

Dark matter from a complex scalar singlet and the role of its discrete symmetries

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Results:

Introduction:

We consider a model where dark matter (DM) emerges from a complex scalar field, $S = \rho + \theta i$, charged under a U(1) global symmetry. This is the simplest model beyond having the SM plus a real scalar singlet, which is too restrained.

From the U(1) the following sym. breaking (SB) terms develops:

- Spontaneously SB (SSB) $[v_s]$ avoids being over-restrained and from suffer instability problems.
- **Explicit SB** gives mass to the Goldstone boson (θ) allowing it to be a DM candidate.

Explicit SB develops from the subset of discrete symmetries of the kinetic term:

 $Z_2: \quad S \to -S \quad Z_3: \quad S \to e^{i2\pi/3} S$ $Z_4: S \to iS \quad DCP: S \to S^*$

DCP can be identified with a "dark" CP and it is the only **remaining sym.** after SSB. An easy long-lived particle means preserving a discrete symmetry and DCP needs for all the couplings to be real.

Objectives:

Consider the following minimal models (MM) containing one single SB term which satisfies different discrete sym. (except V_1) and find if they lead to different subsets in the parameter space for a DM.

$$V_{1} = \frac{1}{2}\mu^{3}S + h.c.; \quad V_{Z_{2}} = \frac{1}{2}\mu_{S}^{2}S^{2} + h.c.$$
$$V_{Z_{3}} = \frac{1}{2}\mu_{3}S^{3} + h.c.; \quad V_{Z_{4}} = \frac{1}{2}\lambda_{4}S^{4} + h.c.$$





Conclusion \rightarrow MM are only distinguished in the FDM and SDM scenarios.