



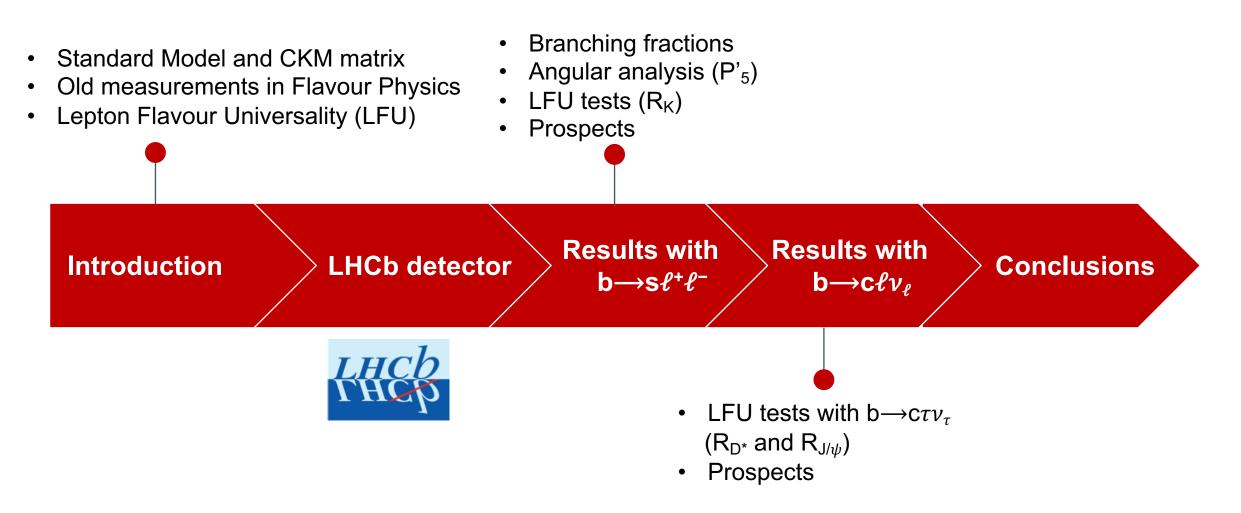
Lepton Flavour Universality tests with semileptonic b-hadron decays

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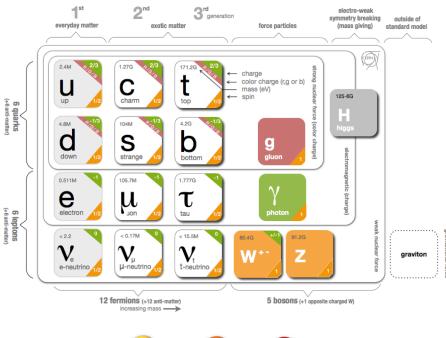
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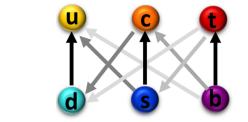
LIP Seminar, 11th July 2019

Outline



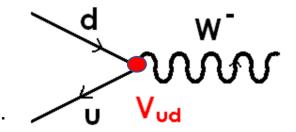
Standard Model and CKM matrix





 $V_{\rm CKM} = \begin{pmatrix} V_{\rm ud} & V_{\rm us} & V_{\rm ub} \\ V_{\rm cd} & V_{\rm cs} & V_{\rm cb} \\ V_{\rm td} & V_{\rm ts} & V_{\rm tb} \end{pmatrix} \approx \begin{bmatrix} 0.974 & 0.225 & 0.003 \\ 0.225 & 0.973 & 0.041 \\ 0.009 & 0.040 & 0.999 \end{bmatrix}$

- In the Standard Model (SM), quarks and leptons are divided in 3 families (or generations).
- Interactions described by the exchange of gauge bosons: photon (EM), W/Z (weak), gluon (strong).
- Flavour Physics: study of transitions between quarks of different flavour.
 - described by the CKM matrix
 ⇒ mediated by a W boson in the SM.



- Transitions between quarks of different families suppressed: $|V_{tb}| \sim 1$, $|V_{cb}| \sim 0.04$, $|V_{ub}| \sim 0.004$.
- However, SM does not account for:
 - Matter/antimatter asymmetry in the Universe.
 - Dark matter.
 - Dark energy.

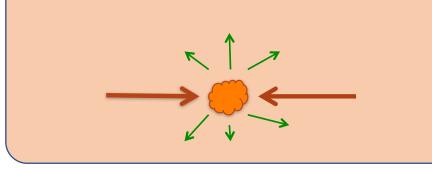
• The SM must be a **low-energy effective theory**.

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Search for NP: direct vs indirect approaches

High energy (direct) approach

- Energy in particle collisions large enough to create new **real** particles.
- Particles appear as "peaks" in a given distribution.
- Approach followed by ATLAS and CMS.



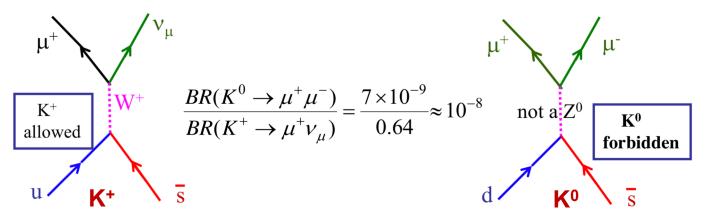
High precision (indirect) approach

- Precision of the measurements high enough to detect New Physics (NP) effects due to virtual particles.
- Indirect measurements can access higher energy scales.
- Approach followed by LHCb and Bfactories.

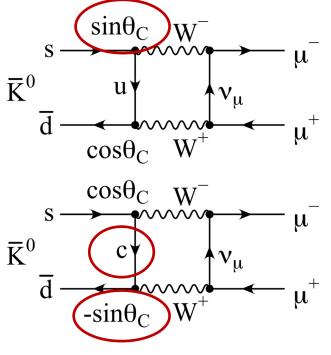


Flavour physics in the past: GIM mechanism

- In the past, flavour physics lead to several **indirect** discoveries:
 - Branching fraction of leptonic K decay **BR**(K⁰ $\rightarrow \mu^+\mu^-$) very suppressed.



cancellation between loop diagrams



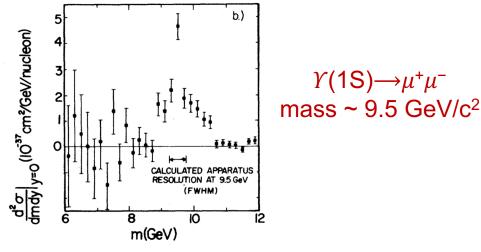
(exact cancellation in the massless quark limit) 5

- **1970**: Glashow, Iliopoulos and Maiani proposed a solution (**GIM** mechanism, only 3 quarks (u,d,s) known at the time) [PRD 2, 1285] <u>(1970)</u>]:
 - Flavour Changing Neutral Currents (FCNC) forbidden in the • SM at tree level.
 - FCNC suppressed: only through **loop** diagrams. •
 - Prediction of the charm quark \Rightarrow Observed in 1973 (J/ ψ). • [PRL 33, 1404, (1974) & PRL 33, 1406(1974)]

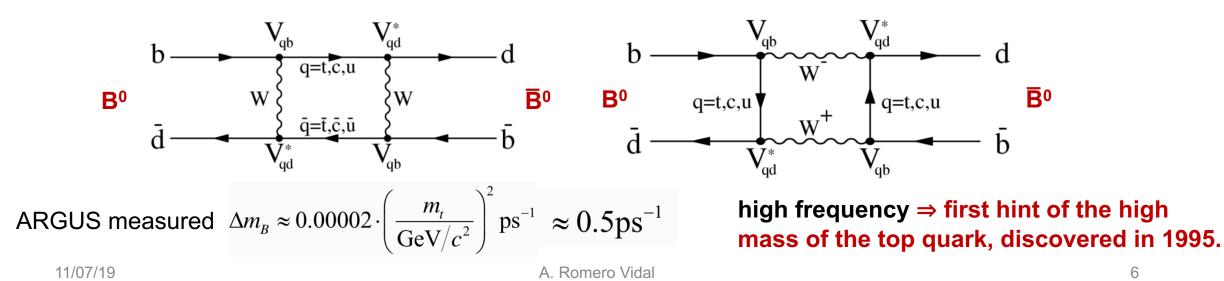
Flavour physics in the past: CPV and B⁰ mixing

 1964: observation of (indirect) CP violation (CPV) in kaons : K⁰_L(CP=−1)→π⁺π⁻(CP=+1) [<u>PRL 13, 138 (1964)</u>]

⇒ 1973: Kobayashi and Maskawa demonstrated that this could be explained if there are at least 3 generations of quarks [<u>PTP</u> 49, 2 (1973) 652-657] ⇒ bottom quark discovered in 1977 [PRL 39, 252 (1977)].



1987: Observation of B⁰-B⁰ meson mixing (property for which a neutral meson (B⁰) transform (oscillates) into its antiparticle (B⁰) over time) [ARGUS collaboration, PLB 192, 1-2, 245-252].



Present: B anomalies

 \circ In recent years, some interesting set of tensions with the SM predictions have arisen:

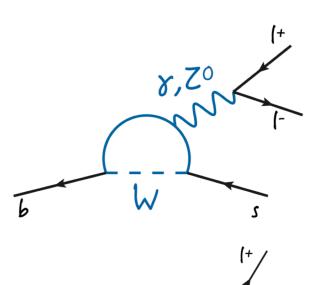
1. In $b \rightarrow s\ell^+\ell^-$ transitions (loop, FCNC):

- Branching fractions of $b \rightarrow s \mu^+ \mu^-$ decays.
- Angular observables of $b \rightarrow s\mu^+\mu^-$ decays.
- Lepton Flavour Universality (LFU) tests in μ/e ratios.

2. In **b** \rightarrow **c** $\ell^+\nu_{\ell}$ transitions (tree, FCCC):

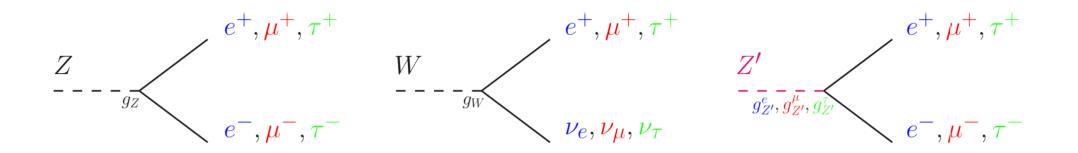
• Lepton Flavour Universality tests in τ/μ ratios.

 $_{\odot}$ In this talk: experimental situation.



Lepton Flavour Universality

- In the SM, couplings of the gauge bosons to charged leptons are universal (Lepton Flavour Universality, LFU).
- Branching fractions involving e, μ and τ differ only due to their different masses (phase space and helicity suppressions)
- Some extensions of the SM predict new particles that can break LFU, i.e. Z', W', leptoquarks ...

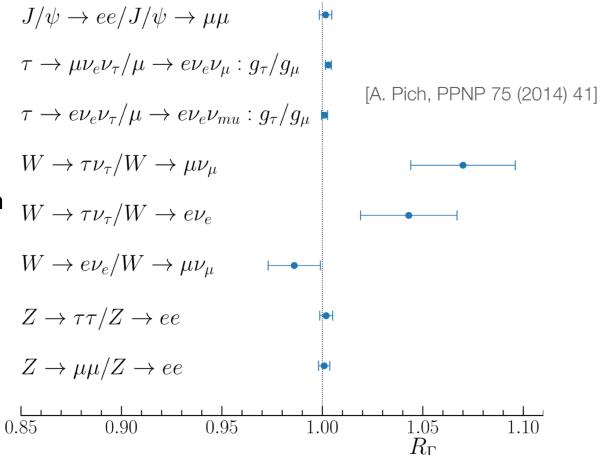


• Any significant deviation of LFU is a **clear sign of NP**.

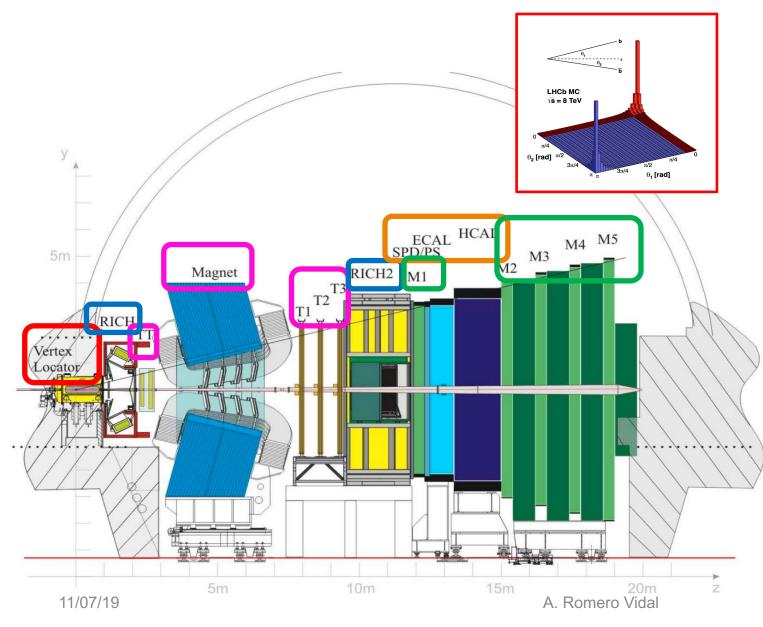
Testing LFU

- Thoroughly tested in the past:
 - Electroweak sector: $Z \rightarrow \ell^+ \ell^-$ and $W^+ \rightarrow \ell^+ \nu_\ell$.
 - **Pseudoscalar mesons**: semileptonic decays of *π*, K and D mesons.
 - **Purely leptonic decays**: $\tau^- \rightarrow \mu^- \nu_\mu \nu_\tau$ and $\tau^- \rightarrow e^- \nu_e \nu_\tau$.
 - Quarkonia decays: $J/\psi \rightarrow e^+e^-$, $J/\psi \rightarrow \mu^+\mu^-$.
- W coupling to τv_{τ} in tension with SM at **2.6** σ level (precision needs to be improved).
- In recent years, some interesting set of tensions with the SM predictions have arisen in the B sector:
 - In b→sℓ+ℓ⁻ transitions (R_K, R_{K*}, BR's, angular distributions ...)
 - In **b** \rightarrow **c** $\tau \nu_{\tau}$ transitions (R_D, R_{D*} and R_{J/ ψ})

[PDG, PRD98, 030001 (2018)]



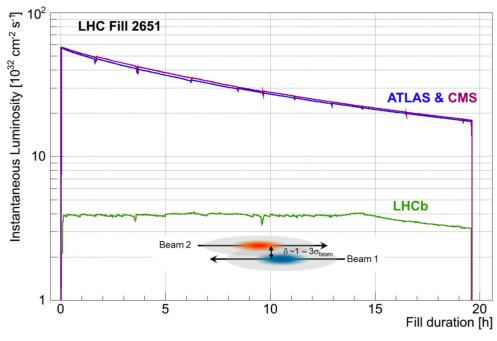
The LHCb detector

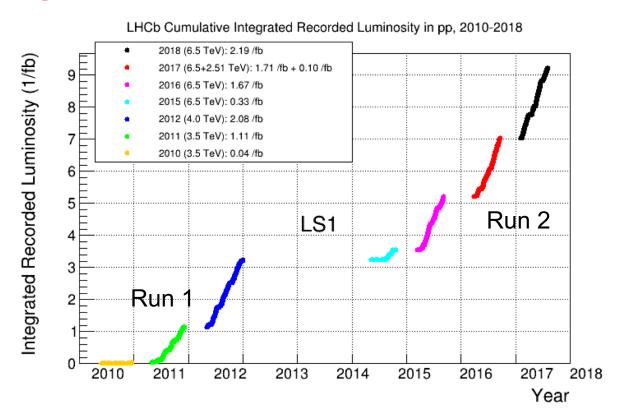


- In pp collisions b/b pairs produced with very small opening angle. LHCb detector is a forward spectrometer.
- Vertex detector (VELO):
 - Excellent vertex resolution: 20µm resolution on impact parameter.
 - Decay time resolution ~45ps.
- Tracking system (plus a 4T magnet):
 - Momentum resolution ⊿p/p~0.4%-0.6%
- **RICH detectors:**
 - Excellent K/ π /p separation.
- Calorimeter systems:
 - Energy measurement (i.e: π^0 , γ).
- Muon system:
 - Very high efficiency for muons.

Detector operation

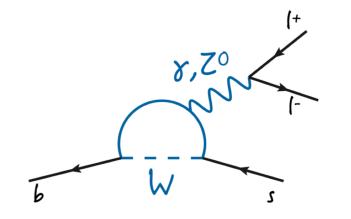
- LHCb designed to run at **lower instantaneous luminosity** \mathcal{L} than ATLAS and CMS.
- Mean number of interactions per bunch crossing ~1.
- pp **beams displaced** to reduce \mathcal{L} .





- **3 fb**⁻¹ of pp collisions at **7-8 TeV** in **Run 1** (2011+2012)
- 6 fb⁻¹ of pp collisions at **13 TeV** in **Run 2** (2015-2018)
- 9 fb⁻¹ in total at the end or Run 2.

Measurements of loop b→sℓ⁺ℓ⁻ transitions

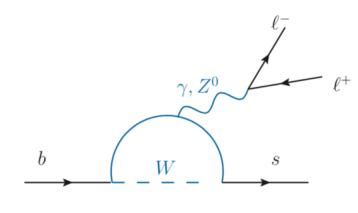


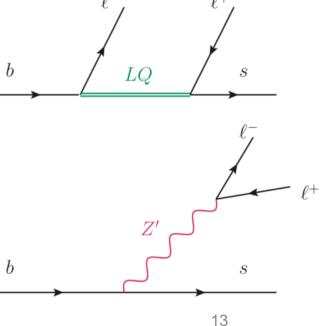
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Flavour Changing Neutral Currents: $b \rightarrow s \mu^+ \mu^-$

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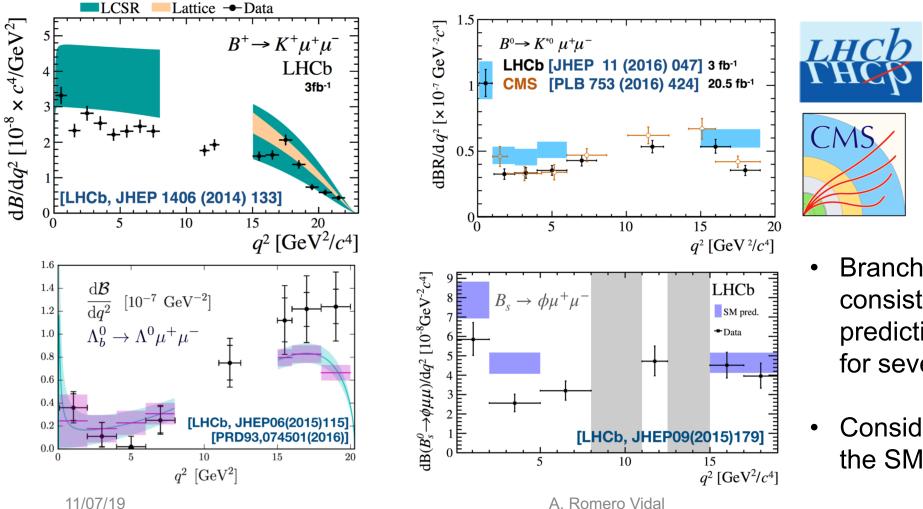
- Flavour Changing Neutral Currents (FCNC) suppressed in the SM:
 - only occur at loop level. 1.
 - 2. GIM suppressed.
 - left-handed chirality. 3.
- This is not necessarily true in a NP scenario.
- Important to study different kind of observables (BR's, angular observables and LFU tests) with different sensitivities to NP.





$b \rightarrow s \mu^+ \mu^-$ branching fractions

• Several measurements performed in different b-hadron decays: $B^+ \rightarrow K^+ \mu^+ \mu^-$, $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $\Lambda_b^0 \rightarrow \Lambda^0 \mu^+ \mu^-$, $B_s^0 \rightarrow \Phi \mu^+ \mu^-$, by different experiments.

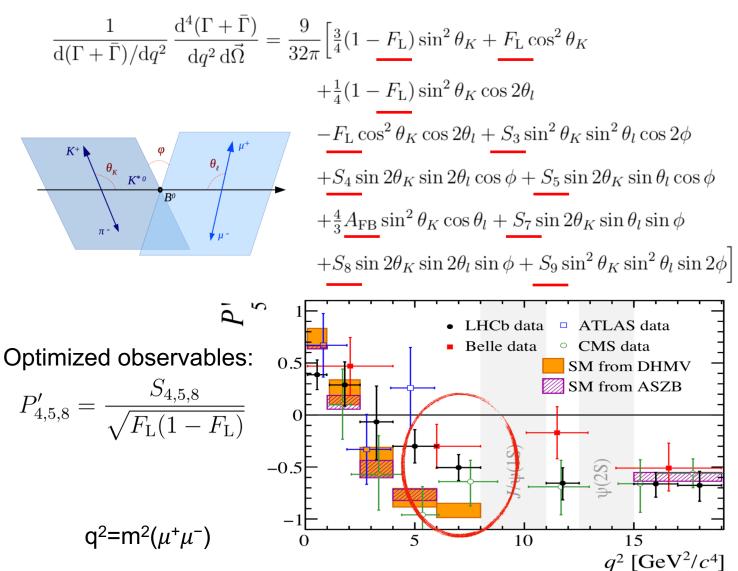


- Branching fractions consistently below the SM prediction at low q²=m²(l⁺l⁻) for several processes.
- Considerable uncertainty in the SM prediction.

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$B^0 \rightarrow K^{*0} \mu^+ \mu^-$ angular analysis

- B⁰→K^{*0}µ⁺µ⁻ is a b→sµ⁺µ⁻ transition: excellent place to search for NP in FCNC process.
- Rates, angular distributions and asymmetries sensitive to NP.
- Experimentally very clean signature.
- Decay described by three angles and dimuon invariant mass squared q² (with parameters depending on q²: F_L, A_{FB}, S_i).
- Many observables with clean theoretical predictions (i.e. P'_{4,5,8}: partial cancellation of hadronic uncertainties) [JHEP 1204(2012) 104].
- LHCb finds in P₅' a deviation from the SM prediction at the level of 3.4σ.



[JHEP 02 (2016) 104]

Theoretical framework: Effective Field Theory

C.(11) • Wilson coefficients (nerturbative short-

• Effective Field Theory (EFT) can describe $b \rightarrow s\ell^+\ell^-$ transitions in terms of an effective Hamiltonian that describes the full theory at low energies (μ).

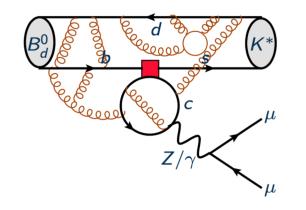
 Contributions from NP will modify the measured values of the Wilson coefficients present in the SM or introduce new operators.

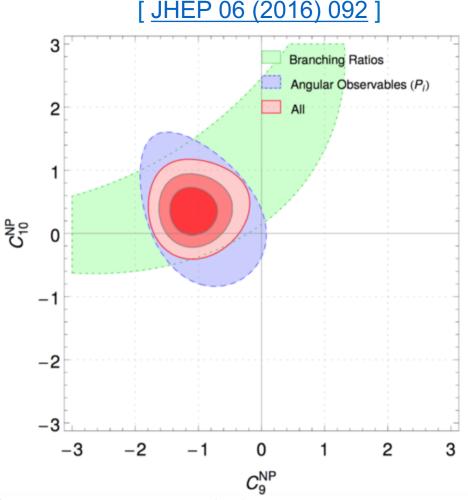
Global fits to $b \rightarrow s \mu^+ \mu^-$ observables

 A large number of measurements performed in b→sℓ⁺ℓ⁻ and b→sγ transitions are used as input in a global fit.

 $\mathcal{H}_{\text{eff}} \sim \sum_{i} C_i(\mu) \mathcal{O}_i(\mu)$

- Wilson coefficients split as C_i = C_iSM + C_i^{NP} with i=7,9,10,7',9',10'.
 Primed C_i': right handed currents, suppressed in the SM.
- Best fit prefers a non-zero value for the vector coupling C₉^{NP} (or C₉^{NP} and axial-vector C₁₀^{NP})
 - ... could QCD effects mimic vector-like NP contribution?





[W. Altmannshofer et al. Phys. Rev. D96 (2017) 055008,
B. Capdevila et al. JHEP 01 (2018) 093, T. Hurth et al. Phys. Rev. D96 (2017) 095034,
G. DAmico et al. JHEP 09 (2017) 010, L.-S. Geng et al. Phys. Rev. D96 (2017) 093006, M. Ciuchini et al. Eur. Phys. J. C77 (2017) 688,
S. Jäger and J. Martin Camalich, Phys. Rev. D93 (2016) 014028 and many others]

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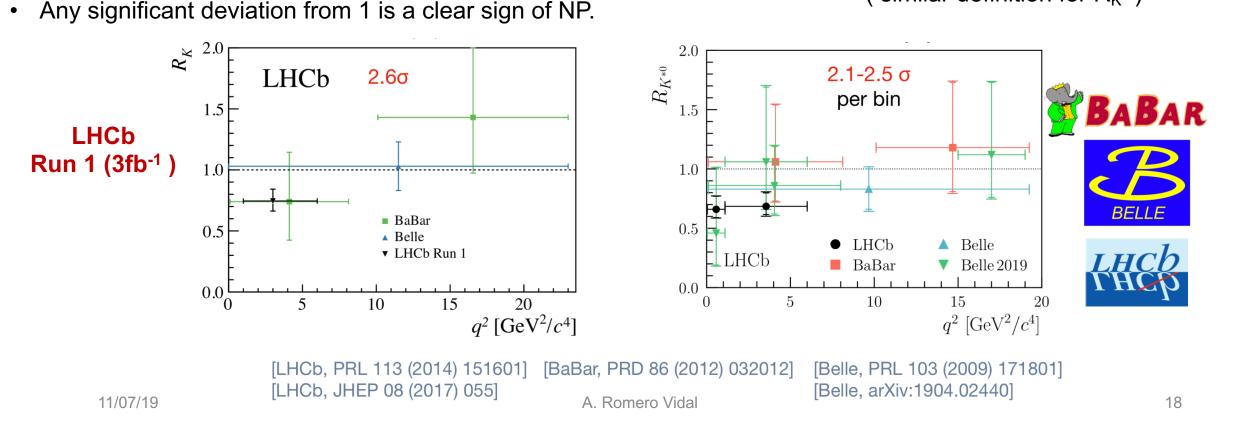
Lepton flavour universality in b \rightarrow s $\mu^+\mu^-$ decays

- Ratios of muons/electrons are **extremely well predicted** in the SM: ٠
 - Hadronic uncertainties of $O(10^{-4})$ ([<u>JHEP 0712 (2007) 040</u>]). ٠
 - QED uncertainties of $O(10^{-2})$ ([EPJC 76 (2016), 8, 440]). ٠

٠

 $R_K = \frac{BR(B^+ \to K^+ \mu^+ \mu^-)}{BR(B^+ \to K^+ e^+ e^-)} \stackrel{\text{SM}}{\cong} 1$

(similar definition for R_{K^*})



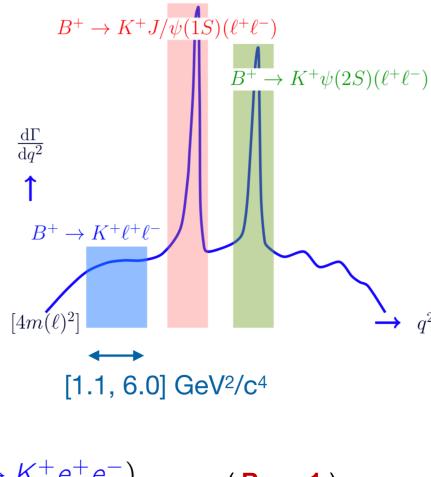
New measurement of R_{κ} at LHCb

- In total, this update uses ~twice as many B's as previous analysis.
 - Re-optimized analysis of **Run 1** dataset (**3fb**⁻¹). ٠
 - Added 2015 and 2016 datasets from Run 2 (2fb⁻¹).
- Details of the analysis:
 - Electrons and muons behave very differently in the LHCb ٠ detector due to larger Bremsstrahlung for electrons.
 - Worse mass and q² resolution.
 - Lower reconstruction efficiency.
 - Measurement perform ٠ $B^- \rightarrow K^- \ell^+ \ell^-$ and B^- systematics.

ement performed as a **double ratio** between

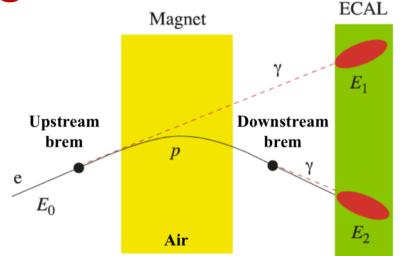
$$\ell^+\ell^-$$
 and $\mathbf{B}^- \to \mathbf{K}^- \mathbf{J}/\psi(\to \ell^+\ell^-)$ modes to cancel most
atics.

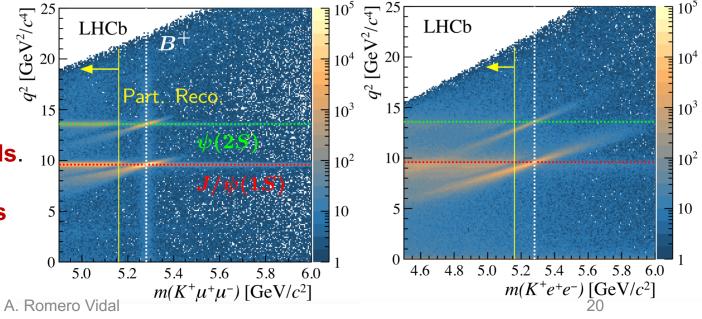
$$R_K = \frac{\mathcal{B}(B^+ \to K^+\mu^+\mu^-)}{\mathcal{B}(B^+ \to K^+J/\psi(\mu^+\mu^-))} / \frac{\mathcal{B}(B^+ \to K^+e^+e^-)}{\mathcal{B}(B^+ \to K^+J/\psi(e^+e^-))} \quad (\mathbf{R}_{J/\psi} = \mathbf{1})$$



Bremsstrahlung

- Electrons lose a fraction of their energy through Bremsstrahlung radiation.
- Bremsstrahlung **recovery procedure** to improve momentum measurement for electrons.
 - Look for photon clusters in the calorimeter (E_T>75 MeV) compatible with electron direction.
- Even after bremsstrahlung recovery, electrons still have degraded momentum and mass/q² resolution.
- Very different trigger signatures: lower trigger efficiency for electrons.
 - Muons identified by muon detectors.
 - Electrons rely on signal in the calorimeters: higher occupancy ⇒ higher trigger thresholds.
- **Critical aspect** of the analysis: get the **differences** between electrons and muons fully under control.





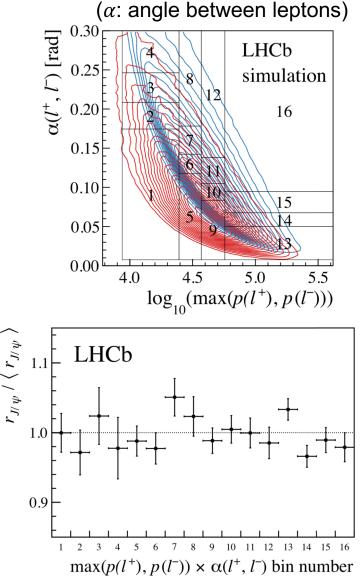
R_K systematics and cross-checks

- Efficiencies computed using simulation carefully calibrated using control channels selected from data:
 - Applied corrections due to B⁺ kinematics, particle-ID, trigger efficiency...
 - Small systematic due to good cancellation in double ratio.
- Numerous cross-checks to ensure good understanding of the efficiencies, i.e check:

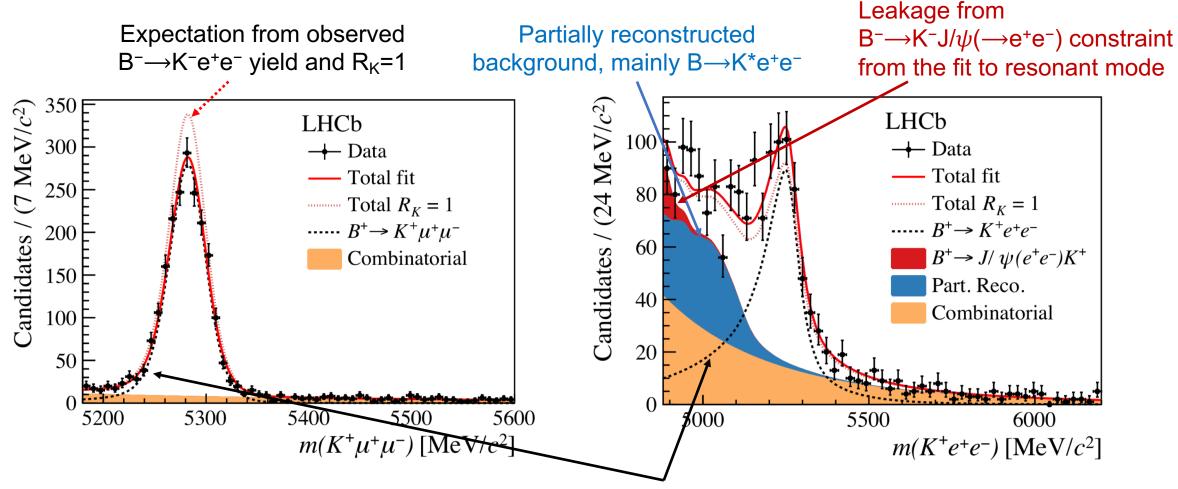
$$r_{J/\psi} = \frac{\mathcal{B}(B^+ \to K^+ J/\psi(\mu^+ \mu^-))}{\mathcal{B}(B^+ \to K^+ J/\psi(e^+ e^-))} = 1$$

$$r_{J/\psi} = 1.014 \pm 0.035 \,(\text{stat} + \text{syst})$$

- Checked that efficiencies are understood in all kinematic regions \Rightarrow $R_{J/\psi}$ flat for all variables examined.
- Cross-checks done independently for Run 1 and Run 2 samples and excellent agreement found.

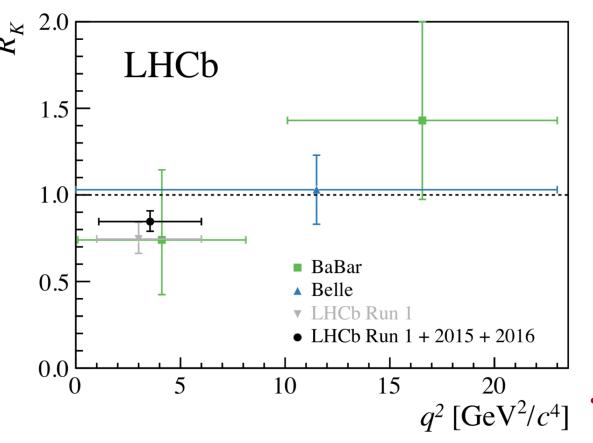


R_{K} simultaneous fit



- Different signal shape between muons and electrons (different x-scales!):
 - worse mass resolution (recovered photons)
 - longer radiative tails (bremsstrahlung)

New R_K result



• Using Run 1 (2011+2012) data:

 $R_{K} = 0.745 \, {}^{+0.090}_{-0.074} \, (\text{stat}) \, \pm 0.036 \, (\text{syst})$

compatible with SM expectation at 2.6σ .

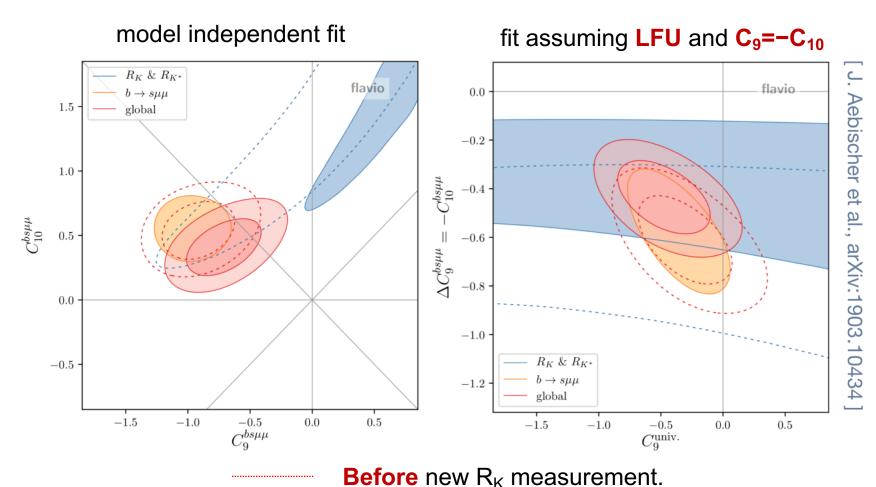
Re-analysing 2011+2012 data and **adding** 2015+2016:

$$R_{K} = 0.846 \, {}^{+0.060}_{-0.054} \, (\text{stat}) \, {}^{+0.014}_{-0.016} \, (\text{syst})$$

compatible with SM expectation at 2.5σ .

- Systematic small due to good cancellation in double ratio:
 - Uncertainty on the model shape.
 - Calibration of B⁺ kinematics and trigger efficiency.

Impact on global fits



After new R_K measurement.

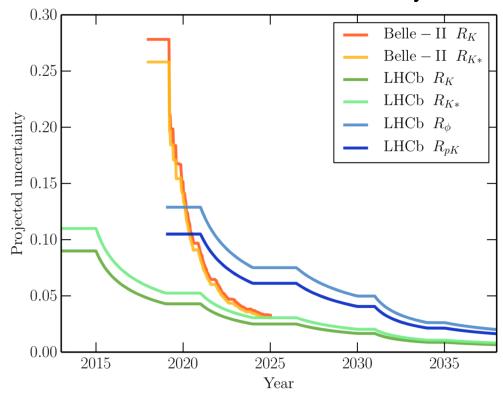
- Best fit point still in tension with the SM.
- Worse compatibility between $R_{K}(R_{K^*})$ and other $b \rightarrow s\mu^+\mu^-$ observables.
- Muonic NP: best fit closer to the SM, C₉=-C₁₀ still preferred.
- Adding LFU NP: slight preference for universal shift in C₉.

[M. Algueró et al., arXiv:1903.09578, A. K. Alok et al., arXiv:1903.09617, M. Ciuchini et al., arXiv:1903.09632, Guido D'Amico et al., arXiv:1704.05438, and more]

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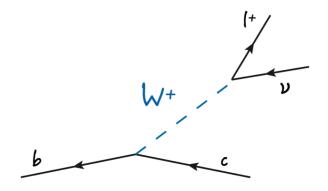
Prospects on b \rightarrow s $\ell^+\ell^-$

- LHCb full Run 2 dataset ~4 times number of B's available in Run 1.
 - Updates of R_K and R_{K^*} and other LFU ratios: R_{ϕ} , R_{pK} ...
 - Angular analyses of $b \rightarrow s\ell^+\ell^-$ transitions also underway
- CMS has collected a sample of ~10¹⁰ B decays.
 - With an effective low p_T electron reconstruction, should get a very competitive number of B⁺→K⁺e⁺e⁻ candidates.
 - Expected **systematics** will be **different** to those at LHCb, i.e. no trigger effect and very different material distribution.
- **ATLAS** pursuing a similar strategy.
- Belle II already started data-taking this year.



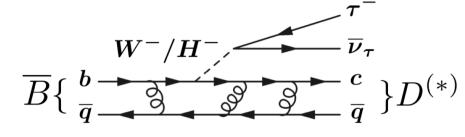
Evolution of uncertainties vs year

Measurements of tree b $\rightarrow c\tau v_{\tau}$ transitions



LFU in semitauonic B decays

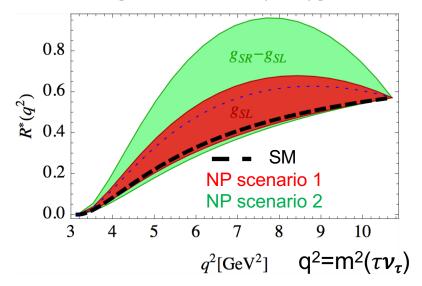
- In some NP scenarios, new particles couple preferentially to the third family ⇒ important to study semileptonic B decays into this family (τ).
- Comparison between semitauonic (τ) and semimuonic (μ) decays sensitive to NP, which could modify branching fractions and angular distributions.



• Perform LFU test through:
$$R(D^{(*)}) = \frac{\mathcal{B}(B^0 \to D^{(*)}\tau\nu)}{\mathcal{B}(B^0 \to D^{(*)}\mu\nu)}$$

- R(D^(*)) very clean SM prediction due to partial cancellation of hadronic form-factors uncertainties in the ratio. Deviation from 1 due to the different lepton masses (i.e. phase space...)
- Observation of violation of LFU would be a sign for NP.

[PRD 85 094025 (2012)]



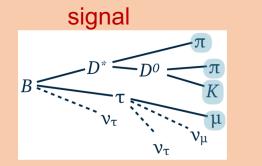
	R(D) _{SM}	R(D*) _{SM}
PRD94 (2016) 9, 094008	0.299 ± 0.003	
PRD95 (2017) 11, 115008	0.299 ± 0.003	0.257 ± 0.003
JHEP 1711 (2017) 061		0.260 ± 0.008
JHEP 1712 (2017) 060	0.299 ± 0.004	0.257 ± 0.005

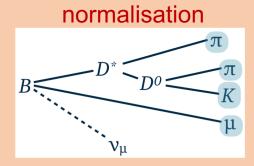
Reconstruction of tau decays

Leptonic decays

Hadronic decays

Signal and normalisation channels share the same visible final state.





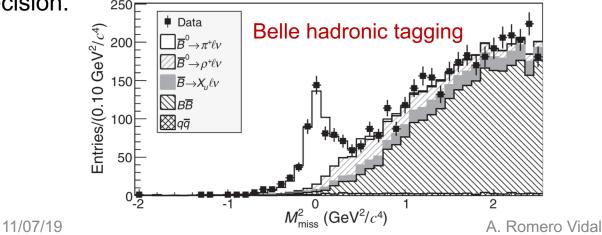
- Part of the systematics cancel in ratio.
- Important background from inclusive semileptonic decays (B→D**ℓν(X)), with many unknowns (form-factors, BR's...)

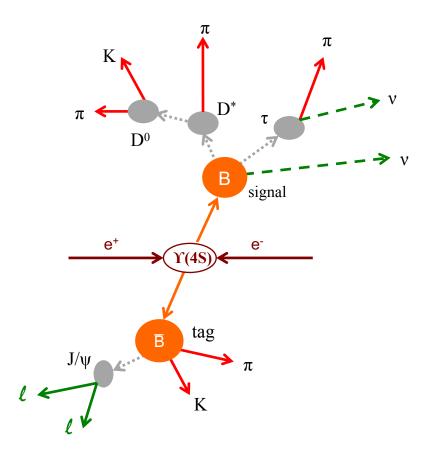
- No background from inclusive semileptonic decays.
- Visible final state is not the same. Systematics (at LHCb) do not cancel in the ratio ⇒ measure with respect to other decay with similar final state.

tau decay	BR(%)	Used by	
leptonic			
$\tau^- \longrightarrow \mu^- \nu_\mu \nu_\tau$	17.39 ± 0.04	B-factories and LHCb	
$\tau^- \rightarrow e^- \nu_e \nu_\tau$	17.82 ± 0.04	B-factories	
hadronic			
$\tau^- \longrightarrow \pi^- \pi^0 \nu_{\tau}$	25.49 ± 0.09	Belle	
$\tau^- \longrightarrow \pi^- \nu_{\tau}$	10.82 ± 0.05	Belle	
$\tau^- \longrightarrow \pi^- \pi^+ \pi^- \nu_{\tau}$	9.02 ± 0.05	LHCb	
$\tau^- \longrightarrow \pi^- \pi^+ \pi^- \pi^0 \nu_{\tau}$	4.49 ± 0.05	LHCb	

Semitauonic B decays at B-factories

- At B-factories (Belle(II) and BaBar), e⁺/e⁻ collisions producing Y(4S)→BB. e⁺ and e⁻ with different energies to produce boosted B's.
- **B-tagging** allows to constrain the momentum of the **B-signal**:
 - Hadronic B-tag: precise measurement of p_{B,sig}. Good determination of q² and m_{miss}² (eff. ~0.3%)
 - SL B-tag: weaker constraint on p_B (eff. ~1%)
- The missing mass (neutrinos) can be measured with high precision.





B-factories vs LHCb

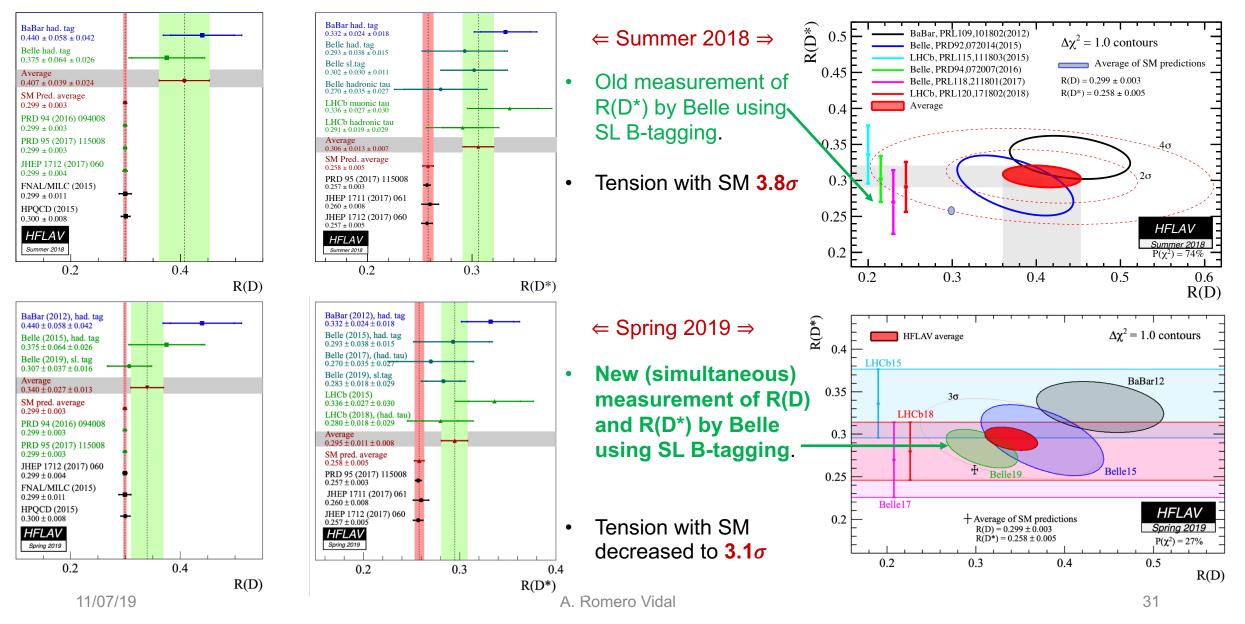
B-factories

- B momentum is known ($\Upsilon(4S) \rightarrow B\overline{B}$)
- B-tag algorithms use the other B in the event:
 - Hadronic B-tag: 0.3% efficient, very pure: all backgrounds are fully reconstructed.
 - **SL** B-tag: 1% efficient, **less pure**.
- Charm mesons reconstructed in multiple decay modes:
 - $D^{*-} \rightarrow D^0 \pi^-$, $D^+ \pi^0$.
 - $D^{*0} \rightarrow D^0 \pi^0$.
 - $D^0 \rightarrow K^- \pi^+ \pi^0$, $K^- \pi^+ \pi^- \pi^-$, $K^- \pi^+$, $K_S^0 \pi^+ \pi^-$, $K_S^0 \pi^0$, $K_S^0 K^+ K^-$, $K^+ K^-$, $\pi^+ \pi^-$. (30% of D⁰ BR's)
 - $D^+ \rightarrow K^- \pi^+ \pi^-$, $K_S^0 \pi^+ \pi^0$, $K_S^0 \pi^+ \pi^- \pi^+$, $K_S^0 \pi^+$, K⁻K⁺ π^+ , K_S^0 K⁺. (22% of D⁺ BR's)

LHCb

- Use the **B flight direction** to measure transverse component of missing momentum.
- Cannot measure longitudinal component, so use some approximation to access rest frame variables.
- Until now only used $D^{*-} \rightarrow D^0 \pi^-$ with $D^0 \rightarrow K^- \pi^+$ for $R(D^*)$ measurements.
- Studies using A_b⁰, B_c⁺ and B_s⁰ hadrons only possible at the LHC (not produced at B-factories).

Status of R(D) and R(D*) measurements



Semitauonic measurements at LHCb

Muonic tau decays

- Measurement using $\tau^- \rightarrow \mu^- \nu_\mu \nu_\tau$ decays.
- **3 neutrinos** present in $B^0 \rightarrow D^* \tau^+ \nu_{\tau}$
- Important **contributions** from exclusive $B^0 \rightarrow D^{*-}$ $\mu^+\nu_{\mu}$ (normalisation) and inclusive $B \rightarrow D^{*-}\mu^+\nu_{\mu}X$ decays, i.e. $(D^{**} \rightarrow D^*\pi\pi)$.
- Important contribution from **doubly-charmed** $B \rightarrow D^{*-}D(X)$ decays with $D \rightarrow X \mu^{+} \nu_{\mu}$.
- R(D*) directly measured from the same data sample (D*-μ*).

Published analyses use only Run 1 data (3fb⁻¹)

Hadronic tau decays

- Measurement using $\tau^- \rightarrow \pi^- \pi^+ \pi^- (\pi^0) \nu_{\tau}$ decays.
- Only 2 neutrinos present in $B^0 \rightarrow D^{*-} \tau^+ \nu_{\tau}$
- tau vertex is known (3π vertex). This gives access to the tau decay time, which can be used to discriminate between signal and background.
- No contamination from semileptonic $B \rightarrow D^{*-}\mu^+\nu_{\mu}(X)$ decays.
- Large contribution from doubly-charmed $B \rightarrow D^{*-}D(X)$ decays with $D \rightarrow \pi^{+}\pi^{-}\pi^{+}X$.
- **Measure BR(B⁰ \rightarrow D^{*-}\tau^{+}\nu_{\tau})** with respect to, i.e. B⁰ $\rightarrow D^{*-}\pi^{+}\pi^{-}\pi^{+}$. Then, use WA BR(B⁰ $\rightarrow D^{*-}\mu^{+}\nu_{\mu})$ to obtain R(D*).

R(D*) at LHCb using $\tau^- \rightarrow \mu^- \nu_\mu \nu_\tau$ decays

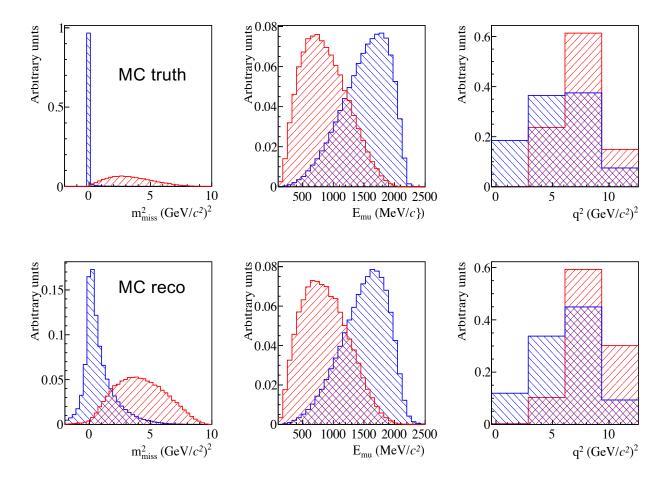
 At the LHC, B momentum cannot be constrained from beam energy (no Y(4S)→BB decay as in B-factories).

$$p \xrightarrow{z \xrightarrow{p_v}} p^{v}$$

 Approximation used to estimate B momentum: B boost along z >> boost of decay products in B rest frame:

$$(\gamma \beta_z)_B = (\gamma \beta)_{D^* \mu} \Longrightarrow (p_z)_B = \frac{m_B}{m(D^* \mu)} (p_z)_{D^* \mu}$$

- ~18% resolution on p_B still good enough to preserve signal and background discrimination.
- 3D template fit to m_{miss}^2 , E_{μ}^* and q^2 (= $m^2(\tau^-\nu_{\tau})$).



Muonic R(D*): Fit model

- $B^0 \rightarrow D^{*-} \mu^+ \nu_{\mu}$. Normalisation mode.
- $B^0 \rightarrow D^{*-} \tau^+ \nu_{\tau}$. Signal.
- Inclusive semileptonic $B \rightarrow D^{*-}\mu^+\nu_{\mu}X$ (and $B \rightarrow D^{*-} \tau^+ \nu_{\tau} X$).
- **Doubly-charmed** $B \rightarrow D^{*-}D(\rightarrow X\mu^{+}\nu_{\mu})X'.$
- Combinatorial $D^0\pi^-$ and $D^{*-}\mu^+$ • background.
- $h \rightarrow \mu$ mis-identification.

Control samples fits to constrain background shapes

- Control samples created adding extra **pions** and **kaons** pointing to B vertex.
- LHCb LHCb MeV 12000 Ce⁷ LHCb_ GeV²/c⁴ 1400 1600 5 1200 1400E 3.25 1200 1000 8000 1000E 6000 800F 600 4000 2000 $\frac{6}{m_{miss}^2} \frac{8}{(\text{GeV}^2/\text{c}^4)}$ 500 1000 1500 2000 250 E_u* (MeV) 0 2 10 2 6 $10 12 q^2 (GeV^2/c^4)$ LHCb 'LHCb MeV 20000 Data 4000E § 18000 2000 $B \rightarrow D^* \tau \nu$ 3500E $B \rightarrow D^*H_c (\rightarrow h_V X') X$ 1500 3000 2000 $B \rightarrow D^{**}hv$ Candidat 10000 1000 $B \rightarrow D^* uv$ 8000 6000 Combinatorial 500 4000 Misidentified u 2000 $\frac{6}{m_{miss}^2} \frac{8}{(\text{GeV}^2/\text{c}^4)}$ 1500 2000 250 E_u* (MeV) ${\begin{array}{*{20}c} 8 & 10 & 12 \\ q^2 \, (\text{GeV}^2/c^4) \end{array}}$ 0 500 1000 0 2 4 2 4 6 34

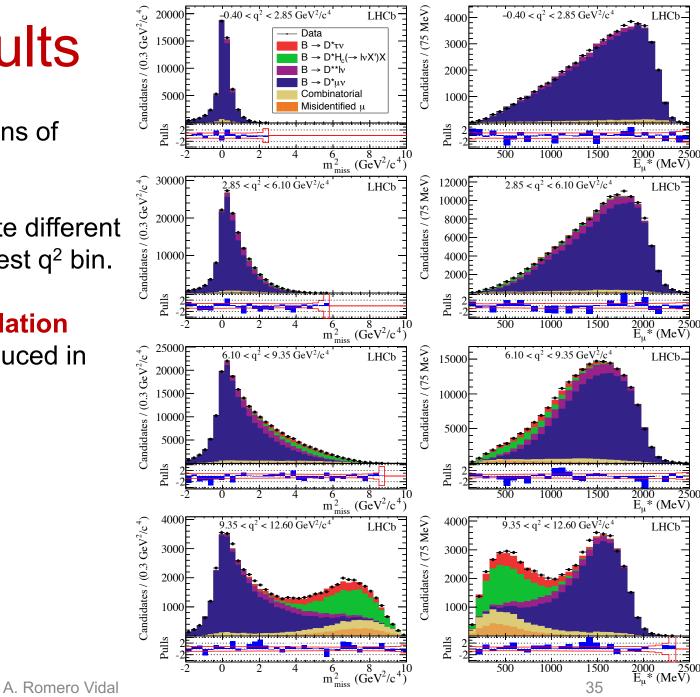
MC re-weighted to match data.

Muonic R(D*): results

- Projections of (left) m²_{miss} and (right) E_μ* in bins of increasing q² from top to bottom.
- Signal clearly visible in highest q² bin. Note different y scales, most signal actually in second-highest q² bin.
- Systematics dominated by the size of simulation sample and h→µ mis-ID. Expected to be reduced in future analyses.
- 3D template fit to m_{miss}^2 , E_{μ}^* and q^2 gives:

R(D*) = 0.336 ± 0.027(stat) ± 0.030(syst)

• Result is 2.1σ above SM (R(D^{*})_{SM} \approx 0.26).



$R(J/\psi)$ at LHCb using $\tau^- \rightarrow \mu^- \nu_\mu \nu_\tau$ decays

• LFU test through the measurement of: R(J/

$$(\psi) = \frac{\mathcal{B}(B_c^+ \to J/\psi \tau \nu)}{\mathcal{B}(B_c^+ \to J/\psi \mu \nu)}$$

- This measurement only possible at the LHC (B_c⁺ not produced at B-factories).
- B_c⁺ form-factors not known precisely ⇒ Theoretical prediction not precise. R(J/ψ) predicted to be in range [0.25,0.28] ([PLB452 (1999) 129], [arXiv:hep-ph/0211021], [PRD73 (2006) 054024], [PRD74 (2006) 074008])
- Improvements in form-factors calculation needed.

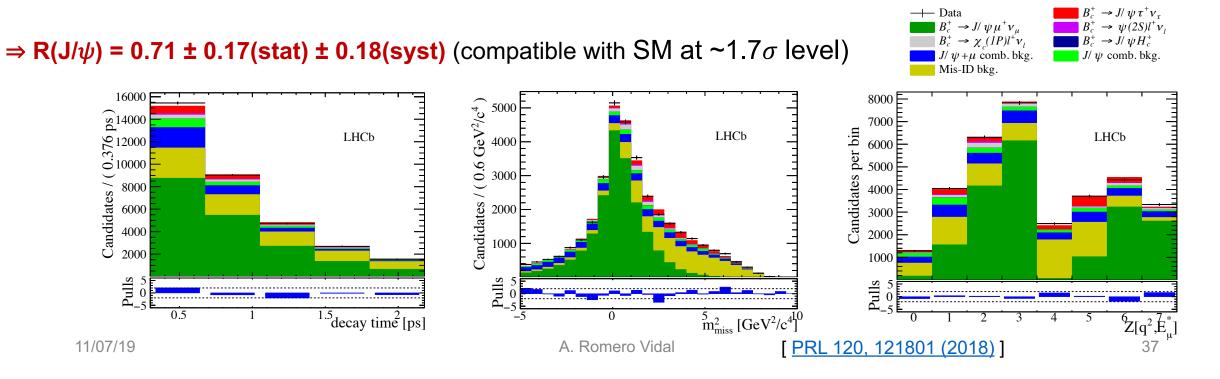


- Main backgrounds:
 - $B_c^+ \rightarrow J/\psi \mu^+ \nu_\mu$ (normalisation mode), $B_c^+ \rightarrow \psi(2S) \mu^+ \nu_\mu$ and $B_c^+ \rightarrow J/\psi D(\rightarrow X \mu^+ \nu_\mu) X'$.
 - Hadron mis-identified as a muon.
 - Combinatorial background (J/ ψ and μ not from same B)

$$(p_{B_c})_z = \frac{m_{B_c}}{m(J/\psi\mu)} \times (p_{J/\psi\mu})_z$$

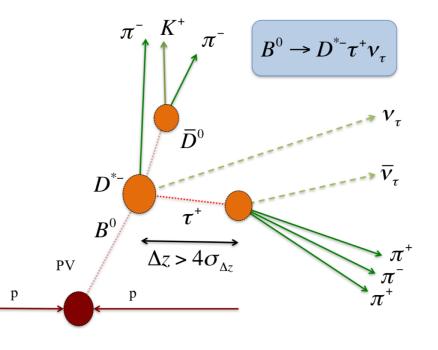
$R(J/\psi)$ results

- B_c⁺ decay time used in addition to q², m_{miss}² and E_μ^{*} in fit model (q² and E_μ^{*} combined in a single variable, Z) in a 3D template fit.
- Form-factors constrained from a control sample enriched in normalisation decays.
- Systematics dominated by the knowledge of the form-factors and size of simulation samples.
- **First evidence** of the decay $B_c^+ \rightarrow J/\psi \tau^+ \nu_\tau$ (**3** σ).

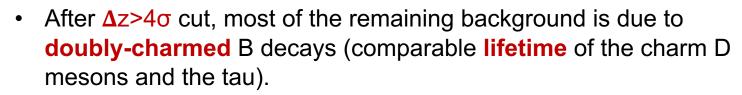


 $\mathbf{R}(D^{*}) \text{ using hadronic } \tau^{-} \longrightarrow \pi^{-}\pi^{+}\pi^{-}\nu_{\tau} \text{ decays}$ $\mathcal{R}(D^{*}) = \left(\frac{\mathcal{B}(\overline{B}^{0} \to D^{*+}\tau^{-}\overline{\nu}_{\tau})}{\mathcal{B}(\overline{B}^{0} \to D^{*+}\pi^{-}\pi^{+}\pi^{-})}\right)_{\text{meas}} \times \left(\frac{\mathcal{B}(\overline{B}^{0} \to D^{*+}\pi^{-}\pi^{+}\pi^{-})}{\mathcal{B}(\overline{B}^{0} \to D^{*+}\mu^{-}\overline{\nu}_{\mu})}\right)_{\text{external}}^{-4\% \text{ precision}}$ $\underset{\text{external}}{\overset{\mathcal{K}(D^{*})}{\longrightarrow}}$

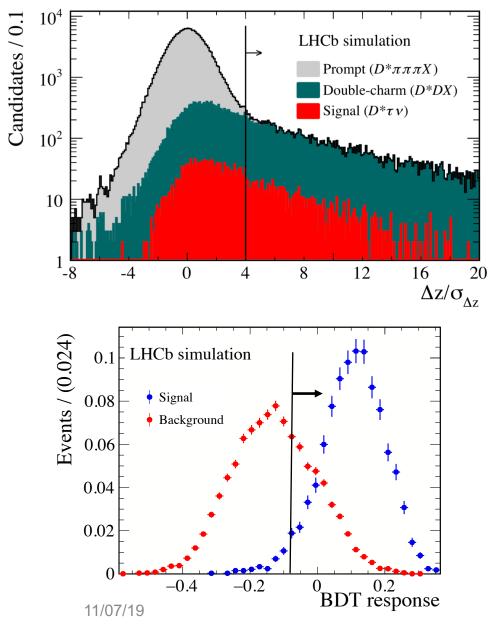
- Signal and normalisation share same visible final state $(D^{*-}\pi^{+}\pi^{-}\pi^{+})$.
- Most of the systematic uncertainties cancel in the ratio (PID, trigger ...).
- Use **topology** of the decay to suppress large $B \rightarrow D^{*-}\pi^{+}\pi^{-}\pi^{+}X$ ("prompt") background (where the 3 pions come from the B vertex).
- Minimum distance Δz between B⁰ and τ vertices >4σ: 35% efficient on signal and reduces prompt background by a factor >100.
- Possible due to the **excellent LHCb vertex resolution**.



Selection



- B→D*⁻D_s⁺X : ~ 10 x signal
- B→D*-D+X : ~ 1 x signal
- B→D*⁻D⁰X : ~0.2 x signal
- Largest background due to $B \rightarrow D^{*-}D_s^+X$ decays with inclusive $D_s^+ \rightarrow \pi^+\pi^-\pi^+X'$ decays.
 - Use a multivariate analysis algorithm (Boost Decision Tree, BDT) to suppress them.
 - Remaining background still large: modelling of the D_s⁺ (and D⁰ and D⁺) using data control samples.
- BDT trained with background MC vs signal MC, including:
 - 3π dynamics (i.e. m($\pi^+\pi^-$) ...),
 - $D^{*-}3\pi$ dynamics (i.e m($D^{*-}3\pi$) ...),
 - Deposit energy in the electromagnetic calorimeter.
- **BDT cut** applied to suppress main $B \rightarrow D^{*-}D_{s}^{+}X$ background.



Signal reconstruction

 B_0

- Due to the presence of 2 neutrinos in the final state: ٠
 - **2 solutions for** $|\mathbf{p}_{\tau}|$ due to the missing neutrino $|\vec{p}_{\tau}| = \frac{(m_{3\pi}^2 + m_{\tau}^2)|\vec{p}_{3\pi}|\cos\theta \pm E_{3\pi}\sqrt{(m_{\tau}^2 m_{3\pi}^2)^2 4m_{\tau}^2|\vec{p}_{3\pi}|^2\sin^2\theta}}{2(E_{3\pi}^2 |\vec{p}_{3\pi}|^2\cos^2\theta)}$ ٠ from the tau decay.
 - **2 solutions for |\mathbf{p}_{\mathsf{B}}|** due to the missing neutrino $|\vec{p}_{B^0}| = \frac{(m_{D*\tau}^2 + m_{B^0}^2)|\vec{p}_{D*\tau}|\cos\theta' \pm E_{D*\tau}\sqrt{(m_{B^0}^2 m_{D*\tau}^2)^2 4m_{B^0}^2|\vec{p}_{D*\tau}|^2\sin^2\theta'}}{from the B decay}$ ٠ from the B decay.

 \Rightarrow 4-fold ambiguity

Approximation: Set the argument of the 2 squared roots to zero:

$$\theta_{max} = \arcsin\left(\frac{m_{\tau}^2 - m_{3\pi}^2}{2m_{\tau}|\vec{p}_{3\pi}|}\right) \qquad \theta'_{max} = \arcsin\left(\frac{m_{B^0}^2 - m_{D*\tau}^2}{2m_{B^0}|\vec{p}_{D*\tau}|}\right)$$

- Possible to reconstruct rest frame variables such as tau decay time and q^2 .
- These variables have negligible biases, and sufficient resolution to preserve **good discrimination** between signal • and background.

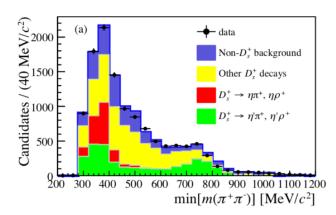
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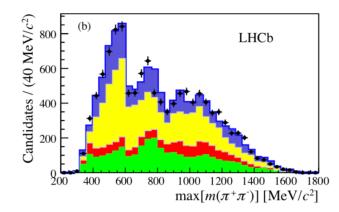
 $2(E_{D*\tau}^2 - |\vec{p}_{D*\tau}|^2 \cos^2 \theta')$

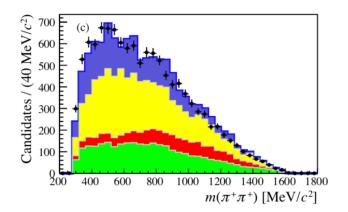
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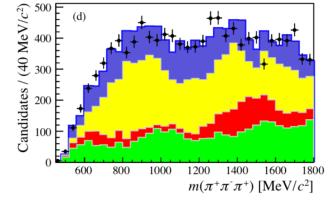
The $D_s^+ \rightarrow \pi^+ \pi^- \pi^+ X$ decay model: low-BDT fit

- Events removed by the BDT cut are used to model the inclusive $D_s^+ \rightarrow \pi^+ \pi^- \pi^+ X$ decay.
- Low BDT region (not used for signal extraction) is used to measure the D_s composition.
 - D_s decay modes with 3 pions + neutrals not very well measured.
 - $D_s \rightarrow 3\pi$ is only 1/15 of the inclusive $D_s \rightarrow 3\pi X$.
- Model obtained from **simultaneous fit** to:
 - Min[m($\pi^+\pi^-$)], Max[m($\pi^+\pi^-$)], m($\pi^+\pi^+$) and m(3π).
- Fit components:
 - ηπ⁺, ηρ⁺,
 - 2. $\eta' \pi^+, \eta' \rho^+,$
 - Other components including: ωπ⁺, ωρ⁺, φπ⁺, φρ⁺, K⁰3π, η3π, η'3π, ω3π, φ3π.
 - 4. Non-D_s component.
- Fit results are used to describe the D_s model at high BDT (signal sample).





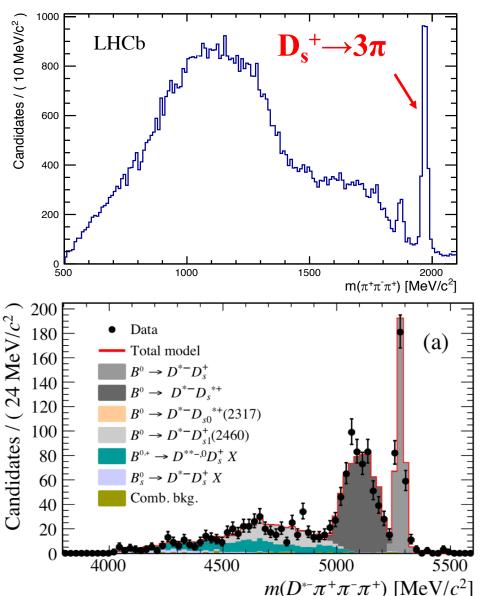




A. Romero Vidal [PRL 120, 171802 (2018), PRD 97, 072013 (2018)] 41

The $B \rightarrow D^{*-}DX$ control samples

- Different control samples are used to improve the description of background B→D*⁻DX components:
 - $D_s^+ \rightarrow 3\pi$
 - $D^0 \rightarrow K3\pi$ (kaon recovered by isolation tools)
 - $D^+ \rightarrow K^- \pi^+ \pi^+$ (mis-ID kaon/pion)
- Monte Carlo corrected using data-driven approach.
- A pure $B \rightarrow D^{*-}D_s X$ control sample obtained by selecting exclusive $D_s \rightarrow 3\pi$ decays.
- Allows to measure the different B→D*-D_s+X contributions from a fit to m(D*-D_s+).
- Uncertainties in the fit parameters **propagated** to final analysis.
- Similar strategy followed for $B \rightarrow D^{*-}D^{+}X$ and $B \rightarrow D^{*-}D^{0}X$ decays.



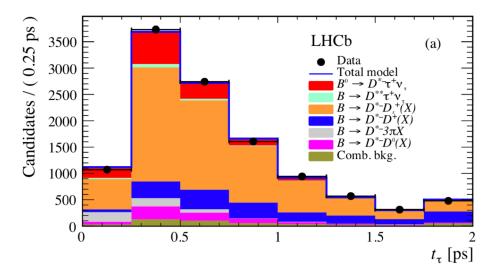
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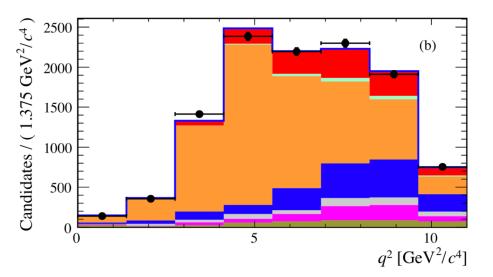
3D template fit

- 3D binned fit to data in (q^2 , τ decay time, BDT) with (8,8,4) bins.
- Model:
 - **1.** Signal $\mathbb{B}^0 \to \mathbb{D}^{*-}\tau^+\nu_{\tau}$: $\tau^- \to \pi^-\pi^+\pi^-\nu_{\tau}$ and $\tau^- \to \pi^-\pi^+\pi^-\pi^0\nu_{\tau}$ (ratio constrained according to known BF's). Free in the fit.
 - 2. $B \rightarrow D^{**} \tau v_{\tau}$. Fixed to the expected yield, 11% of signal (assign syst. uncertainty).
 - 3. Doubly-charmed B decays:
 - a) $B \rightarrow D^{*-}D_s^{+}X$. Includes $B^0 \rightarrow D^{*-}D_s^{+}$, $B^0 \rightarrow D^{*-}D_s^{*+}$, $B^0 \rightarrow D^{*-}D_{s0}^{*+}$, $B^0 \rightarrow D^{*-}D_{s1}$, $B_s^0 \rightarrow D^{*-}D_s^{+}X$, $B \rightarrow D^{**}D_s^{+}X$. Shape constrained from control sample.
 - b) $B \rightarrow D^{*-}D^{0}X$. Yield constrained from control sample.
 - c) $B \rightarrow D^{*-}D^{+}X$. Free in the fit.
 - 4. Prompt $B \rightarrow D^{*-}3\pi X$. Yield constrained from control sample.
 - 5. Comb. background. Yield constrained from control sample.

 $\Rightarrow \mathsf{BR}(\mathsf{B}^{0} \rightarrow \mathsf{D}^{*-} \tau v) = (1.42 \pm 0.094(\mathsf{stat}) \pm 0.129(\mathsf{syst}) \pm 0.054(\mathsf{ext}))\%$ $\Rightarrow \mathsf{R}(\mathsf{D}^{*}) = 0.291 \pm 0.019(\mathsf{stat}) \pm 0.026(\mathsf{syst}) \pm 0.013(\mathsf{ext})$

 1σ agreement with the SM (R(D^{*})_{SM} ≈ 0.26).

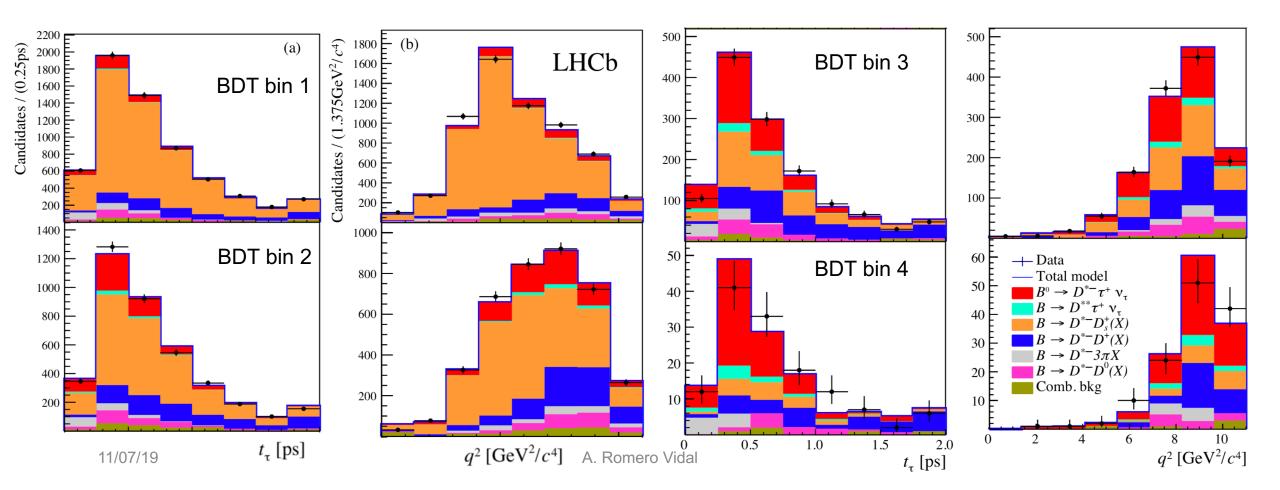




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Fit projections in BDT bins

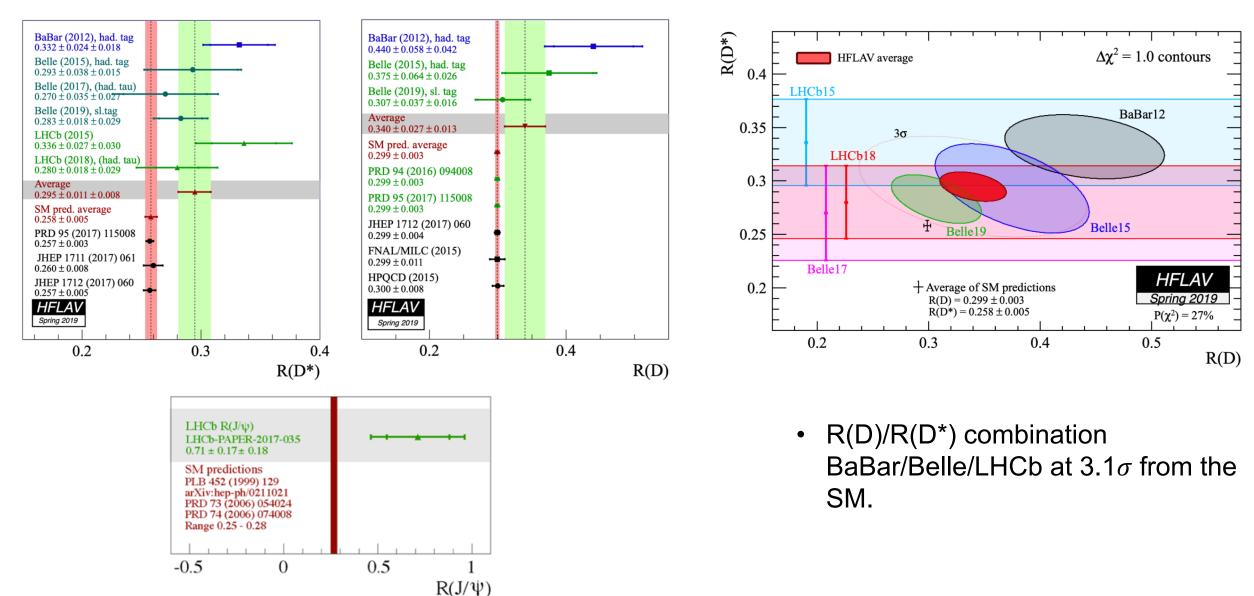
- Quality of the fit good in all bins of BDT.
 - Plots ordered from lowest BDT (bin 1) to highest BDT (bin 4)
 - High signal purity at high BDT.



Systematic uncertainties

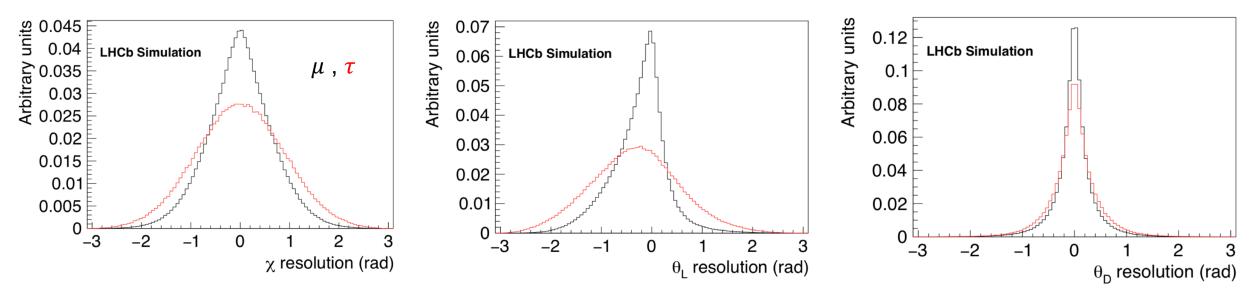
Source	$\delta R(D^{*-})/R(D^{*-})[\%]$		
Simulated sample size	4.7		
Empty bins in templates	1.3	These errors will be reduced with more data.	
Signal decay model	1.8		
$D^{**}\tau\nu$ and $D^{**}_s\tau\nu$ feeddowns	2.7		
$D_s^+ \to 3\pi X$ decay model	2.5		
$B \to D^{*-}D^+_s X, B \to D^{*-}D^+X, B \to D^{*-}D^0X$ backgrounds	3.9		
Combinatorial background	0.7	New external	
$B \to D^{*-} 3\pi X$ background	2.8	measurements can	
Efficiency ratio	3.9	help to reduce this	
Normalization channel efficiency (modeling of $B^0 \to D^{*-}3\pi$)	2.0	uncertainty!!!	
Total uncertainty	9.1		

Summary on R(X_c)



Future: Angular analyses

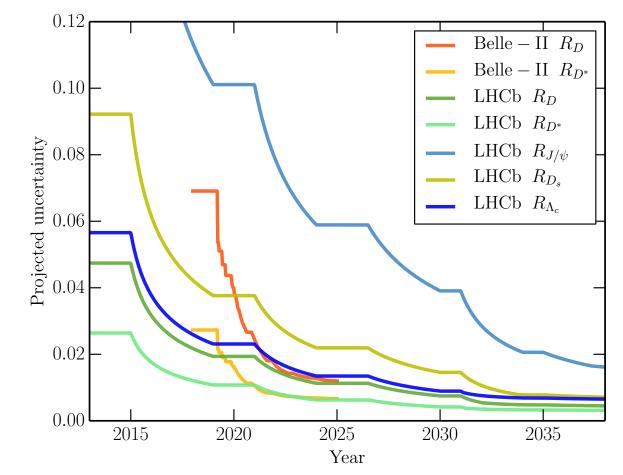
Resolutions using leptonic tau reconstruction:



- A future step is to look at the **angular distributions** in $B^0 \rightarrow D^{*-}\tau^+\nu_{\tau}$, very sensitive to NP contributions.
- Resolution in the D* **polarisation** angle good enough to measure D* polarisation.
- Resolution in *χ* and θ_L not so good, but some sensitivity is expected. This is compensated by the huge amount of statistics (to be) collected by LHCb.
- Prospects on the precision of angular analyses still to be studied.

LHCb prospects on R(X_c)

 Production fractions and efficiencies used to extrapolate the uncertainties.



- LHCb can perform tests of LFU not accessible at Belle II:
 - $R(\Lambda_{c}^{(*)}), R(J/\psi), R(D_{s}^{(*)}).$
- Precision in R(X_c) about 2-3% at the end of the Upgrade II.
- Sensitivity to angular observables need to be studied.

Conclusions

- Intriguing **anomalies** found in measurements of B-hadron decays:
 - loop-level $b \rightarrow s\ell^+\ell^-$ transitions.
 - tree-level $b \rightarrow c\ell^+ \nu_{\ell}$ transitions.
- Lepton Flavour Universality tests are theoretically clean probes for New Physics.
- \circ Precision in measurements still needs to improve to provide a definite picture.
- Upcoming measurements using Run 2 full statistics will help to resolve the current situation.
- Ongoing and future experiments upgrades and the start-up of Belle II open the door to many improvements in precision, so interesting times are ahead.

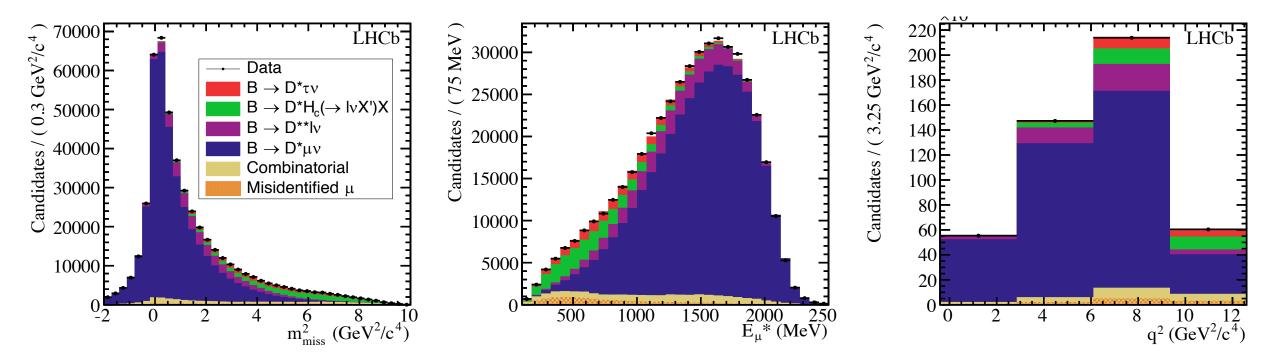


Future: R(D) vs R(D*) with hadronic tau decays

- Hadronic R(D*) measurement performed using $B^0 \rightarrow D^{*-}3\pi$ as normalisation mode.
- $B^0 \rightarrow D^{*-} 3\pi$ branching fraction measured with ~4% precision (PDG).
- For R(D) measurement ($B^- \rightarrow D^0 \tau^+ \nu_{\tau}$), an alternative normalisation mode needed:
 - $B^- \rightarrow D^0 3\pi$ poorly measured (37%)
 - $B^- \rightarrow D^0 D_s^+$ measured with 10% precision. In addition, $D_s^+ \rightarrow 3\pi$ measured with 5% precision.
 - Need improvements in these branching fractions: Belle II, BES III.
- In the LHCb Upgrade-I, the use of software trigger will allow to measure R(D^(*)) combining muonic and hadronic tau decays.

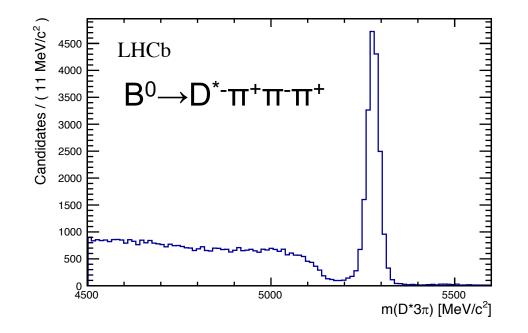
Decay	Branching fraction (PDG)	Precision
$B^0 \rightarrow D^- 3\pi$	(6.0 ± 0.7) x 10 ⁻³	11.7%
$B^0 \rightarrow D^- D_s^+$	(7.2 ± 0.8) x 10 ⁻³	11.1%
$B^0 \rightarrow D^{*-}3\pi$	(7.21 ± 0.29) x 10 ⁻³	4.0%
$B^0 \rightarrow D^{*-}D_s^+$	(8.0 ± 1.1) x 10 ⁻³	13.8%
$B^- \rightarrow D^0 3 \pi$	(5.6 ± 2.1) x 10 ⁻³	37.5%
$B^- \rightarrow D^0 D_s^+$	(9.0 ± 0.9) x 10 ⁻³	10.0%
$B^{-} \rightarrow D^{*0}3\pi$	(1.03 ± 0.12) x 10 ⁻³	11.6%
$B^{-} \rightarrow D^{*0}D_{s}^{+}$	(8.2 ± 1.7) x 10 ⁻³	20.7%
$D_s^+ \longrightarrow 3\pi$	(1.09 ± 0.05) x 10 ⁻²	4.6%

R(D*) muonic: fit projection



The normalisation mode

- Normalization channel as similar as possible to the signal $\rightarrow B^0 \rightarrow D^{*-}\pi^+\pi^-\pi^+$.
- This cancels production yield, BR uncertainties and systematics linked to trigger, PID and selection.
- In PDG 2014, BR(B⁰→D^{*-}π⁺π⁻π⁺) known with 11% precision.
- New BaBar measurement 4.3% (stat+syst) precision. [Phys. Rev. D94 (2016) 091101]
- In this analysis ~17000 events (1% precision)



Isolation

• Signal candidates are required to be well isolated.

K⁺

 D^0

р

D

Δz>4σ

В

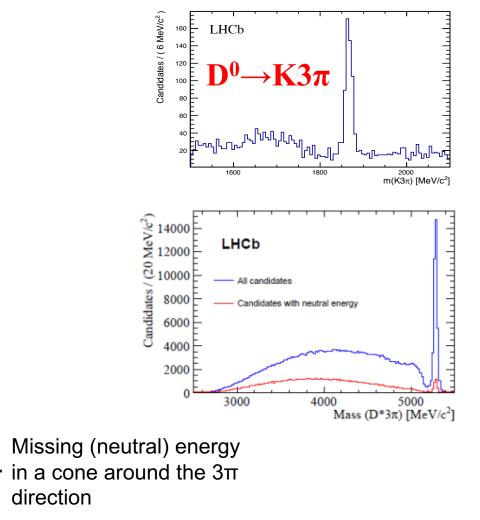
PV

Π

• Events with extra particles pointing to the B and/or tau vertices are vetoed.

Π

 $B^0 \rightarrow D^{*-}DX$



Ζ

р

y

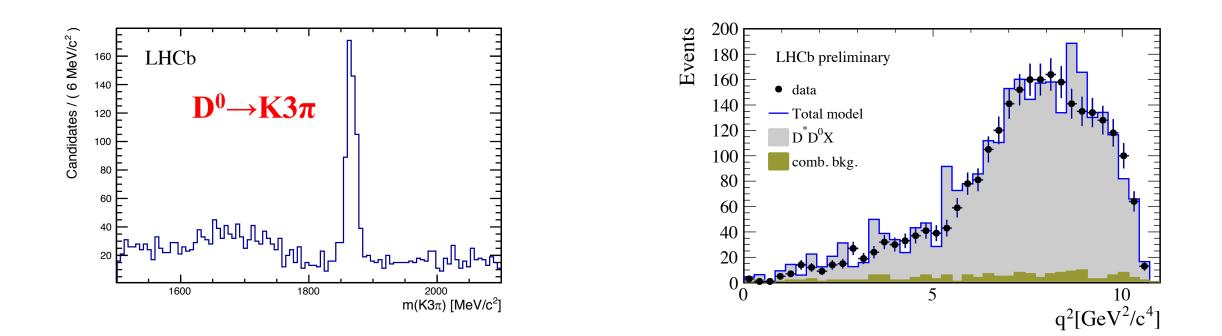
К,π

Κ, π

3π

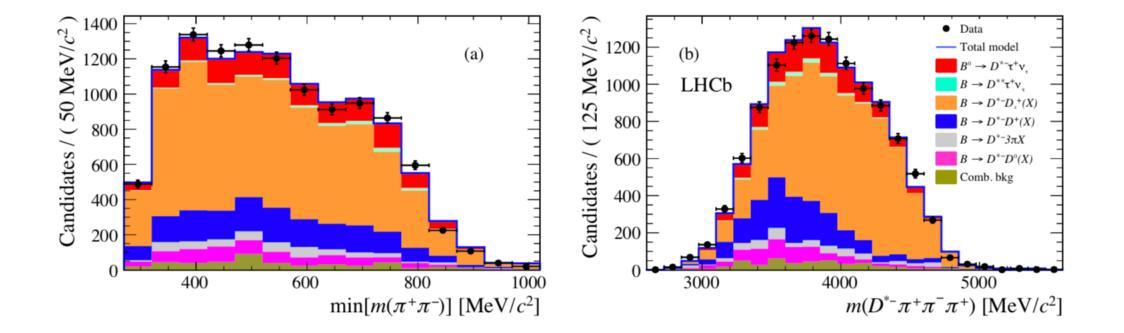
D⁰ control sample

- $X_b \rightarrow D^*D^0X$ decays can be isolated by selecting exclusive $D^0 \rightarrow K3\pi$ decays (kaon recovered using isolation tools).
- A correction to the q² distributions is applied to the Monte Carlo to match data.



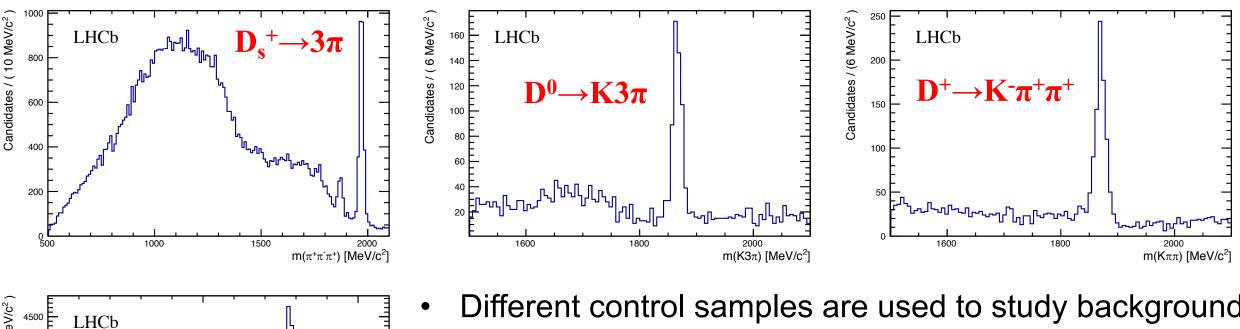
Fit projections on M(D*3 π) and min[M(2 π)]

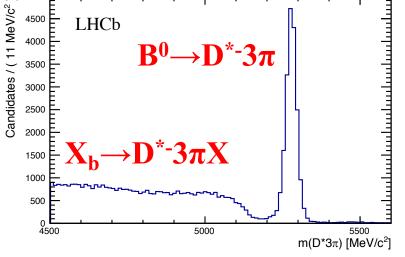
• Important variables in the BDT.



• Excellent agreement with data.

Control samples

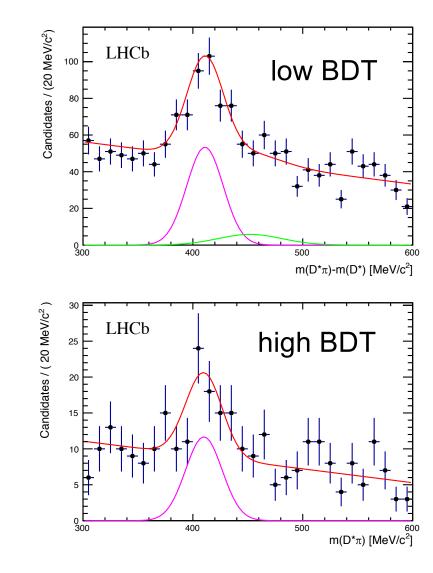




- Different control samples are used to study background components:
 - $D_s^+ \rightarrow 3\pi$
 - $D^0 \rightarrow K3\pi$ (kaon recovered by isolation tools)
 - $D^+ \rightarrow K^- \pi^+ \pi^+$ (mis-ID kaon/pion)
- Monte Carlo corrected using data-driven approach. ullet

Additional cross-checks: $B \rightarrow D^{**} \tau v$

- $B^0 \rightarrow D^{**} \tau v$ and $B^+ \rightarrow D^{**0} \tau v$ constitute potential feed-down to the signal.
- D**(2420)⁰ is reconstructed using its decay to D*+π⁻ as a crosscheck.
- The observation of the D**(2420)⁰ peak allows to compute the D** 3π BDT distribution and to deduce a D** τv upper limit. This upper limit is consistent with the theoretical prediction.
- Subtraction in the signal of 0.11±0.04 due to D**τν events leading to a systematic uncertainty of 2.3%.



Belle II prospects on R(D) and R(D*) [arXiv:1808.10567]

- Improve the precision on R(D) and R(D*) to the 2-3% level.
- Better control on backgrounds like $B \rightarrow D^{**} \ell \nu$, very important for these measurements.

