

Neutrino-less Double Beta Decays

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Sep. 5-10, 2021 @Zoom

22nd edition

PANIC Lisbon Portugal

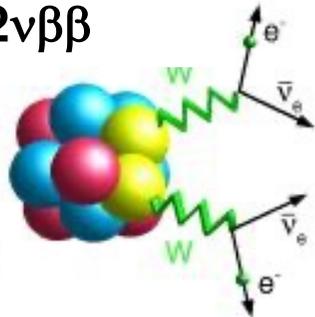
Particles and Nuclei International Conference



Double Beta Decay: SM and BSM

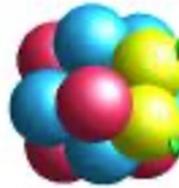
SM allowed process:

$2\nu\beta\beta$



- $2n \rightarrow 2p$ and $Z \rightarrow Z+2$
- hidden behind single β -decays

BSM
add \rightarrow

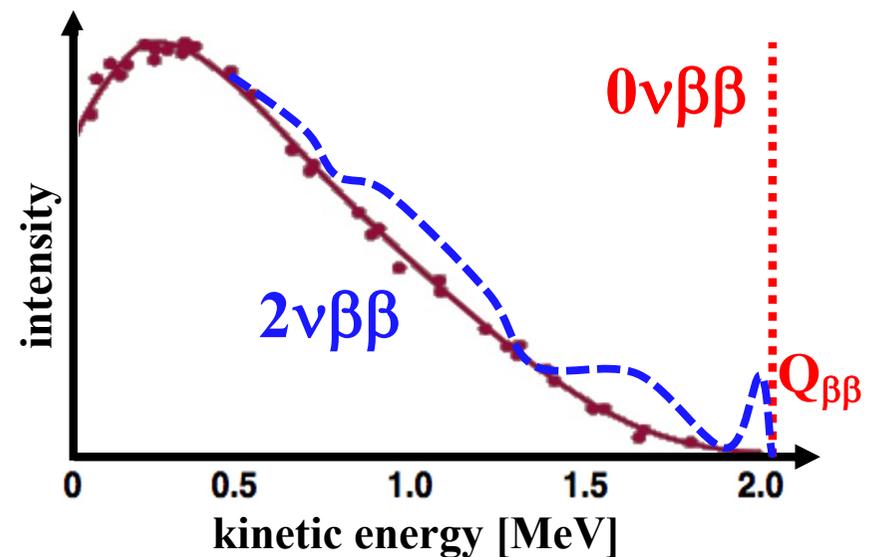
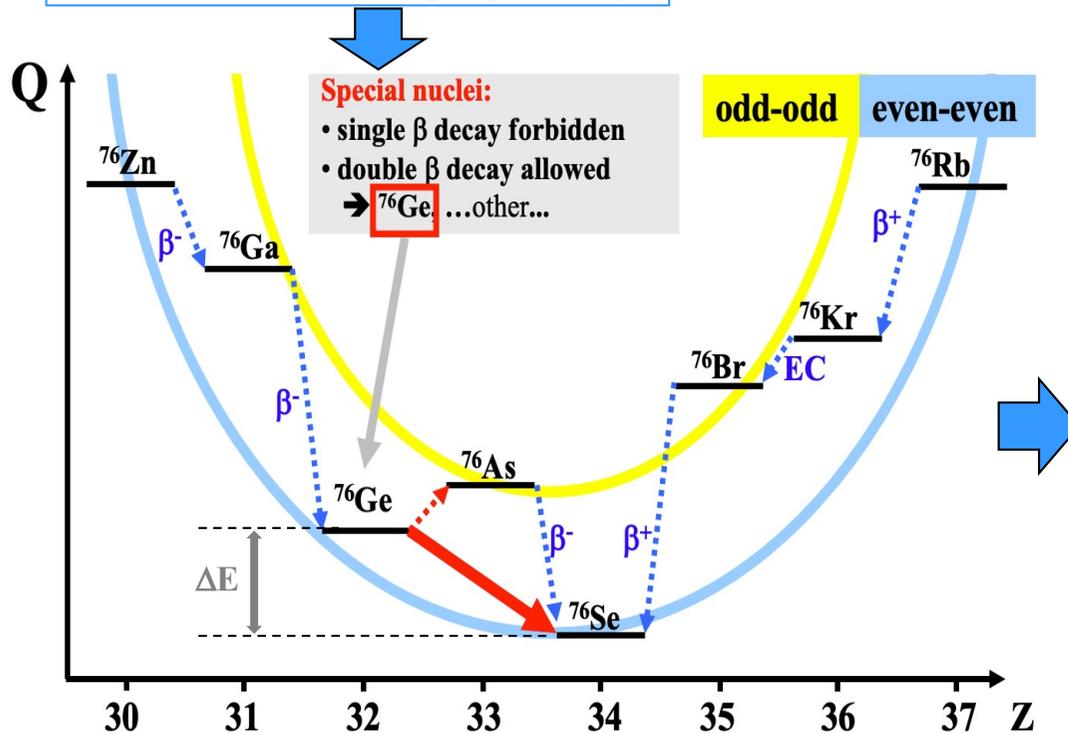


some
BSM
decay

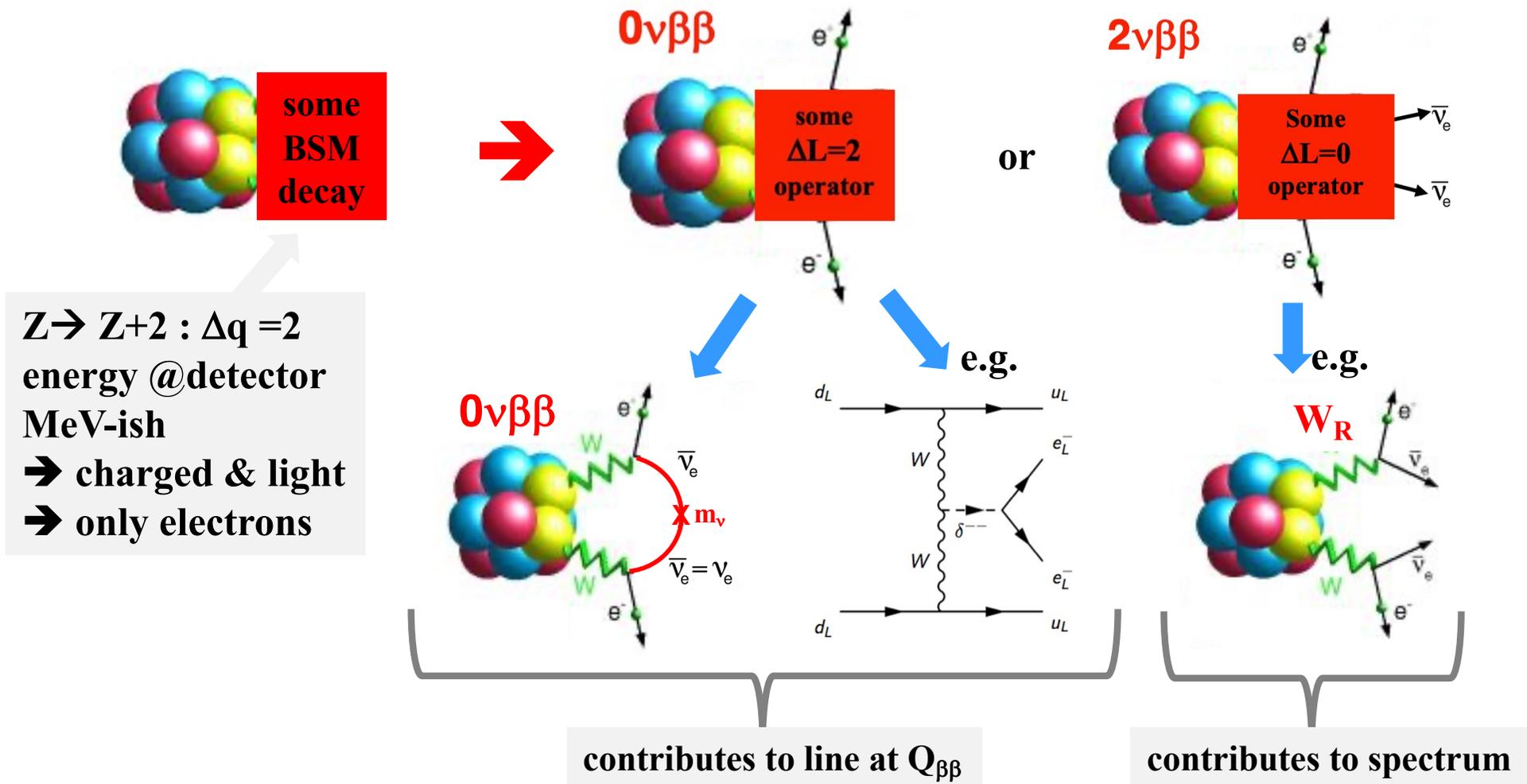
- SM: $2e+2\nu \rightarrow$ spectrum

- BSM, $Z \rightarrow Z+2$, $Q_{\beta\beta} \sim \text{MeV-ish}$
 \rightarrow requires 2 electrons:

- $2e + \text{no other particles}$
 \rightarrow line @ $Q_{\beta\beta}$ and $\Delta L=2$
- $2e + X$
 \rightarrow modified spectrum $\leq Q_{\beta\beta}$



Double Beta Decay and BSM Physics

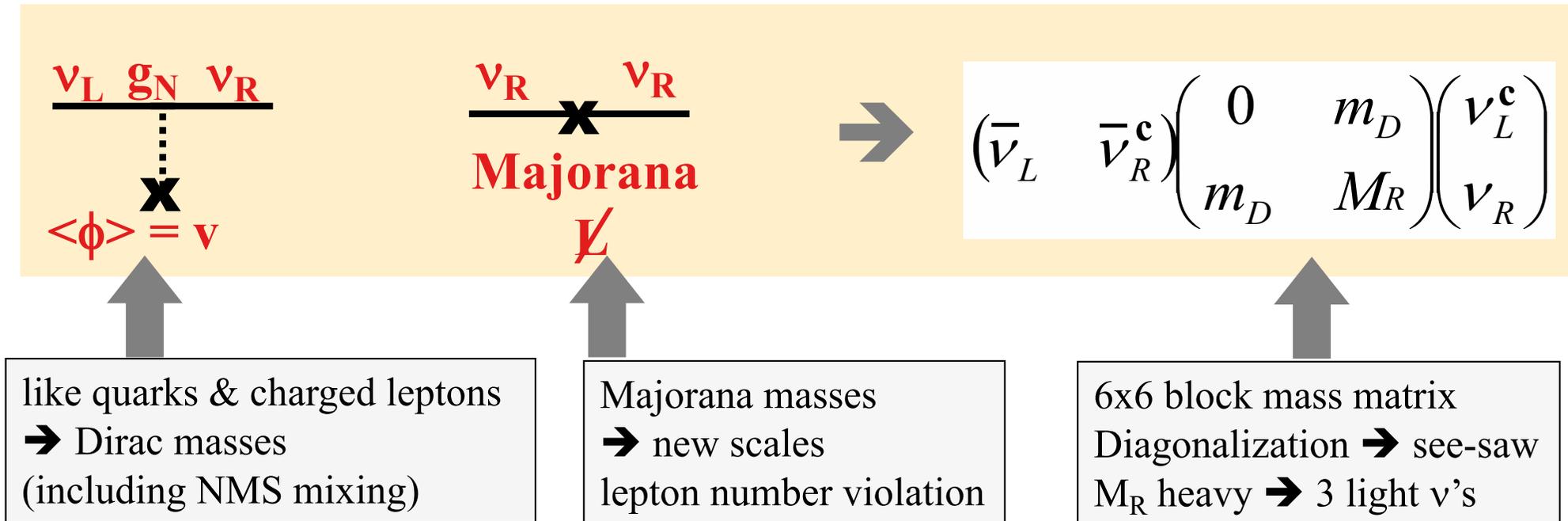


- many reasons for BSM physics with and w/o $\Delta L=2$
- if one ignores all other reasons → neutrino masses
- many BSM models affect both $0\nu\beta\beta$ and $2\nu\beta\beta$

→ why think only about neutrinos?

Bottom-Up from a Neutrino Perspective

Simplest: Another copy of SM field \rightarrow new representation \rightarrow extra gauge group \rightarrow ...
 \rightarrow Simplest and suggestive possibility: right handed neutrinos



Assumptions:

- 3 right handed neutrino states \rightarrow complete gauge singlet ($1_c, 1_L, Y=0$)
- no extra symmetry which forbids L violation

Consequences:

- new parameters
- modified symmetries: L violation
- nice explanation of light neutrinos
- multi-scale theory

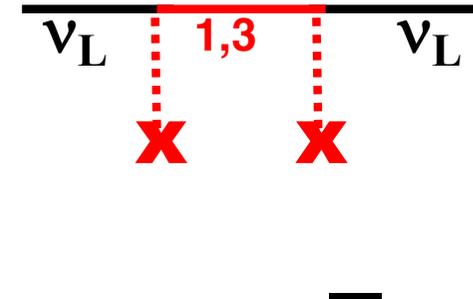
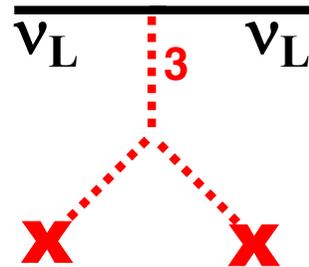
\rightarrow SM+

Many other Possibilities

- add scalar triplets (3_L) or add fermionic (1_L) or (3_L)

→ Majorana masses without ν_R :

$$M_L L L^c$$



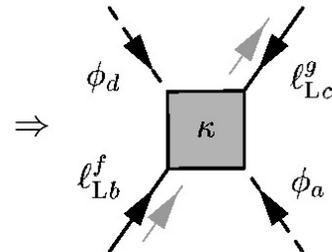
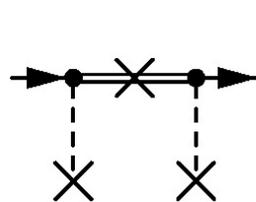
- add both ν_R and new singlets / triplets:

→ see-saw type II, III

$$m_\nu = M_L - m_D M_R^{-1} m_D^T$$

- add higher dimensional operators: $d=5, \dots$

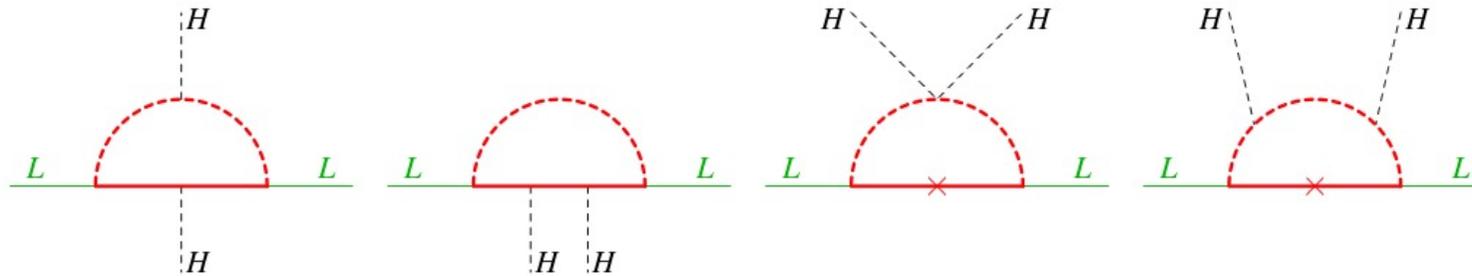
→ model independence



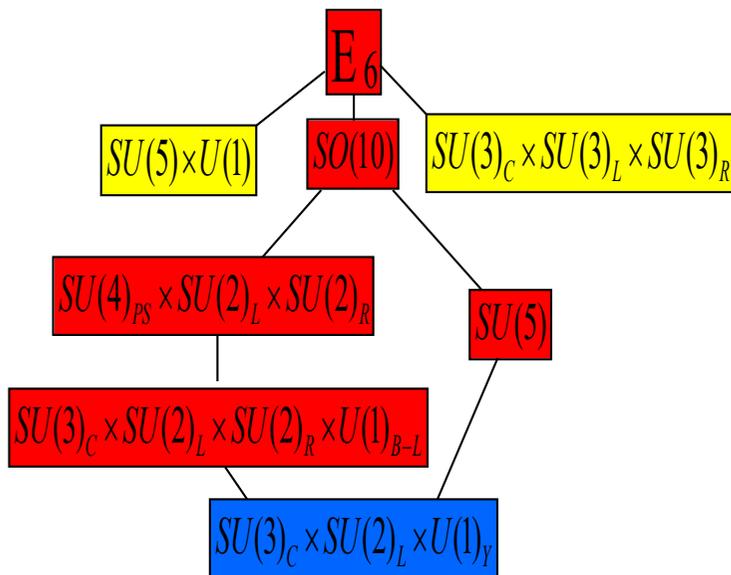
$$\Leftrightarrow \mathcal{L}_{mass} = \kappa \cdot \bar{\nu}_L^C \nu_L \Phi^T \Phi$$

$$\Rightarrow M_L \bar{L} L^c$$

■ Radiative neutrino mass generation \leftrightarrow BSM



■ Gauge unification, GUTs and Flavour symmetries



Generations are only copies
 → horizontal symmetries?
 A4, S3, D5, T', ...
 breaking...
 → explain regularities

more options \leftrightarrow more
 constrains, precision

Leptons	2/3	2/3	2/3
	u ~5	c ~1350	t 175000
	-1/3	-1/3	-1/3
	d ~9	s ~175	b ~4500
Quarks	0?	0?	0?
	ν_1	ν_2	ν_3
	e 0.511	μ 105.66	τ 1777.2
	1.	2.	3.
	generation		

■ More: further ν 's, SUSY, gauge groups, extra d, strings,...

What is realized in Nature?

A huge number of options and a huge parameter space...

... but we know little: $2 \times \Delta m^2$, mixings, \sim mass, unitarity bounds, astro/cosmo

- **choose a ‘minimal model’ and explain all there is (so far...)**



- three ν_R to explain only masses and mixings
- ν MSM \rightarrow keV neutrinos as DM, BAU (Shaposhnikov et al.)
- often leaves other questions out (e.g. unification, inflation, DE, flavour, ...)

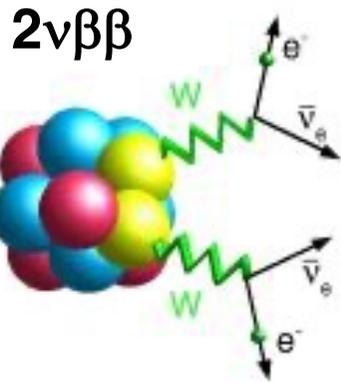
- **many wider and very well motivated BSM options & hints**

- often connected to other fields \rightarrow **involves well motivated new physics**
- must explain all existing data in neutrino physics **and** in other fields
 - \rightarrow connections to other fields: **LFV, LHC, DM**
- must be consistent with diverse theoretical constraints
 - \rightarrow BBN, cosmology, SN explosions, ...
 - \rightarrow leptogenesis \leftrightarrow variants, other BAU
 - \rightarrow keV sterile ν 's as warm dark matter \leftrightarrow other DM candidates

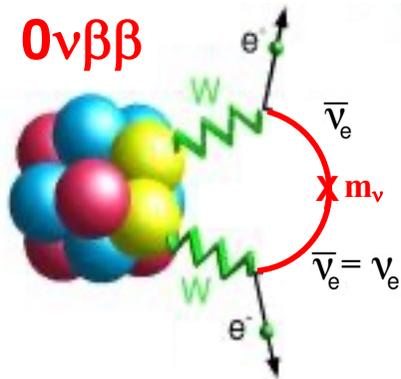
WE DON'T KNOW \rightarrow be open minded!

Majorana Neutrino Interpretation of $0\nu\beta\beta$

SM

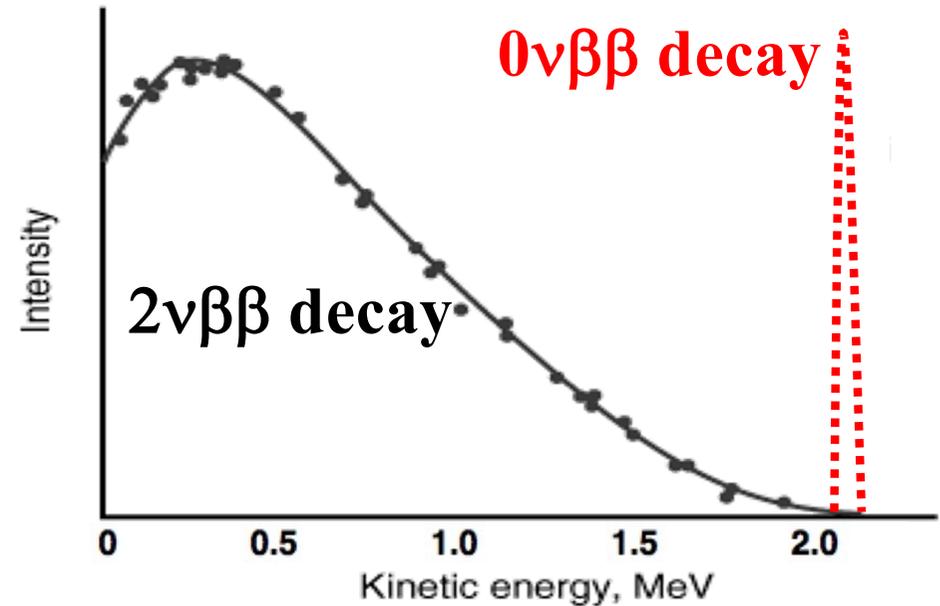


assume
Majorana
mass



$2\nu\beta\beta$ decay seen for different isotopes

$T^{1/2} = O(10^{18} - 10^{21} \text{ years}) \rightarrow \text{up to } 10^{11} \otimes T_{\text{Universe}}$



$$1/\tau = G(Q,Z) |M_{\text{nucl}}|^2 \langle m_{ee} \rangle^2$$

lifetime τ

NMEs have
uncertainties...

- observe expected $2\nu\beta\beta$ signal
- look for a line at $Q_{\beta\beta}$

interpret signal \rightarrow **Majorana mass**

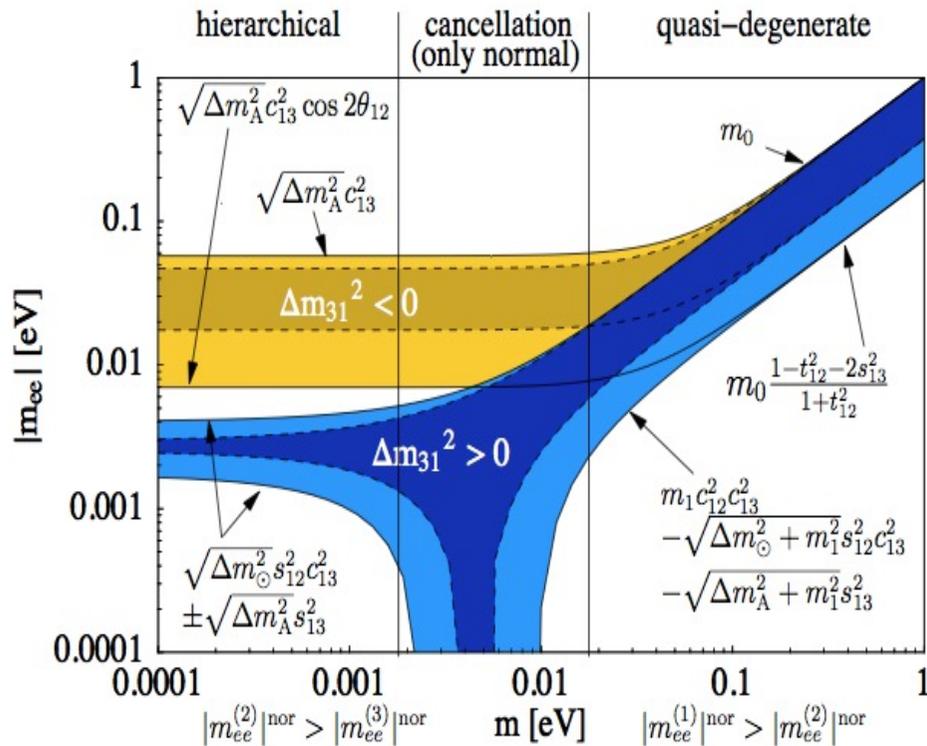
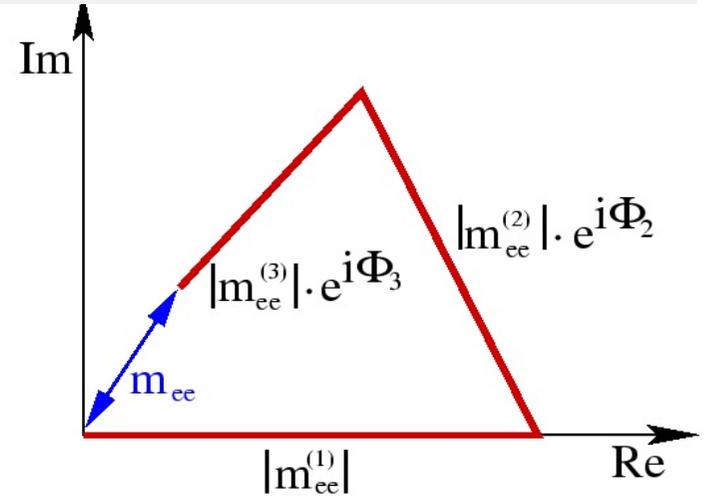
m_{ee} : The Effective Neutrino Mass

$$m_{ee} = |m_{ee}^{(1)}| + |m_{ee}^{(2)}| \cdot e^{i\Phi_2} + |m_{ee}^{(3)}| \cdot e^{i\Phi_3}$$

$$|m_{ee}^{(1)}| = |U_{e1}|^2 m_1$$

$$|m_{ee}^{(2)}| = |U_{e2}|^2 \sqrt{m_1^2 + \Delta m_{21}^2}$$

$$|m_{ee}^{(3)}| = |U_{e3}|^2 \sqrt{m_1^2 + \Delta m_{31}^2}$$

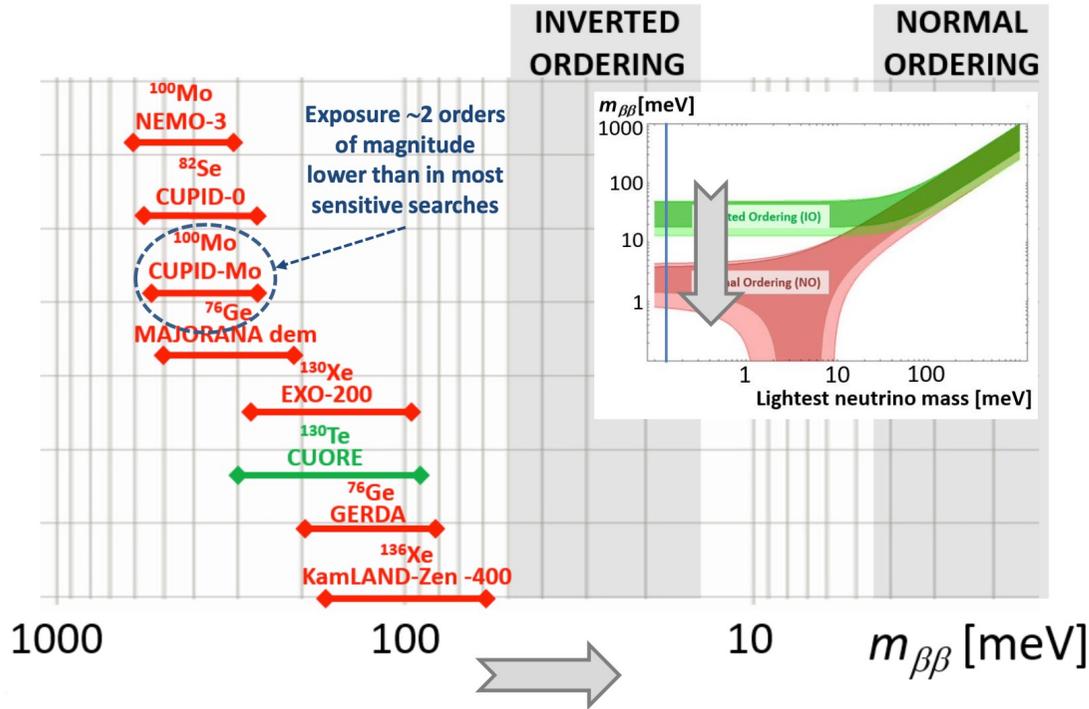


Comments:

- NMEs \rightarrow unavoidable theory errors
- known Δm^2 from oscillations
TH + EX errors \rightarrow bands
- ton-scale projects aim at covering IH
- global fits ...NH?
- cosmology...
- **Assumption:** No other BSM physics...!

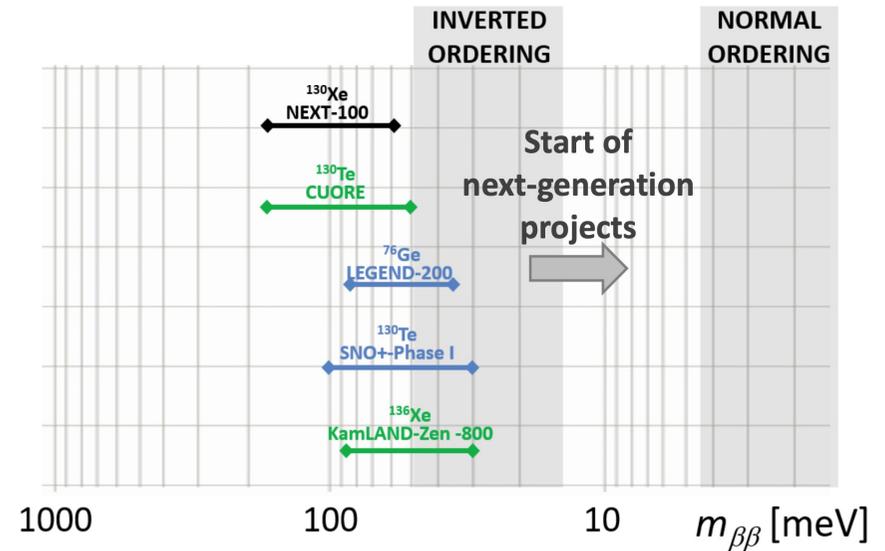
Experimental Status & Outlook

A. Giuliani, talk at TAUP2021, <https://indico.ific.uv.es/event/6178/>

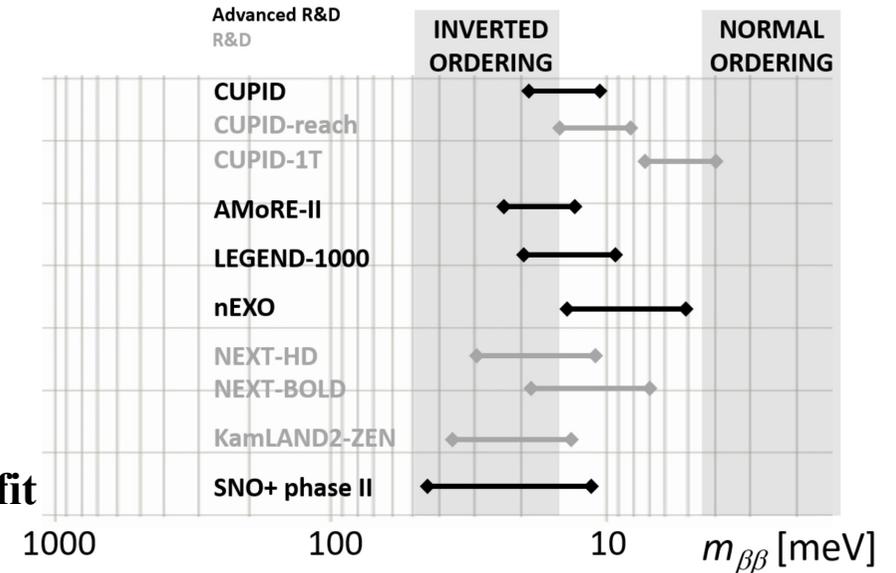


- next years: start to cover all IH
- in the long run:
 - will cover all of IH
 - improved BSM tests
 - so far no convincing project to test (most) NH
- while waiting: watch NH preference
- now: global fits prefer NH weakened by T2K/NoVA CP fit

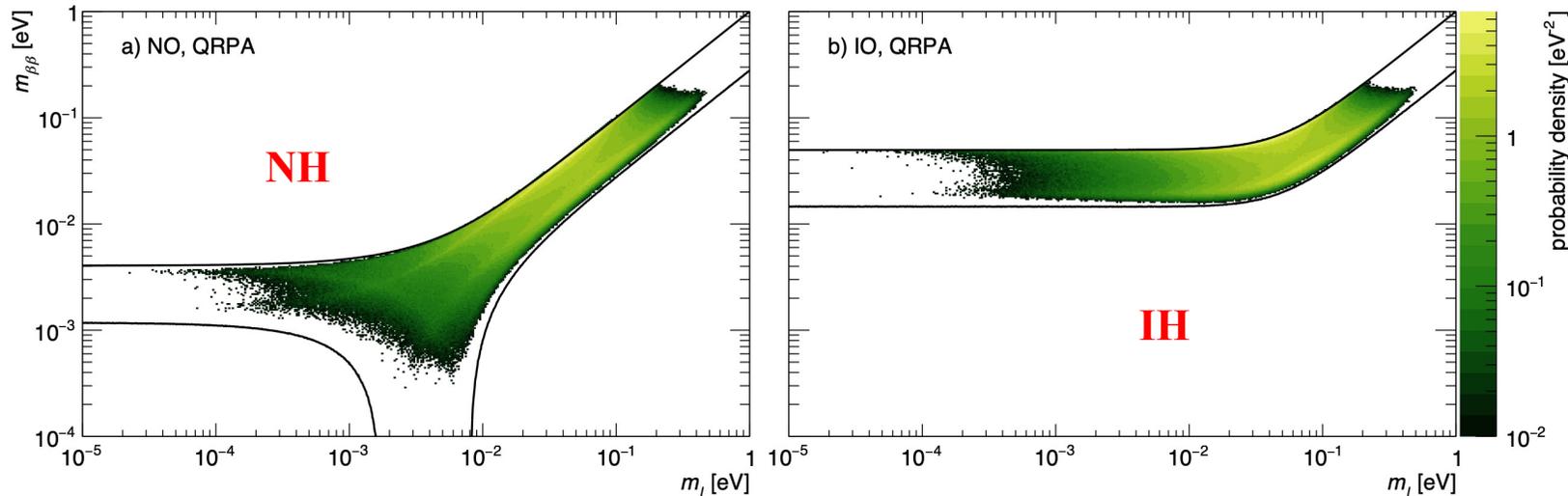
possible ~5y situation



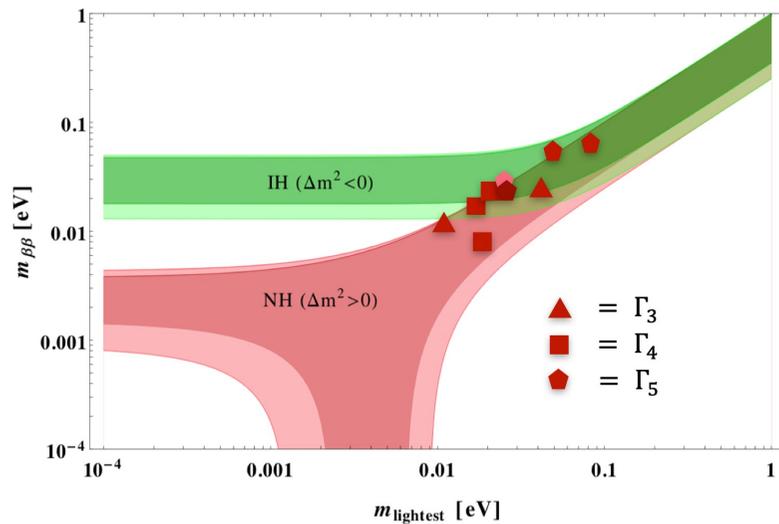
long term goals of next generation projects



A lower or Limit?



Agostini, Benato, Detwiler PRD 96 (2017) 053001 ; similar: Caldwell, Ettengruber, Merle, Schulz, Totzauer, PRD 96 (2017) 7, 073001

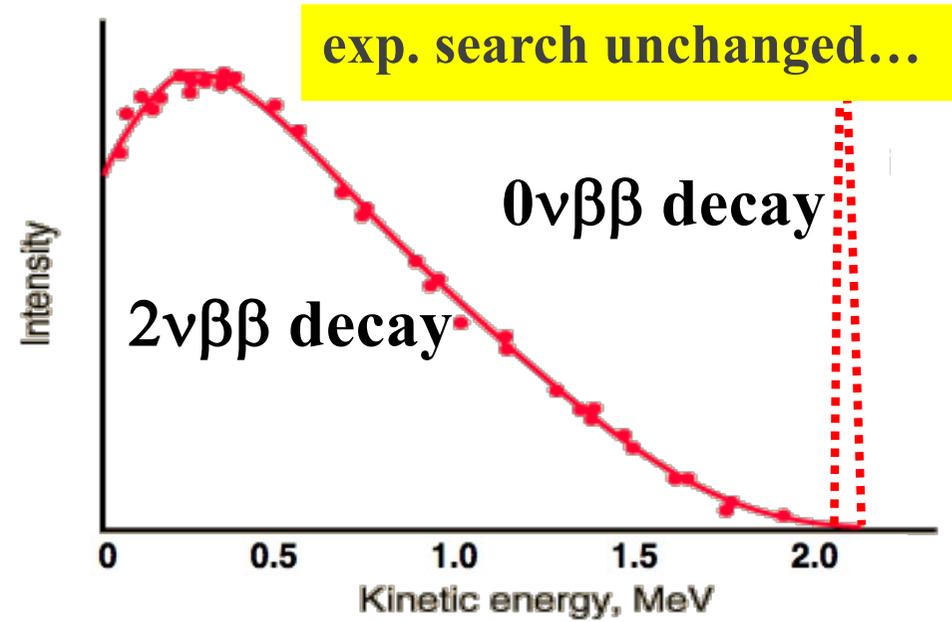
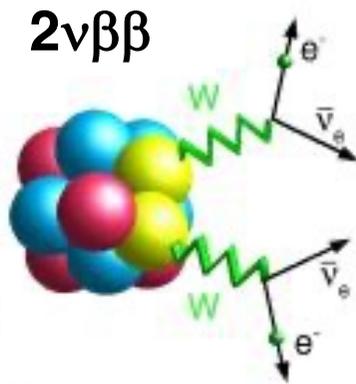


Feruglio, FLASY 2019, <http://indico.cern.ch/event/795851>

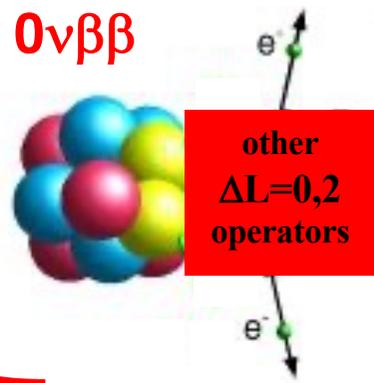
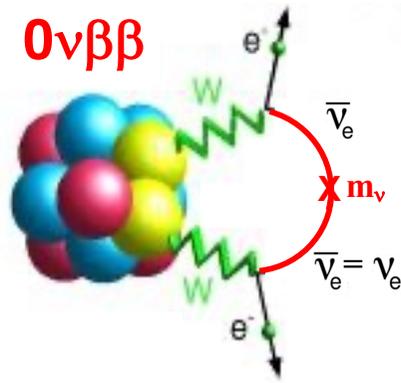
- **global fits of neutrino parameters**
 - ➔ constrain model parameters (NH \leftrightarrow IH)
 - ➔ weak preference for NH (2.7σ in NuFit 5.0)
 - ➔ **data: no lower limit**
 - ...and $0\nu\beta\beta$ may be other physics ➔ $m_{\beta\beta} = 0$, tiny
- **theory: there is no probability**
 - ➔ SM+ m_ν : any number per se equally good
 - ➔ no theory measure (linear, log, ...)
 - ➔ discrete symmetries: added assumptions
 - ➔ **depends on symmetry ➔ no generic preference**
- **same methods do not work for quarks, e, μ , τ**

More general: L Violating Processes

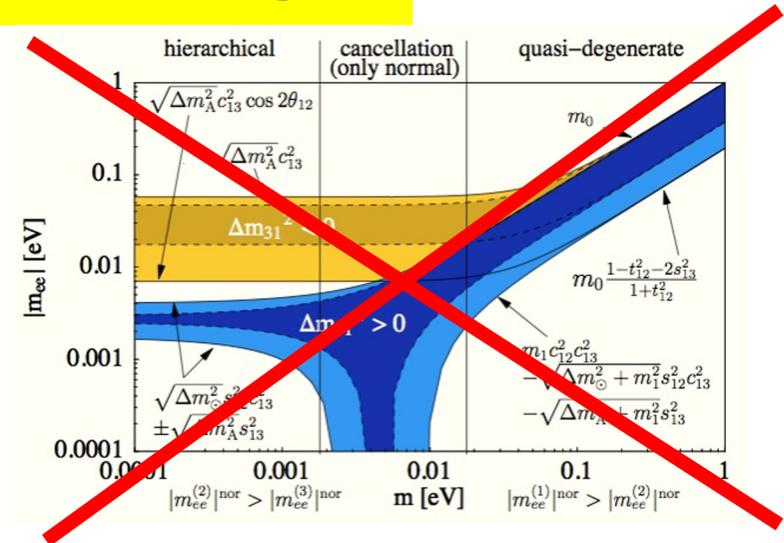
SM



BSM

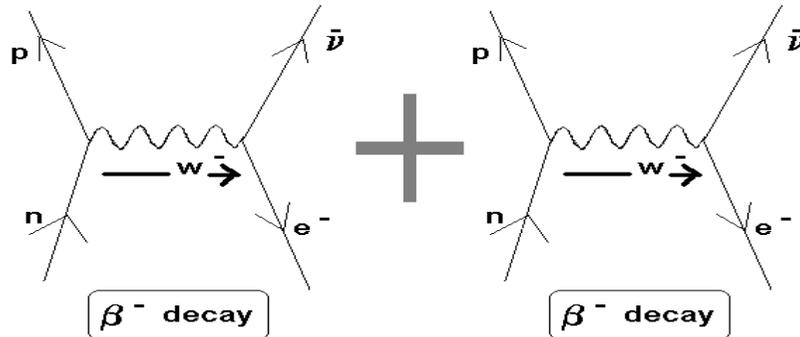


...interpretation changes:



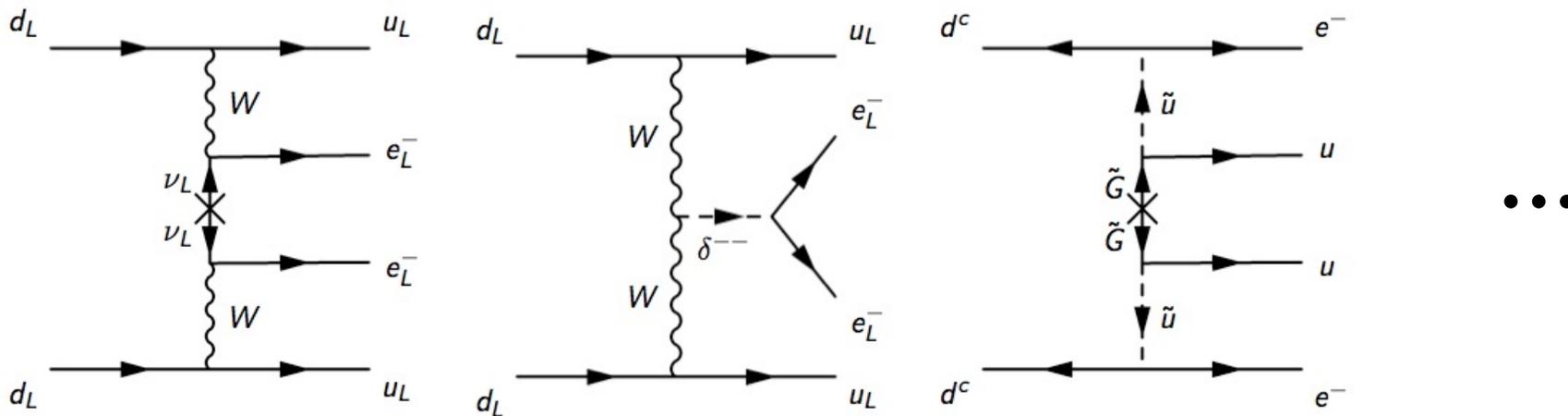
Other Double Beta Decay Processes

Standard Model:



→ 2 electrons + 2 neutrinos
 $2\nu\beta\beta$

Majorana ν -masses or other $\Delta L=2$ physics: → 2 electrons



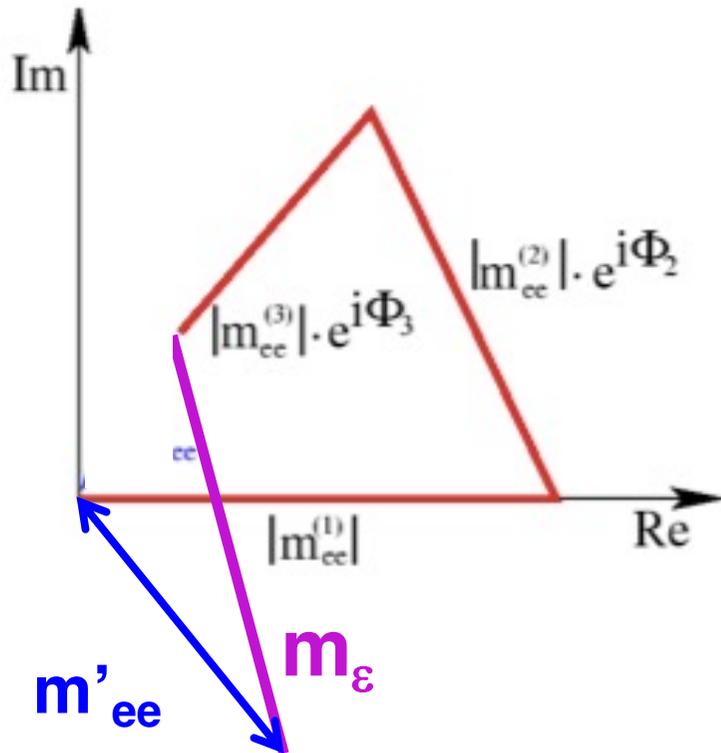
$0\nu\beta\beta$

Majorana
 neutrino masses
 \leftrightarrow Dirac?

SM + Higgs triplet

SUSY

important connections to LHC and LFV ...
 sub eV Majorana mass \leftrightarrow TeV scale physics



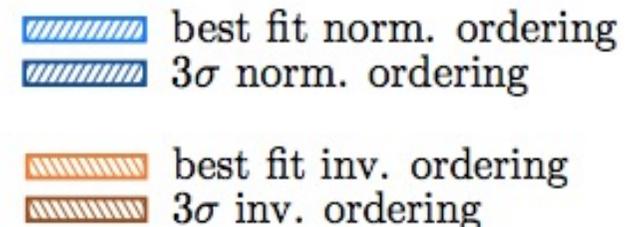
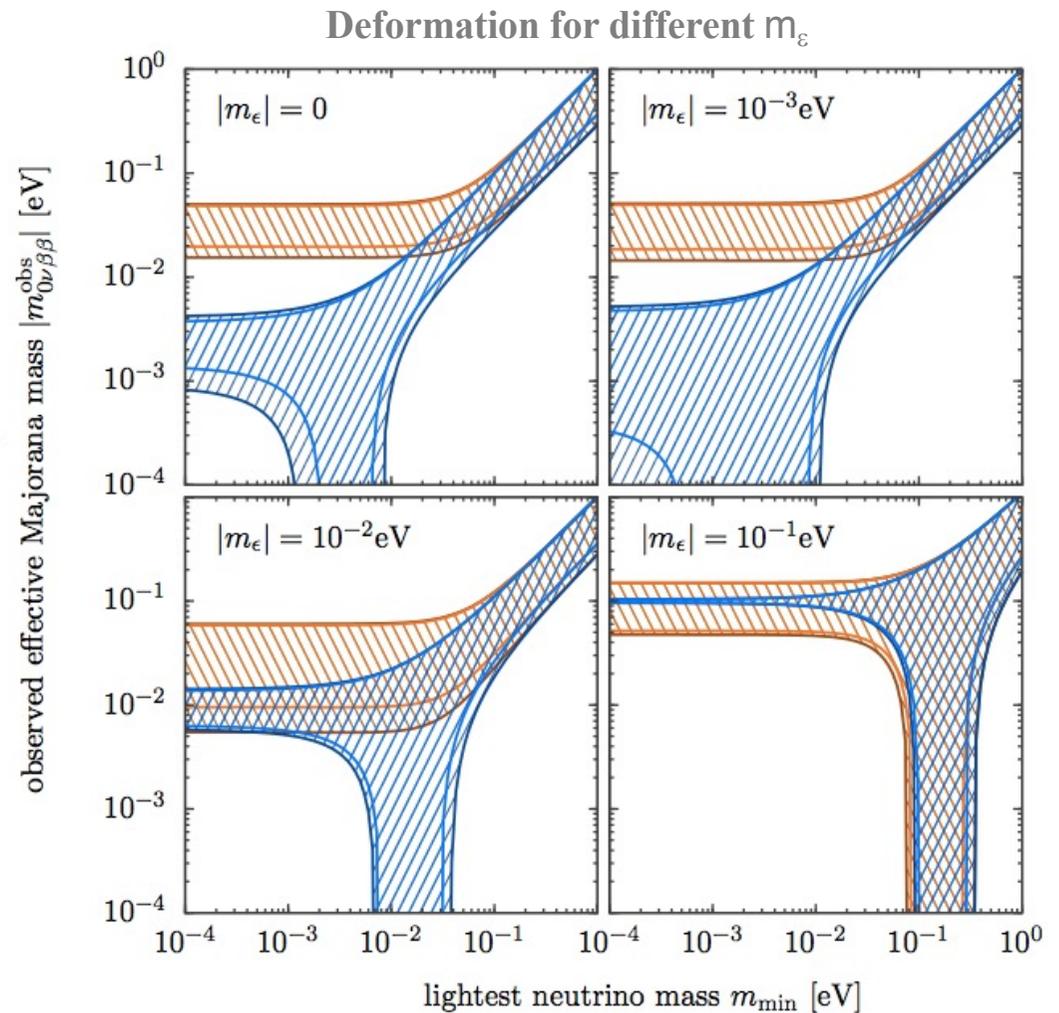
Interferences:

total $0\nu\beta\beta$

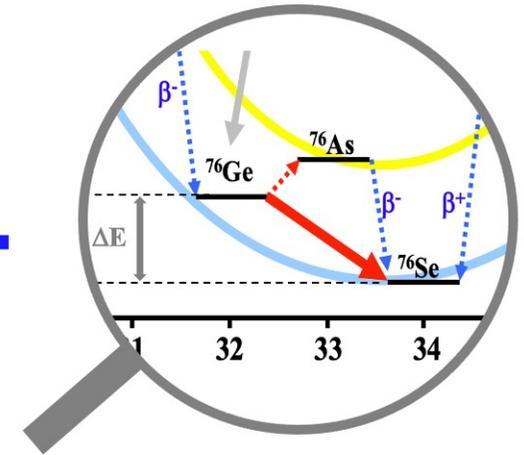
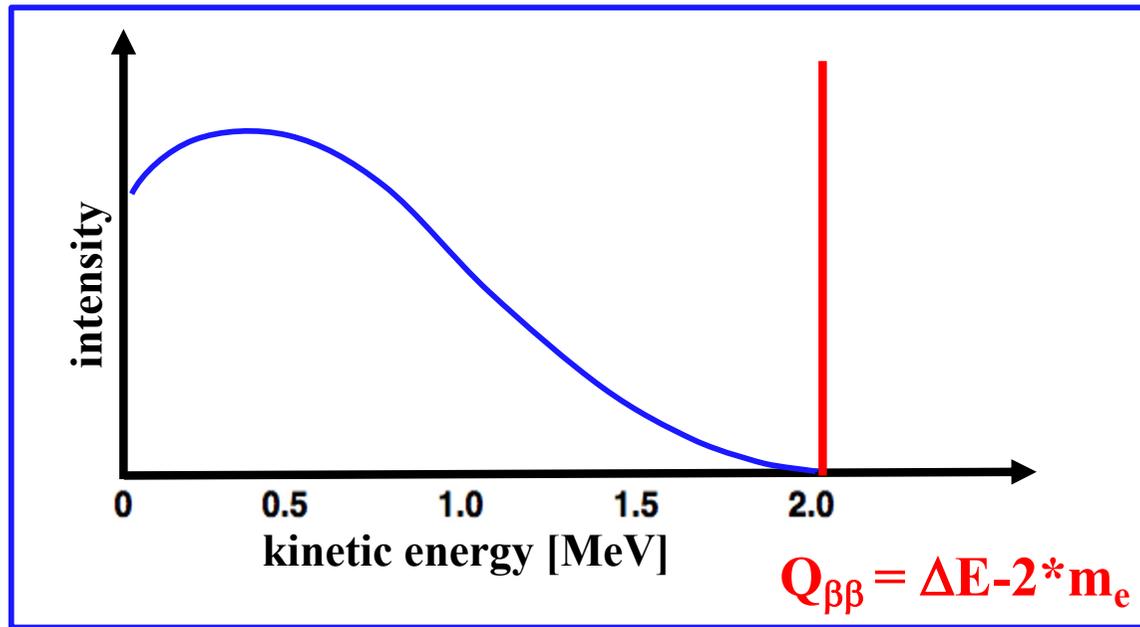
- neutrino contribution
- $m_{\epsilon} \leftrightarrow$ other $\Delta L=2$ physics
- shifts masses, mixings, CP phases
- deformation due to TeV physics

Extreme case:

- Dirac neutrino masses
- all due to other $\Delta L=2$ physics



Modifications of the Energy Spectrum



BSM can affect:

- amplitudes of both $2\nu\beta\beta$ and $0\nu\beta\beta \leftrightarrow$ NME uncertainties
- shape
- line position

backgrounds in a real experiment \leftrightarrow what can be detected

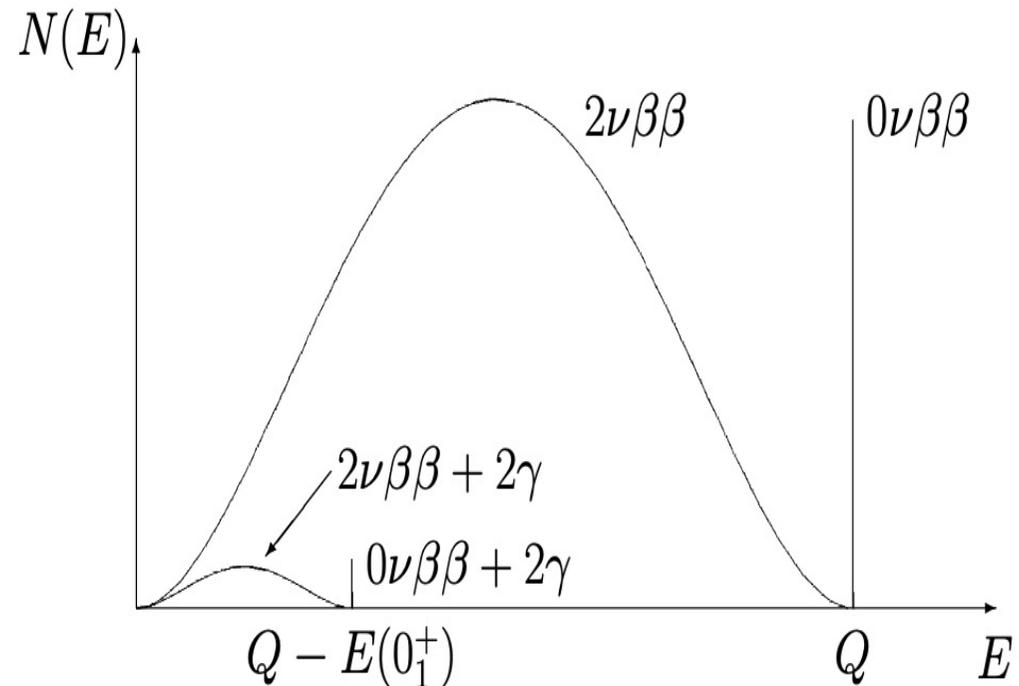
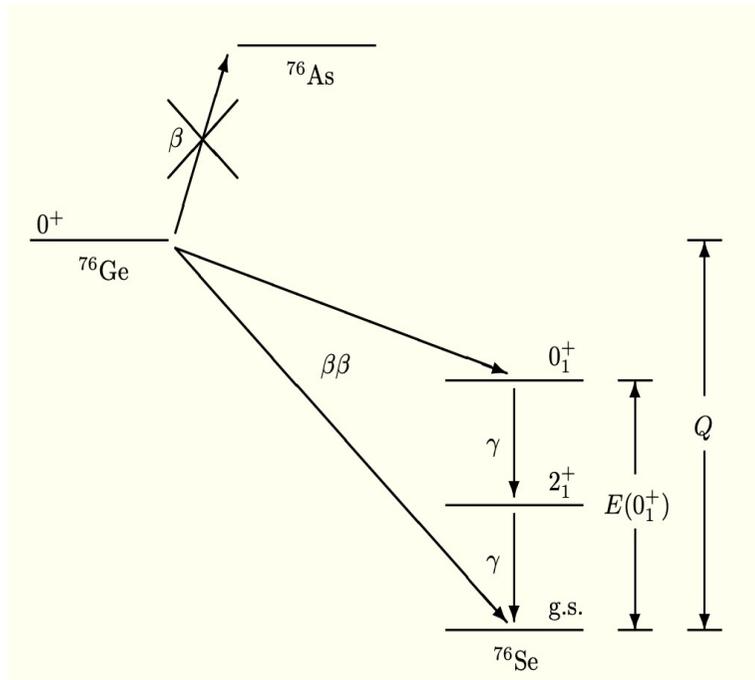
Examples:

- sterile neutrinos with masses \simeq MeV and $2\nu\beta\beta$
- right-handed currents $\rightarrow M_{\text{WR}}$
- ...

Consistency Test of $0\nu\beta\beta$ with one Isotope

$0\nu\beta\beta$ can also go to excited states

Duerr, ML, Zuber, PRD 84 (2011) 093004



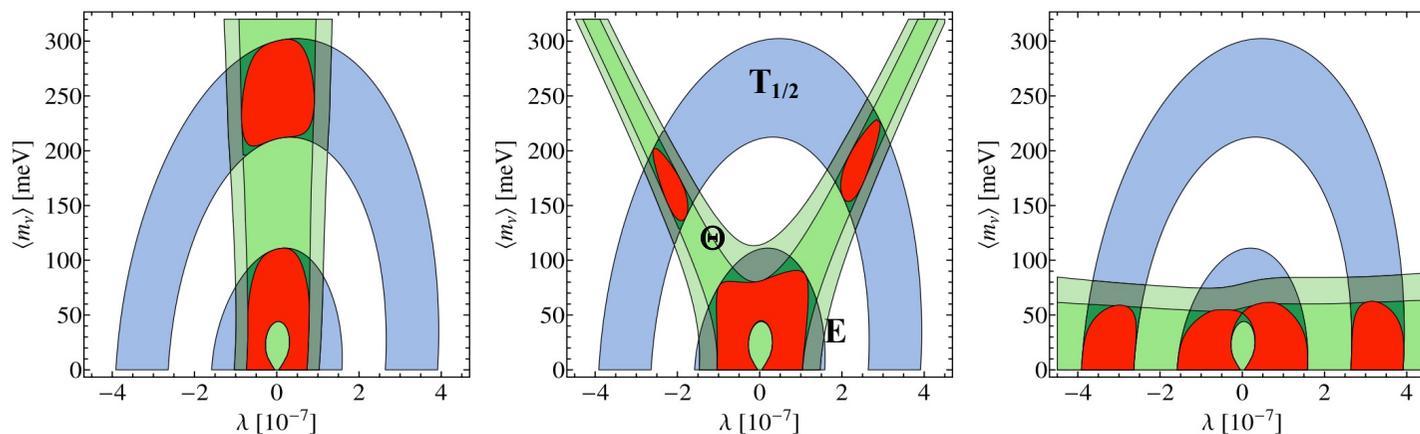
→ small corrections to the $0\nu\beta\beta$ spectrum
from sub-leading SM allowed transitions – not BSM

A nice consistency test with one isotope

$\Delta E \leftrightarrow$ known energy levels → does not depend on NME uncertainties

Resolving the Origin of $0\nu\beta\beta$

- **Over-constraining neutrino parameters**
 - m_ν from $0\nu\beta\beta \leftrightarrow$ other determinations: KATRIN, cosmology, ...
 - **so far only bounds and no discrepancy \rightarrow future:**
 - perfect match \rightarrow just Majorana neutrinos?
 - mismatch \rightarrow something else must contribute
 - finite mass and no $0\nu\beta\beta \rightarrow$ Dirac mass?
- **Further information in $0\nu\beta\beta$ experiments**
 - microscopic physics \leftrightarrow angular distribution (SuperNEMO, NEXT)



Arnold et al., Eur. Phys. J. C (2010) 70: 927

left-right symmetry
 \rightarrow right handed currents

$$\lambda \approx \left(\frac{M_{WL}}{M_{WR}} \right)^2 \sqrt{\frac{m_\nu}{M_R}}$$

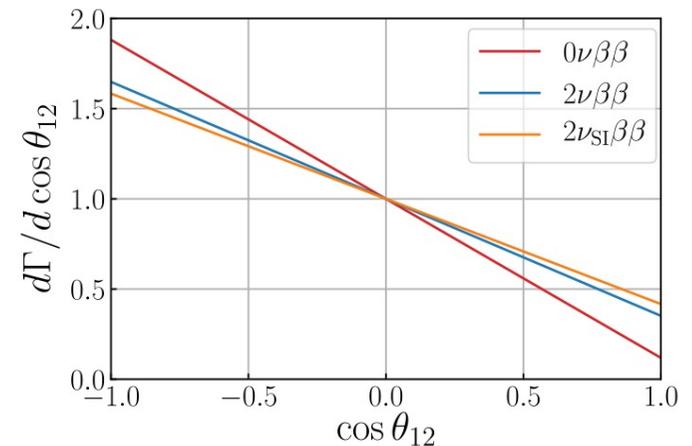
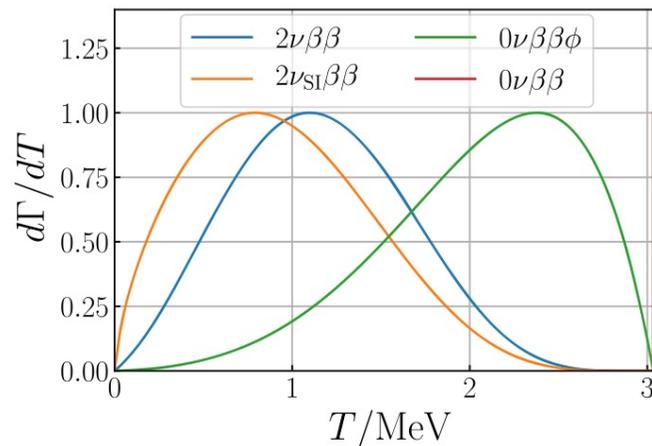
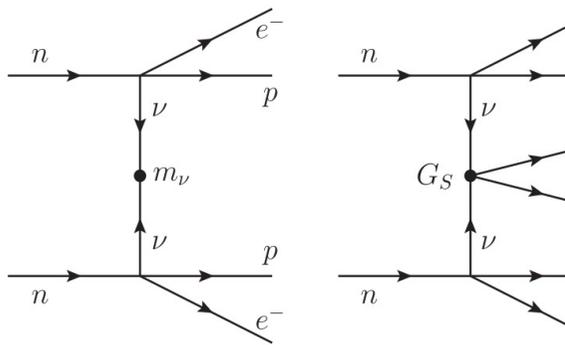
- a) pure MM contribution
- b) 30% RHC $_\lambda$ admixture
- c) pure RHC $_\lambda$

Combine $0\nu\beta\beta$ and $2\nu\beta\beta$

Future experiments will measure $2\nu\beta\beta$ spectrum very precisely
→ combine $0\nu\beta\beta$ and $2\nu\beta\beta$ (effects on top of high statistics SM signal)

- **Lepton-number conserving right-handed currents**
 Deppisch, Graf, Simkovic, PRL 125 (2020) 17, 171801
- **Sterile neutrino and light fermion searches through energy end point**
 Bolton, Deppisch, Graf, Simkovic, PRD 103 (2021) 5, 055019, Agostini, et al., PLB 815 (2021) 136127
- **Neutrino self-interactions**
 Deppisch, Graf, Rodejohann, Xu, PRD 102 (2020) 5, 051701

• ...



High-Energy Collider Experiments

testing long- and short range operators

LEFT \leftrightarrow SMEFT \rightarrow more information

dimension-9 short-range operators
scalar O_i and vector O_i^μ four-quark operators

\leftrightarrow Complete set of d=9 Operators in SMEFT
Li, Ren, Xiao, Yu, Zheng, PRD104 (2021) 1, 015025

$$\begin{aligned}
 \mathcal{O}_1 &= (\bar{u}_L^\alpha \gamma_\mu d_L^\alpha) (\bar{u}_L^\beta \gamma^\mu d_L^\beta), & \mathcal{O}'_1 &= (\bar{u}_R^\alpha \gamma_\mu d_R^\alpha) (\bar{u}_R^\beta \gamma^\mu d_R^\beta), \\
 \mathcal{O}_2 &= (\bar{u}_R^\alpha d_L^\alpha) (\bar{u}_R^\beta d_L^\beta), & \mathcal{O}'_2 &= (\bar{u}_L^\alpha d_R^\alpha) (\bar{u}_L^\beta d_R^\beta), \\
 \mathcal{O}_3 &= (\bar{u}_R^\alpha d_L^\beta) (\bar{u}_R^\beta d_L^\alpha), & \mathcal{O}'_3 &= (\bar{u}_L^\alpha d_R^\beta) (\bar{u}_L^\beta d_R^\alpha), \\
 \mathcal{O}_4 &= (\bar{u}_L^\alpha \gamma_\mu d_L^\alpha) (\bar{u}_R^\beta \gamma^\mu d_R^\beta), \\
 \mathcal{O}_5 &= (\bar{u}_L^\alpha \gamma_\mu d_L^\beta) (\bar{u}_R^\beta \gamma^\mu d_R^\alpha), \\
 \mathcal{O}_6^\mu &= (\bar{u}_L \gamma^\mu d_L) (\bar{u}_L d_R), & \mathcal{O}_6^{\mu'} &= (\bar{u}_R \gamma^\mu d_R) (\bar{u}_R d_L), \\
 \mathcal{O}_7^\mu &= (\bar{u}_L \lambda^A \gamma^\mu d_L) (\bar{u}_L \lambda^A d_R), & \mathcal{O}_7^{\mu'} &= (\bar{u}_R \lambda^A \gamma^\mu d_R) (\bar{u}_R \lambda^A d_L), \\
 \mathcal{O}_8^\mu &= (\bar{u}_L \gamma^\mu d_L) (\bar{u}_R d_L), & \mathcal{O}_8^{\mu'} &= (\bar{u}_R \gamma^\mu d_R) (\bar{u}_L d_R), \\
 \mathcal{O}_9^\mu &= (\bar{u}_L \lambda^A \gamma^\mu d_L) (\bar{u}_R \lambda^A d_L), & \mathcal{O}_9^{\mu'} &= (\bar{u}_R \lambda^A \gamma^\mu d_R) (\bar{u}_L \lambda^A d_R),
 \end{aligned}$$

d=7

$\psi^2 \phi^4$		$\psi^4 \phi$	
$\mathcal{O}_{eH^3H^1}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{im} \epsilon^{jn} (l_{pi} C l_{rj}) H_m H_n (H^1 H)$	$\mathcal{O}_{d_{1u}H}$	$\epsilon^{ij} (\bar{d}_p^i l_{rj}) (\epsilon_p C u_{ra}) H_i$
	$\psi^2 \phi^3 D$	$\mathcal{O}_{d_{1q}H}^{(1)}$	$\mathcal{Y}[\overline{[3]}] \epsilon^{ik} \epsilon^{jm} (\bar{d}_p^i l_{rj}) (l_{ri} C q_{tak}) H_m$
\mathcal{O}_{eH^3D}	$\epsilon^{ij} \epsilon^{km} (l_{pi} C \gamma^\mu e_r) H_j H_n D_\mu H_k$	$\mathcal{O}_{d_{1q}H}^{(2)}$	$\mathcal{Y}[\overline{[3]}] \epsilon^{ij} \epsilon^{km} (\bar{d}_p^i l_{rj}) (l_{ri} C q_{tak}) H_m$
	$\psi^2 \phi^2 D^2$	$\mathcal{O}_{d_{1q}H}^{(3)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{d}_p^i l_{rj}) (l_{ri} C q_{tak}) H_m$
$\mathcal{O}_{eH^2}^{(1)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ik} \epsilon^{jm} (l_{pi} C \sigma^{\mu\nu} l_{rj}) D_\mu H_k D_\nu H_m$	$\mathcal{O}_{d_{1q}H}^{(4)}$	$\mathcal{Y}[\overline{[333]}] \epsilon^{ij} \epsilon^{km} (\bar{d}_p^i l_{rj}) (l_{ri} C q_{tak}) H_m$
$\mathcal{O}_{eH^2}^{(2)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ik} \epsilon^{jm} (l_{pi} C l_{rj}) D_\mu H_k D^\mu H_m$	$\mathcal{O}_{eH}^{(1)}$	$\mathcal{Y}[\overline{[333]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
	$F \psi^2 \phi^2$	$\mathcal{O}_{eH}^{(2)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ik} \epsilon^{jm} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
$\mathcal{O}_{W^2H^2}^{(1)}$	$\mathcal{Y}[\overline{[3]}] \epsilon^{ik} \epsilon^{jn} (\tau^I)_n W_{\mu\nu}^I (l_{pi} C \sigma^{\mu\nu} l_{rj}) H_k H_m$	$\mathcal{O}_{eH}^{(3)}$	$\mathcal{Y}[\overline{[3]}] \epsilon^{ik} \epsilon^{jm} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
$\mathcal{O}_{W^2H^2}^{(2)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ik} \epsilon^{jn} (\tau^I)_n W_{\mu\nu}^I (l_{pi} C \sigma^{\mu\nu} l_{rj}) H_k H_m$	$\mathcal{O}_{eH}^{(4)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
\mathcal{O}_{BH^2}	$\mathcal{Y}[\overline{[3]}] \epsilon^{ik} \epsilon^{jm} B_{\mu\nu} (l_{pi} C \sigma^{\mu\nu} l_{rj}) H_k H_m$	$\mathcal{O}_{eH}^{(5)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
	$\psi^4 D$	$\mathcal{O}_{eH}^{(6)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
$\mathcal{O}_{d_{1u}D}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} (\bar{d}_p^i \gamma^\mu u_{ra}) (l_{ri} C D_\mu l_{rj})$	$\mathcal{O}_{eH}^{(7)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
$\mathcal{O}_{d_{1q}D}(\not{E})$	$\mathcal{Y}[\overline{[33]}] \epsilon^{abc} (\bar{l}_i^a D_\mu d_{ra}) (d_{pa} \gamma^\mu C l_{rc})$	$\mathcal{O}_{eH}^{(8)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$
$\mathcal{O}_{d_{1q}D}(\not{E})$	$\mathcal{Y}[\overline{[333]}] \epsilon^{abc} (\bar{l}_i^a \gamma^\mu d_{pa}) (d_{ra} C D_\mu d_{sc})$	$\mathcal{O}_{eH}^{(9)}$	$\mathcal{Y}[\overline{[33]}] \epsilon^{ij} \epsilon^{km} (\bar{e}_p^i l_{rj}) (l_{ri} C l_{tk}) H_m$

d=9

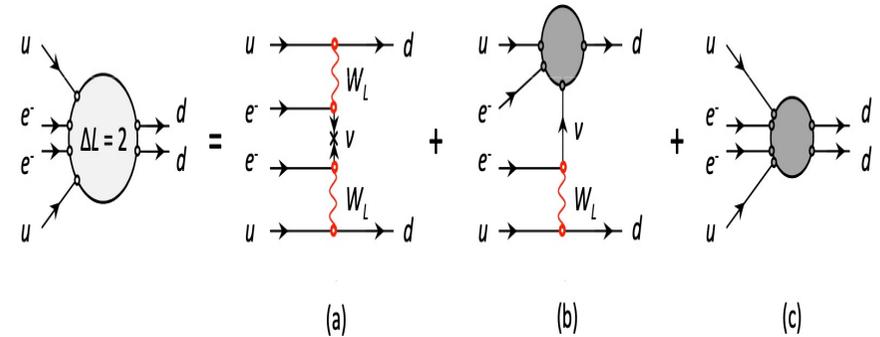
...many more...

- **Same-Sign Dilepton Signals:** $p + p \rightarrow 2e^\pm + 2\text{jets}$
- **Missing Transverse Energy (MTE):** $p + p \rightarrow e^\pm + 2\text{jets} + \text{MTE}$
- **Dijet with Missing Transverse Energy:** $p + p \rightarrow 2\text{jets} + \text{MTE}$

Study / combine with other Channels

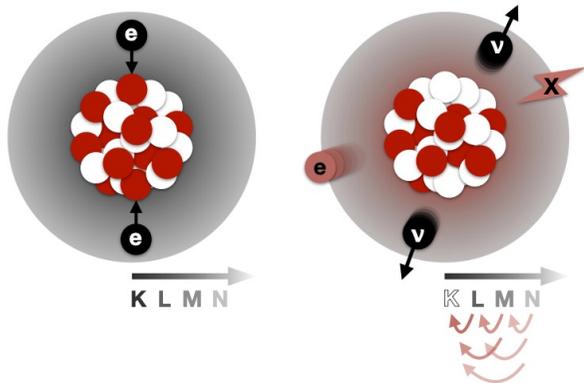
Other $\Delta L = 2$ channels:

- $0\nu\beta^+\beta^+ : A, Z \rightarrow A, Z-2 + 2e^+$
- $0\nu\beta^+EC : A, Z + e^- \rightarrow A, Z-2 + e^+$
- $0\nu ECEC : A, Z + 2e^- \rightarrow A, Z-2$

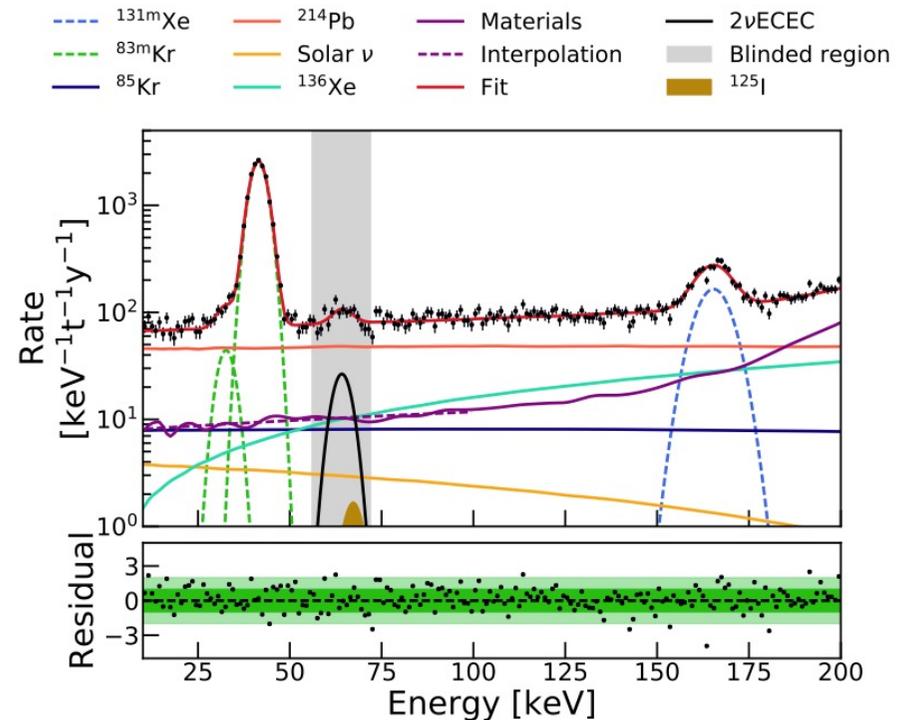


Observation of $2\nu ECEC$ by XENON1T

E. Aprile et al., Nature 568 (2019) 7753, 532.



$T_{2\nu ECEC} = (1.8 \pm 0.5_{\text{stat}} \pm 0.1_{\text{sys}}) \times 10^{22} \text{ y}$
 longest ever measured process
 \simeq one trillion times the age of the Universe



Neutrino Properties from Other Experiments

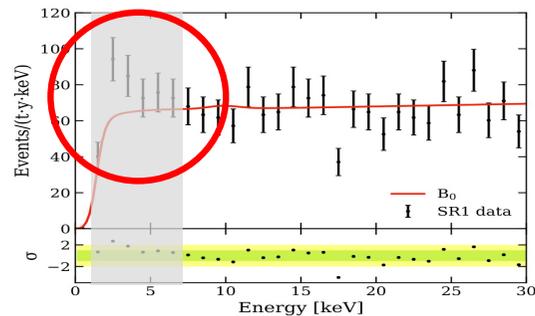
You don't have to observe neutrinos to learn about them

↔ impact of ν 's on other observables

- cosmology
- astrophysics
- $0\nu\beta\beta$ decay
- ...

interconnections of deviations from SM: A recent example

Excess at low E_R in XENON1T



→ triggered many potential explanations

one of them: **a large neutrino magnetic moment**

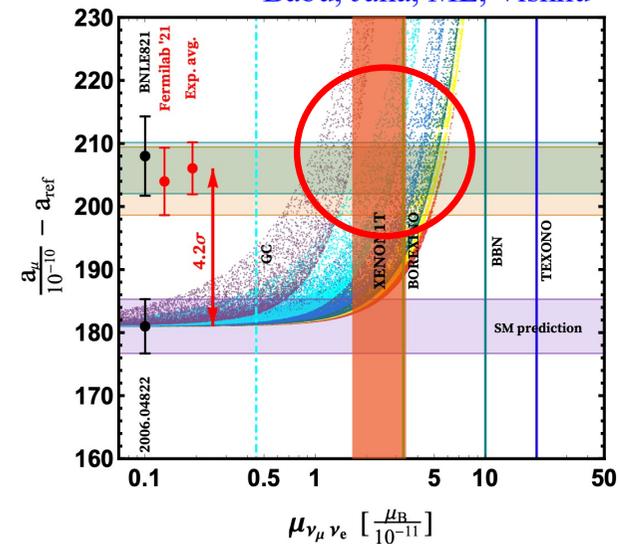
Usually not expected ↔ the same diagrams induce radiatively to high neutrino masses

→ symmetries which decouple μ and m_ν

→ leads interestingly also to the required corrections to $g-2$

deviation of $g-2$ from SM

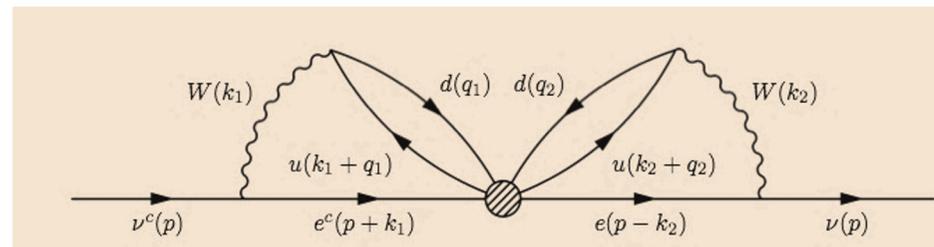
Babu, Jana, ML, Vishnu



Majorana: Pro's and Con's

QFT rules the world:

- include in \mathcal{L} any operator which is not forbidden by symmetry
... will anyway be generated by quantum corrections
- GR anomalies: Any global symmetry like L should be gauged or broken
- L: SM+ = SM+ ν_R allows $\Delta L=2$ terms \rightarrow Majorana terms must exist
- Schechter-Valle theorem: *Any $\Delta L=2$ operator which mediates $0\nu\beta\beta$ decay induces via loops Majorana mass terms \rightarrow Majorana masses unavoidable!*



- Calculate: 4 loops \rightarrow induces $\delta m_\nu = 10^{-25}$ eV Dürr, ML, Merle
 \rightarrow very tiny Majorana mass contribution
 \rightarrow academic interest: cannot explain observed neutrino masses and splitting's
 \rightarrow neutrinos might be predominantly Dirac with a tiny Majorana contribution

- **Global symmetries expected to be broken gravitational effects** $\rightarrow M_{\text{Planck}}$

lepton number $L \rightarrow$ expect at least:

$$m_\nu \approx \frac{v^2}{M_{\text{Planck}}} \approx 10^{-5} \text{ eV}$$

Too small to explain oscillations but too large as subdominant splitting

Connection to matter-antimatter asymmetry

- **Another hierarchy problem**

– Dirac masses, not explained, small radiative corrections natural \leftrightarrow chiral symmetry

– Majorana, e.g. type I see-saw \rightarrow

• naturalness: $M \leq 10^7\text{-}10^8 \text{ GeV}$

• vanilla leptogenesis; $M \geq 10^9 \text{ GeV}$

• similar arguments for other types of see-saw

$$\delta m_H^2 \simeq \frac{y_\nu^2}{16\pi^2} M^2$$

- **A henn and egg problem**

Arguments rest on the fact that n_R does not carry any other charge

\rightarrow well motivated extensions where ν_R has some charge (e.g. horizontal S...)

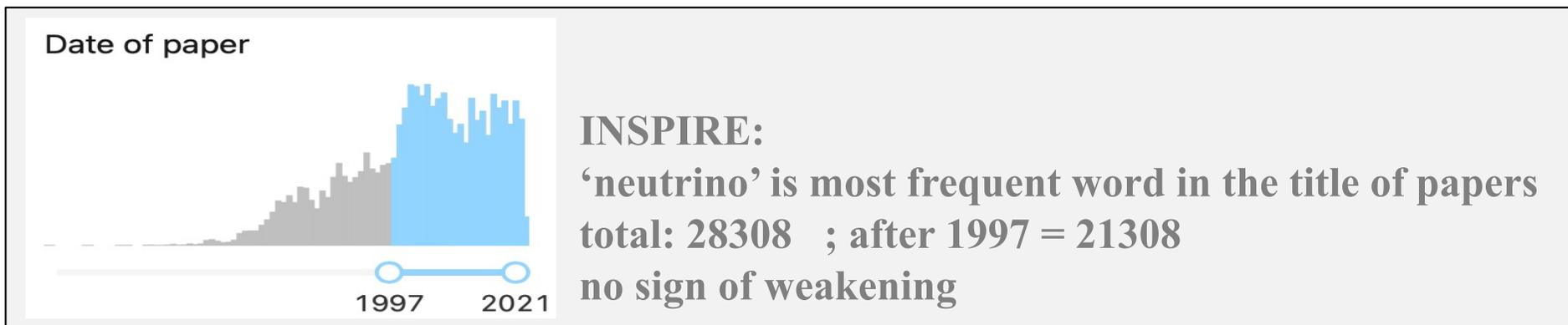
\rightarrow Majorana masses forbidden if symmetry is not explicitly broken (SSB is OK)

\rightarrow models with Dirac masses – not spoiled by loops \leftrightarrow symmetry

Summary

Neutrino physics:

- will remain a hot topic in fundamental physics



- **double beta decay is and will be a very important sub-topic**

To find out if $0\nu\beta\beta$ exists and what BSM mechanism is behind it:

- need better data from different types of experiments
- combine information:
 - $0\nu\beta\beta$ and $2\nu\beta\beta$ (more than one isotope, angular, ...)
 - from other ν -physics (masses, mixings, oscillations)
 - from other fields (e.g. magnetic moment & $g-2$, ...)
 - cosmology (leptogenesis, ...)
- ➔ a lot of experimental work: R&D, ..., construction
- ➔ further theoretical ideas and studies