

# New Target and Probes of Sterile Neutrino Dark Matter

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# Outline

Introduce a gauge singlet fermion, mix it with SM neutrinos

$$\nu_4 = \cos \vartheta \nu_s + \sin \vartheta \nu_a$$

Flavor eigenstates:  $\nu_a$  active, weakly interacting,  $\nu_s$  pure singlet.  
 $\vartheta$  is vacuum mixing angle.

[DM relic abundance target](#)

[Experimental probes](#)

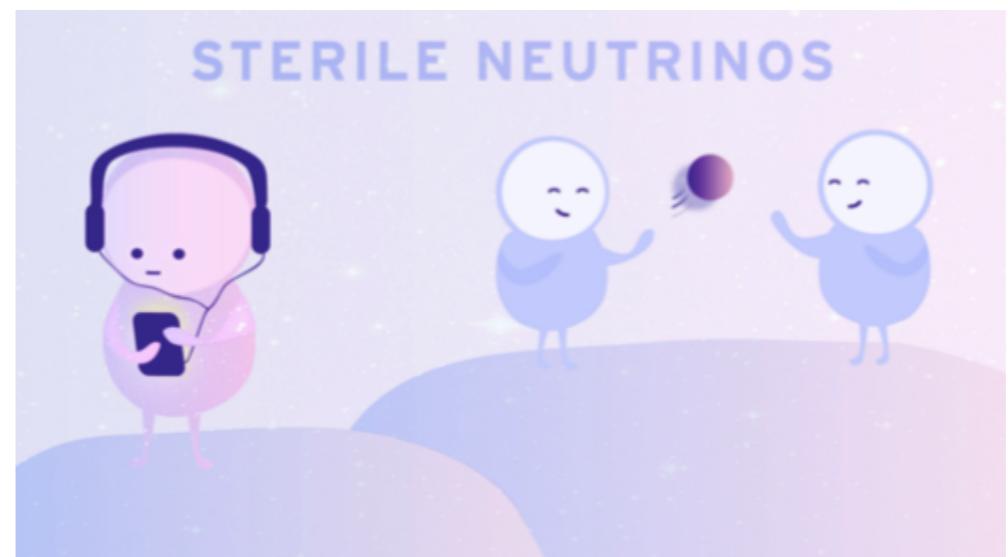
# Non-Thermal

Fully thermalizing  $\nu_4$  with SM sector overclose the universe:

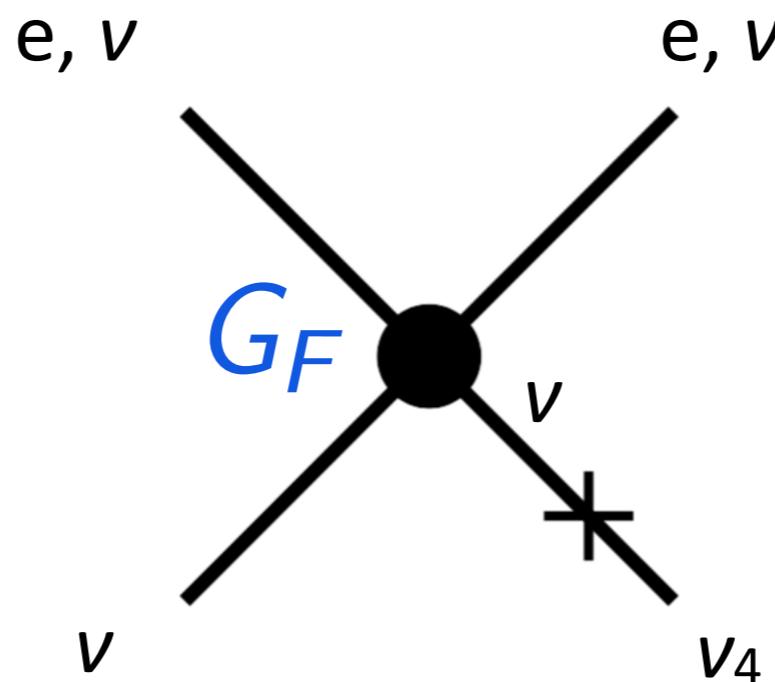
$$\Omega_4 \sim 10 \left( \frac{m_4}{\text{keV}} \right)$$

$\nu_4$  heavier than keV from dwarf galaxies (Tremaine, Gunn 1979).

Must be produced in a non-thermal way with a small  $\vartheta \ll 1$ .



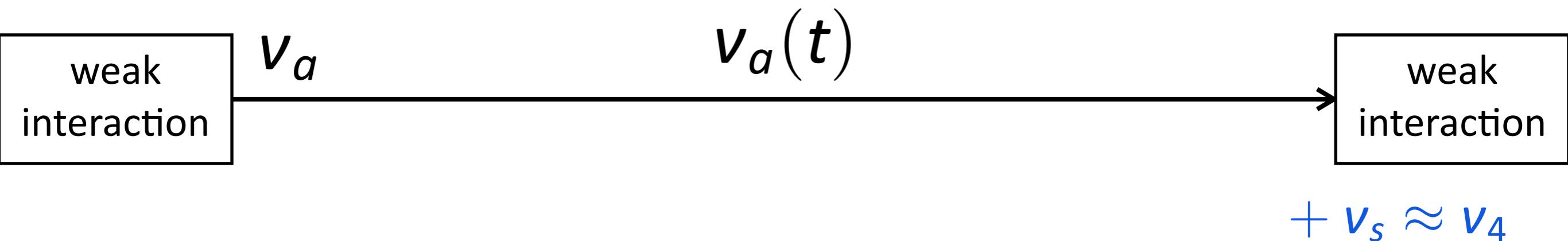
# Dodelson-Widrow Mechanism



Tiny mixing angle  $\vartheta$  controls the relic density.

hep-ph/9303287

# Neutrino Oscillation in Early Universe



Two time scales:

In the thermal bath, neutrino after produced remains coherent state until destroyed.

In between, active-sterile neutrino oscillation occurs.

# The Key Equation

$$\frac{df_4}{d \log(1/T)} = \frac{\Gamma}{2H} P_{\nu_a \rightarrow \nu_4} f_a$$

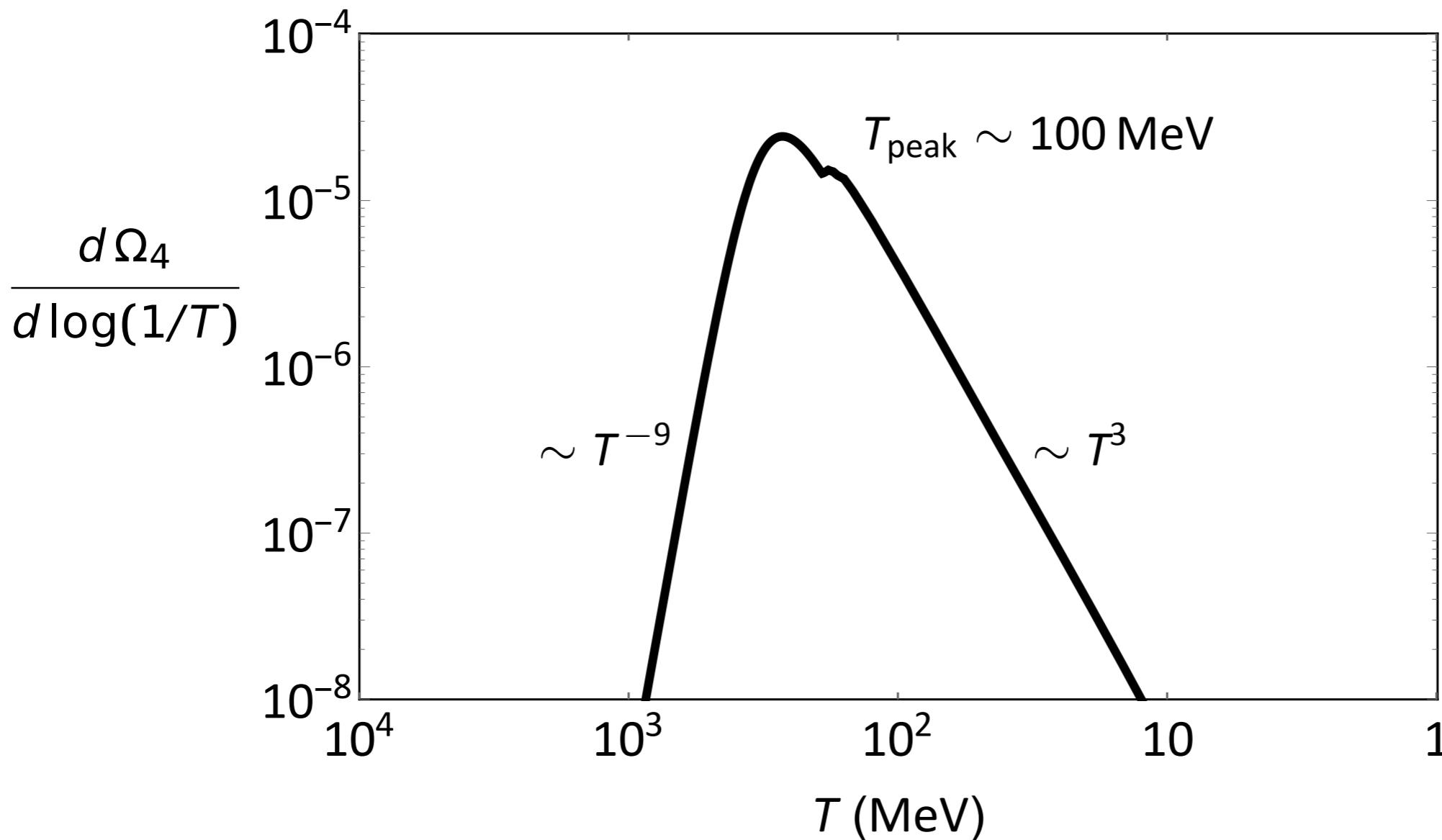
$\Gamma/H$  : Counts number of cycles for the above process to repeat before neutrino decoupling.

Roughly

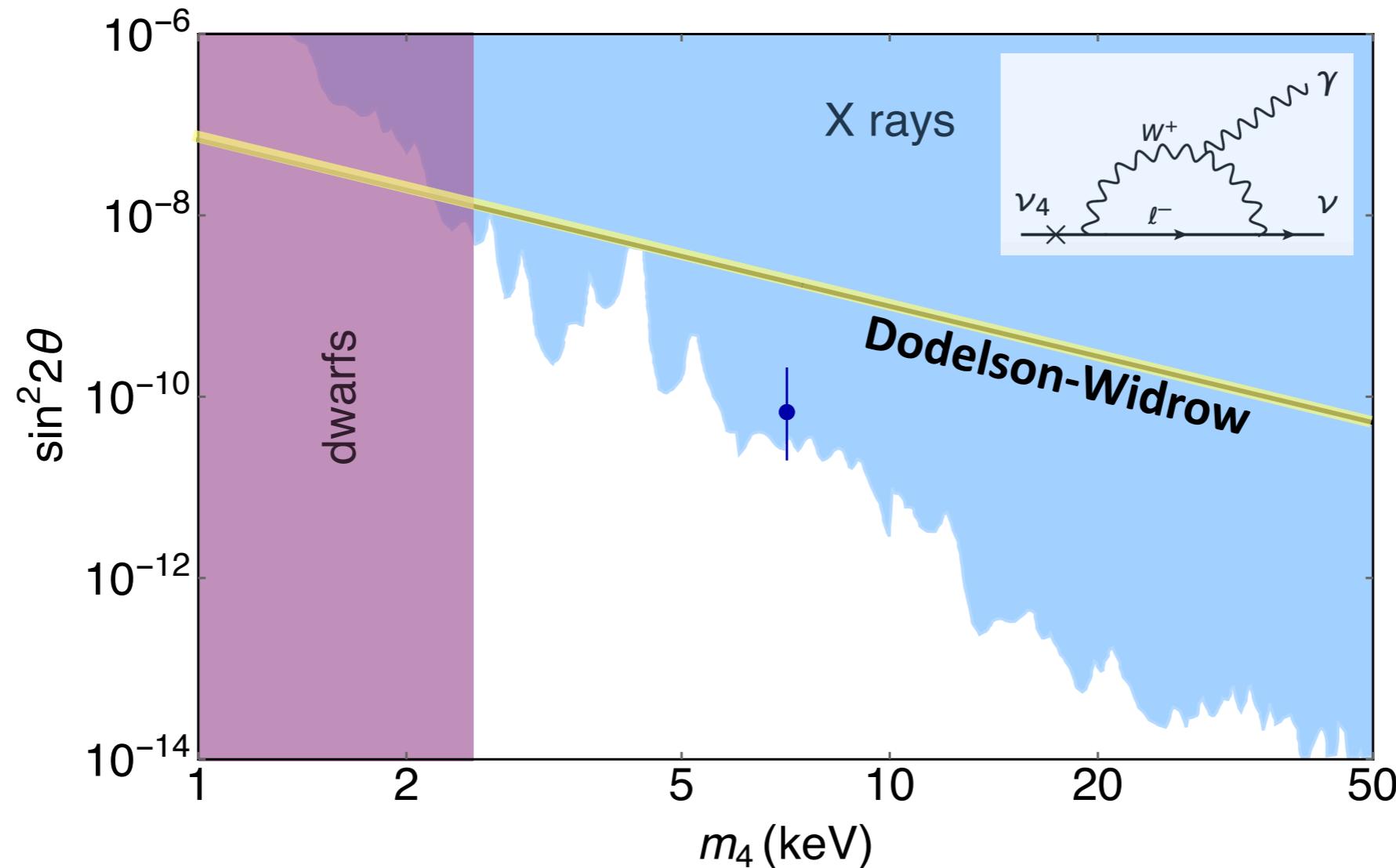
$$P_{\nu_a \rightarrow \nu_4} \sim \sin^2 2\theta$$

Subtlety,  $P_{\nu_a \rightarrow \nu_4} \sim \sin^2 2\theta$  is highly suppressed for very large  $\Gamma$ .

# Production Time Window



# Severely Constrained



Abazajian (1705.01837)

# A Simple Idea

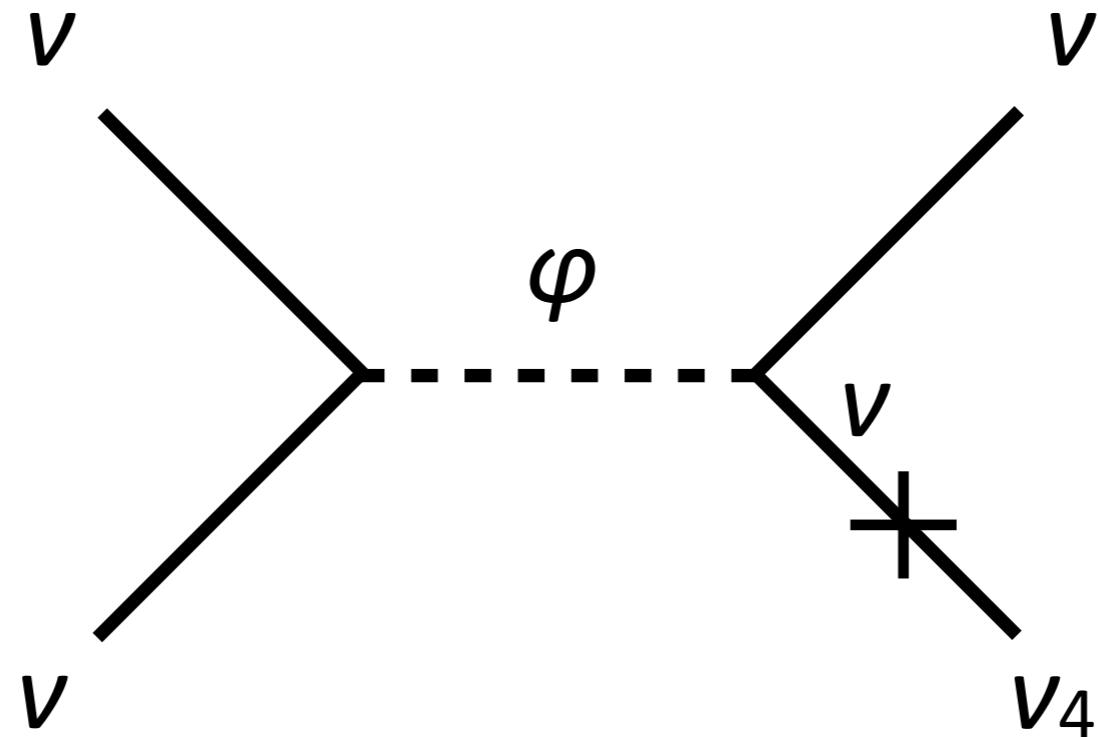
$$\Omega_4 [\text{DW}] \sim \Gamma_w \times \sin^2 2\vartheta$$

$\Gamma_w$  given by weak interaction; X-ray constrains  $\vartheta$ .

Intuition: compensate smaller mixing with larger reaction rate.

BSM neutrino interactions with known particles (e.g. electron, photon) tightly constrained. **Resort to neutrino self interaction: never directly measured, allowed to be much stronger.**

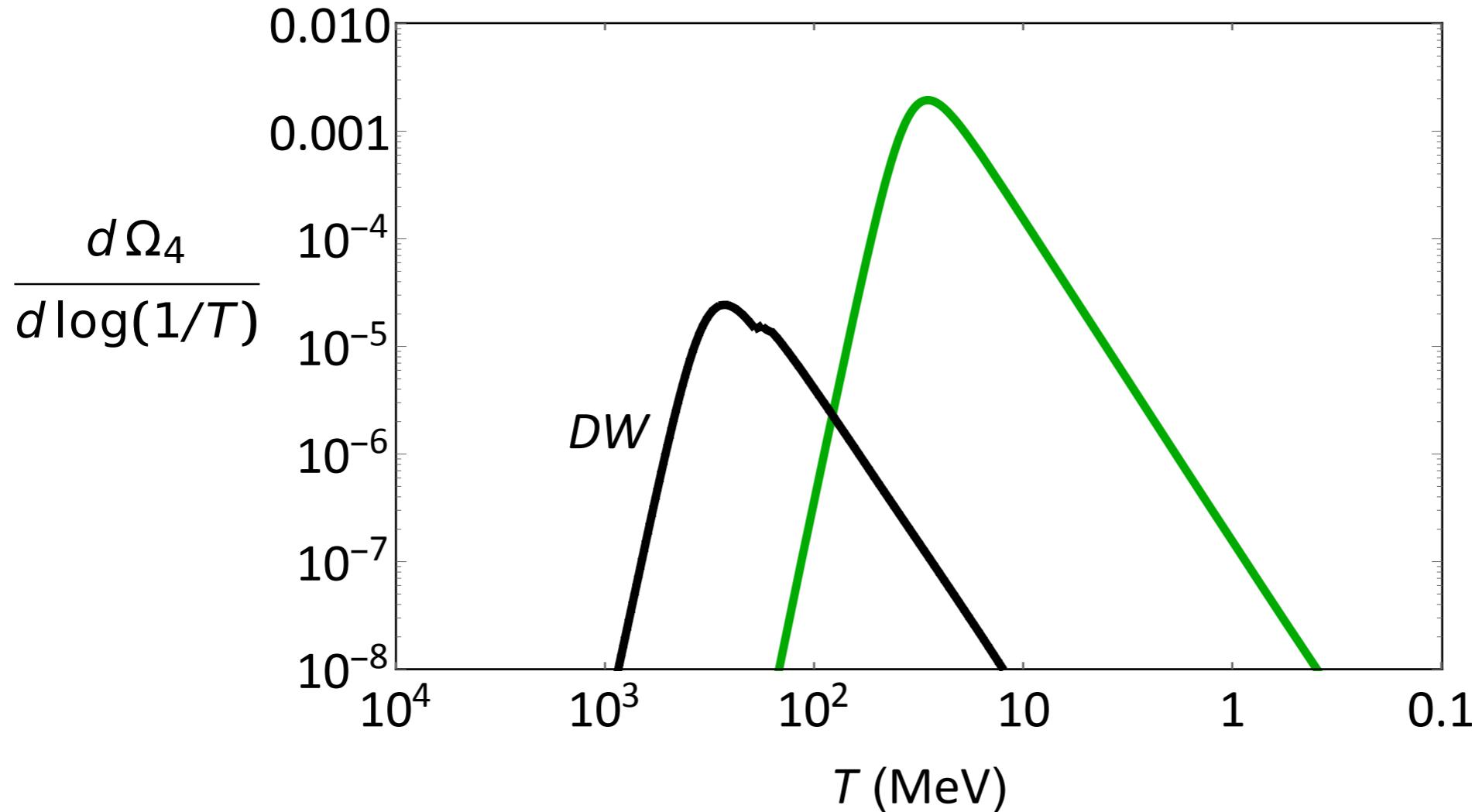
# Majoron-like Scalar Mediator



New dark matter production model

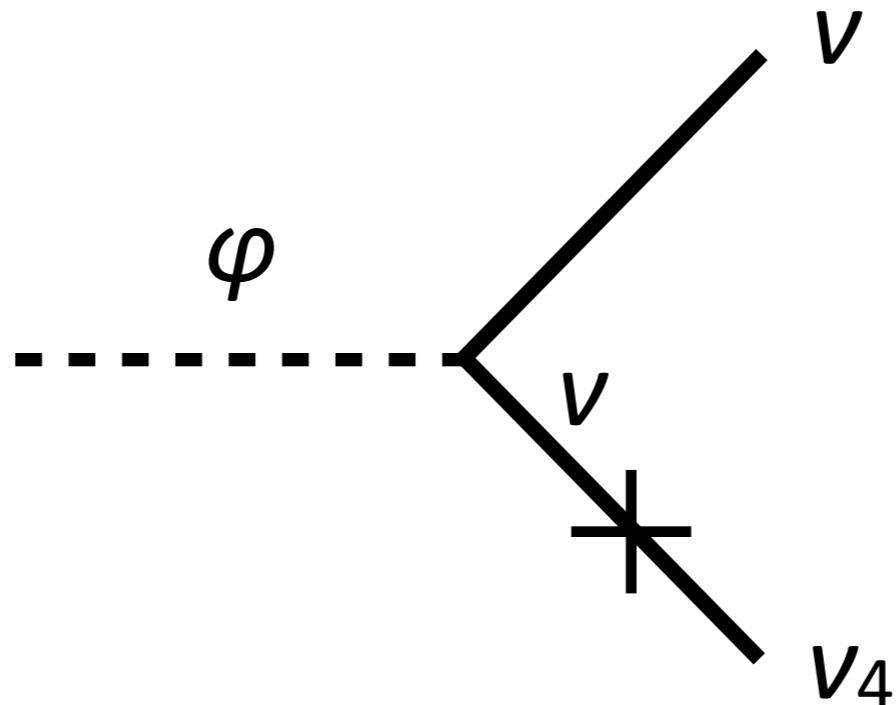
de Gouvêa, Sen, Tangarife, YZ (1910.04901)

# Case of Heavy Mediator



Final relic density:  $\Omega_4 \propto \frac{\lambda^3}{m_\phi^2} \gg \frac{g^3}{M_W^2}$

# Case of Light Mediator

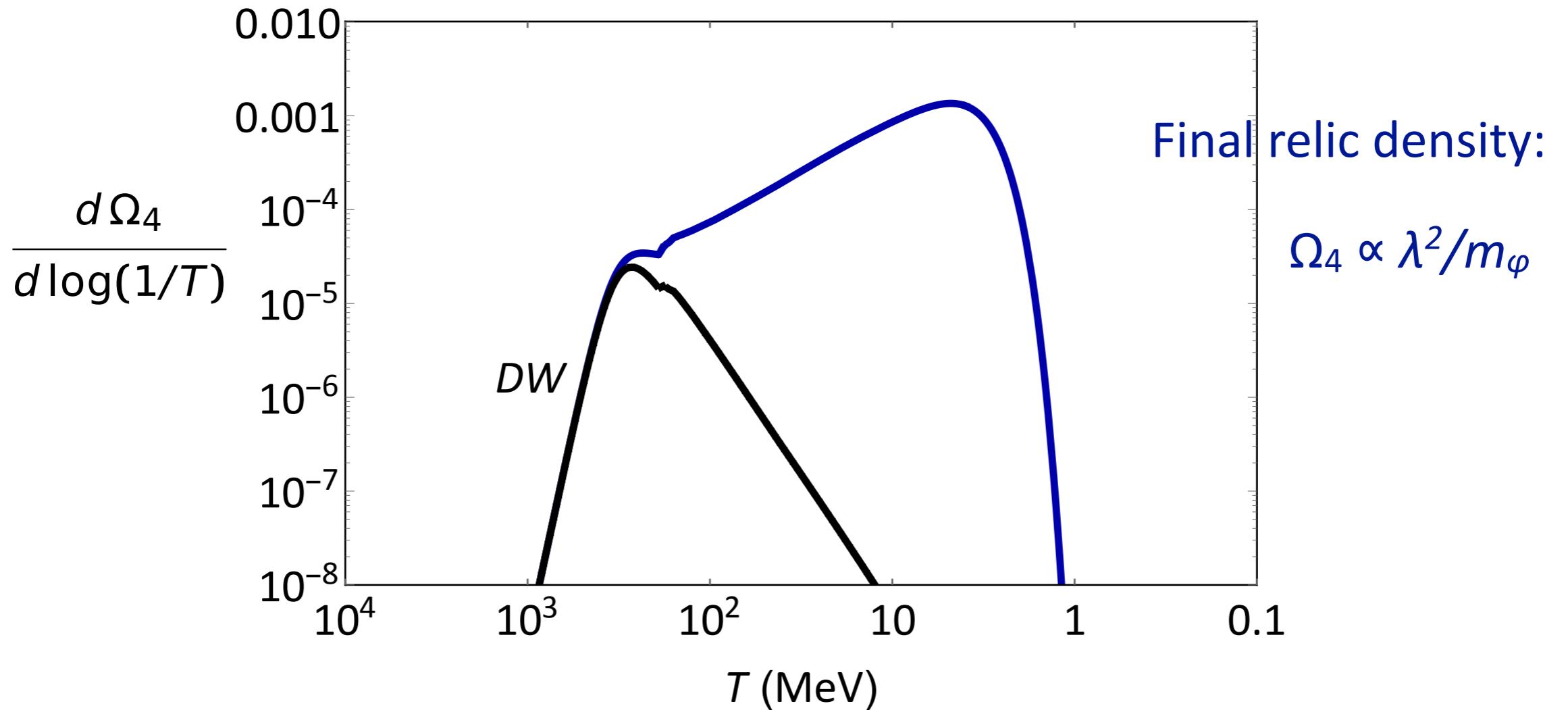


When  $T > m_\varphi$ ,  $\varphi$  exists in thermal bath, decays to  $\nu_4$ .

$\Gamma_{\text{decay}} \sim \lambda^2$ , more important than scattering for  $\lambda \ll 1$ .

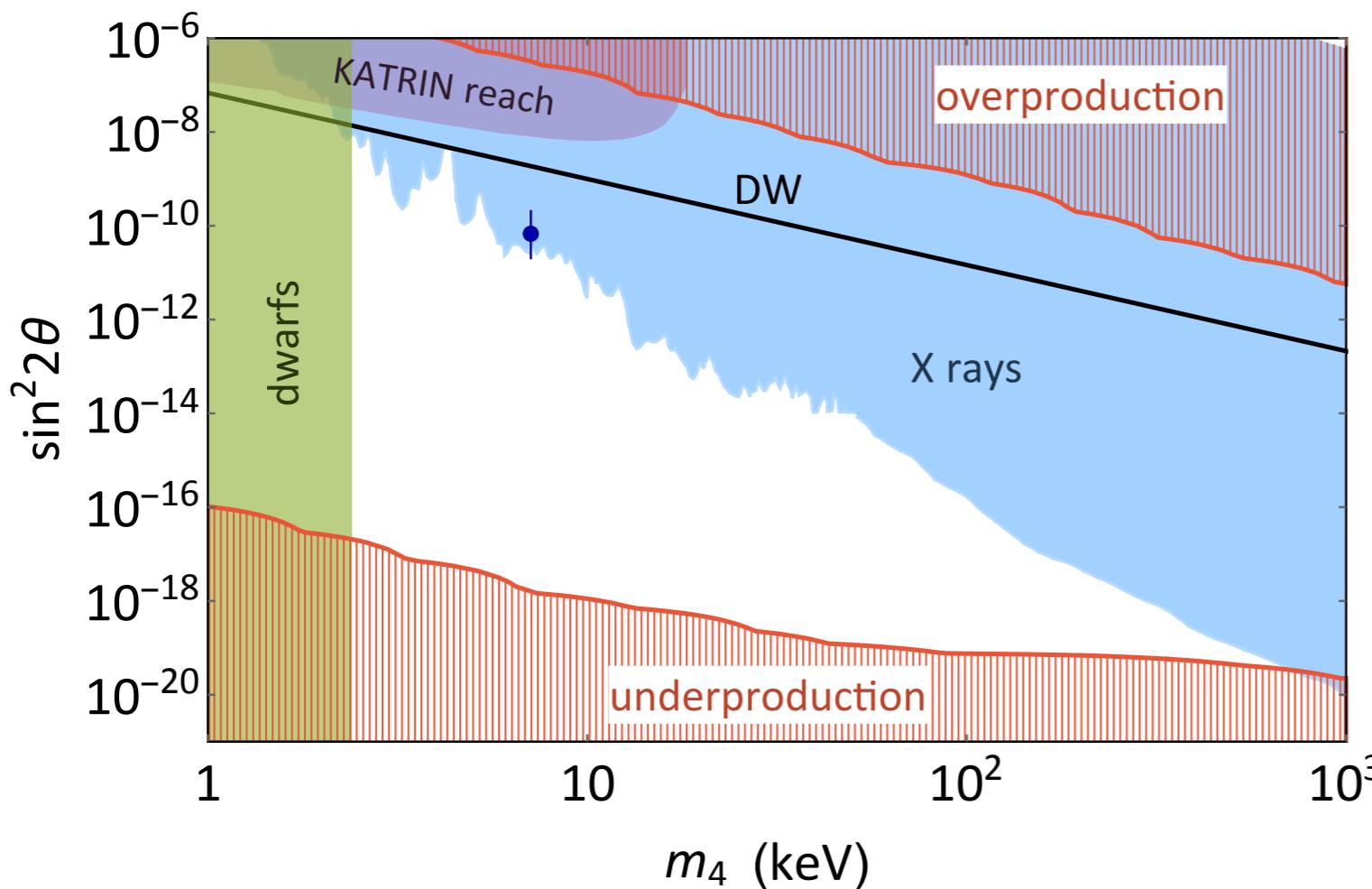
**Opens up new parameter space.**

# Case of Light Mediator



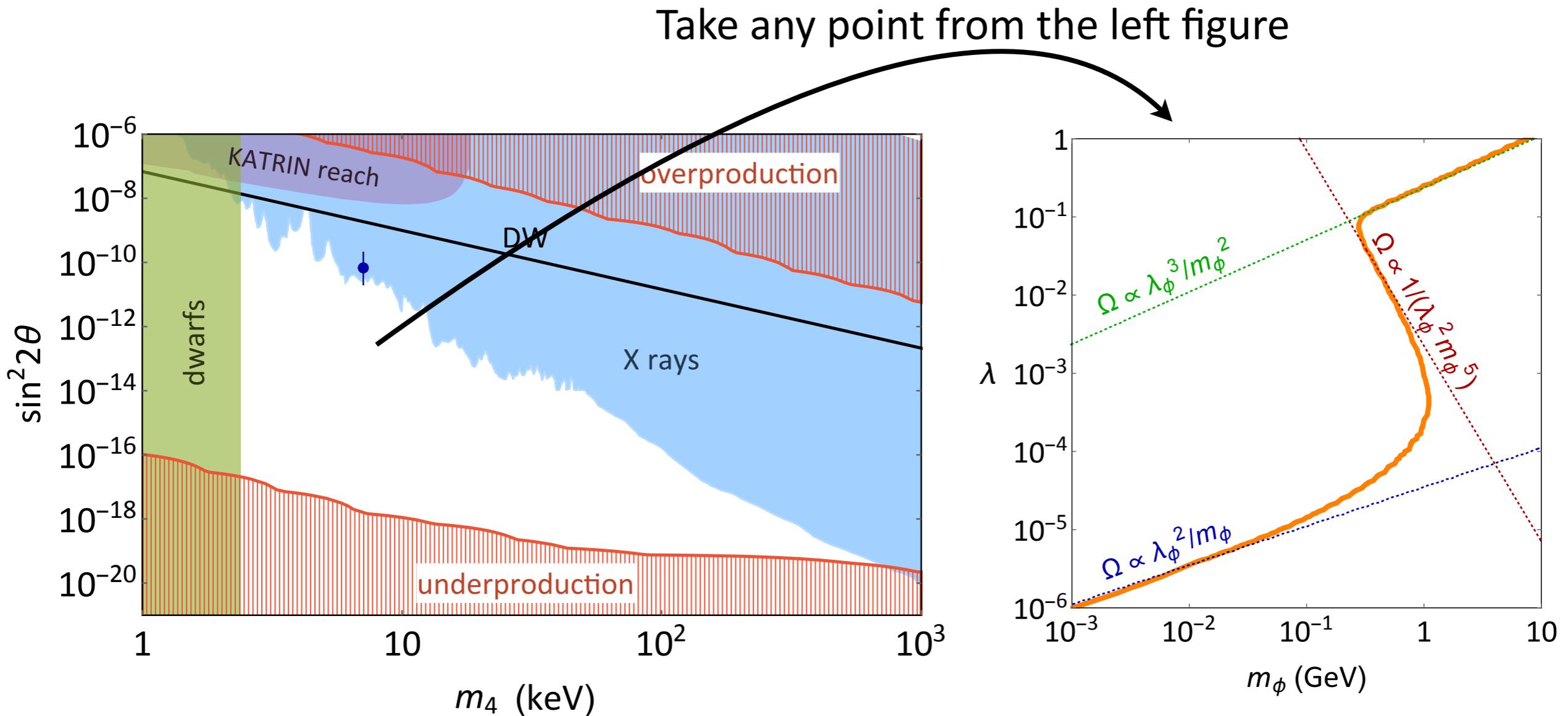
Additional scenario: new interaction can also impact mixing angle

# Wide Window for Relic Density



de Gouvêa, Sen, Tangarife, YZ (1910.04901)

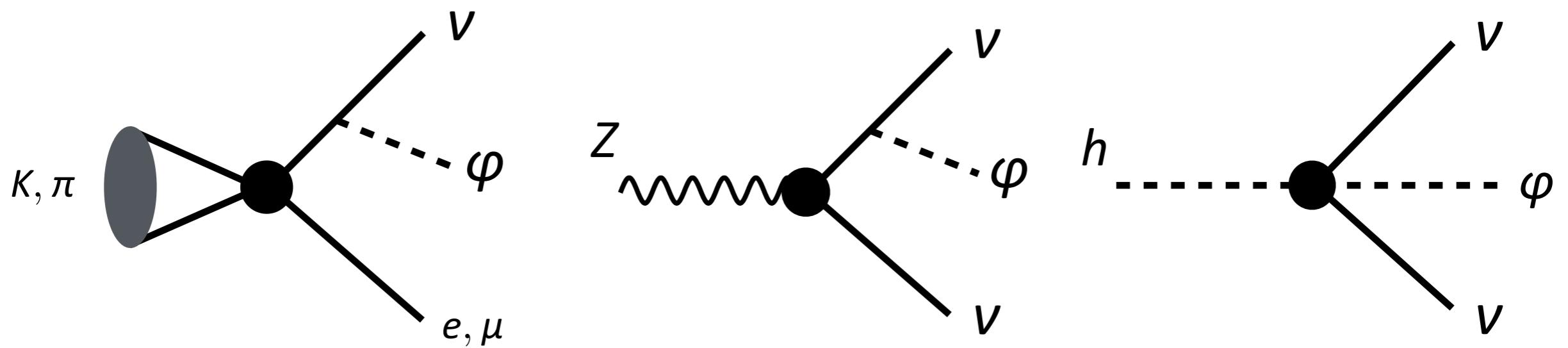
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de Gouvêa, Sen, Tangarife, YZ (1910.04901)

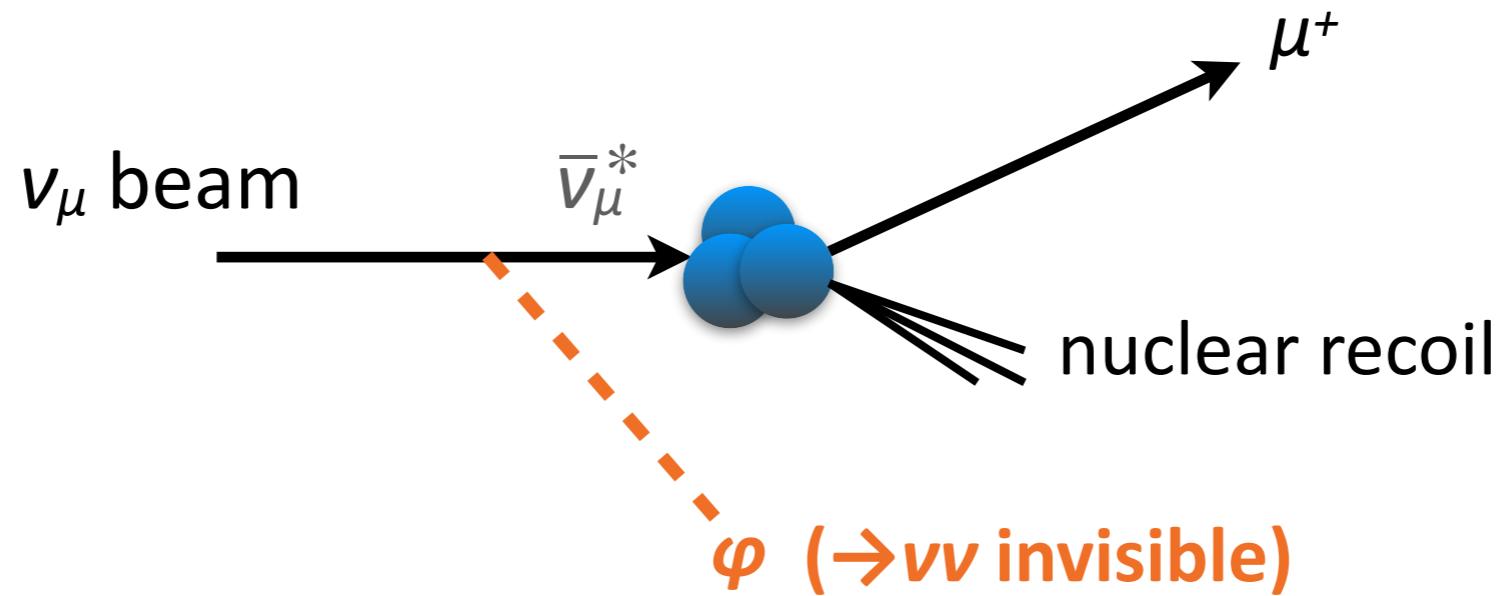
# Probing Neutrino Self Interaction

# Meson, Z, Higgs Invisible Decays



Berryman, de Gouvêa, Kelly, YZ (1802.00009)  
Brdar, Lindner, Vogl, Xu (2003.05339)  
Pasquini, Perez (1511.01811)  
Barger, Keung, Pakvasa (PRD 1982)

# Mono-Neutrino Signal

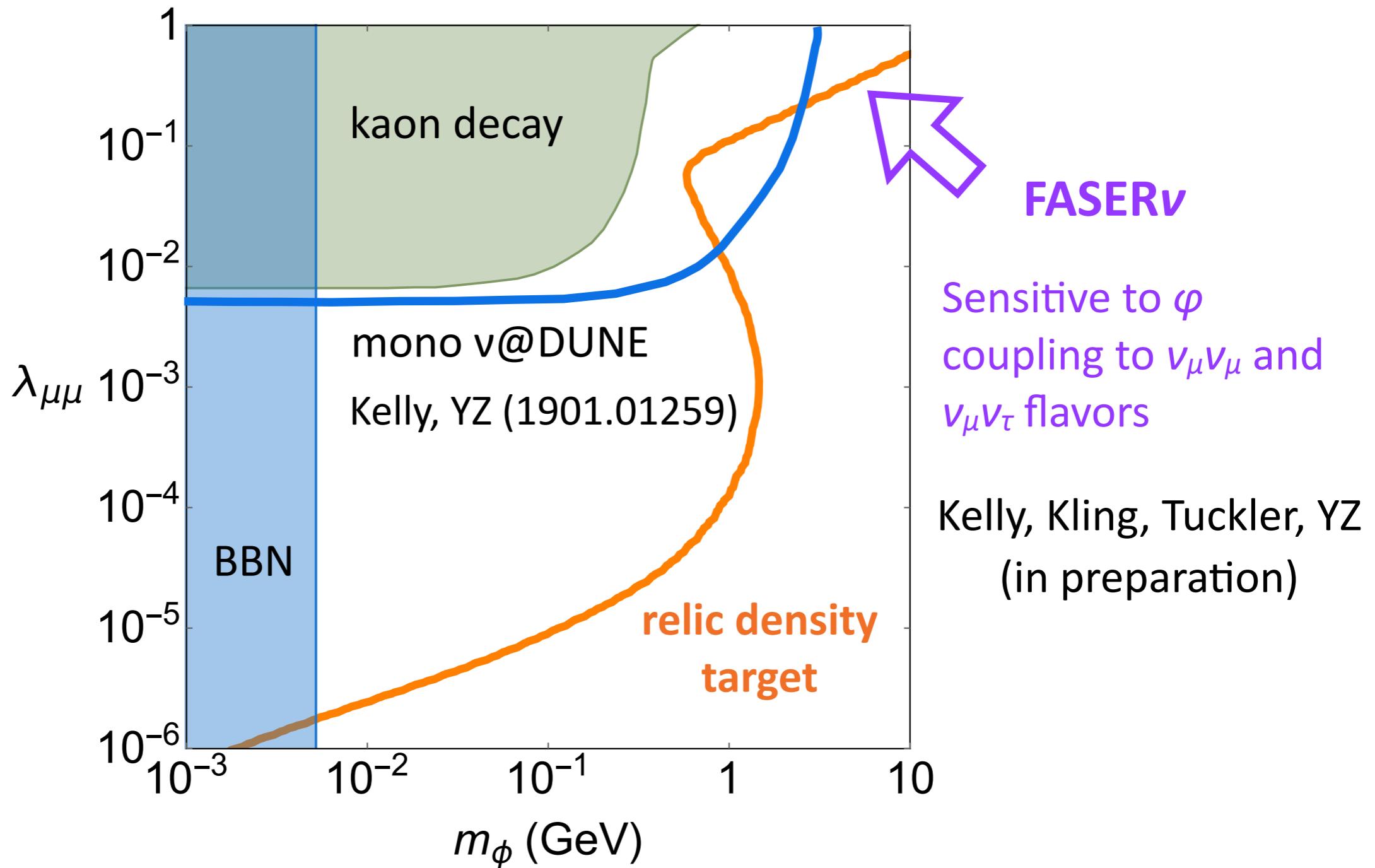


Beamstrahlung process:  $\nu_\mu + N \rightarrow \mu^+ + N' + \varphi$ , features

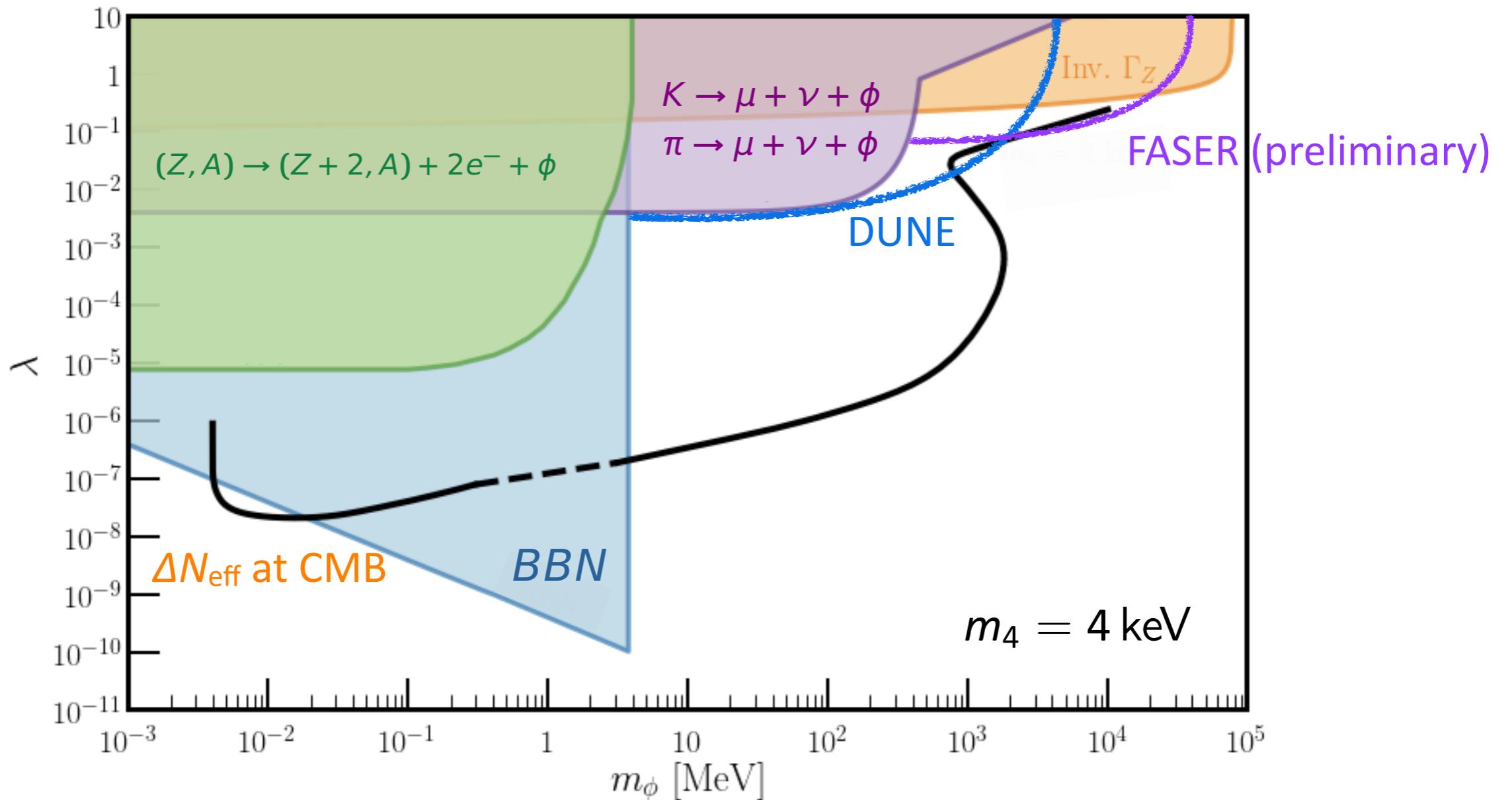
- “Wrong-sign” outgoing muon
- Missing transverse momentum  $p_T$

Kelly, YZ (1901.01259)

# Coverage



# The Big Picture



Kelly, Sen, YZ (2011.02487)

# Conclusion

Neutrino self interactions can have a novel impact on the origin of sterile neutrino dark matter.

Broad open parameter space relevant for DM relic density.

Excitingly, there are a number of ways to test this scenario with upcoming experiments.

**thanks!**