





Theoretical issues of the SM and physics at the TeV scale

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PANIC - Lisbon - 2021 - On line





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SM@TeV Thanks to all contributors

| Status and future prospects of precision computations for Higgs Physics at the LHC | Prof. Fabrizio Caola |
|---|--------------------------|
| Online | 13:00 - 13:25 |
| Measurement of Higgs differential distributions at the LHC | Arun Kumar |
| Online | 13:25 - 13:45 |
| Higgs couplings to fermions and bosons (including inclusive cross sections, coupling combine Serhat Ördek | nations, spin/CP measure |
| Probing Higgs couplings to light quarks via Higgs pair production | Lina Alasfar |
| Online | 14:05 - 14:25 |
| Combined SMEFT interpretation of Higgs, diboson, and top quark data from the LHC | Juan Rojo 🥝 |
| Online | 14:45 - 15:05 |
| Higgs rare and exotic decays at the LHC | Miha Muskinja |
| Online | 15:05 - 15:25 |
| Full NLO QCD corrections to Higgs-pair production in the Standard Model and beyond | Dr Julien Baglio |
| Online | 15:25 - 15:45 |
| Double Higgs production at the LHC | Louis Portales |
| Online | 15:45 - 16:05 |
| Perspectives for Higgs measurements at Future Colliders | Ang Li |
| Online | 16:20 - 16:40 |
| ILC Higgs Physics Potential | Dr Shin-ichi Kawada |
| Online | 16:40 - 17:00 |
| Nailing Higgs Couplings at Future Colliders | Ayan Paul |
| Online | 17:00 - 17:20 |
| Single boson production overview (W, Z, \vee) at the LHC | Mario Pelliccioni |
| Online | 17:20 - 17:40 |

| Status of NNLO QCD corrections for process with one or more jets in the final state at the LHC | João Pires |
|---|----------------------------|
| Online | 13:00 - 13:25 |
| V+jets/+heavy flavour production at the LHC | Alexandre Laurier |
| Online | 13:25 - 13:45 |
| Jet substructure and fragmentation (including TOP) at the LHC | Andy Buckley |
| Online | 13:45 - 14:05 |
| Modelling the data at the LHC: status and issues (overview including soft QCD and TOP) | Efe Yazgan |
| Online | 14:05 - 14:25 |
| Status of VBS measurements at the LHC | Shuli |
| Online | 14:40 - 15:00 |
| Four-lepton production in gluon fusion at NLO matched to parton showers | Dr Silvia Ferrario Ravasio |
| Online | 15:00 - 15:20 |
| Precise predictions for photon pair production | Alessandro Broggio |
| Online | 15:20 - 15:40 |
| Multibosons production at the LHC (diboson, triboson) | Oleg Kuprash |
| Online | 15:40 - 16:00 |
| Overview of precision measurements (angular coefficients, charge asymmetry, sin20, mW, etc) Vladislav Shalaev | at the LHC |
| Top quark Cross sections overview (including re-interpretation) at the LHC | |
| Online | 16:35 - 16:55 |
| Status of single top measurements at the LHC | Víctor Rodríguez Bouza |
| Online | 16:55 - 17:15 |
| Top quark mass measurements at the LHC | Christoph Garbers |
| Online | 17:15 - 17:35 |
| Top quark properties overview (asymmetries, CP violation, spin correlations, FCNC) at the LHC | Dr Jacob Kempster |
| Online | 17:35 - 17:55 |
| Associated productions with top (t+X, tt+X with X=W,Z,×, heavy-flavours, tt) at the LHC | Tomas Dado |
| Online | 17:55 - 18:15 |



The Standard Model Simplicity





• $SU(3)_c \times SU(2)_L \times U(1)_Y$ gauge symmetries.

• Matter is organised in chiral multiplets of the fund. representation.

• The SU(2) x U(1) symmetry is spontaneously broken to U(1)_{EM}.

Yukawa interactions lead to fermion masses, mixing and CP violation.

• Matter+gauge group => Anomaly free

Neutrino masses can be easily accommodated.

Renormalisable = valid to "arbitrary" high scales.

A number of accidental symmetries,

 $U(1)_{I}, U(1)_{R}, SU(2)_{I} \times SU(2)_{R}, SU(N)_{f}, GIM, \dots$

which allow to explain what we see and we don't see in our exps.



The Standard Model Simplicity



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 $\mathscr{L}_{SM}^{(4)} = -\frac{1}{\Delta} F^{\mu\nu} F_{\mu\nu} + \bar{\psi} i D \psi + (y_{ij} \bar{\psi}_L^i \phi \psi_R^j + h.c.) + |D_\mu \phi|^2 - V(\phi)$

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 $\mathscr{L}_{SM}^{(4)} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \bar{\psi} i D \psi + (y_{ij} \bar{\psi}_L^i \phi \psi_R^j + h.c.) + |D_{\mu} \phi|^2 - V(\phi)$







 $\mathscr{L}_{SM}^{(4)} = -\frac{1}{\Lambda} F^{\mu\nu} F_{\mu\nu} + \bar{\psi} i D \psi + (y_{ij} \bar{\psi}_L^i \phi \psi_R^j + h.c.) + |D_{\mu} \phi|^2 - V(\phi)$







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$$\mathscr{L}_{SM}^{(4)} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \bar{\psi} i D \psi + (y_{ij} \bar{\psi}_L^i \bar{\psi}_L^j \bar{\psi}_L^j$$



Observations:

- No EWBG (Higgs too light, not enough CPV,...)
- No Dark Matter ?
- (g-2)µ ?
- FUV ?



$\phi \psi_R^j + \text{h.c.}) + |D_\mu \phi|^2 - V(\phi)$



} low energy

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$$\mathscr{L}_{SM}^{(4)} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \bar{\psi} i D \psi + (y_{ij} \bar{\psi}_L^i \bar{\psi}_L^j \bar{\psi}_L^j$$



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$\phi \psi_R^j + \text{h.c.}) + |D_\mu \phi|^2 - V(\phi)$







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- Theory predictions seem adeguate. (The key role of MCs is hidden in this plot).







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Precision observables do not point to any clear deviation either.

The most puzzling experimental "issue" of the SM is that we don't really understand why it works so well...

Whatever New Physics might exist to address the SM theoretical shortcomings, its effects must be "small" so that have gone undetected so far.

The main path ahead is twofold

1] Explore the unexplored

2] Increase the precision of TH and EXP to identify possible deviations.





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1] Explore the unexplored

2] Increase the precision of TH and EXP to identify possible deviations.

 $\begin{array}{|c|c|} \alpha_s(M \\ \Delta \alpha_{hs}^{(5)} \\ M_Z \\ M_Z \\ m_t & [0 \\ m_H \\ \hline M_W \\ \hline m_H \\ \hline M_W \\ \hline \Gamma_W & [\\ BR_W \\ BR_W \\ BR_W \\ BR_W \\ \hline P_{\tau}^{pol} \\ \sin^2 \theta \\ \hline \Gamma_Z & [0 \\ \sigma_h^0 & [r \\ R_\ell^0 \\ A_{FB}^0 \\ \hline A_\ell & (S \\ R_b^0 \\ R_c^0 \\ A_{FB}^0 \\ A_{FB} \\ A_b \\ A_c \\ \hline \sin^2 \theta \\ \sin^2 \theta \\ \hline \end{array}$

| | Measurement | Posterior | Prediction | Pull | -3 -2 | 1 0 |
|--|--|--|---|--|--|--------------------------------------|
| $egin{array}{llllllllllllllllllllllllllllllllllll$ | $\begin{array}{c} 0.1177 {\pm} 0.0010 \\ 0.027611 {\pm} 0.000111 \\ 91.1875 {\pm} 0.0021 \\ 172.59 {\pm} 0.45 \\ 125.30 {\pm} 0.13 \end{array}$ | 0.1179 ± 0.0009 0.027572 ± 0.000106 91.1880 ± 0.0020 172.76 ± 0.44 125.30 ± 0.13 | $0.1197 {\pm} 0.0028$ $0.027168 {\pm} 0.000355$ $91.2038 {\pm} 0.0087$ $175.97 {\pm} 1.98$ $112.68 {\pm} 12.89$ | -0.7 1.2 -1.8 -1.7 0.98 | $lpha_S \left(M_Z^2 ight) \ \Delta lpha_{ m had}^{(5)} \left(M_Z^2 ight) \ m_t \ [{ m GeV}] \ m_H \ [{ m GeV}]$ | |
| [GeV] | $80.379 {\pm} 0.012$ | $80.360 {\pm} 0.005$ | $80.355 {\pm} 0.006$ | 1.8 | $M_W [\text{GeV}]$ | |
| [GeV] | $2.085{\pm}0.042$ | $2.0883 {\pm} 0.0006$ | $2.0883 {\pm} 0.0006$ | -0.08 | M_{Z} [GeV] | |
| $V \rightarrow had$ $V \rightarrow \ell \nu$ | $0.6741 {\pm} 0.0027$ $0.1086 {\pm} 0.0009$ | $\begin{array}{c} 0.67486 {\pm} 0.00007 \\ 0.10838 {\pm} 0.00002 \end{array}$ | $0.67486 {\pm} 0.00007$ $0.10838 {\pm} 0.00002$ | -0.28 0.24 | $\Gamma_Z ~[{ m GeV}] \ \sigma^0_{ m had} ~[{ m nb}]$ | |
| $=A_{\ell}$ | $0.1465{\pm}0.0033$ | $0.1473 {\pm} 0.0004$ | $0.1473 {\pm} 0.0005$ | -0.23 | R_ℓ | |
| $	heta_{ m eff}^{ m lept}(Q_{ m FB}^{ m had})$ | $0.2324{\pm}0.0012$ | $0.23149 {\pm} 0.00006$ | $0.23149 {\pm} 0.00006$ | 0.91 | $A_{FB}^{0,\ell}$ $ppol$ | |
| GeV] ıb] | 2.4955 ± 0.0023 41.4802 ± 0.0325 20.7666 ± 0.0247 0.0171 ± 0.0010 | 2.4945 ± 0.0006 41.4910 ± 0.0076 20.750 ± 0.0080 0.01627 ± 0.00010 | 2.4943 ± 0.0007 41.4930 ± 0.0080 20.7460 ± 0.0087 0.01626 ± 0.00010 | 0.50 -0.38 0.79 0.84 | $A_{\ell} (SLD)$ A_{c} A_{b} A_{b} | |
| SLD) | $\begin{array}{c} 0.1513 {\pm} 0.0021 \\ 0.21629 {\pm} 0.00066 \\ 0.1721 {\pm} 0.0030 \\ 0.0992 {\pm} 0.0016 \\ 0.0707 {\pm} 0.0035 \\ 0.923 {\pm} 0.020 \\ 0.670 {\pm} 0.027 \end{array}$ | $\begin{array}{c} 0.14727 {\pm} 0.00045 \\ 0.21588 {\pm} 0.00010 \\ 0.17221 {\pm} 0.00005 \\ 0.1032 {\pm} 0.0003 \\ 0.0738 {\pm} 0.0002 \\ 0.93475 {\pm} 0.00004 \\ 0.6679 {\pm} 0.0002 \end{array}$ | $\begin{array}{c} 0.14731 {\pm} 0.00047 \\ 0.21587 {\pm} 0.00010 \\ 0.17221 {\pm} 0.00005 \\ 0.10327 {\pm} 0.00033105 \\ 0.0738 {\pm} 0.0002 \\ 0.93475 {\pm} 0.00004 \\ 0.6679 {\pm} 0.0002 \end{array}$ | 1.9 0.63 -0.04 -2.5 -0.88 -0.59 0.08 | $\begin{array}{c c} A_{FB}^{0,c} \\ A_{FB}^{0,b} \\ R_{c}^{0} \\ R_{b}^{0} \\ \sin^{2} \theta_{\text{eff}}^{\ell}(Q_{FB}^{\text{had}}) \\ \sin^{2} \theta_{\text{eff}}^{\text{lept}} (\text{Tev/LHC}) \end{array}$ | |
| $	heta_{ m eff}^{ m lept}(m Tev/LHC)$ | $0.23137 {\pm} 0.00022$ | 0.23149 ± 0.00006 | $0.23150 {\pm} 0.00006$ | -0.57 | HEPfit | $Pull = \frac{O_{exp}}{\sigma_{ex}}$ |

[Courtesy of De Blas et al., work in progress]







$$i m_f / v$$

$$igm_W g_{\mu
u} = 2i v g_{\mu
u} \cdot m_W^2 / v^2$$

$$g \frac{m_Z}{\cos \theta_W} g_{\mu\nu} = 2i v g_{\mu\nu} \cdot m_Z^2 / v^2$$

Unique mass generation mechanism for fermions and vectors.

| | | [AT | LAS | 2020 | L | | |
|--|---|------|--------|---------------|---------------------------------------|-----------------|----------------------|
| | | | | | | | |
| ATLAS Pre √s = 13 TeV, 24.5 m _µ = 125.09 GeV | liminary - 139 fb⁻¹ , v_, < 2.5 | ⊢⊷IT | otal 🕻 | Sta | ıt. 💳 S | Syst. | I SI |
| p _{SM} = 87% | Ϋ́Η' | | | | Total | Stat. | Syst. |
| ggF γγ | è. | | | 1.03 | 3 ± 0.11 (| $\pm \; 0.08$, | $^{+0.08}_{-0.07}$) |
| ggF ZZ | eļ 👘 | | | 0.94 | 4 ^{+0.11} _{-0.10} (| ±0.10, | ± 0.04) |
| ggF WW | ÷ | | | 1.08 | B ^{+0.19} _{-0.18} (| ±0.11, | ± 0.15) |
| ggF ττ | H ata h | | | 1.02 | 2 ^{+0.60} _{-0.55} (| +0.39 -0.38, | $^{+0.47}_{-0.39}$) |
| ggF comb. | ė. | | | 1.00 |) ± 0.07 (| ± 0.05 , | ± 0.05) |
| VBF γγ | I ser i | | | 1.3 | 1 ^{+0.26} _{-0.23} (| +0.19 -0.18, | $^{+0.18}_{-0.15})$ |
| VBF ZZ | | | | 1.2 | 5 ^{+0.50} _{-0.41} (| +0.48 -0.40, | $^{+0.12}_{-0.08}$) |
| VBF WW | | | | 0.60 | -0.36 - 0.34 | +0.29 -0.27, | ±0.21) |
| VBF ττ | H inter H | | | 1.1 | 5 ^{+0.57} _{-0.53} (| +0.42 -0.40, | $^{+0.40}_{-0.35}$) |
| VBF bb | | - | | — 3.03 | 3 ^{+1.67} _{-1.62} (| +1.63 -1.60, | +0.38 -0.24) |
| VBF comb. | | | | 1.1 | 5 ^{+0.18} _{-0.17} (| ±0.13, | $^{+0.12}_{-0.10})$ |
| VH γγ | | | | 1.32 | 2 ^{+0.33} _{-0.30} (| +0.31 -0.29, | $^{+0.11}_{-0.09})$ |
| VH ZZ | | | | 1.53 | 3 ^{+1.13} _{-0.92} (| +1.10 -0.90, | +0.28 -0.21) |
| VH bb | | | | 1.02 | 2 ^{+0.18} _{-0.17} (| ±0.11, | +0.14 -0.12) |
| VH comb. | I | | | 1.1(|) ^{+0.16} _{-0.15} (| ±0.11, | +0.12 -0.10) |
| ttH+tH γγ | e p | | | 0.90 | -0.27 (-0.24) | +0.25 -0.23, | $^{+0.09}_{-0.06}$) |
| ttH+tH VV | + | - | | 1.72 | 2 ^{+0.56} _{-0.53} (| +0.42 -0.40, | $^{+0.38}_{-0.34}$) |
| ttH+tH ττ ⊢ | ÷ | - | | 1.20 |) ^{+1.07} _{-0.93} (| +0.81 -0.74, | $^{+0.70}_{-0.57}$) |
| ttH+tH bb 🛏 | | | | 0.79 | $9 + 0.60 \\ - 0.59$ (| ±0.29, | +0.52 -0.51) |
| <i>ttH+tH</i> comb. | ÷. | | | 1.10 |) ^{+0.21} _{-0.20} (| +0.16 -0.15, | +0.14 -0.13) |
| | | | | | | | |
| -2 0 | 2 | 2 | 4 | | 6 | | 8 |









$$i m_f / v$$

$$igm_W g_{\mu
u} = 2i v g_{\mu
u} \cdot m_W^2 / v^2$$

$$g \frac{m_Z}{\cos \theta_W} g_{\mu\nu} = 2i v g_{\mu\nu} \cdot m_Z^2 / v^2$$

Unique mass generation mechanism for fermions and vectors.

| A | TLAS Pr | eliminar | ່ | Total | Stat | | Svet | |
|------|-------------------|--|-----|-------|--------------|----------------------|-----------------|----------------------|
| √s = | = 13 TeV, 24 | .5 - 139 fb ⁻¹ | | Total | Jiai | | <i>y</i> or. | 0 |
| | = 125.09 Ge | eV, y _H < 2. | 5 | | | | | |
| PSM | = 07 /0 | | | | | Total | Stat. | Syst. |
| gg | F γγ | e | | | 1.03 | ±0.11 (| $\pm \; 0.08$, | $^{+0.08}_{-0.07}$) |
| gg | F <i>ZZ</i> | e | | | 0.94 | +0.11 -0.10 (| ± 0.10 , | ± 0.04) |
| gg | F WW | ÷ | | | 1.08 | +0.19 -0.18 (| ±0.11, | ±0.15) |
| gg | F ττ | - | 1 | | 1.02 | + 0.60 - 0.55 (| +0.39 -0.38, | +0.47 -0.39) |
| gg | F comb. | ę | | | 1.00 | ± 0.07 (| $\pm \; 0.05$, | ± 0.05) |
| VE | BF γγ | H | | | 1.31 | +0.26 -0.23 (| +0.19 -0.18, | +0.18 -0.15) |
| VE | 3F <i>ZZ</i> | ⊢ ■− | 9 | | 1.25 | +0.50 -0.41 (| +0.48 -0.40, | +0.12 -0.08) |
| VE | BF WW | H | | | 0.60 | +0.36 -0.34 (| +0.29 -0.27, | ±0.21) |
| VE | 3F ττ | - | 4 | | 1.15 | +0.57 -0.53 (| +0.42 -0.40, | $^{+0.40}_{-0.35}$) |
| VE | 3F <i>bb</i> | E | = | - | 3 .03 | +1.67 -1.62 (| +1.63 -1.60, | +0.38 -0.24) |
| VE | 3F comb. | I | | | 1.15 | +0.18 -0.17 (| ±0.13, | +0.12 -0.10) |
| Vł | Η γγ | |) | | 1.32 | +0.33 -0.30 (| +0.31 -0.29, | +0.11 -0.09) |
| VH | + ZZ | | | | 1.53 | +1.13 -0.92 (| +1.10 -0.90, | +0.28 -0.21) |
| Vł | l bb | | | | 1.02 | +0.18 -0.17 (| ±0.11, | +0.14 -0.12) |
| Vł | d comb. | • | | | 1.10 | +0.16 -0.15 (| ±0.11, | +0.12 -0.10) |
| ttŀ | I+tH γγ | ÷ | | | 0.90 | +0.27 -0.24 (| +0.25 -0.23, | $^{+0.09}_{-0.06}$) |
| ttŀ | I+tH VV | H H | • • | | 1.72 | +0.56 -0.53 | +0.42 -0.40, | $^{+0.38}_{-0.34}$) |
| ttF | l+tH ττ | - | | | 1.20 | +1.07 -0.93 (| +0.81 -0.74, | $^{+0.70}_{-0.57}$) |
| tt⊦ | l+tH bb | Here in the second seco | | | 0.79 | $^{+0.60}_{-0.59}$ (| $\pm \; 0.29$, | +0.52 -0.51) |
| ttŀ | <i>I+tH</i> comb. | ÷ | | | 1.10 | +0.21 -0.20 (| +0.16 -0.15, | +0.14 -0.13) |
| | | | | | | | | |
| 2 | 0 | | 2 | Λ | | 6 | | Q |
| -2 | 0 | | 2 | 4 | | 0 | | 0 |

[ATLAS 2020]









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V(H

 $V^{\rm SM}$







V(H

 $V^{\rm SM}$

$$-3iv \cdot m_h^2/v^2$$







V(H

 V^{SM}

$$\begin{aligned} T(\Phi) &= \frac{m_H^2}{2} H^2 + \lambda_3 v H^3 + \frac{\lambda_4}{4} H^4 + \dots \\ T(\Phi) &= -\mu^2 (\Phi^{\dagger} \Phi) + \lambda (\Phi^{\dagger} \Phi)^2 \implies \begin{cases} v^2 &= \mu^2 / \lambda \\ m_H^2 &= 2\lambda v^2 \end{cases} \quad \begin{cases} \lambda_3^{\rm SM} &= \lambda \\ \lambda_4^{\rm SM} &= \lambda \end{cases} \end{aligned}$$



 $-3 iv \cdot m_h^2/v^2$





$$\begin{aligned} f(\Phi) &= \frac{m_H^2}{2} H^2 + \lambda_3 v H^3 + \frac{\lambda_4}{4} H^4 + \dots \\ f(\Phi) &= -\mu^2 (\Phi^{\dagger} \Phi) + \lambda (\Phi^{\dagger} \Phi)^2 \implies \begin{cases} v^2 &= \mu^2 / \lambda \\ m_H^2 &= 2\lambda v^2 \end{cases} \quad \begin{cases} \lambda_3^{\rm SM} &= \lambda \\ \lambda_4^{\rm SM} &= \lambda \end{cases} \end{aligned}$$

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Reicher



One of the flagship measurements foreseen for the HL-LHC. [Di Micco et al., 1910.00012]



The precision path Two questions

1. What is the expected experimental precision on key SM measurements at the TeV scale and a reasonable goal for the corresponding TH predictions? Are we there yet?

2. How to frame and interpret our results to maximise the sensitivity to New Physics?



Towards the HL-LHC

- 20-fold data sample
- 1/5 statistical uncertainties
- Comparable reduction of systematic uncertainties?
- Definition of tails and access to rare processes















$$\frac{\kappa_i^2\cdot\kappa_f^2}{\kappa_H^2}$$













 \rightarrow













Currently limits on k_{λ} from H and HH are comparable and will stay so at the HL-LHC. Borderline sensitivity to say something about EW baryogenesis...



Precision calculations for the LHC Status



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Precision calculations for the LHC The path

"Rules of thumb at the LHC":

- Predictions must be calculated at least to **NLO QCD** to control the central value at 10-20%.
- **N2LO QCD** provides control at 5% level and on the uncertainties stabilizing \bullet the perturbative expansion.
- **N2LO QCD** is expected to be of the same order as NLO EW $\alpha_S^2 \sim \alpha_W$, yet • **EW** corrections grow large and negative at high energies (Sudakov logs).
- **N3LO QCD** is the frontier of precision aiming ~1% of MHO uncertainties.
- **Resummation** Universal, all-order terms that are potentially large for some \bullet observables (logs or 1PI loops for propagators) need to be resummed. They might refer to global or non-global observables. Resummation leads to mprovements in precision and accuracy.











Precision calculations for the LHC N3LO revolution

[Anastasiou et al., 1602.00695]



Table 2: Gluon fusion Higgs boson production cross sections and uncertainties as a function of the pp collider energy.

| \sqrt{s} | σ | δ (theory) | $\delta(ext{PDF})$ | $\delta(\alpha_s)$ |
|------------|-------------|-----------------------------------|--|---|
| 13 TeV | √ 48.61 pb | +2.08 pb +4.27% -3.15 pb (-6.49%) | $\left(\pm 0.89 \mathrm{pb} (\pm 1.85\%)\right)$ | +1.24 pb $(+2.59%)$ $-1.26 pb$ $(+2.62%)$ |
| 14 TeV | √ 54.72 pb | +2.35 pb $+4.28%-3.54pb -6.46\%$ | $\left(\pm 1.00 \text{pb} (\pm 1.85\%)\right)$ | +1.40 pb $+2.60%$ $-2.62%$ |
| 27 TeV | V 146.65 pb | +6.65 pb +4.53% -9.44 pb +4.53% | $\left(\pm 2.81 \text{pb} \left(\pm 1.95\%\right)\right)$ | +3.88 pb $(+2.69%)$ $-3.82 pb$ $(-2.64%)$ |

• Drell-Yan now available [Duhr, Dulat and Mistelberger, 2001.07717] [Duhr, Dulat and Mistelberger, 2007.13313]

Collider Energy / TeV

δ(EW)

 Very significant reduction of MHO uncertainties.

 Differential distributions are available.

 $\delta(\text{PDF}+\alpha_s)$

60

 δ (PDF–TH)

80

• Uncertainty budget points to PDF as the main source of error.





40

 $\delta(1/m_t)$

 $\delta(t,b,c)$

 δ (scale)

20

 $\delta_i/\delta_{\rm f}$

Precision calculations for the LHC Fully exclusive simulations



the 2012 Sakurai Prize for Theoretical Particle Physics by the American Physical Society, along with the late Guido Altarelli.

Brvan Webber (left) and

Torbjörn Sjöstrand (right).

Credit: Lund University, T

Siöstrand

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[Mazzitelli et al. , 2012.14267]



Precision calculations for the LHC Fully exclusive simulations



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Brvan Webber (left) and

Torbjörn Sjöstrand (right).

Credit: Lund University, T

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Precision calculations for the LHC Status: PDF's



- Complete N3LO PDF's evolution not available yet. Non-singlet evolution available at 4 loops already.
- Error budget with many sources. MHO uncertainties yet to be included in the final assessment.
- Reaching 1% will be very challenging.
- Room for a breakthrough from lattice?







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- Very fast progress in conceptual as well as technical aspects.
- Tight and consolidated community, with high momentum.
- Considering the status of 20 years ago seems clear that NNLO will be completed and N3LO will start to become available for $2\rightarrow 2$ (see 3-loop $q\bar{q} \rightarrow \gamma\gamma$ results)
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 Analytically historically matching the FO accuracy.

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 Having a NLL and beyond PS, is being explored now. To be seen.

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The precision path Two questions

1. What is the expected experimental precision on key SM measurements at the TeV scale and a reasonable goal for the corresponding TH predictions? Are we there yet?

2. How to frame and interpret our results to maximise the sensitivity to New Physics?





Three key properties of the SM:

- Mass generation with gauge invariance
- Unitarity (up to a predefined Λ)
- Perturbativity/renormalizability



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Rattazzi® adapted



UCLouvain Fabio Maltoni





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 $\mathscr{L}^{(4)}$ $\mathscr{L}^{(2)}$ \mathscr{L} + $m_{v} = 0$ $U(1)_L^3 \times U(1)_B$ GIM $Y_u, Y_d, Y_l \Rightarrow$ Flayor & \mathcal{P} Rattazzi® adapted







Rattazzi® adapted





Rattazzi® adapted

$$\frac{1}{\Lambda} \mathscr{L}^{(5)} + \frac{1}{\Lambda^2} \mathscr{L}^{(6)} + \dots$$

$$U(\Lambda)_L \to m_\nu \neq 0 \qquad \Rightarrow \Lambda \ge 10^{14}$$
Flavor $\Rightarrow \mu \to e\gamma, \Delta m_K, \dots$

$$CP' \Rightarrow edm's \qquad \Rightarrow \Lambda \ge 10^6$$
Dipoles $\Rightarrow (g - 2)_\mu$

$$U(1)_B \Rightarrow p \to \pi^0 e^+ \qquad \Rightarrow \Lambda \ge 10^{15}$$

$$\Rightarrow \Lambda \ge 10^{15}$$







| Λ_{UV} | |
|----------------|--|
| | |

One can satisfy all the previous requirements, by building an EFT on top of the SM that respects the gauge symmetries:

$$\mathscr{L}_{\text{SMEFT}} = \mathscr{L}_{\text{SM}}^{(4)} + \frac{1}{\Lambda^2} \sum_{i}^{N_6} c_i \mathcal{O}_i^{(6)} + \frac{1}{\Lambda^4} \sum_{j}^{N_8} \mathcal{O}_j^{(6)}$$

With the "only" assumption that all new states are heavier than energy probed by the experiment $\sqrt{s} < \Lambda$.

The theory is renormalizable order by order in $1/\Lambda$, perturbative computations can be consistently performed at any order, and the theory is predictive, i.e., well defined patterns of deviations are allowed, that can be further limited by adding assumptions from the UV. Operators can lead to larger effects at high energy (for different reasons).

- $C_i \mathcal{O}_i^{(8)}$





Energy helps precision





The master equation of an EFT approach has three key elements:

$$\Delta Obs_n = Obs_n^{\mathsf{EXP}} - Obs_n^{\mathsf{SM}} = \frac{1}{\Lambda^2} \sum_i a_{n,i}^{(6)}(\mu) c_i^{(6)}(\mu) + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$



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Most precise/accurate experimental measurements with uncertainties and correlations



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Most precise EFT predictions



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increased NP Sensitivity

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Most precise EFT predictions

increased NP Sensitivity \Rightarrow increased UV identification power

The way of SMEFT A simple approach

EFT bounds translate to constraints on parameters of UV models

Simplest case: single-field extensions of the SM

Mass limits (in TeV)

[Ellis et al. 2012.02779]

A powerful approach Is this easy?

It's as exciting as challenging. Pattern of deformations enter many observables in a correlated way.

Needs to manage complexity, uncertainties and correlations.

Needs coordinated work among analysis groups in collaborations traditionally working separately (top, Higgs, EW,...)

Needs coordinated work between theorists and experimentalists (model dependence, validity, interpretations, matching to the UV).

A new paradigm: shifting value from "the best single measurement" to "the best combinable measurement"!

A powerful approach What are we going to learn?

[Peskin, ICHEP2020]

Global fits First explorations: EWPO+H+EW+Top

- Already now and without a dedicated experimental effort there is considerable information that can be used to set limits:
- Fitmaker [J. Ellis, M. Madigan, K. Mimasu, V. Sanz, T. You 2012.02779]
- SMEFIT [J. Either, G. Magni, F. M., L. Mantani, E. Nocera, J. Rojo, E. Slade, E. Vryonidou, C. Zhang, 2105.00006]
- SFitter [Biekötter, Corbett, Plehn, 2018] + [I. Brivio, S. Bruggisser, F. M., R. Moutafis, T. Plehn, E. Vryonidou, S. Westhoff, C. Zhang, 1910.03606] (separated)
- HEPfit [de Blas, et al. 2019]
- 30+ operators, linear and/or quadratic fits, Higgs/Top/EW at LHC, WW at LEP and EWPO.

Global fits Workflow

Methodology

EleniVryonidou®

Data 317 data points: Top: ttbar, single-top, associated top production, distributions. Higgs production and decay, differential distributions, STXS. Diboson production, distributions Global EW/Top/Higgs SMEFT fit Fit results can be used to bound specific UV complete models

New data can be straightforwardly added

Output

Global fits Operators vs processes

Global EW(PO)+H+Top **Examples**

[Ellis et al. 2012.02779]

34 operators, $SU(2)^2 \times SU(3)^3$

EWPO fitted, 341 data points

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36 operators, $SU(2)^2 \times SU(3)^3$

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Global fits: now vs future EWPO+EW+Higgs

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The SM at the TeV Conclusions

- SM at the TeV scale is in an extremely good shape. No signs of significant deviations have been detected.
- Tremendous improvements in the accuracy/precision of SM predictions have been achieved, opening a new realm of opportunities.
- The LHC campaign of precision measurements is entering a new phase measuring at unprecedented precision a large number of channels and accessing for the first time rare final states.
- A far reaching approach to interpreting SM measurements is to constrain the SM interactions at the TeV scale (and beyond) by employing the SMEFT, maximising sensitivity to heavy new physics.
- Considerable theory effort going on, being matched by the experimental work.
- EFT's are also being used to gauge sensitivity to NP at future colliders.
- Busy future ahead with even more integrated TH/EXP activities.

