\kappa}^{lep} (1 + \gamma^5) e^\kappa) + \\ & \kappa (1 + \gamma^5) \nu^\kappa) - m_\nu^\kappa (e^\kappa U_{\lambda\kappa}^{lep} (1 - \gamma^5) \nu^\kappa) - \frac{g}{2M} H (\bar{\nu}^\lambda \nu^\lambda) - \\ & \frac{1}{2} \phi^0 (\bar{u}_j^\lambda \gamma^5 \nu^\lambda) + \frac{1}{2} \frac{m_d^\lambda}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 \nu^\lambda) + \frac{1}{4} \frac{m_u^\lambda}{M} \phi^0 (\bar{u}_j^\lambda \gamma^5 \nu^\lambda) + \\ & \frac{ig}{2M\sqrt{2}} \phi^+ (-m_d^\kappa (\bar{u}_j^\lambda C_{\lambda\kappa} (1 - \gamma^5) d_j^\kappa) + m_u^\lambda (\bar{u}_j^\lambda C_{\lambda\kappa} (1 + \gamma^5) d_j^\kappa) + \\ & \kappa (1 + \gamma^5) u_j^\kappa) - m_u^\kappa (\bar{d}_j^\lambda C_{\lambda\kappa}^\dagger (1 - \gamma^5) u_j^\kappa) - \frac{g}{2} \frac{m_u^\lambda}{M} H (\bar{u}_j^\lambda u_j^\lambda) - \\ & (\bar{u}_j^\lambda \gamma^5 u_j^\lambda) - \frac{ig}{2} \frac{m_d^\lambda}{M} \phi^0 (\bar{d}_j^\lambda \gamma^5 d_j^\lambda) + \bar{G}^a \partial^2 G^a + g_s f^{abc} \partial_\mu \bar{G}^a G^b g_\mu^c + \\ & \partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \frac{M^2}{c_w^2}) X^0 + \bar{Y} \partial^2 Y + igc_w W_\mu^+ (\partial_\mu \bar{X}^0 X^- - \\ & s_w W_\mu^+ (\partial_\mu \bar{Y} X^- - \partial_\mu \bar{X}^+ Y) + igc_w W_\mu^- (\partial_\mu \bar{X}^- X^0 - \\ & \partial_\mu \bar{X}^- X^-) + igc_w W_\mu^0 (\partial_\mu \bar{X}^+ X^+ - \\ & \partial_\mu \bar{X}^- X^-) + ig_s W_\mu^+ (\partial_\mu \bar{X}^+ X^+ - \\ & \partial_\mu \bar{X}^- X^-) - \frac{1}{2}gM (\bar{X}^+ X^+ H + \bar{X}^- X^- H + \frac{1}{c_w^2} \bar{X}^0 X^0 H) + \frac{1-2c_w^2}{2c_w} igM (\bar{X}^+ X^0 \phi^+ - \bar{X}^- X^0 \phi^-) + \\ & \frac{1}{2c_w} igM (\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-) + igM s_w (\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^-) + \\ & \frac{1}{2}igM (\bar{X}^+ X^+ \phi^0 - \bar{X}^- X^- \phi^0) \end{aligned}$$

A great triumph of 20th century science.

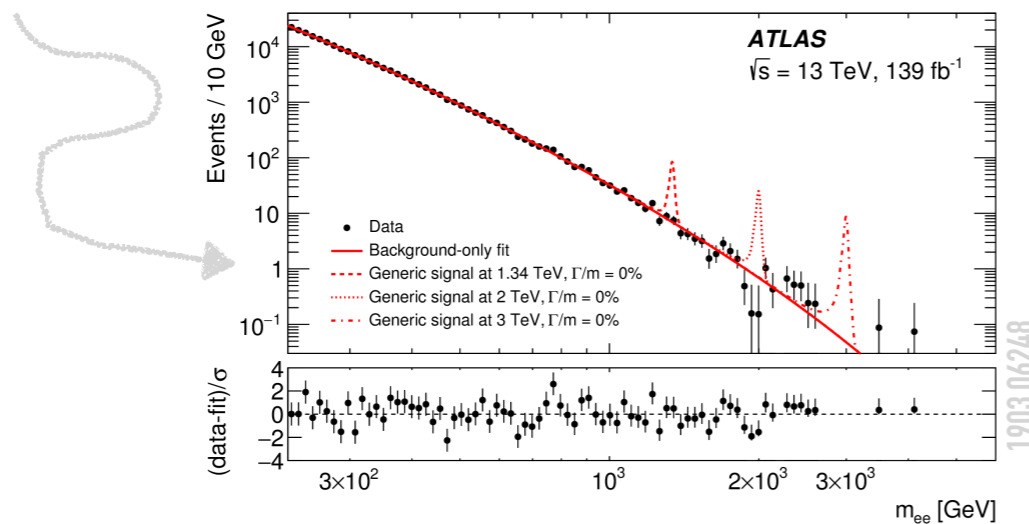
# complementary paths to NP @LHC



## Direct

*bump hunting*

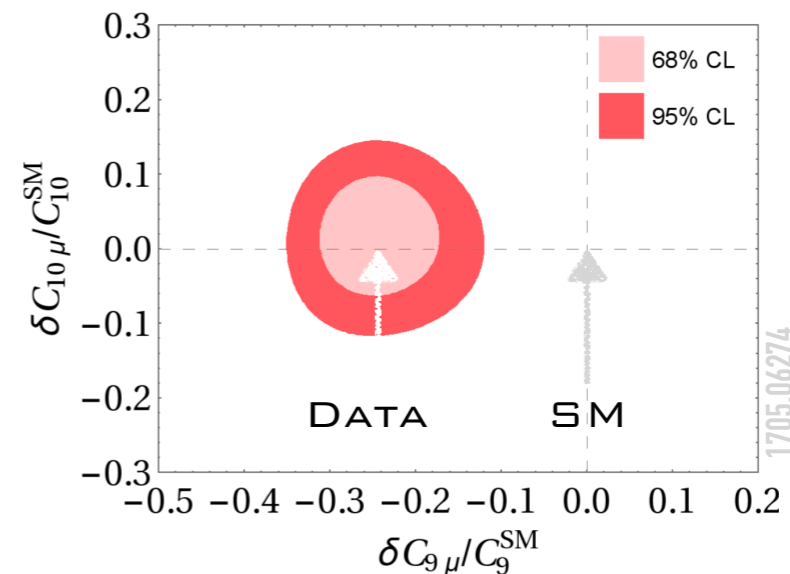
- ▶ searching for the decay products of potentially produced NP particles



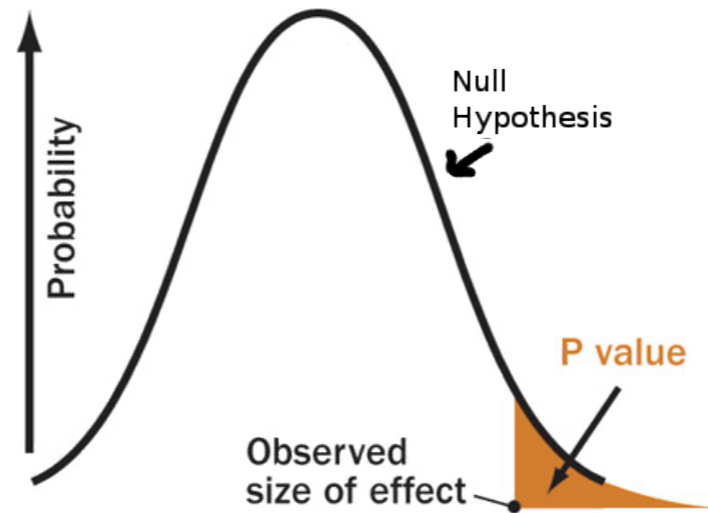
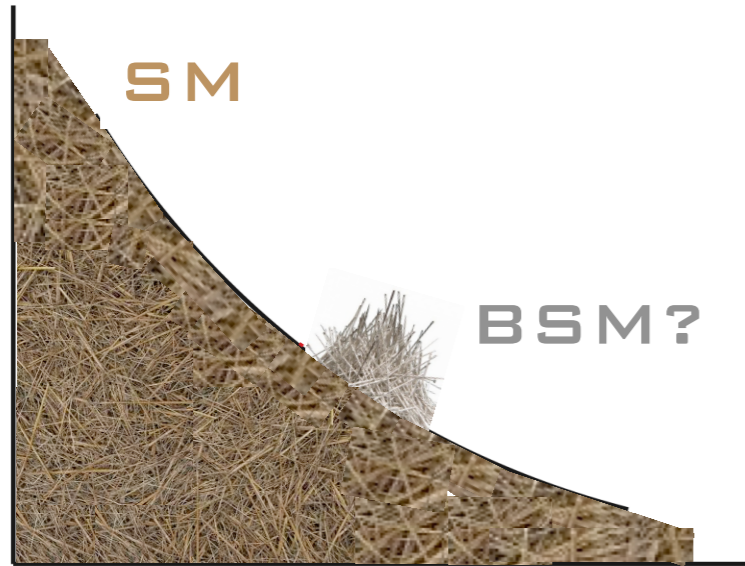
## Indirect

*precision & rareness*

- ▶ searching for NP particles running in quantum loops (virtual)



# Statistics! significance of a bump



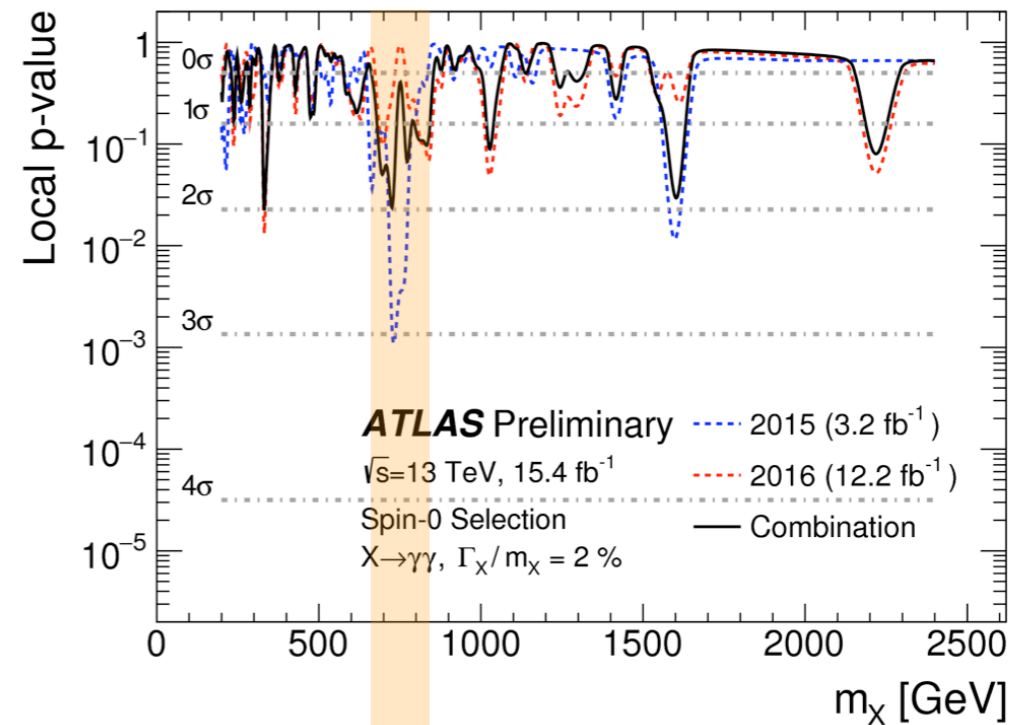
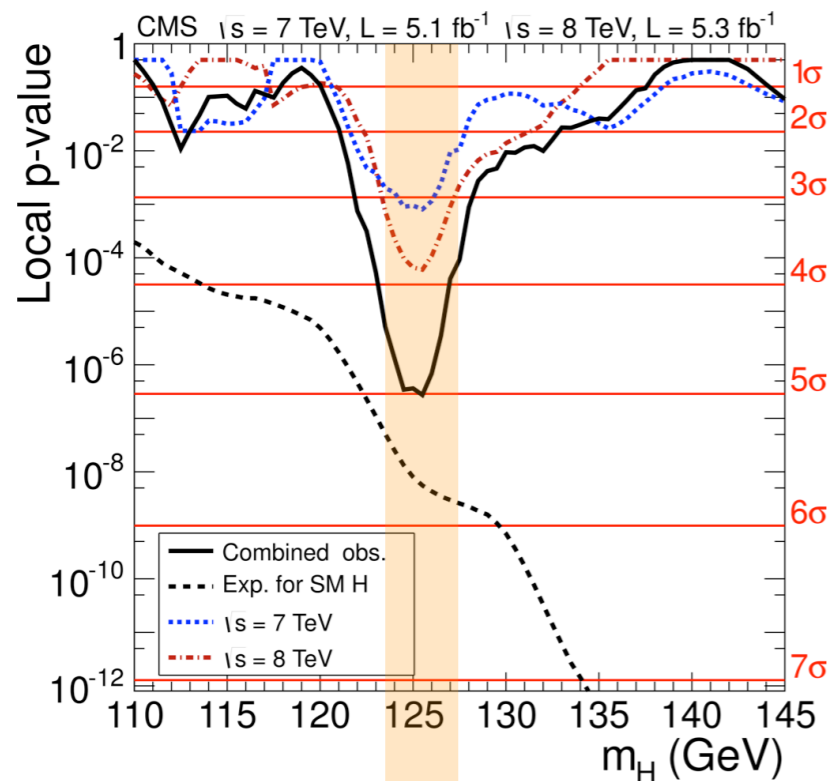
**5 $\sigma$  (discovery)**  
P-VAL: 0.0000003

**3 $\sigma$  (evidence)**  
P-VAL: 0.0013

@125 GeV: the Higgs!

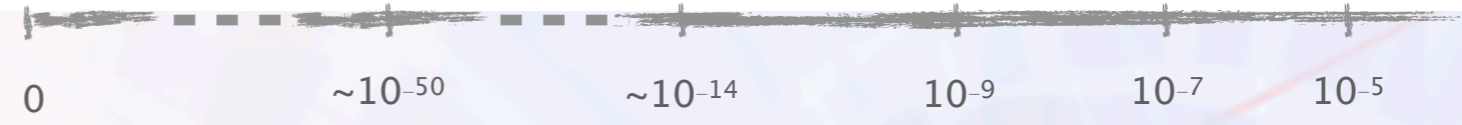
@750 GeV: a fluctuation!

$X \rightarrow \gamma\gamma$



# precise & rare

baryon number violation  
 lepton flavour violation  
 GIM suppressed e.g.  $t \rightarrow c/u$   
 helicity suppressed e.g.  $B \rightarrow \mu\mu$   
 EW penguins e.g.  $b \rightarrow sll$   
 CKM suppressed e.g.  $b \rightarrow u$



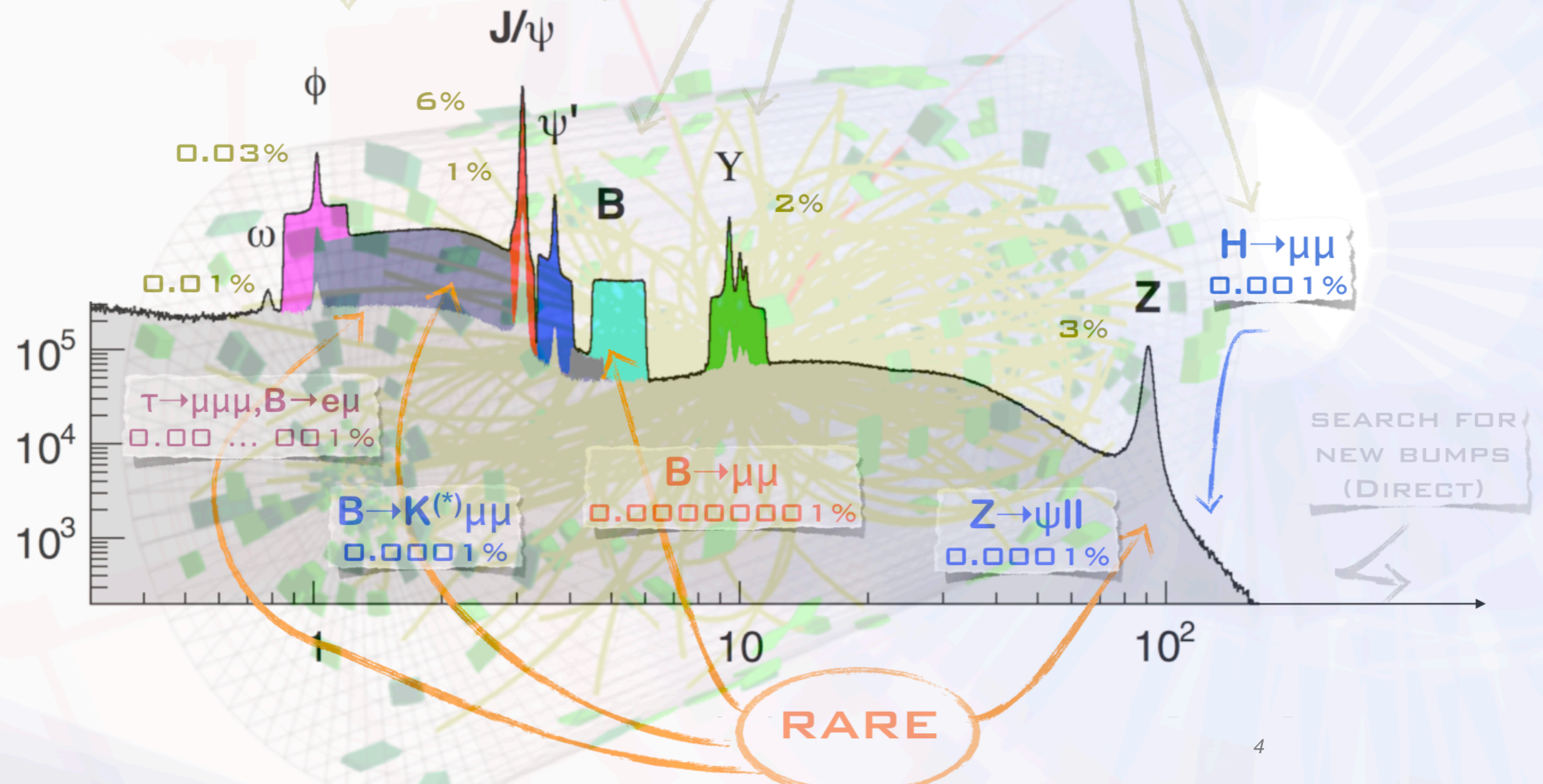
NOT-SO-RARE → PRECISION!

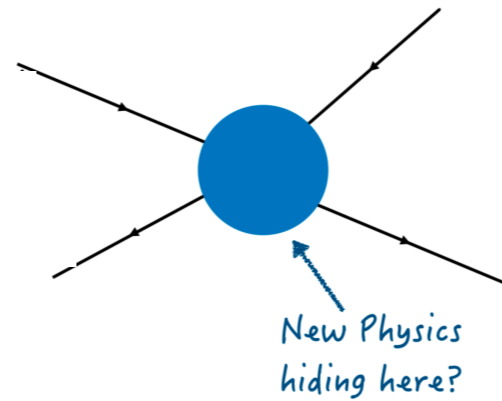
MEDIUM RARE ← EWK PENGUINS

VERY RARE ← FCNC/GIM+HELICITY

ULTRA RARE ← LFV

PRECISION





# Rare Decays of SM particles, towards NP

beauty  
charm  
strangeness  
top  
W,Z  
Higgs  
leptons ( $\tau$ )

SUSY  
Z',W'  
leptoquarks  
unexpected

?

# rare beauty | $B_s \rightarrow \mu\mu$

a milestone discovery of the LHC physics program

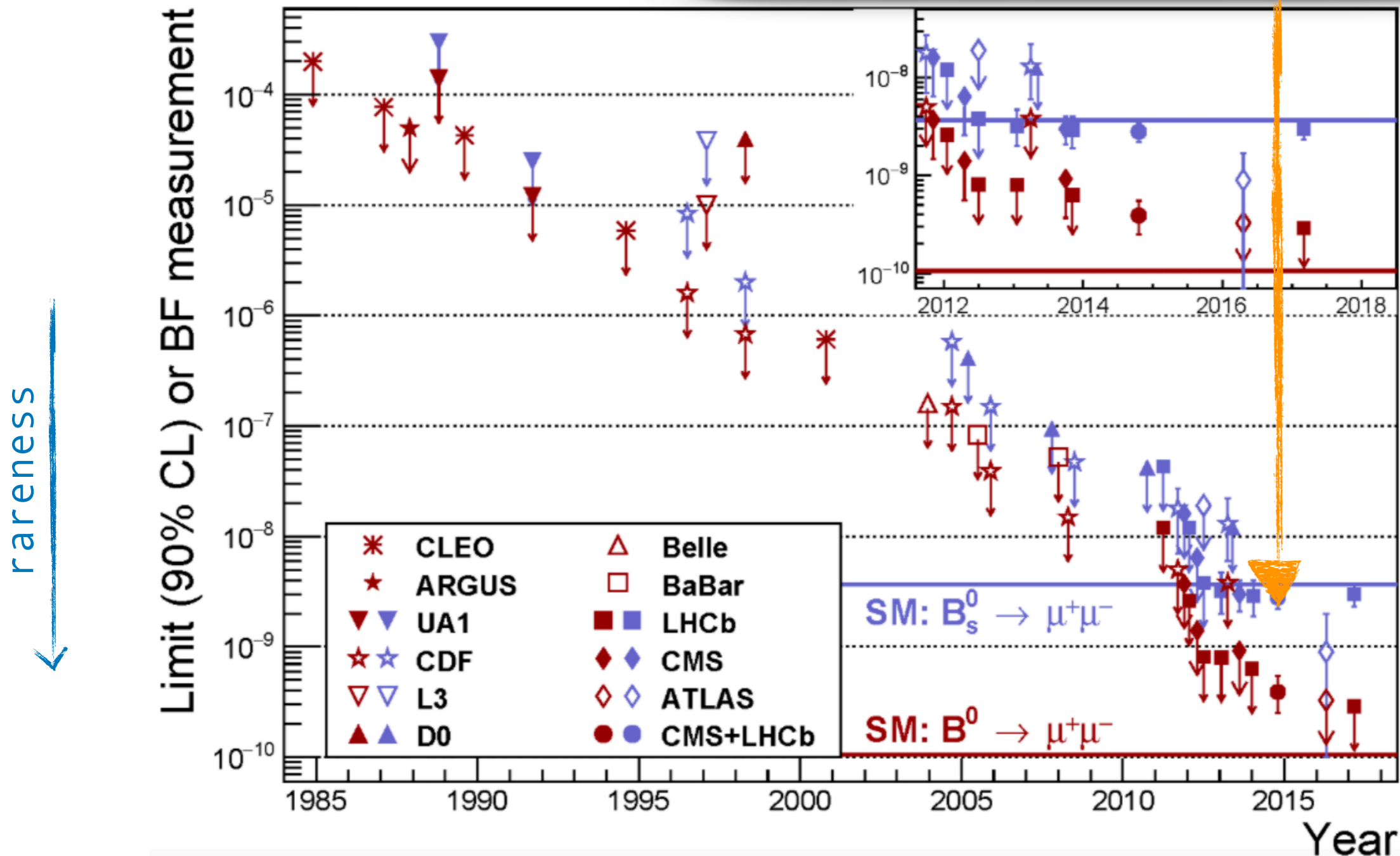
LETTER

OPEN

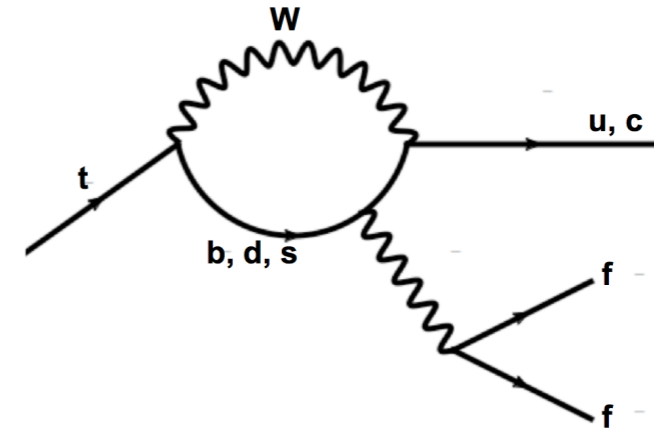
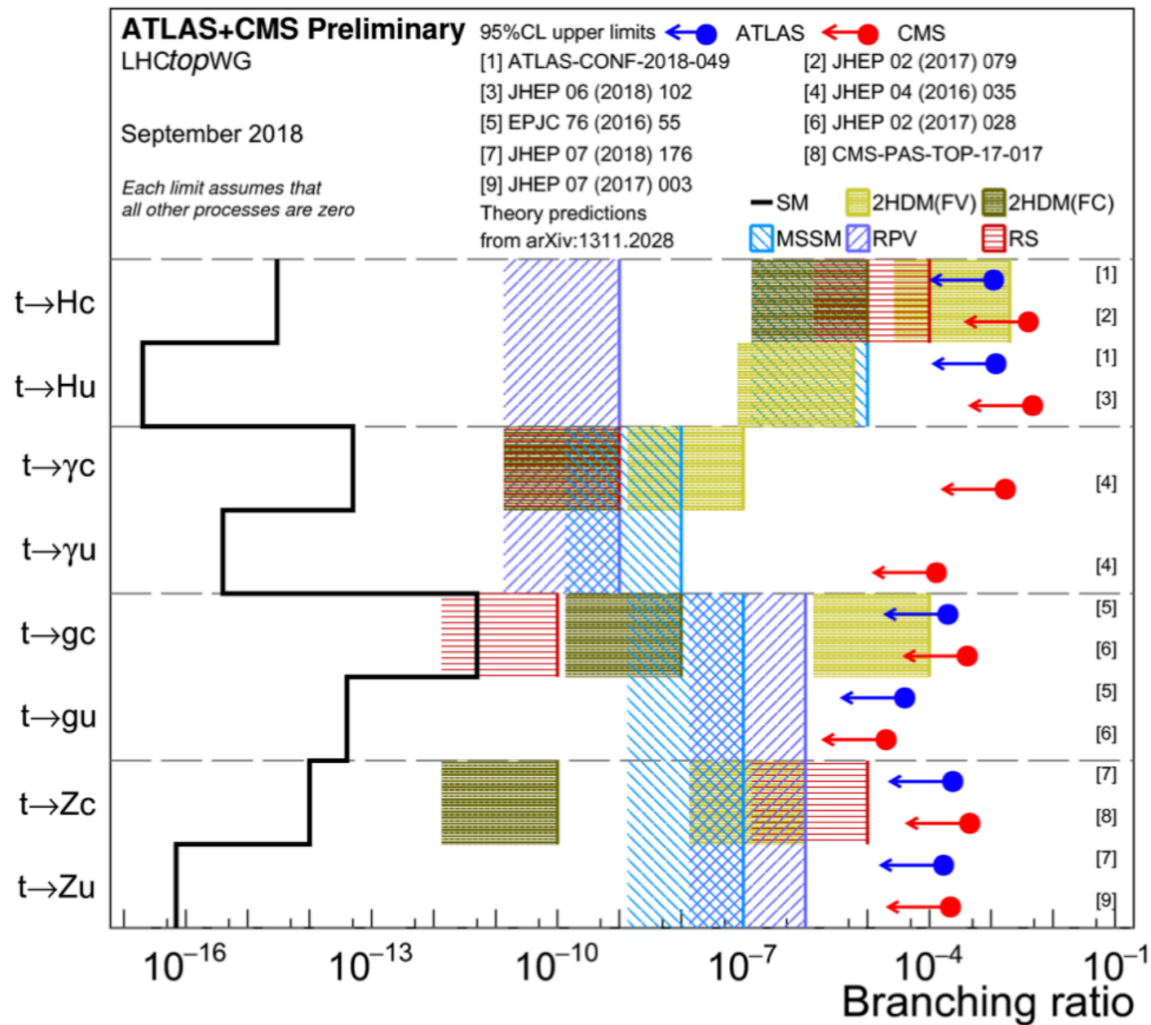
doi:10.1038/nature14474

Observation of the rare  $B_s^0 \rightarrow \mu^+ \mu^-$  decay from the combined analysis of CMS and LHCb data

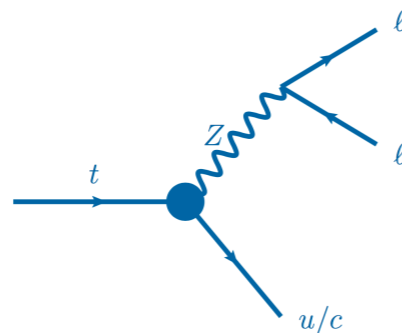
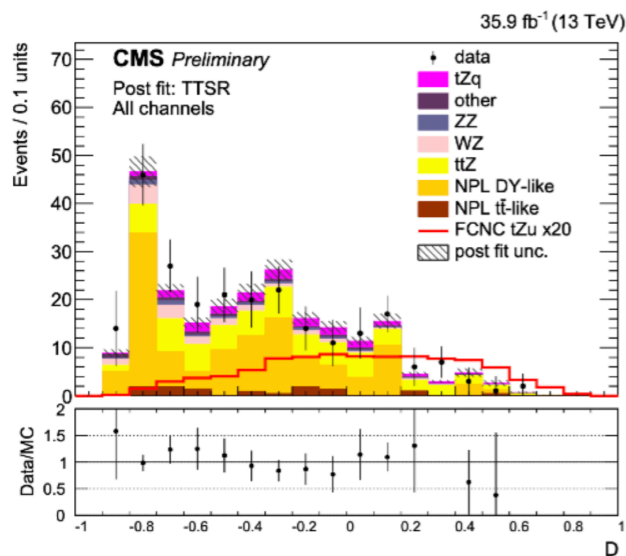
The CMS and LHCb collaborations\*



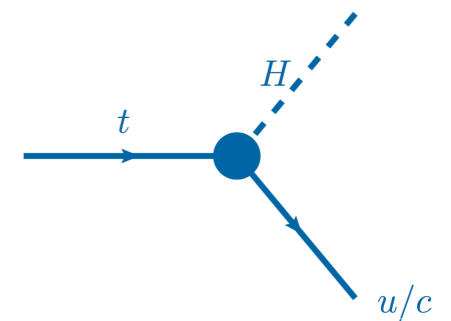
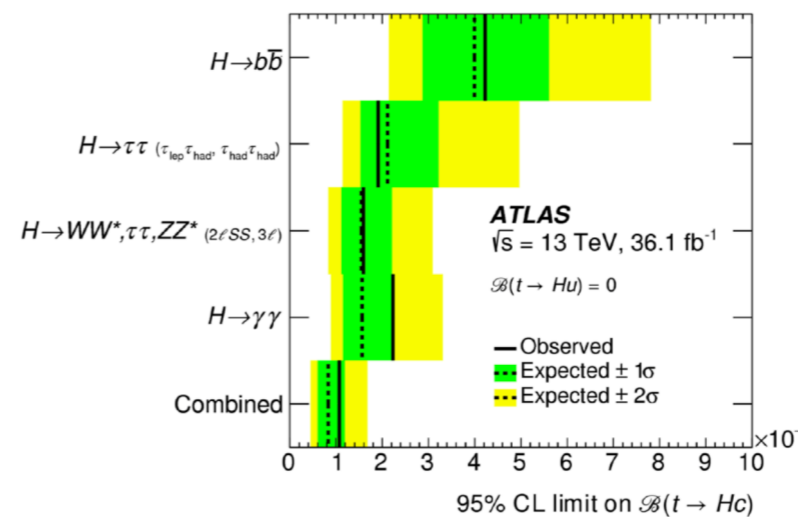
# rare top | $t \rightarrow u/c$



- FCNC/GIM in top sector lead to very rare processes
  - $BF \sim 10^{-14}$
- rates enhanced in NP models
  - MSSM ( $10^{-7}$ ), 2HDM ( $10^{-6}$ ), RS ( $10^{-5}$ )
- current limits  $\sim 10^{-4}$

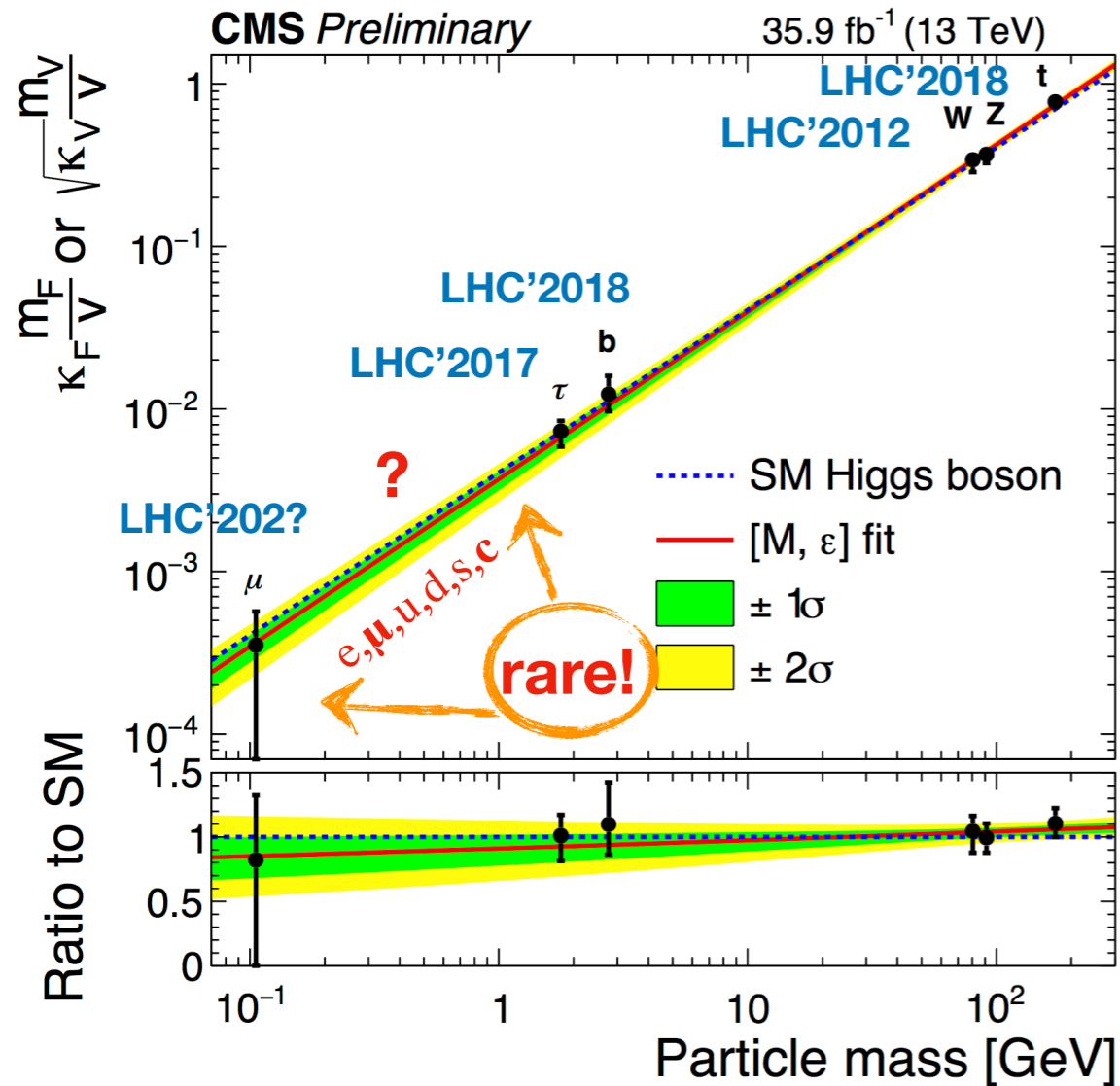


CMS-TOP-17-017

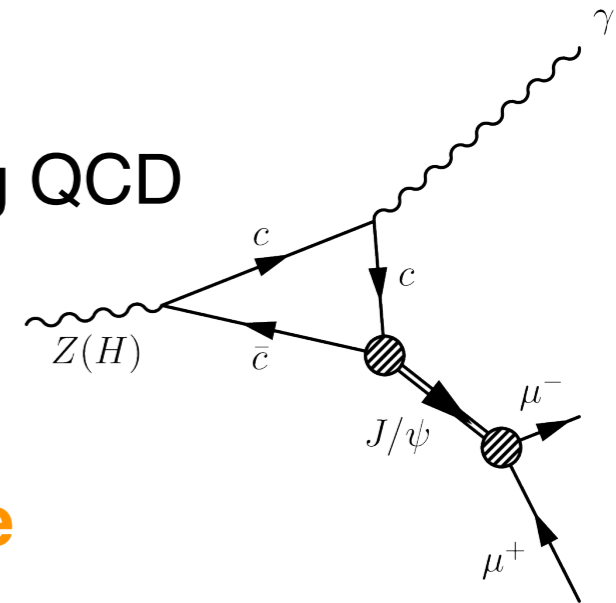


1803.09923

# rare Higgs

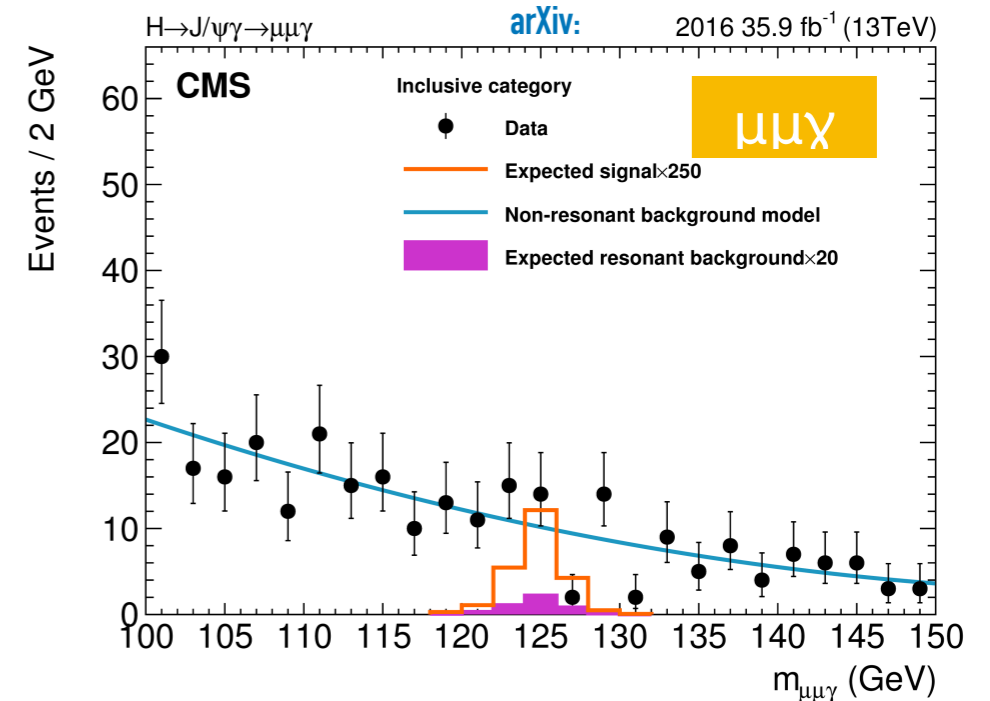


- $H \rightarrow qq$ 
  - overwhelming QCD background
- $H \rightarrow Q\gamma$ 
  - clean but **rare**
  - $H \rightarrow Y/\psi/\phi/\rho + \gamma$



## Higgs couplings:

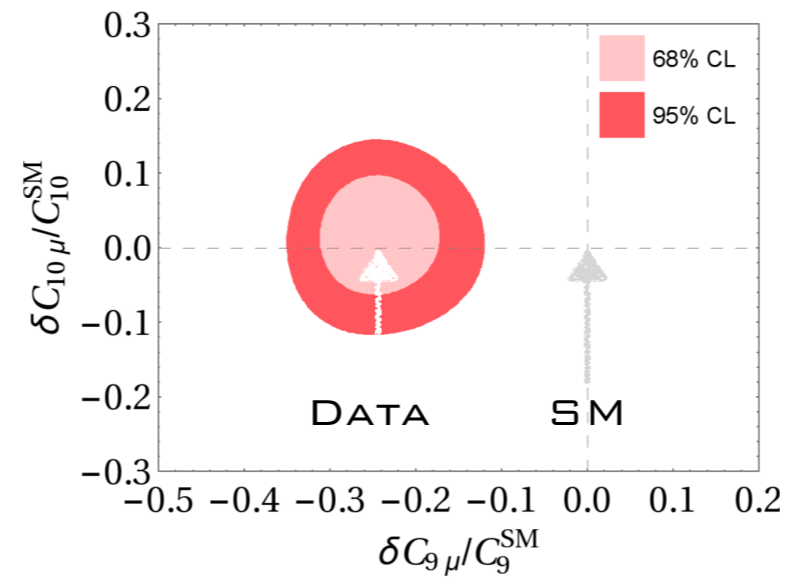
- H to **W,Z,t,b,τ**: done
- H to **γ**: no mass → no coupling
- H to **μ**: clean signature; expect Run2(+Run3)
- H to **c**: challenging, in reach @HL-LHC
- H to **u,d,s,e**: almost hopeless @LHC but NP!



Currently @CMS

$\mu(H \rightarrow cc) < 70$  |  $\mu(H \rightarrow J/\psi\gamma) < 220$





# Lepton Flavour Universality & Flavour Anomalies

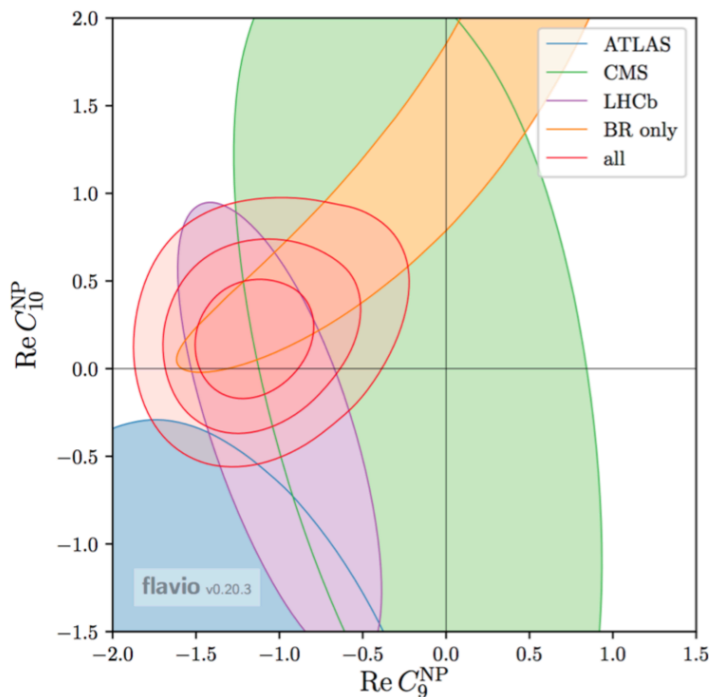
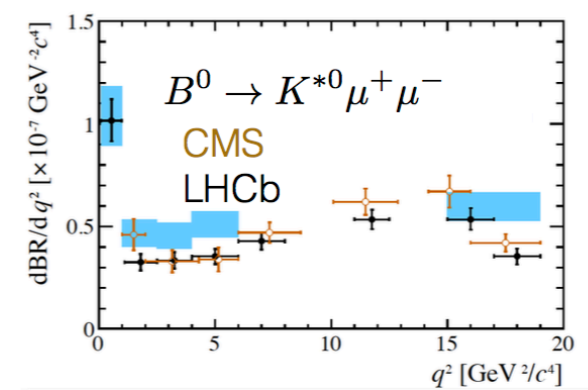
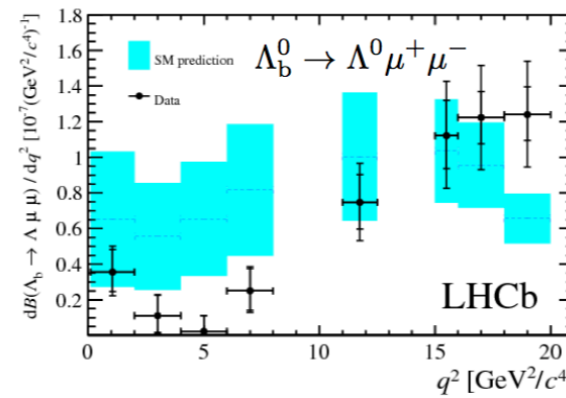
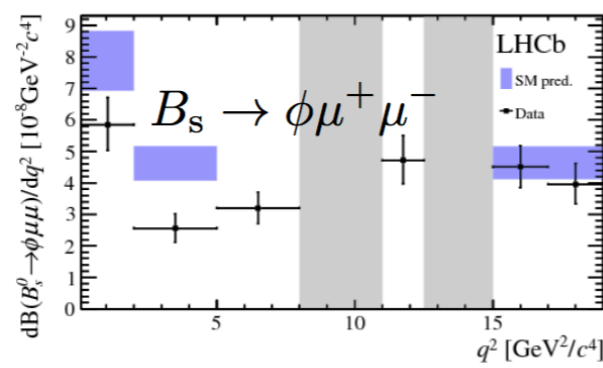
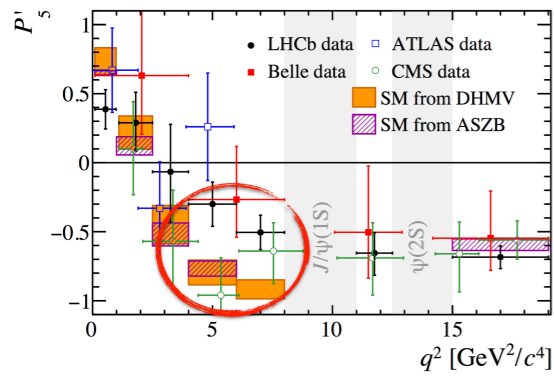
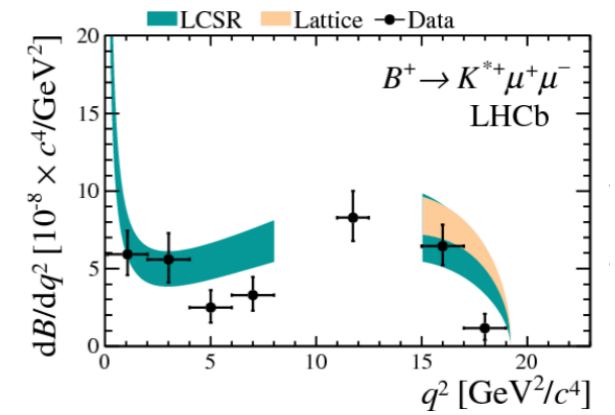
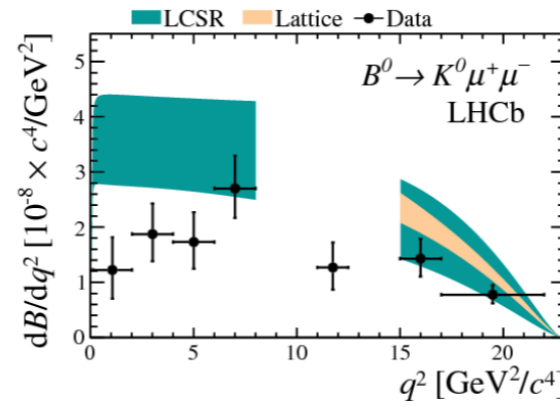
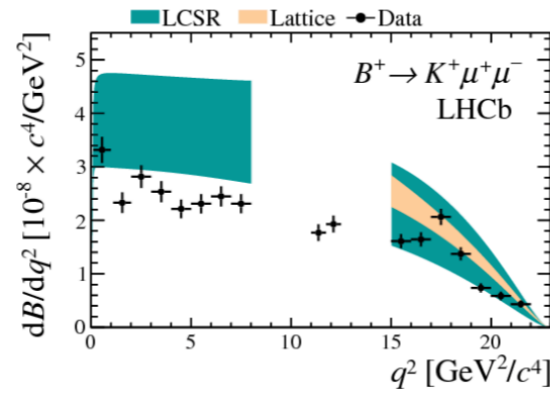
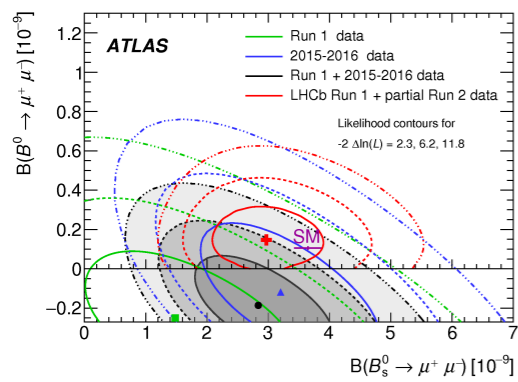


$b \rightarrow sll$   
 $b \rightarrow clv$



$Z', W'$   
leptoquarks  
?

# $b \rightarrow s \mu \mu$ | global fits



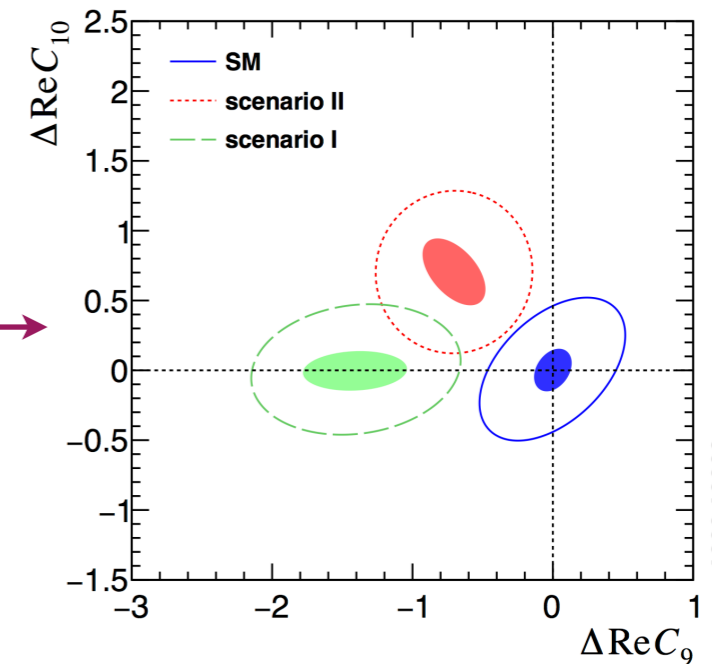
## Effective Field Theory

$$H_{\text{eff}} \propto \sum_i (C_i^{\text{SM}} + C_i^{\text{NP}}) \cdot O_i$$

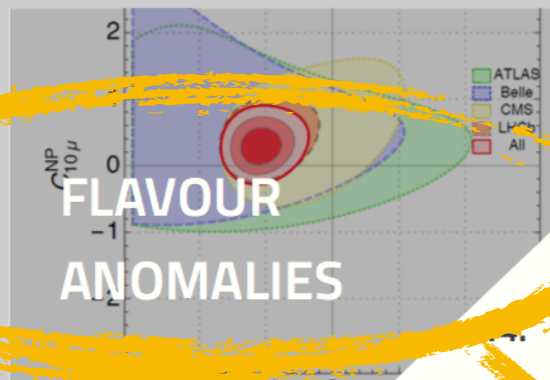
← Now

HL-LHC →

complemented by lepton flavour universality tests!



# LIP NEWS



## Flavour Anomalies First hints of New Physics at the LHC?

Nuno Leonardo

Over the last few years, a persistent set of deviations from the Standard Model (SM) predictions has emerged from the data. These have been detected in decays of b-quark hadrons. While the deviations are not sufficiently significant if considered individually, when taken together they are. These so-called “flavour anomalies” stand currently as a most exciting indication of New Physics (NP) and a hottest topic in the field of HEP at the moment.

New phenomena beyond the standard theory of particle physics are pursued in a multitude of paths. At the LHC, a main path, which explores the energy frontier, aims at directly detecting new heavy particles, beyond those of the SM. These NP particles may be produced in the collisions, and their presence detected through the products of their decay. Another path, which explores the luminosity frontier, aims at detecting the presence of NP indirectly, through precision measurements. Here, NP particles may virtually contribute to the amplitude of SM-allowed processes, and be revealed through measured deviations relative to the SM expectation, in observable particle properties. The two approaches are complementary and each is actively pursued by exploring a large variety of processes.

Hints of the presence of NP may accordingly be revealed through excesses in distributions (e.g. a bump in the mass spectrum) or measured deviations (e.g. on a particle's decay rate). And as it happens, several such hints, of both kinds, have turned up in the LHC data. However, so far, none of sufficiently high statistical significance, so as to unequivocally exclude possible background fluctuations as their source. Nonetheless, in the case of certain b-hadron decays, several such deviations from theory expectation seem to conspire together – while each individual deviation is still not significant *per se*, the coherent pattern displayed by their ensemble is.

Each deviation is associated to one of two underlying b-quark transitions: (i)  $b \rightarrow sll$ , i.e. bottom to strange quark plus pair of opposite-charge leptons, and (ii)  $b \rightarrow cl\nu$ , i.e. bottom to charm quark plus charged lepton and neutrino. The former can occur only at loop level in the SM (flavor changing neutral current, that is forbidden in SM, at tree level), with high sensitivity to NP (where NP particles can run in the loops). The latter (charged current) occurs at tree level.

The neutral-current transitions,  $b \rightarrow sll$ , are realised in various rare B decays, both leptonic, e.g.  $B_s \rightarrow \mu^+ \mu^-$ , and semileptonic, e.g.  $B \rightarrow S \mu^+ \mu^-$ , where S stands for a strange-quark hadron (e.g. K,  $K^*$ ,  $\Phi$ ,  $\Lambda$ ). In addition to decays of the latter class, there are many NP-sensitive observables associated to the angular distributions of the decay products. Deviations are detected to a varying degree in many of these. The departure from theory was initially detected by LHCb in one such angular observable, denoted  $P'_{\mu\mu}$ , in the decay  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ . It should be remarked here that for this decay a challenge arises in calculating the theory predictions – specifically, going from the underlying quark-level transition  $b \rightarrow sll$  to the

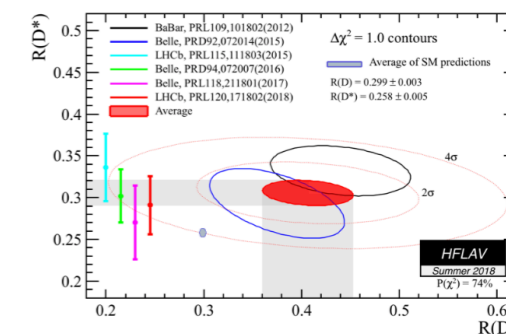
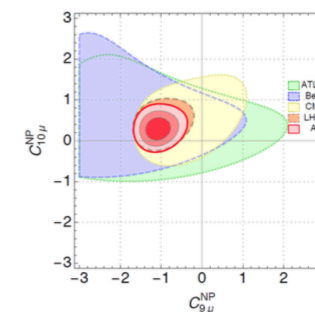
experimentally observed B-meson decay, there are QCD contributions involved whose estimation is non-trivial. And while the  $P'_{\mu\mu}$  observable is constructed in such a way as to be more robust in terms of such QCD ( $B \rightarrow S$ ) form-factor determinations, some debate persists on the theory front.

There is another major chapter in the saga of flavor anomalies. And this time perhaps even more dramatic: it involves violation of lepton flavor universality (LFU). Apart from the differences in their masses, the SM interactions do not distinguish between the different leptons. This means, for example, that the rates of the decays  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  and  $B^0 \rightarrow K^{*0} e^+ e^-$  involving muons and electrons should be comparable. The LHCb data has however revealed that their ratio,  $R_{K^*}$ , seems to display a noticeable departure from unit. Important to remark here is that the above-mentioned form-factor uncertainties cancel in the ratios, rendering these observables rather robust theoretically. Indications of LFU violation had actually been also detected earlier at the B factories (BaBar and Belle experiments), between taus and muons, in the decays  $B \rightarrow D^{(*)} \tau \nu$  and  $B \rightarrow D^{(*)} \mu \nu$ , where the corresponding ratios,  $R_D$  and  $R_{D^*}$ , exhibit departures from their SM expectations (see figure). These were quite unexpected, with the underlying transitions  $b \rightarrow cl\nu$  occurring at tree level.

Naturally, the anomalies have raised a large excitement amongst both experimentalists and theorists. After all, the ensemble of anomalies when interpreted collectively appear to indicate a departure from the SM, with a significance above the  $5\sigma$  mark (see figure). Theorists have been actively putting forward classes of models that attempt to explain the anomalies, along with other tensions in the flavor sector, e.g.  $(g-2)_\mu$ , while simultaneously accommodating other experimental constraints, e.g. from  $B_s$  mixing and dilepton mass spectra. Among these, models with extra gauge bosons ( $Z'$ ) or leptoquarks (LQ) appear to be favoured.

From the experimental side, a clarification will be sought by thoroughly exploiting the LHC Run 2 data. Not only will the LHCb measurements be repeated to reach increased precision, contributions from ATLAS and CMS will offer independent input with orthogonal systematics. For example, during 2018 a large, dedicated dataset has been collected by CMS specifically for this purpose. Belle2 is coming online, and within a few years its data will provide decisive input. Dedicated searches for scenarios addressing the anomalies, including  $Z'$  and LQ, will be pursued at the LHC.

Whether the source of the anomalies turns out to be more mundane statistical fluctuations, underestimations in theory calculations, or genuine NP, it is exciting that a clarification is within reach over the next few years. A confirmation of these flavour anomalies would point to new particles or interactions and have profound implications for our understanding of particle physics.



Current status of the flavor anomalies. Left: Global fit to  $b \rightarrow sll$  observables, with results projected on the plane of two EFT coefficients. Right: Fit to  $b \rightarrow cl\nu$  observables. The red ellipses represent the regions favoured by the data. The SM lies at the origin (0,0) of the left plot and on the small region at about (0.3,0.25) on the right plot. The tension between data and SM is clearly visible.