Search for collective behaviour and multiparton interactions in *ep* scattering at HERA

On behalf of the H1 and ZEUS collaborations

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12^{th} Workshop on MPI at the LHC, October 12 2021

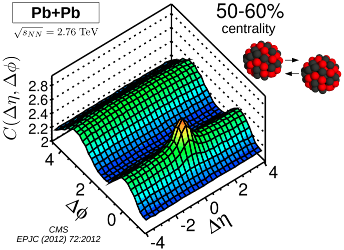
Motivating questions

1. How small can a colliding system be while still exhibiting the collective features typically associated with the quark–gluon plasma in heavy-ion collisions?
2. How many Multiparton Interactions are needed to set the stage for a collective environment?

Recent measurements using the H1 and ZEUS detectors will be presented in neutral current DIS and photoproduction.

New ZEUS publication: arxiv:2106.12377 (submitted to JHEP)
Recent H1 preliminaries: Analysis note

Motivation for the analysis



- Two-particle correlations in heavy-ion collisions show a clear **double ridge**, which is interpreted as a sign of fluid-like behaviour (QGP).
- $C(\Delta\eta, \Delta\varphi) = S(\Delta\eta, \Delta\varphi)/B(\Delta\eta, \Delta\varphi)$,
  S and B are formed from pairs from the same- and mixed-events, respectively.

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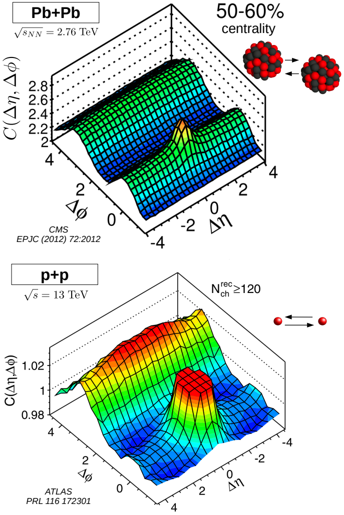

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- Such collisions were thought to be too small to produce a thermally equilibrated QGP.

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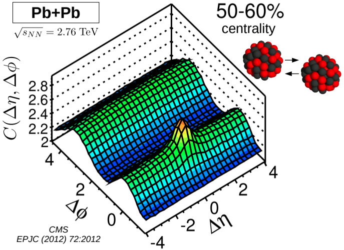

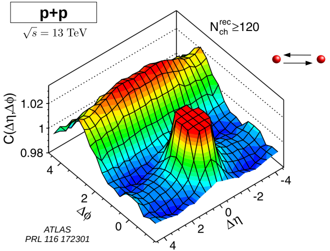

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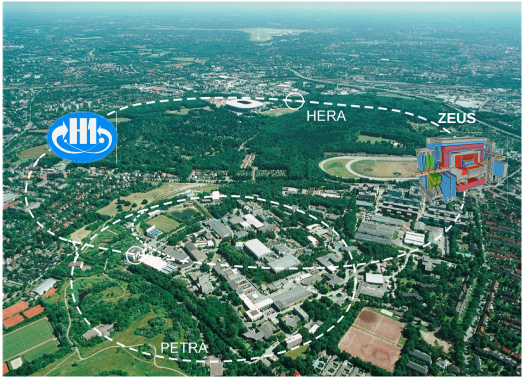
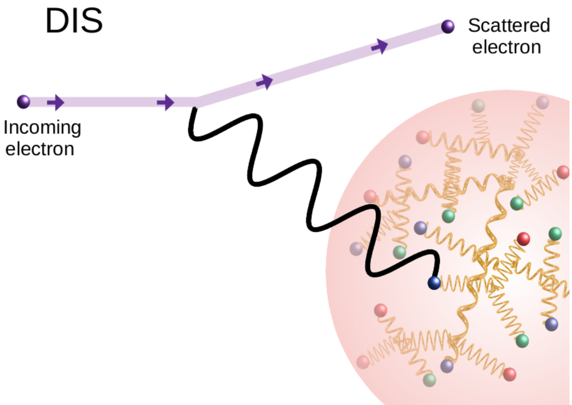
- The start of the LHC revealed that high-multiplicity $p + p$ collisions also have a double-ridge!
- Such collisions were thought to be too small to produce a thermally equilibrated QGP.
- **What about even more fundamental *ep* scattering at HERA??**

The HERA collider and main experiments



- Location: DESY, Hamburg, Germany
- Data taking: 1992 - 2007
- 27.5, 27.6 GeV electrons/positrons
  920 GeV protons
  $\rightarrow \sqrt{s} = 318, 319$ GeV
- HERA I+II:
  500 pb^{-1} per experiment

Deep inelastic scattering (DIS)



- DIS is defined by large virtualities: $Q^2 \gg \Lambda^2_{\mathrm{QCD}}$.
- Transverse radius (R_t) and longitudinal length (L) of the probed region are given by:

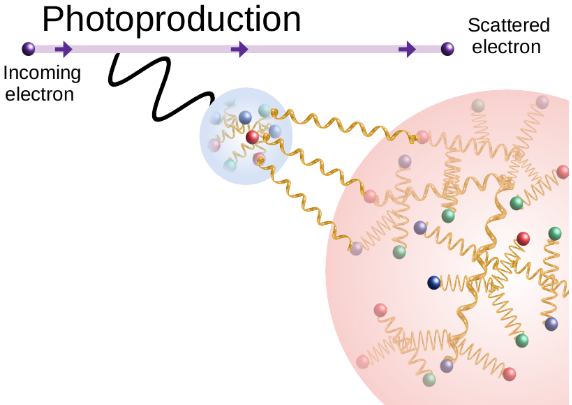
$$R_t \sim \frac{1}{Q}$$

$$L \sim \frac{1}{m_{proton}\, x}$$

PRD 95 114008

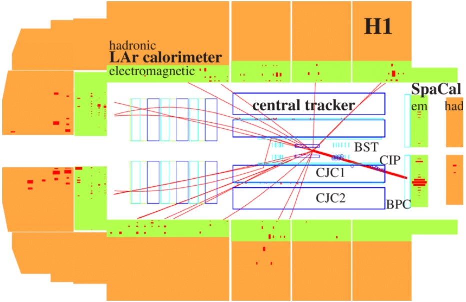
- Neutral current (NC) DIS involves the exchange of photon or Z boson.

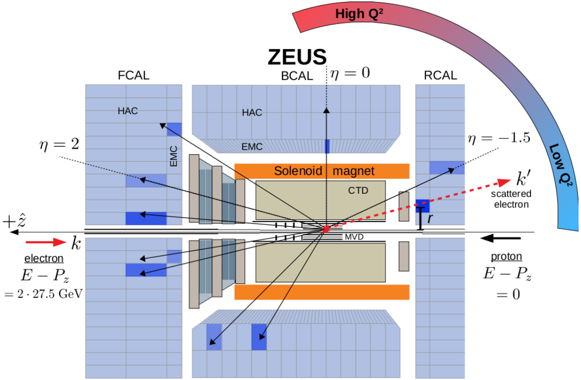
Photoproduction (PhP)



- Photoproduction (γp) is defined by small virtualities: $Q^2 \ll \Lambda^2_{\mathrm{QCD}}$.
- Exchanged photon may fluctuate into quarks and gluons.
- Larger interaction regions are probed.
- **Multiparton Interactions are possible.**
- Scattering is hadron-like.

H1 and ZEUS detectors





Event and track selection (main cuts only)

DIS event selection

	scattered electron	Q^2	$\sum(E_i - P_{z,i})$	N_{ch}	V_Z
H1	in SpalCal	5 to 100 GeV^2	35 to 75 GeV	≥ 2	-35 to +35 cm
ZEUS	in CAL	$\geq$ 20 GeV^2	47 to 69 GeV	≥ 20	-30 to +30 cm

Photoproduction event selection

	scattered electron	$\sum(E_i - P_{z,i})$	N_{ch}	V_Z
H1	in tagger	NA	≥ 2	-30 to +30 cm
ZEUS	absent	$\leq$ 55 GeV	≥ 20	-30 to +30 cm

Track selection

	p_{T}	η	DCA
H1	0.3 to 3 GeV	-1.6 to 1.6 (0 to 5 for DIS in HCM)	$<$ 5 cm in XY
ZEUS	0.1 to 5 GeV	-1.5 to 2.0	$<$ 2 cm in XY and Z

There are several differences between the H1 and ZEUS analyses but compatible results are obtained nevertheless.

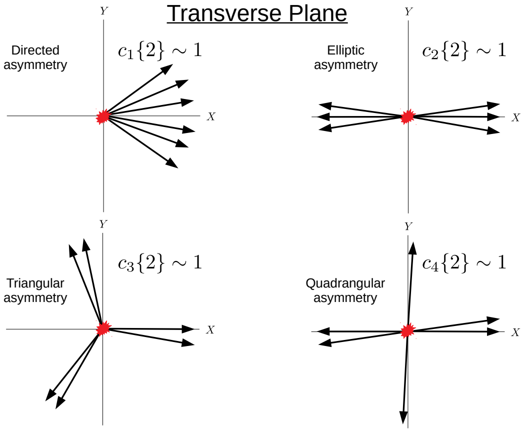
Two-particle correlation functions

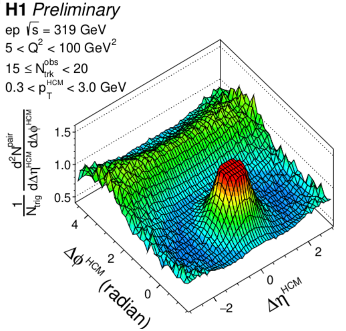
Two-particle azimuthal correlations are measured:

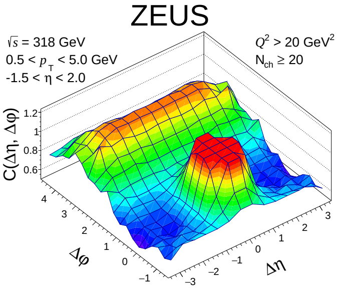
$c_n\{2\} = \langle\langle \cos n(\phi_i - \phi_j) \rangle\rangle$.

φ_i is the azimuthal angle of particle i.

n is the harmonic.





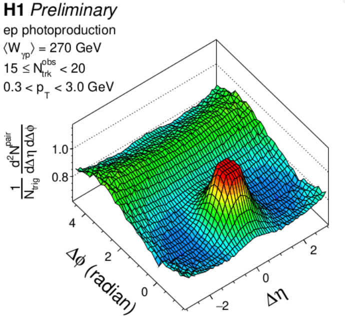


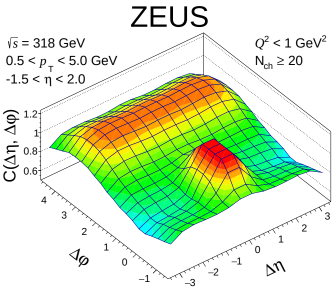
A near-side peak and away-side ridge are clearly visible.

No visible double-ridge.

Note: Kinematic selection differs between H1 and ZEUS.

Results: H1 & ZEUS ridge plots in photoproduction





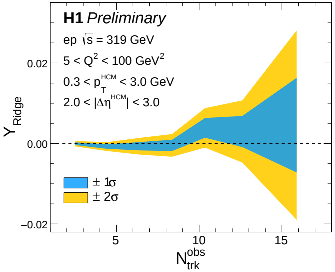
A near-side peak and away-side ridge are clearly visible.

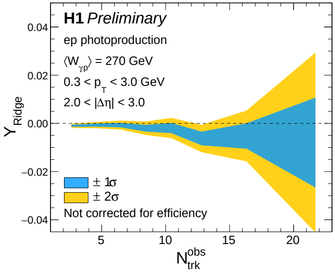
No visible double-ridge.

Correlation strengths are significantly smaller than those in DIS.

Note: Kinematic selection differs between H1 and ZEUS.

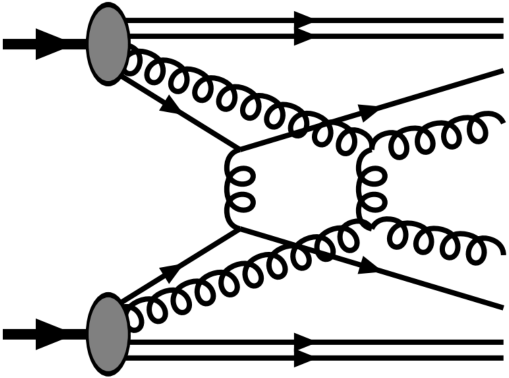
Ridge yields in H1





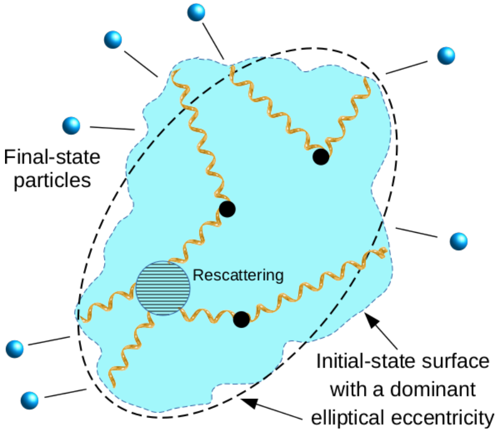
Using a Zero-Yield-At-Minimum assumption, the ridge yields are extracted.
Ridge yields in both DIS and Photoproduction are consistent with zero.

Multiparton Interactions (MPI)



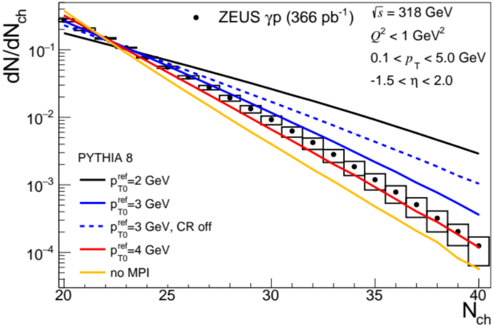
- MPI occur when there's more than one $2 \to 2$ partonic scattering between the beam particles in a given event.
- If the scatterings are sufficiently hard ($p_T \gtrsim 1$ GeV), they can be modeled in an event generator like PYTHIA.
- **Established feature in high-multiplicity hadronic collisions. So far not conclusively observed in *ep* scattering.**

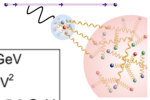
A subsequent rescattering phase is possible



- The initial scattering is shown here with 3 MPIs (black dots)
- Unlike in DIS, the spatial extent of this "initial state" is finite with an irregular shape in general.
- Subsequently, a phase of rescattering may occur, whereby a local thermal equilibrium might form.

Results: dN/dN_{ch}





	$p_{\mathrm{T0}}^{\mathrm{ref}} = 2$	$p_{\mathrm{T0}}^{\mathrm{ref}} = 3$	$p_{\mathrm{T0}}^{\mathrm{ref}} = 4$
$\langle nMPI \rangle$	8.3	3.8	2.2

The number of MPI (nMPI) and IR divergencies are controled by the p_{T0} parameter in PYTHIA.
It is used to regularize the interaction cross section in PYTHIA.

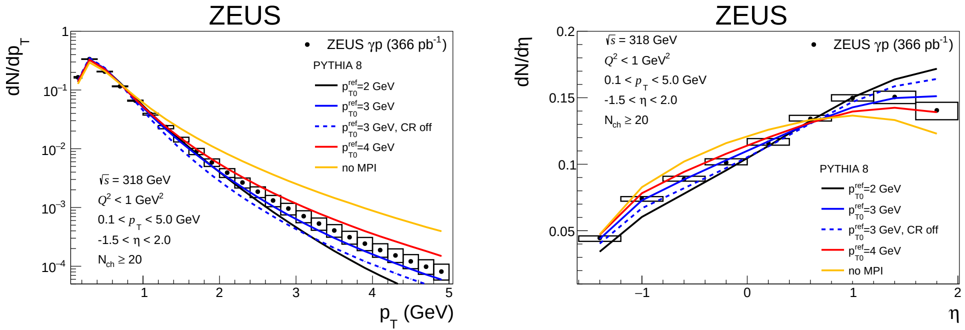
$$\frac{d\sigma}{dp_{\mathrm{T}}^2} \propto \frac{\alpha_s^2(p_{\mathrm{T0}}^2 + p_{\mathrm{T}}^2)}{(p_{\mathrm{T0}}^2 + p_{\mathrm{T}}^2)^2}$$

The energy dependence of this parameter is given by $p_{\mathrm{T0}} = p_{\mathrm{T0}}^{\mathrm{ref}} \, (W/7\,\mathrm{TeV})^{0.215}$, where W is the γp $\sqrt{s}$.

More MPI → lower $p_{\mathrm{T0}}^{\mathrm{ref}}$

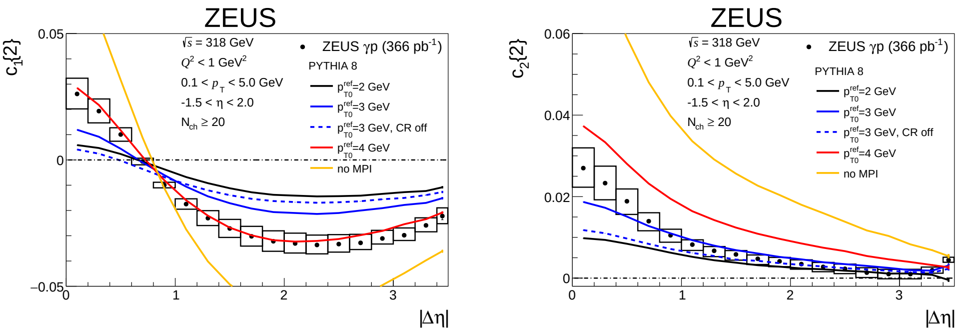
Colour Reconnection (CR) is PYTHIA's modeling of rescattering between partons from different MPIs

Results: dN/dp_T and $dN/d\eta$



- The scenarios of no MPI and very many MPI are disfavored.
- Red and blue lines favored: $2 \lesssim \langle nMPI \rangle \lesssim 4$.

Results: $c_1\{2\}$ and $c_2\{2\}$ versus $\Delta\eta$



- Correlation strengths are diluted by MPI.
- The scenarios of no MPI and very many MPI are disfavored.
- Red and blue lines favored: $2 \lesssim \langle nMPI \rangle \lesssim 4$.

Summary

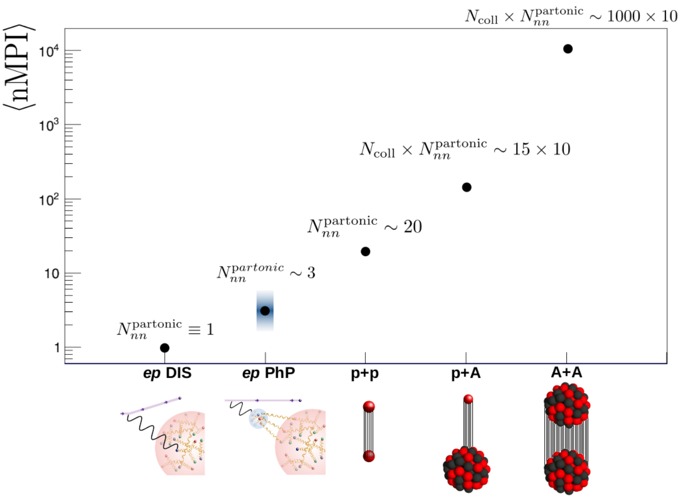
- Measurements of charged-particle azimuthal correlations have been presented using H1 and ZEUS data in *ep* photoproduction (γp) and NC DIS.
- There is no clear indication of a double ridge in either γp or DIS. The observations **do not** reveal significant collective behaviour like that seen in heavy-ions or high-multiplicity hadronic collisions.
- The concept of multiparton interactions provides a useful tool to help understand the emergence of collective behaviour. It sets the stage for a potential rescattering phase.

	nMPI	Collectivity
ep photoproduction	~ 3	No
pp high-multiplicity	~ 20	Yes

The initial states in both systems may be similar in their spatial extent but completely different in the number of MPI.

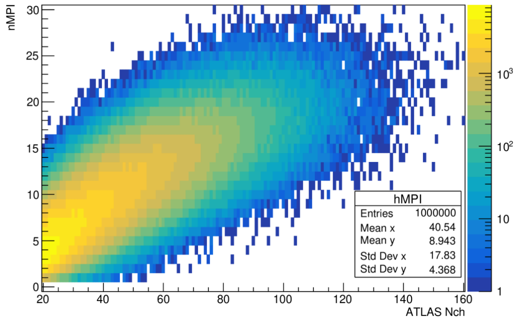
Backup

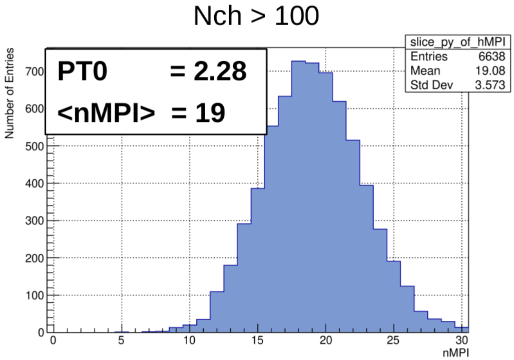
Illustration of MPI growth



- Rough illustration of how nMPI grows from DIS to heavy-ions
- N_{coll}: number of binary nucleon-nucleon collsions
- $N_{nn}^{\mathrm{partonic}}$: number of parton scatterings per binary nucleon-nucleon collision
- Estimates for N_{coll} taken from
  - Ann. Rev. Nucl. Part. Sci. 57, 205 (2007)
  - PRC 97 024905 (2018).
- Estimates for $N_{nn}^{\mathrm{partonic}}$ taken from PYTHIA

nMPI in high-multiplicity $p + p$ PYTHIA at LHC energies





PYTHIA $p + p$ events at $\sqrt{s} = 13$ TeV were generated.
N_{ch} was counted according to the ATLAS acceptance used in PRL 116 172301.
$-2.5 < \eta < 2.5$, $0.4 < p_{\mathrm{T}} < 50$ GeV

Two- and four-particle correlation functions

Two-particle azimuthal correlations are measured:

$c_n\{2\} = \langle\langle \cos n(\phi_i - \phi_j)\rangle\rangle.$

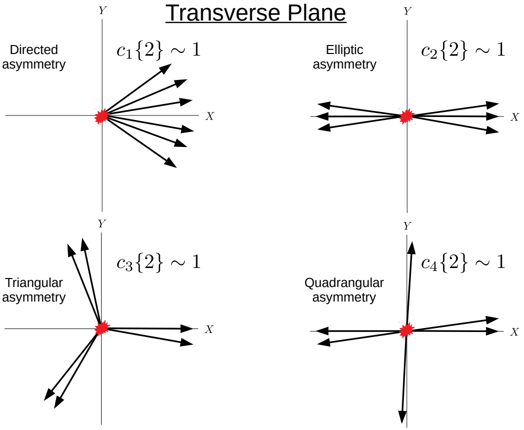
φ_i is the azimuthal angle of particle i.

n is the harmonic.

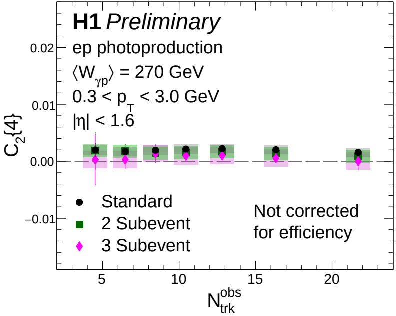
Four-particle cumulant correlations are also measured:

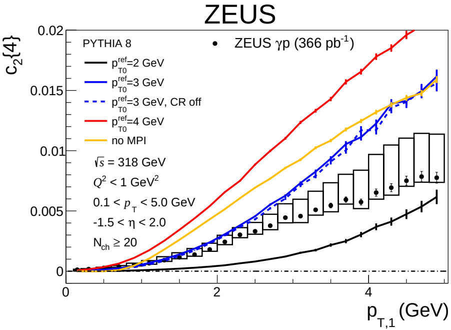
$$C_n\{4\} = \langle\langle \cos n(\phi_i + \phi_j - \phi_k - \phi_l)\rangle\rangle$$
$$c_n\{4\}(p_{\mathrm{T},1}) = C_n\{4\}(p_{\mathrm{T},1}) - 2\,c_n\{2\}(p_{\mathrm{T},1})\,c_n\{2\}$$

where $p_{\mathrm{T},1}$ is the transverse momentum of particle i.



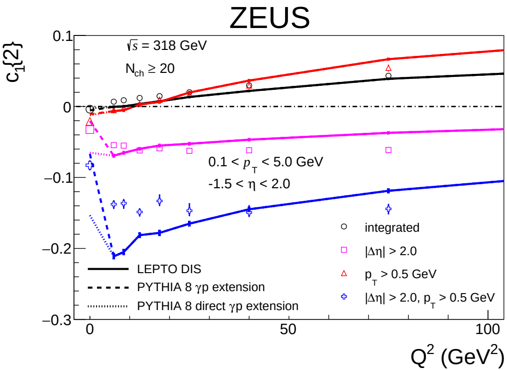
Results: Four-particle cumulants in photoproduction

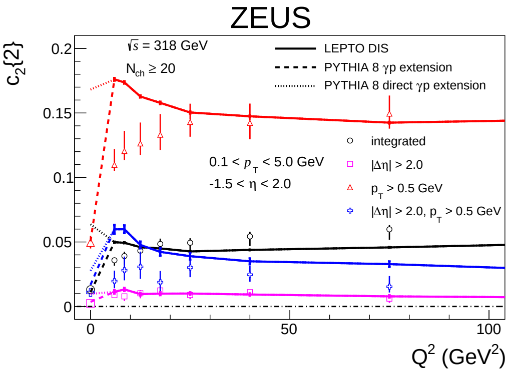




- Four-particle cumulant is positive, which is in contrast to the negative values seen in non-central heavy-ion collisions.
- The scenarios of no MPI and very many MPI are disfavored.

Results: Q^2 evolution of $c_1\{2\}$



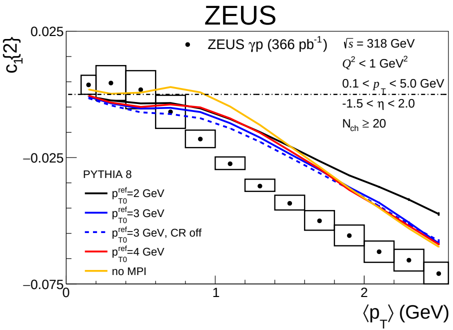


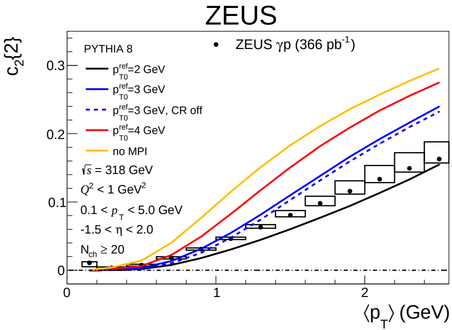
Photoproduction correlation strengths ($Q^2 = 0$) are clearly diminished wrt those in DIS.

The LEPTO model of DIS gives a rough qualitative description of the data.

PYTHIA 8 with only the direct component of γp predicts much stronger correlations than the full calculation (direct + resolved).

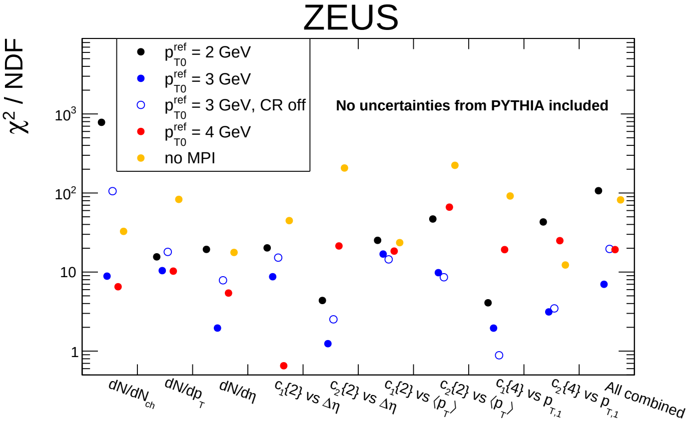
Results: $c_1\{2\}$ and $c_2\{2\}$ versus $\langle p_\mathrm{T} \rangle$



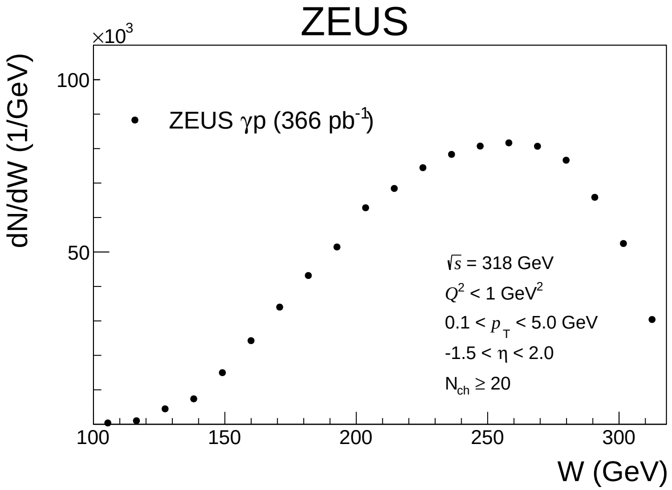


- $c_1\{2\}$ versus $\langle p_\mathrm{T} \rangle$ not sensitive to MPI and not described well by PYTHIA.
- More extreme levels of MPI are favored by $c_2\{2\}$ versus $\langle p_\mathrm{T} \rangle$.

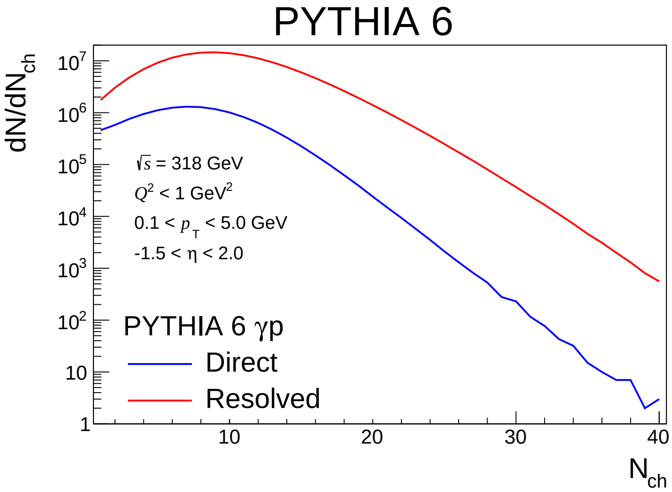
Condensed view of PYTHIA 8 comparisons



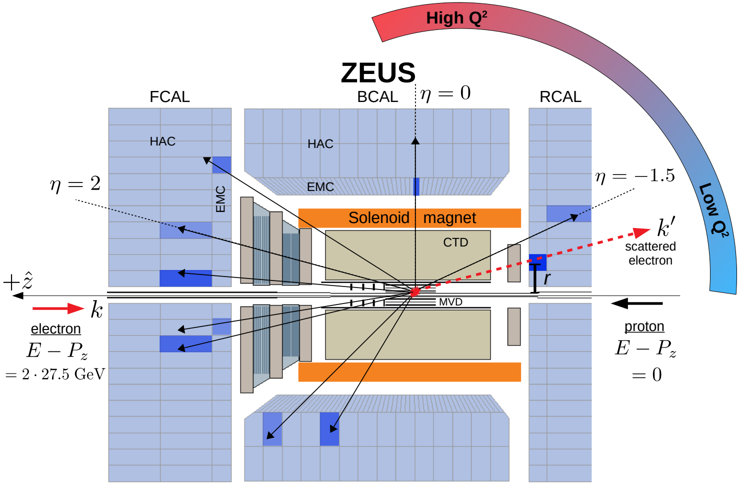
W distribution



Direct and Resolved event distributions



ZEUS track selection

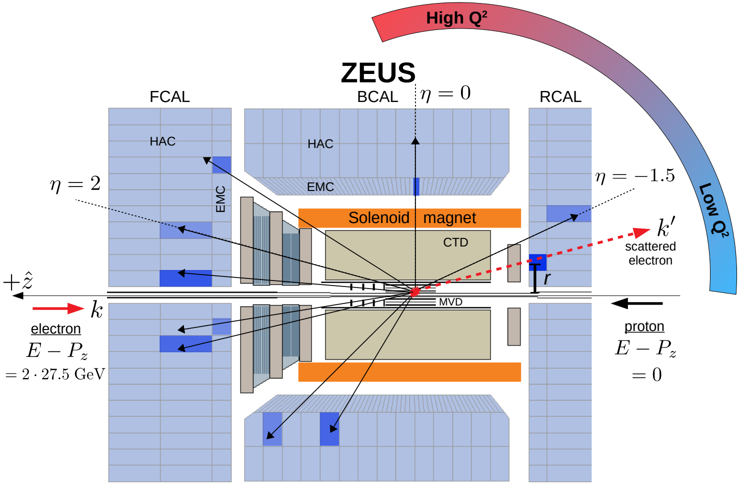


Track selection for correlation analysis

- Reject scattered electron (if detected)
- $-1.5 < \eta < 2.0$
- $0.1 < p_T < 5.0$ GeV
- ≥ 1 MVD hit
- $\mathrm{DCA}_{XY,Z} < 2$ cm
- $\Delta R > 0.4$ (cone around scattered electron)

$$N_{\mathrm{ch}} = \sum_{i}^{N_{\mathrm{rec}}} w_i^{(1)}$$

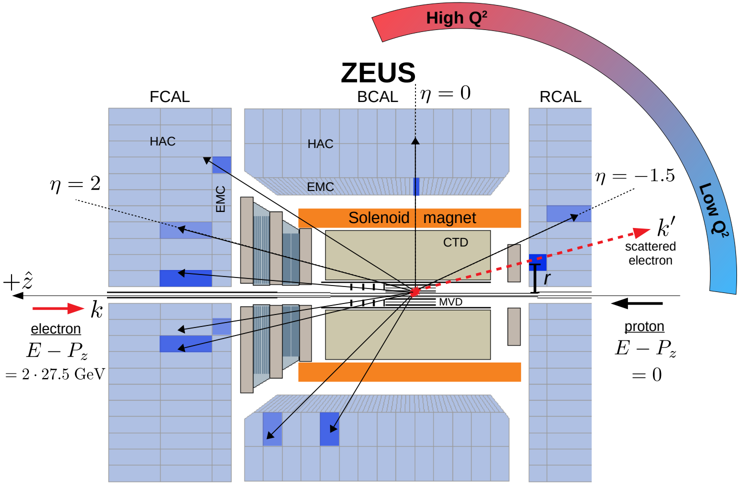
ZEUS DIS event selection



DIS Event selection (0.2 M)

- $N_{\rm ch} \geq 20$
- DIS triggers
- electron probability $> 90\%$
- $Q^2 = -(k - k')^2 > 5\ \mathrm{GeV}^2$
- $k'_0 > 10$ GeV
- $r > 15$ cm
- $\theta_e > 1$ rad
- $47 < \sum(E_i - P_{z,i}) < 69$ GeV
- $|V_z| < 30$ cm

ZEUS photoproduction event selection



Photoproduction event selection (5 M)

- $N_{ch} \geq 20$
- PhP oriented triggers
- electron probabiltiy $< 90\%$
- $k'_0 < 15$ GeV
- $\sum(E_i - P_{z,i}) < 55$ GeV
- $|V_z| < 30$ cm

Tracking efficiency corrections

The efficiency correction weights for 1-, 2-, and 4-particle distributions are defined as:

$$w^{(n)} = \frac{N^n_{gen}(\vec{x})}{N^n_{rec}(\vec{x})}$$

The are computed differentially in Monte Carlo simulations of the ZEUS detector:

dimension of $\vec{x}$	One-particle (n=1)	Two-particle (n=2)	Four-particle (n=4)
x_1	φ	$\varphi_1 - \varphi_2$	$\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4$
x_2	η	$\langle \eta_i - \langle \eta \rangle \rangle$	$\langle \eta_i - \langle \eta \rangle \rangle$
x_3	p_{T}	$\langle p_{T,i} - \langle p_T \rangle \rangle$	$\langle p_{T,i} - \langle p_T \rangle \rangle$
x_4 (charge)	q	$\|q_1 + q_2\|$	$\|q_1 + q_2 + q_3 + q_4\|/2$
x_5	-	N_{rec}	N_{rec}