

Stirren and shaken

Dynamical behavior of boson stars and dark matter cores

Lorenzo Annulli, PhD@CENTRA, Instituto Superior Técnico June 25. 2020

6th IDPASC/LIP PhD Students Workshop

In General Relativity, ultralight bosons can clump to form self-gravitating structures

These solutions, called Newtonian boson stars (NBS), seem to be a good description of dark matter (DM) cores in haloes (fuzzy DM models)

Typical bosons masses are $\sim 10^{-22}~{\rm eV}$

- Stability
- Interaction with surrounding bodies

In which way the presence of a black hole changes the local DM density? When a star crosses a NBS, will change its properties? To which extent?

Impact on scalar and gravitational wave (GW) emission

Nontrivial environmental effect on GW phase?

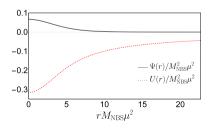
U(1)-invariant, self-interacting, complex scalar field Φ minimally coupled to gravity

$$\mathcal{S} \equiv \int d^4x \sqrt{-g} \left(\frac{R}{16\pi} - \frac{1}{2} g^{\mu\nu} \partial_\mu \Phi \partial_\nu \Phi^* - \frac{\mu^2}{2} |\Phi|^2 \right)$$

Assuming Φ non-relativistic, spherically symmetric and stationary, $\Phi = \Psi(r)e^{-i(\mu-\gamma)t}$, in the Newtonian limit

$$\label{eq:Schrödinger-Poisson:} \begin{aligned} \nabla^2\Psi = 2\mu\left(\mu U + \gamma\right)\Psi\,, \quad \nabla^2U = 4\pi\mu^2\Psi^2 \end{aligned}$$

Fundamental NBSs satisfy the scaling-invariant mass-radius relation $M_{\rm NBS}\simeq 9.1/(R_{\rm NBS}\mu^2)$



Dynamical response to external perturbers, which disturb the spherically symmetric, stationary background,

$$\Phi = \left[\Psi_0(r) + \delta \Psi(t, r, \theta, \varphi)\right] e^{-i\Omega t}$$

Linearized system of equations

$$i\partial_t \delta \Psi = -\frac{1}{2\mu} \nabla^2 \delta \Psi + (\mu U_0 + \gamma) \, \delta \Psi + \mu \Psi_0 \delta U$$
$$\nabla^2 \delta U = 4\pi \left[\mu^2 \Psi_0 \left(\delta \Psi + \delta \Psi^* \right) + P \right]$$

The fluctuation $\delta \Psi$ can be used to calculate physical quantities such as energy, linear and angular momenta radiated, from the flux of certain currents,

$$\left(\dot{E}^{\mathrm{rad}}, \dot{P}^{\mathrm{rad}}_{i}, \dot{L}^{\mathrm{rad}}_{z}\right) = \lim_{r \to \infty} r^{2} \int d\theta d\varphi \sin\theta \, \left(\delta T^{S}_{r\mu} \xi^{\mu}_{t}, \delta T^{S}_{r\mu} e^{\mu}_{i}, \delta T^{S}_{r\mu} \xi^{\mu}_{\varphi}\right)$$

Things we computed

Free perturbations

 Small perturbations aka Quasi Normal Modes

Sourced perturbations

Static

- A super massive black hole sitting in the center of a DM halo
- A black hole eating its host DM halo

Dynamics

- Massive objects plunging into boson stars
- A perturber oscillating at the center
- Low-energy and high-energy binaries within boson stars

Things we computed

Free perturbations

 Small perturbations aka Quasi Normal Modes

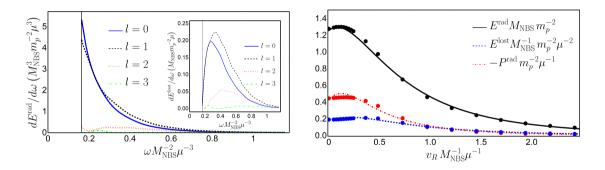
Sourced perturbations

Static

- A super massive black hole sitting in the center of a DM halo
- A black hole eating its host DM halo

Dynamics

- Massive objects plunging into boson stars
- A perturber oscillating at the center
- Low-energy and high-energy binaries within boson stars





First self-consistent calculation of dynamical friction in these backgrounds

Binaries "stir" the NBS core, backreaction affects GW at leading -6PN

Gravitational collapse to a BH is accompanied by only a small change in the DM core

The NBS gets accreted, for realistic parameters only after several Hubble times

Benchmark for numerical relativity simulations involving BS in extreme mass-ratio regime.

Extensions: eccentric motion, self-gravitating vectors or other nonlinearly interacting scalars.