

# Exercises

\* What are  $n, u, P$  for relativistic degenerate fermion gas?

$$dN = gV \frac{d^3k}{(2\pi)^3} \Rightarrow n = \frac{k_F^3}{3\pi^2}$$

$$u = \int \frac{2}{8\pi^3} \underbrace{c k}_{\text{relativistic dispersion}} 4\pi n^2 dk = \frac{k_F^4}{4\pi^2}$$

$$P = \frac{2}{3} \int k \cdot c \frac{k^2}{2\pi^2} dk = \frac{1}{3\pi^2} \frac{k_F^4}{4} = \frac{1}{3} u$$

$$\boxed{P = \frac{1}{3} u}$$

## \* Lane-Emden Solutions

$$\underline{n=0} \quad \frac{1}{\mathcal{F}^2} \frac{d}{d\mathcal{F}} \left( \mathcal{F}^2 \frac{d\vartheta_0}{d\mathcal{F}} \right) = -\vartheta_0 = -1$$

$$\int \frac{d}{d\mathcal{F}} \left( \mathcal{F}^2 \frac{d\vartheta_0}{d\mathcal{F}} \right) d\mathcal{F} = \int -\mathcal{F}^2 d\mathcal{F} \Rightarrow$$

$$\mathcal{F}^2 \frac{d\vartheta_0}{d\mathcal{F}} = C_1 - \frac{\mathcal{F}^3}{3} \Rightarrow$$

$$\frac{d\vartheta_0}{d\mathcal{F}} = \frac{C_1}{\mathcal{F}^2} - \frac{\mathcal{F}}{3}$$

$C_1$  must be zero because otherwise

$$\frac{d\vartheta_0}{d\mathcal{F}} \rightarrow \infty \text{ as } \mathcal{F} \rightarrow 0$$

$$\theta_0 = c_2 - \frac{f^2}{6} \quad \underline{\theta_0(f=0) = 1}$$

$$\theta_0(f) = 1 - \frac{f^2}{6} \quad f_1 = \sqrt{6}$$

$$R = 2f_1$$

$$n=1 \quad \frac{d}{df} \left( f^2 \frac{d\theta_1}{df} \right) = -f^2 \theta_1$$

$$f^2 \theta_1'' + 2f \theta_1' + f^2 \theta_1 = 0$$

$$\theta_1'' + \theta_1 + \frac{2}{f} \theta_1' = 0$$

$$\text{Try } \theta_1 = A \sin f \cdot f^s$$

$$A(1+s) f^{-2+s} (2f \cos f + s \sin f) = 0$$

For  $S = -1$  we get a solution

$$\theta_1 = \frac{\sin J}{J} \quad \lim_{J \rightarrow 0} \frac{\sin J}{J} = 1 \quad J_1 = \pi$$

\* Relate  $M$  and  $R$  without explicit or implicit dependence on  $\rho_c$

$$R = \left[ \frac{(n+1) K \rho_c^{\frac{1-n}{n}}}{4\pi G} \right]^{1/2} \quad J_1 = a J_1$$

$$M = -4\pi a^3 \rho_c J_1^2 \left. \frac{d\theta}{dJ} \right|_{J_1}$$

$$a = \left[ \frac{(n+1) K \rho_c^{\frac{1-n}{n}}}{4\pi G} \right]^{1/2}$$

$$R \sim a \sim \rho_c^{\frac{1-n}{2n}}$$

$$M \sim d^3 \rho_c \sim \rho_c^{\frac{3-3n}{2n} + 1} \sim \rho_c^{\frac{3-n}{2n}}$$

$$K = N_n G M^{\frac{n-1}{n}} R^{\frac{3-n}{n}}$$

This combination is independent of  $\rho_c$

$$N_n = \frac{(4\pi)^{1/n}}{n+1} \left[ -\mathcal{F}_1^{\frac{n+1}{n-1}} \frac{d\mathcal{D}_n}{d\mathcal{F}_1} \right]^{\frac{1-n}{n}}$$

$\rho = K \rho^{5/3}$  non-relativistic  
 $\gamma = 1 + \frac{1}{n} \Rightarrow n = \frac{3}{2}$

$\rho = K \rho^{4/3}$  relativistic  $n = 3$

$$K = N_{3/2} G M^{\frac{1}{2}} R = N_{3/2} G M^{\frac{1}{2}} R$$

$$M^{1/3} R = \text{const.} \Rightarrow \boxed{M \sim \frac{1}{R^3}}$$

$$K = N_3 G M^{\frac{2}{3}} R^0 = N_3 G M^{\frac{2}{3}}$$

$$M = \text{const.} \quad E = \sqrt{c^2 u^2 + m^2 c^4}$$

Chandrasekhar mass limit

## \* Gravitational Collapse of fermions

Non-interacting

$$E = KE + PE = c k_F - \frac{GMm}{r} < 0$$

$$\eta = \frac{k_F^3}{3\pi^2} \Rightarrow k_F = (3\pi^2 \eta)^{1/3} = \left( \frac{3\pi^2 N}{\frac{4}{3}\pi r^3} \right)^{1/3} =$$

$$= \left( \frac{9\pi}{4} \right)^{1/3} \frac{N^{1/3}}{r}$$

$$E = \left( \frac{9\pi}{4} \right)^{1/3} \frac{N^{1/3}}{r} - G \frac{Nm^2}{r} < 0$$

$$G N m^2 > \left(\frac{9\pi}{4}\right)^{1/3} N^{1/3} \Rightarrow$$

$$N^{2/3} > \left(\frac{9\pi}{4}\right)^{1/3} \frac{1}{G m^2} \quad G = \frac{1}{M_{pl}^2}$$

$$N^{2/3} > \left(\frac{9\pi}{4}\right)^{1/3} \frac{M_{pl}^2}{m^2} \Rightarrow$$

$$N > \left(\frac{9\pi}{4}\right)^{3/2} \frac{M_{pl}^3}{m^3} = N_{crit}$$

For  $m = 1 \text{ GeV}$   $M_{crit} = N_{crit} m = \mathcal{O}(M_\odot)$

Maximum NS mass  $\sim 2.4 M_\odot$

TOV  $M = 0.7 M_\odot$

# \* Gravitational Collapse of bosons

Uncertainty principle  $\Delta x \Delta p \geq \hbar$

$$\Delta p \geq \frac{\hbar}{\Delta x} \quad p_{\text{crit}} = \frac{1}{r} \quad \hbar = c = 4.5 \times 10^{-11} \text{ m}$$

$$E = KE + PE = \overset{=1}{c} p_{\text{crit}} - \frac{GMm}{r} =$$

$$= \frac{1}{r} - \frac{GNm^2}{r} < 0$$

$$1 < Nm^2G \Rightarrow N \geq \frac{m_{\text{pl}}^2}{m^2}$$

$$N_{\text{crit}} = \frac{m_{\text{pl}}^2}{m^2}$$

bosons

$$\frac{10^{38}}{m^2} \sim \frac{10^{-19} M_{\odot}}{m^2}$$

$$N_{\text{crit}} \sim \frac{m_{\text{pl}}^3}{m^3}$$

fermions

$$\frac{10^{57}}{m^3} \sim \frac{M_{\odot}}{m^3}$$



\* Astrophysical Black holes

$$M > 3M_{\odot}$$

Mass of supermassive star  $9M_{\odot} \leq M \leq 200M_{\odot}$

NS formation,  $M \gtrsim 20M_{\odot}$  BH formation

\* Primordial Black holes

# Chemical Equilibrium & Charge Neutrality

$$n \rightarrow p + e^- + \bar{\nu}_e$$

$$\mu_n = \mu_p + \mu_e \quad (1) \text{ (neutrinos escape)}$$

$$n_p = n_e \quad \text{(charge neutrality)}$$

$$n_p = \frac{k_{FP}^3}{3\pi^2}, \quad n_e = \frac{k_{Fe}^3}{3\pi^2} \Rightarrow k_{FP} = k_{Fe}$$

At zero temperature

$$\mu_p = E_p = m_p + \frac{k_{FP}^2}{2m_p} \quad \text{proton non-relativistic}$$

$$\mu_e = k_{Fe} c \stackrel{1}{=} k_{FP} \quad \text{electron relativistic}$$

we should check this a posteriori

$$m_n + \frac{k_{Fn}^2}{2m_n} = m_p + \frac{k_{Fp}^2}{2m_p} + \underbrace{k_{Fp}}_{p_e} \quad (1)$$

$$p_{tot} = n_p m_p + n_n m_n \quad (2) \text{ (ignore electrons)}$$

$$\frac{k_{Fn}^2}{2m_n} = k_{Fp} \left( 1 + \frac{k_{Fp}}{2m_p} \right) - Q$$

$$Q = m_n - m_p \approx 1.3 \text{ MeV}$$

we assume  $\frac{k_{Fp}^2}{m_p} \gg Q \ll k_{Fp}$

(must be checked later)

$$\frac{k_{Fn}^2}{2m_n} = k_{Fp} \quad (3)$$

$$\textcircled{2} \quad \rho_{\text{tot}} \approx m_n n_n + m_p n_p \approx m_n n_n$$

$$n_n = \frac{\kappa_{Fn}^3}{3\pi^2} \quad n_p = \frac{\kappa_{Fp}^3}{3\pi^2}$$

(must be checked!)

$$\kappa_{Fn} = (3\pi^2 n_n)^{1/3}$$

$$\textcircled{2} \Rightarrow \rho_{\text{tot}} \approx n_n m_n \Rightarrow n_n \approx \frac{\rho_{\text{tot}}}{m_n}$$

$$= \frac{\rho_{\text{tot}}}{\rho_{\text{nuc}}} \frac{\rho_{\text{nuc}}}{m_n}$$

$$\rho_{\text{nuc}} \approx 2.8 \times 10^{14} \text{ g/cm}^3$$

$$n_n \approx 1.7 \times 10^{38} \text{ cm}^{-3} \left( \frac{\rho_{\text{tot}}}{\rho_{\text{nuc}}} \right)$$

$$\begin{aligned}
 n_p = n_e &= \frac{k_{FP}^3}{3\pi^2} \stackrel{\textcircled{3}}{\approx} \left( \frac{k_{Fn}}{2m_n} \right)^3 \frac{1}{3\pi^2} = \\
 &= \frac{1}{(2m_n)^3} \underbrace{\left( \frac{k_{Fn}}{3\pi^2} \right)^2}_{n_n} \frac{(3\pi^2)^2}{3\pi^2} = \frac{3\pi^2}{8m_n^3} n_n^2
 \end{aligned}$$

$$n_p = n_e \approx 9.5 \times 10^{35} \text{ cm}^{-3} \left( \frac{\rho_{\text{tot}}}{\rho_{\text{nuc}}} \right)^2$$

Indeed for  $\rho_{\text{tot}} \approx \rho_{\text{nuc}}$

$$n_p = n_e \ll n_n$$

$$k_{Fn} = (3\pi^2 n_n)^{1/3} \approx 340 \left( \frac{\rho_{\text{tot}}}{\rho_{\text{nuc}}} \right)^{1/3} \frac{\text{MeV}}{c}$$

$$k_{FP} = k_{Fe} = (3\pi^2 n_p)^{1/3} = 60 \left( \frac{\rho_{\text{tot}}}{\rho_{\text{nuc}}} \right)^{2/3} \frac{\text{MeV}}{c}$$

$$m_p \sim 0.9 \text{ GeV}$$

Indeed protons are  
non-relativistic

$$m_e \sim 0.5 \text{ MeV}$$

$$\mu_{Fe} \sim 60 \text{ MeV}$$

electrons are  
relativistic

$$n \approx \frac{1}{e^{\frac{E-\mu}{k_B T}} + 1}$$

$$\frac{\mu}{T} \gg 1 \quad \text{cold object}$$

$$\frac{\mu}{T} \ll 1 \quad \text{hot object}$$

$$10^{13} \mu \sim 1 \text{ GeV}$$

$$10^7 \mu \sim 1 \text{ keV} = T$$

$$\mu \sim \text{GeV}$$