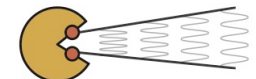




Accessing the coupled-channels dynamics with femtoscopy correlations at LHC

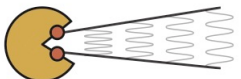
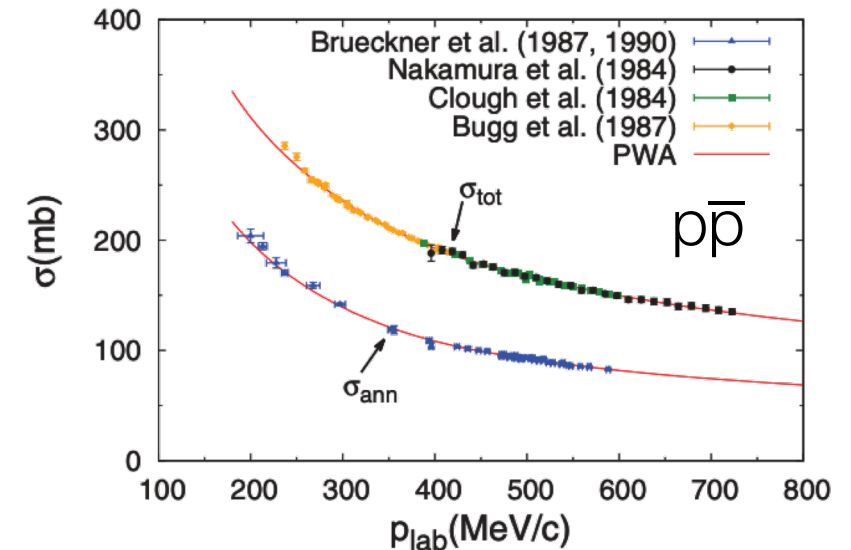
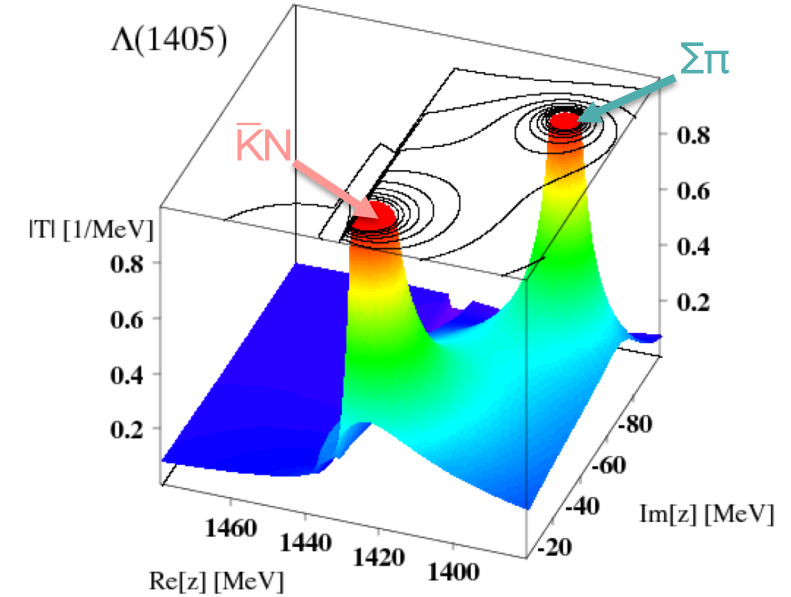
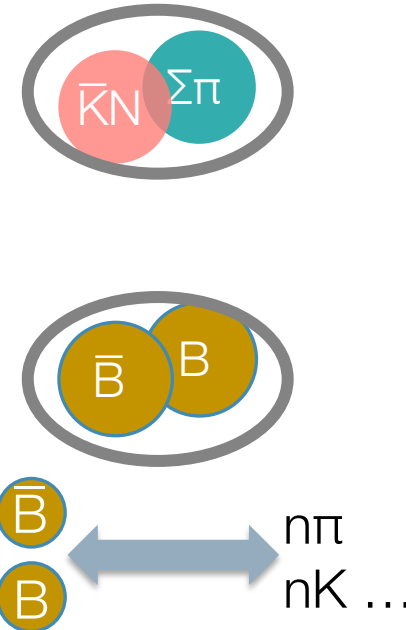
V. Mantovani Sarti (TUM) on behalf of the ALICE Collaboration

PANIC 2021 – 8.09.2021



- Coupled-channel dynamics widely present in hadron-hadron strong interactions
 - Close in mass and same quantum numbers (e.g. B,S,Q)
 - On-shell and off-shell processes from one channel to the other
- Can be at the origin of several phenomena
 - Molecular states as $\Lambda(1405) \rightarrow$ interplay of $\bar{K}N-\Sigma\pi$
- Annihilation dynamics for $B-\bar{B}$ interactions
 - Multi-meson channels below threshold

See A. Ohnishi's talk in previous session



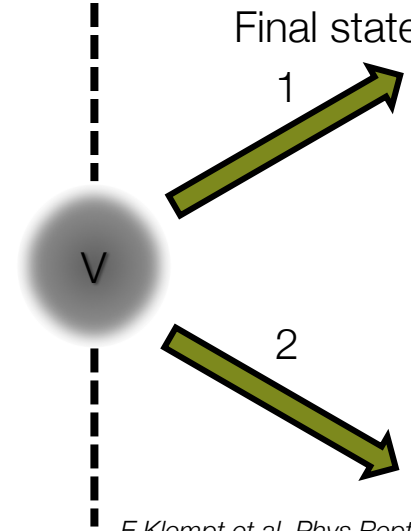
TUM Scattering experiments...

- Fixed initial state (e.g. K^-) and different final states:
 - Elastic scattering: $1 \rightarrow 1$
 - Inelastic scattering: $1 \rightarrow 2, 3, \dots$
- Measurement of the cross section in different channels
- Disadvantages:
 - Not accessible down to zero momenta
 - Large uncertainties
 - Limited to few h-h interactions

Initial state

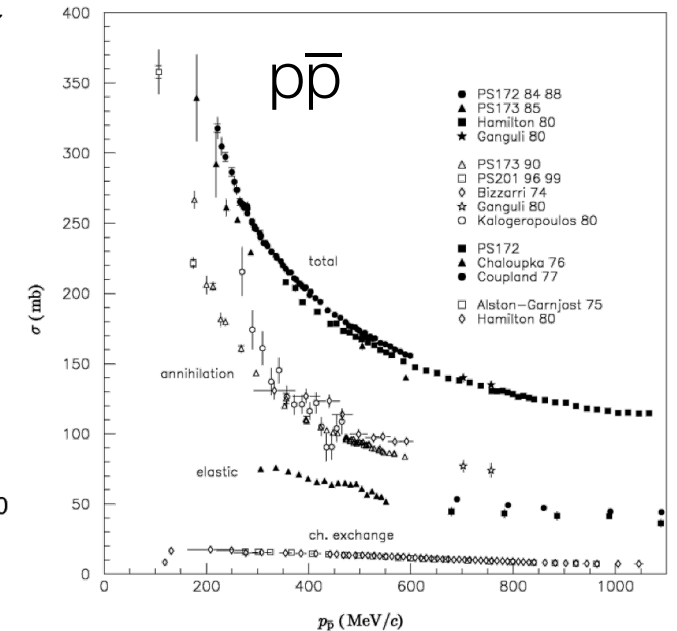
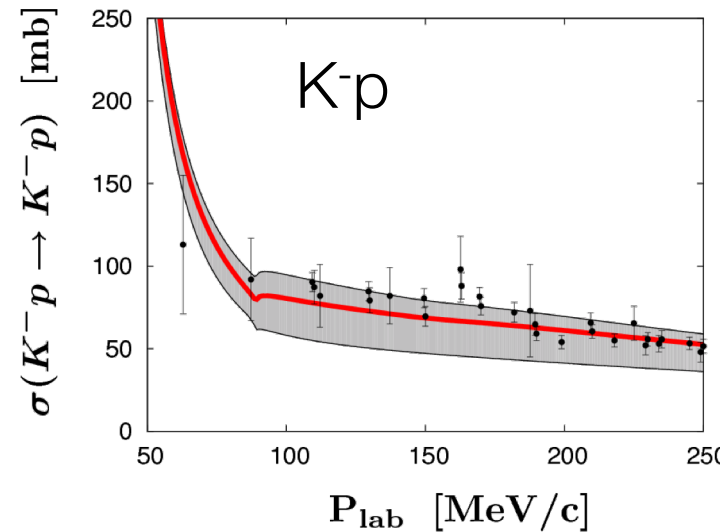


Final state



E.Klempt et al. Phys.Rept. 368 (2002) 119-316

G.S. Abrams et al. Phys.Rev. 139 (1965) B454-B457



TUM ... and now femtoscopy at ALICE

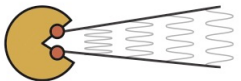
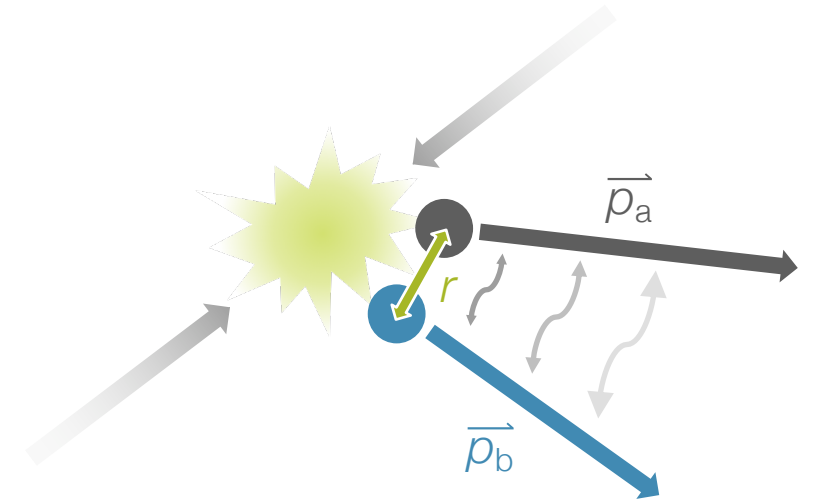
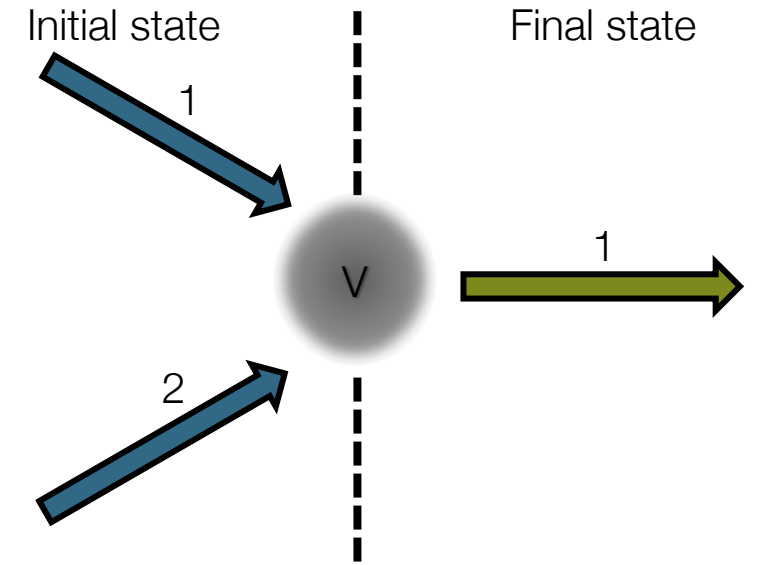


- Fixed final state with the measured pair and all possible initial states
 - inclusive measurement: $1, 2, 3, 4, \dots \rightarrow 1$

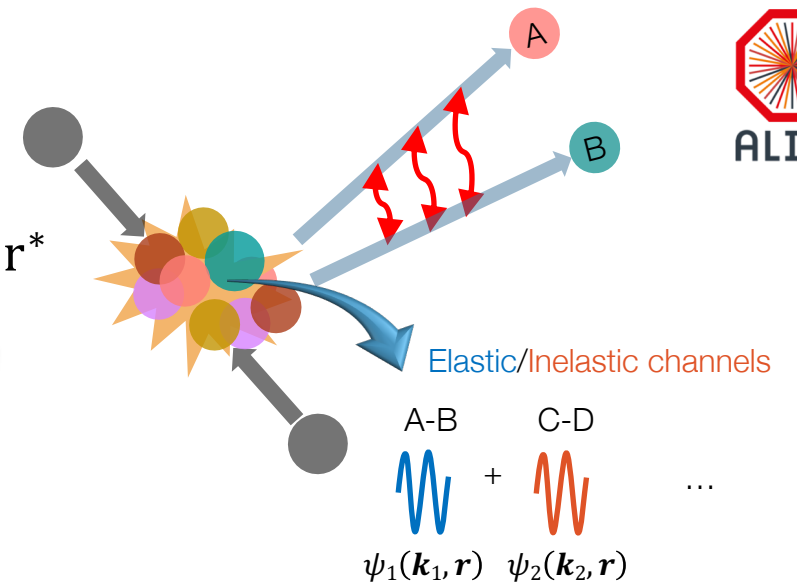
$$C(\mathbf{k}^*) = \zeta(\mathbf{k}^*) \cdot \frac{N_{\text{same}}(\mathbf{k}^*)}{N_{\text{mixed}}(\mathbf{k}^*)} \begin{cases} > 1 \text{ attraction} \\ = 1 \text{ no inter.} \\ < 1 \text{ repulsion} \end{cases}$$

- Access to any particle-pair interaction as long as enough statistics is available
 - Baryon-baryon sector widely investigated in ALICE Run 2

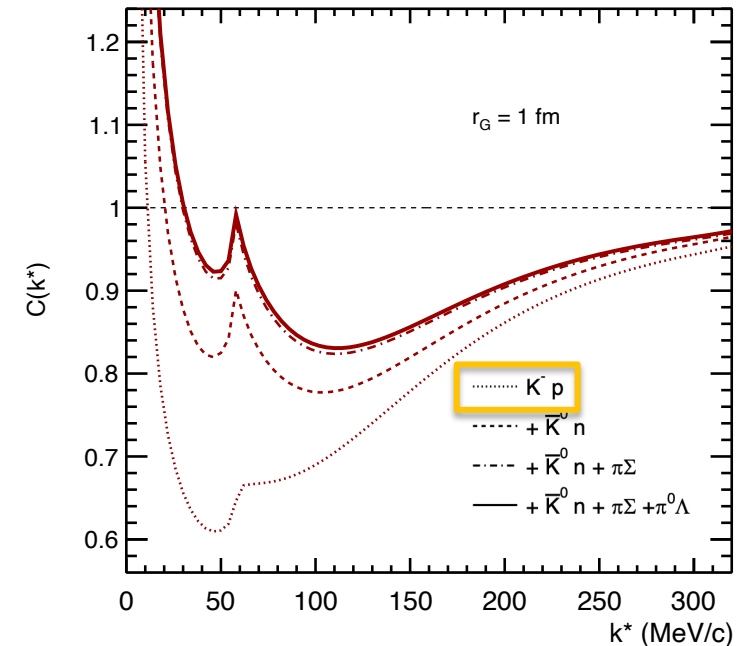
PRC 99 (2019) 024001
PLB 797 (2019) 134822
PRL 123 (2019) 112002
PRL 124 (2020) 09230
PLB 805 (2020) 135419
PLB 811 (2020) 135849
Nature 588 (2020) 232-238
 arXiv:2104.04427
 arXiv:2105.05578
 arXiv:2105.05683
 arXiv:2105.05190



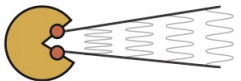
$$C(k^*) = \underbrace{\int S(\vec{r}^*) |\psi_{1 \rightarrow 1}(\vec{k}^*, \vec{r}^*)|^2 d^3r^*}_{\text{elastic A-B} \rightarrow \text{A-B}} + \underbrace{\sum_{j \neq 1} w_j \int S(\vec{r}^*) |\psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)|^2 d^3r^*}_{\text{inelastic C-D} \rightarrow \text{A-B, ...}}$$



- Wave functions $\psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)$
 - Above threshold ($E \gtrsim M_{\text{pair}}$) \rightarrow cusp structure e.g. $\bar{K}^0 n$
 - Below threshold ($E < M_{\text{pair}}$) \rightarrow shift upward of CF e.g. $\Sigma \pi$
- Conversion weights w_j related to the produced pairs as initial states and coupling strength
- Correlation function is sensitive to inelastic channels but how does the emitting source enter?



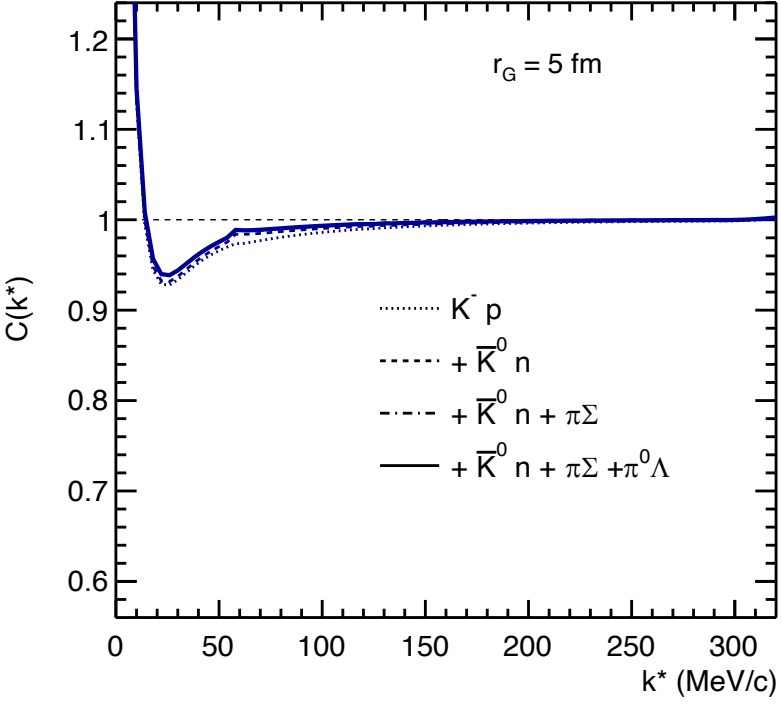
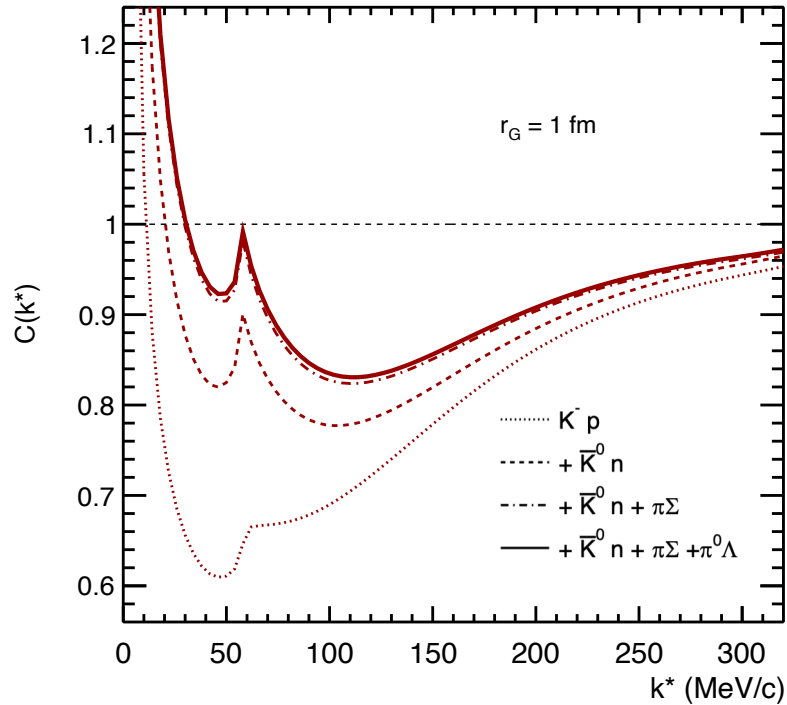
J. Haidenbauer Nucl.Phys.A 981 (2019)
 Y. Kamiya et al. Phys.Rev.Lett. 124 (2020)
 V.M.S., L. Fabbietti, O. Vazquez-Doce 2012.09806



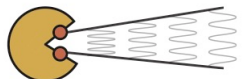
- Inelastic contributions negligible as the source size increases → mostly driven by the elastic interaction

$$C(k^*) = \int S(\vec{r}^*) |\psi_{1 \rightarrow 1}(\vec{k}^*, \vec{r}^*)|^2 d^3 r^* + \sum_{j \neq 1} w_j \int S(\vec{r}^*) |\psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)|^2 d^3 r^*$$

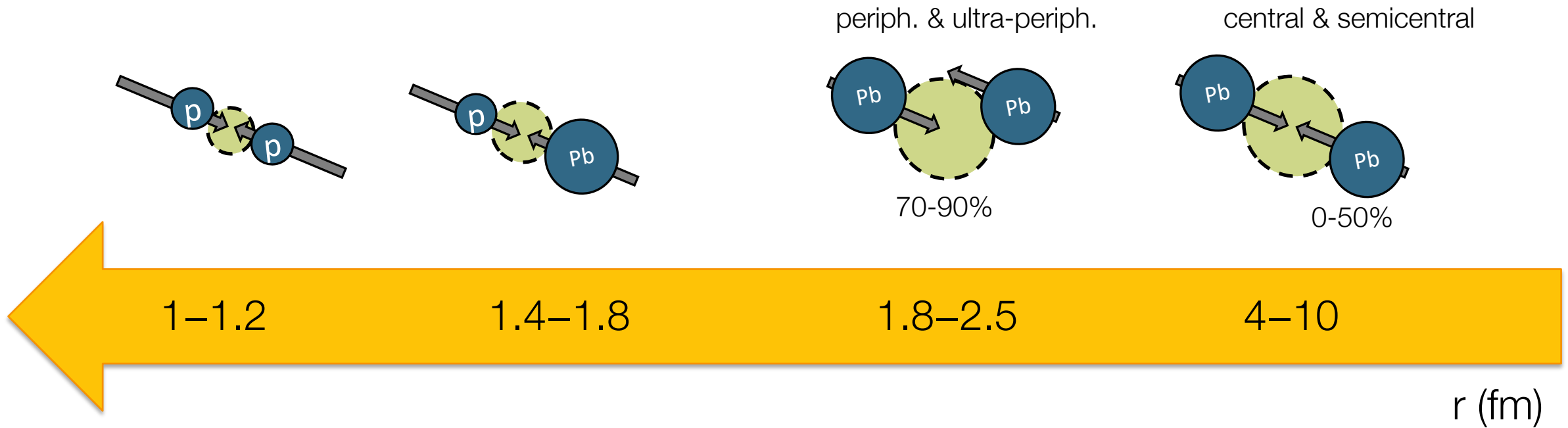
$$C(k^*) = \int S(\vec{r}^*) |\psi_{1 \rightarrow 1}(\vec{k}^*, \vec{r}^*)|^2 d^3 r^* + \sum_{j \neq 1} w_j \int S(\vec{r}^*) |\psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)|^2 d^3 r^*$$



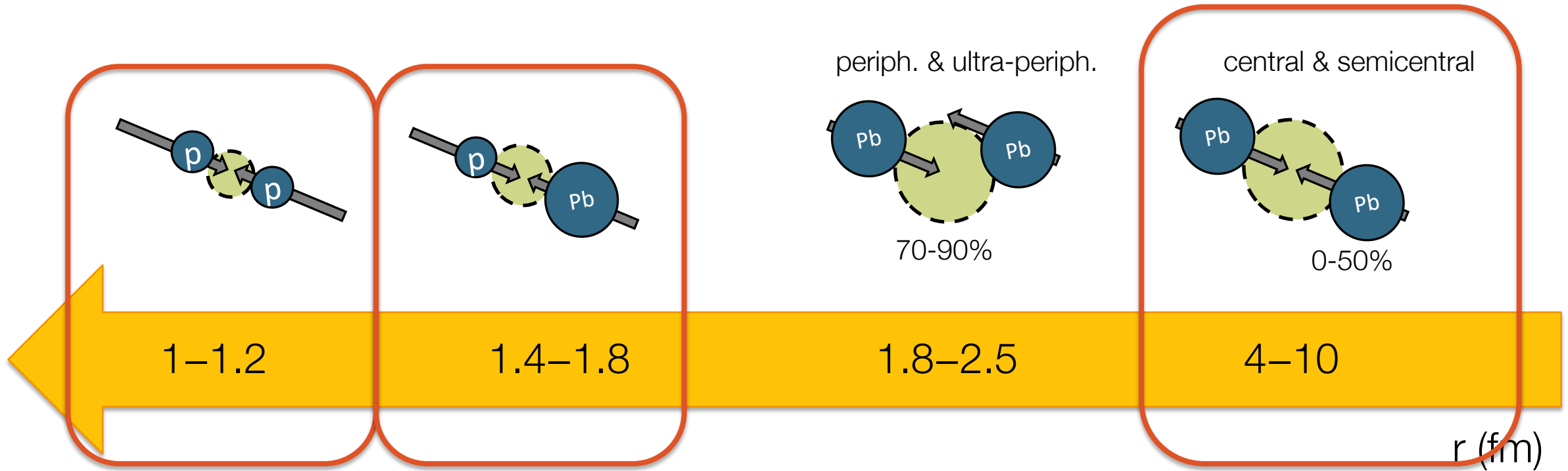
J. Haidenbauer Nucl.Phys.A 981 (2019)
 Y. Kamiya et al. Phys.Rev.Lett. 124 (2020)
 V.M.S., L. Fabbietti, O. Vazquez-Doce 2012.09806



- At ALICE by changing the colliding system we can probe distances ranging from 1 fm up to 10 fm



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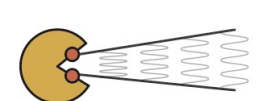
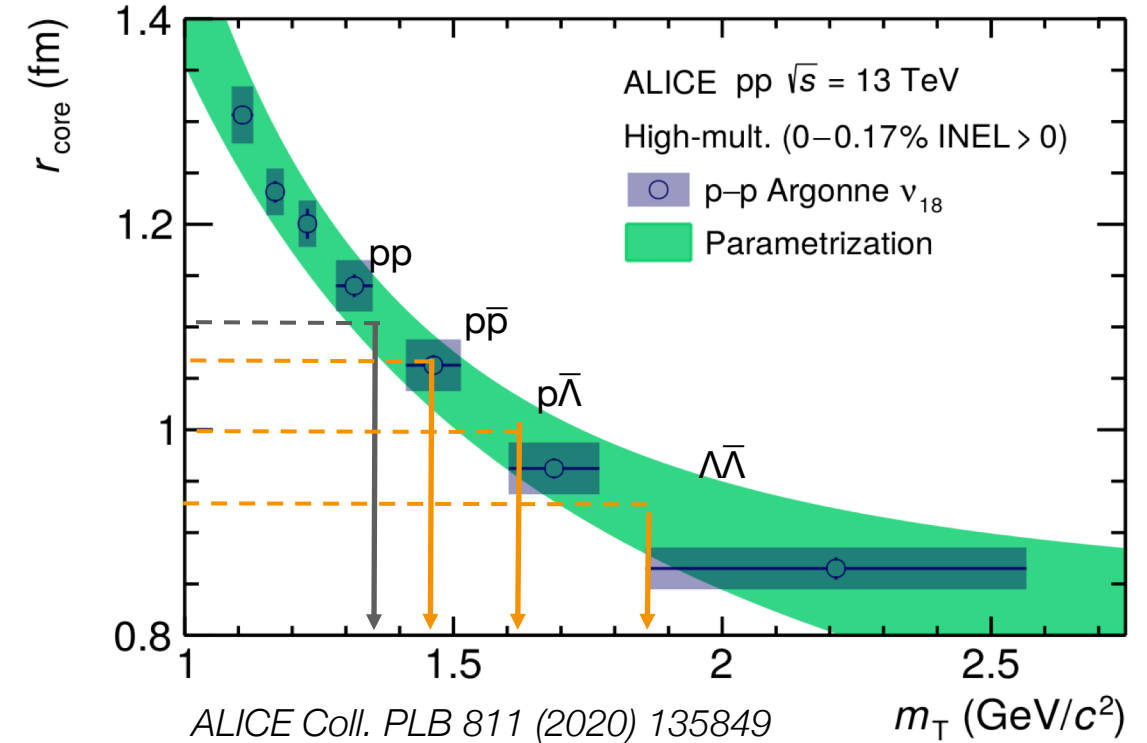
K-p and K+p in pp 5,7,13 TeV, p-Pb and Pb-Pb 0-50% 5.02 TeV
 $p\bar{p}$, $p-\bar{\Lambda}$ and $\Lambda-\bar{\Lambda}$ in pp HM 13 TeV and Pb-Pb 5.02 TeV

In pp collisions:

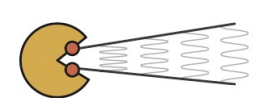
- Two main contributions
 - Common Gaussian core r_{core} from $\langle m_T \rangle$ scaling due to collective effects (p-p, K^+ -p)
 - Strongly decaying resonances contributions
- Effective Gaussian with $r_G \sim 1$ -1.2 fm

In Pb–Pb collisions a gaussian source is assumed:

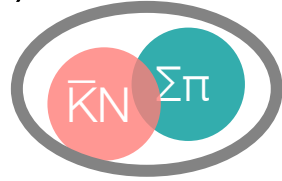
- Fixed from K^+ -p and B–B pairs
- Scaling of the source size with m_T in agreement with previous measurements at 2.76 TeV



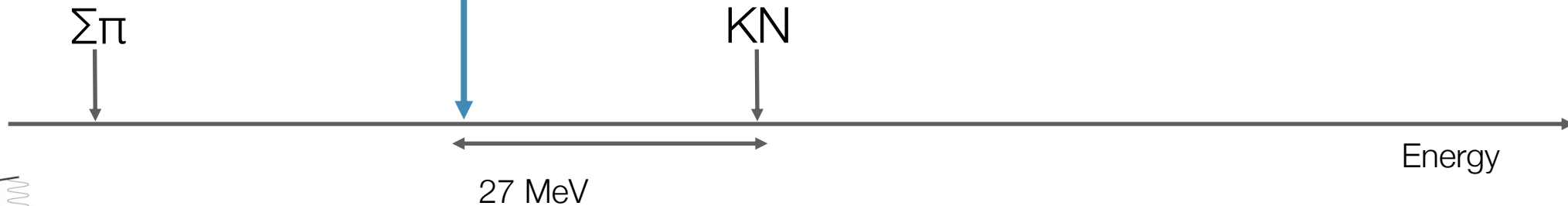
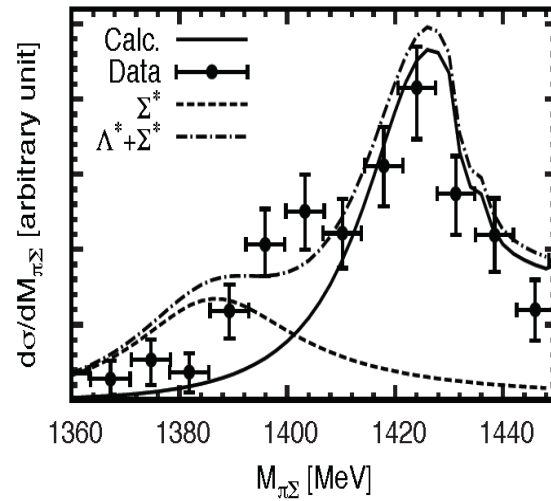
K^- -p and K^+ -p femtoscopy



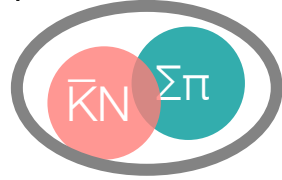
$\Lambda(1405)$ molecular state



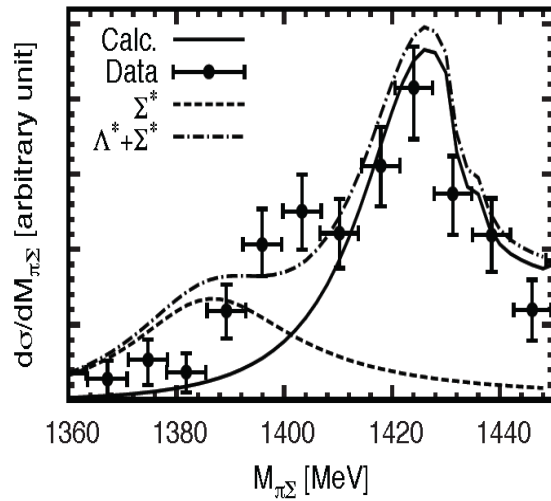
$\Lambda(1405) \rightarrow \Sigma\pi$



$\Lambda(1405)$ molecular state



$$\Lambda(1405) \rightarrow \Sigma\pi$$

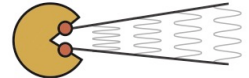
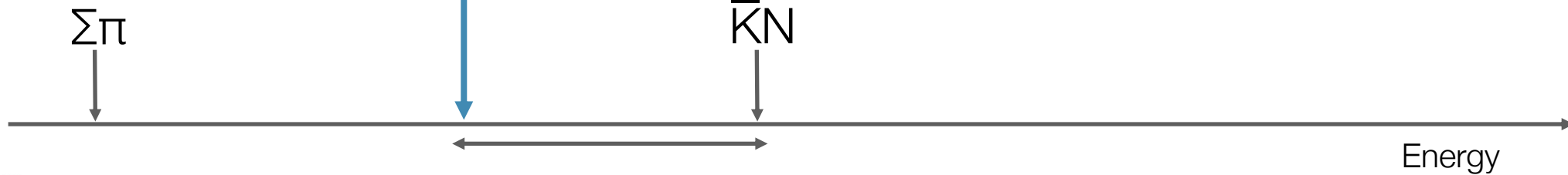
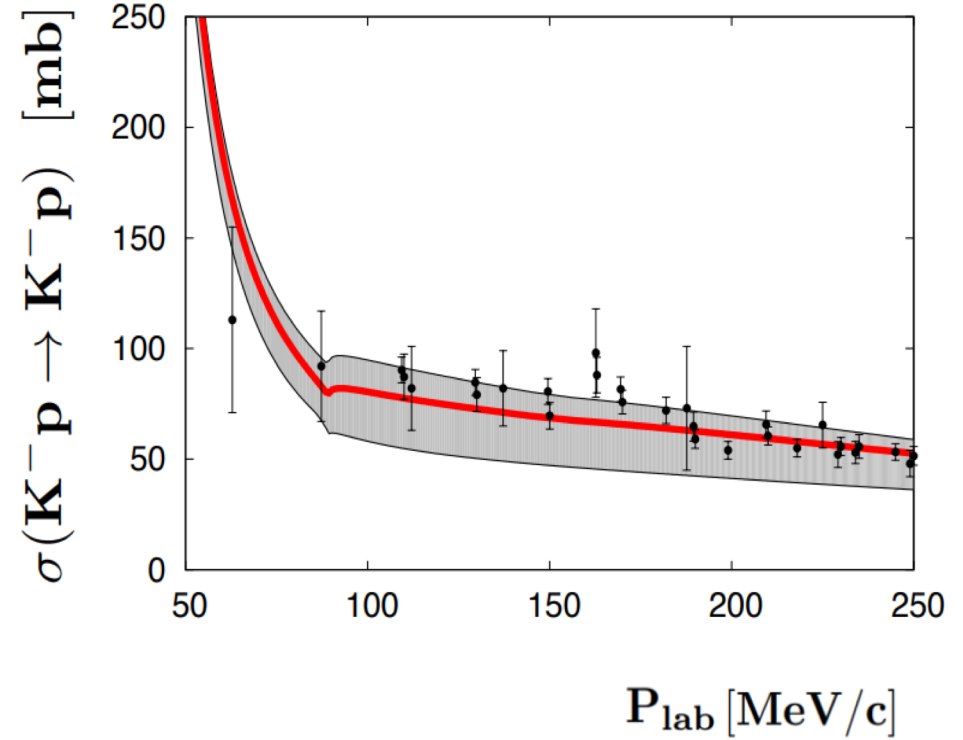


Kaonic hydrogen

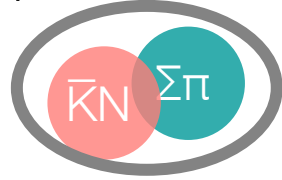
SIDDHARTA Coll. PLB 704 (2011)

Scattering Experiments

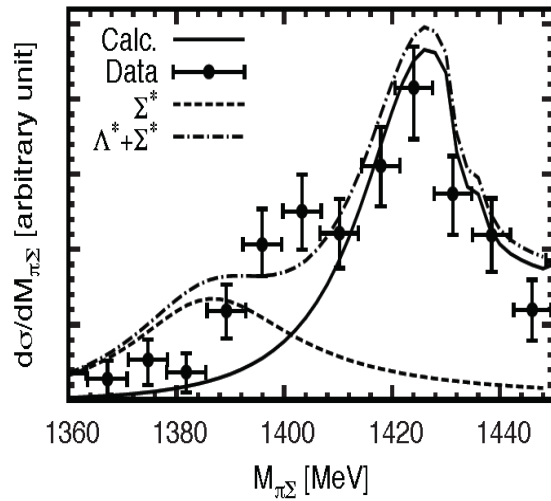
Ikeda et al. Nucl.Phys.A 881 (2012)



$\Lambda(1405)$ molecular state

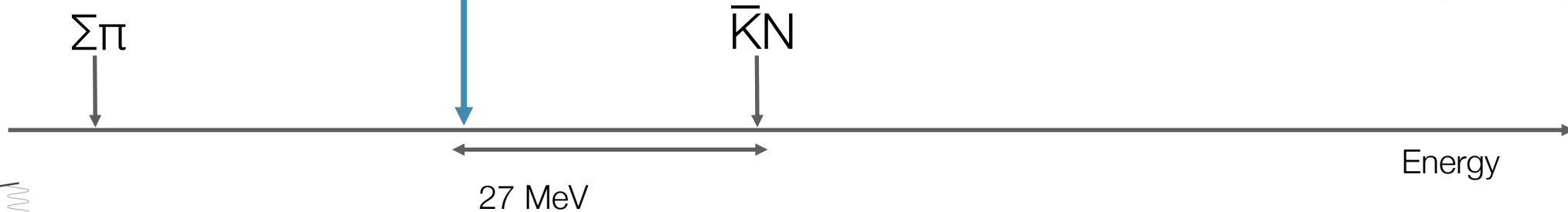
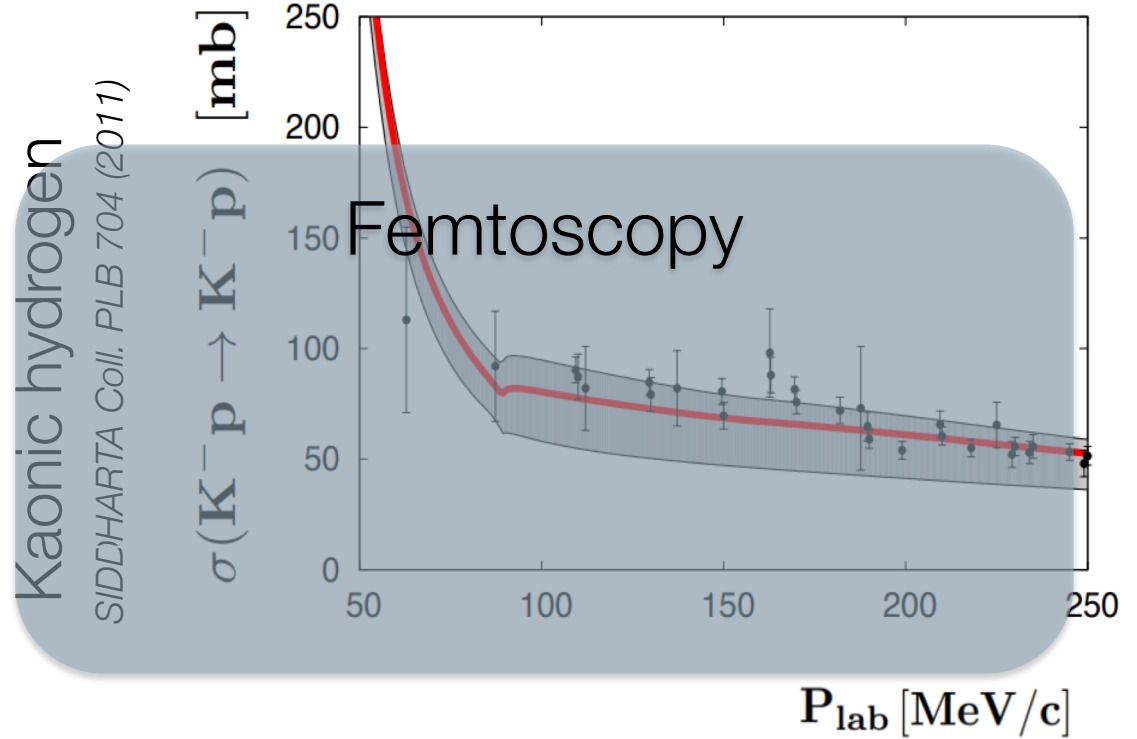


$$\Lambda(1405) \rightarrow \Sigma\pi$$



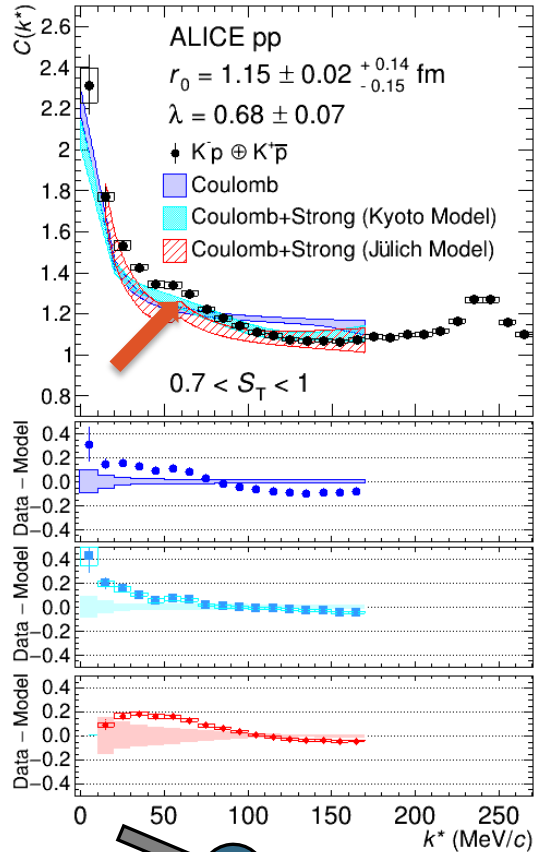
Scattering Experiments

Ikeda et al. Nucl.Phys.A 881 (2012)

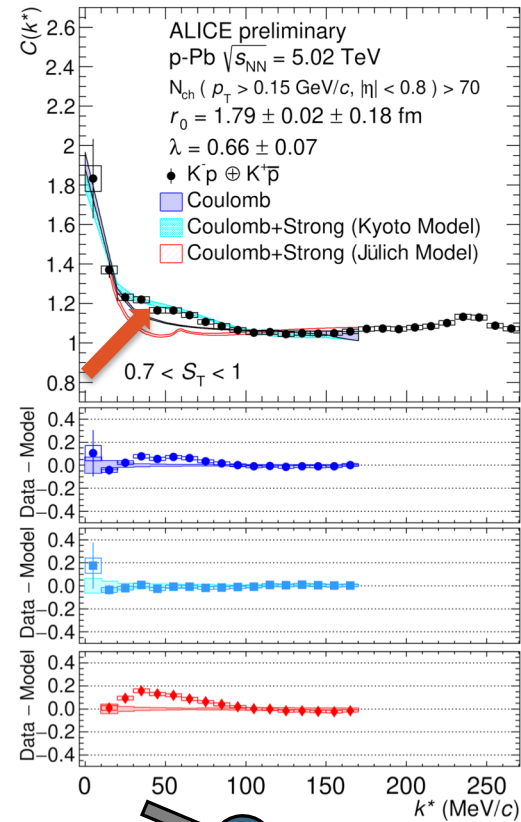


- Disappearance of \bar{K}^0 -n cusp at $k^* \sim 60$ MeV/c as the source increases

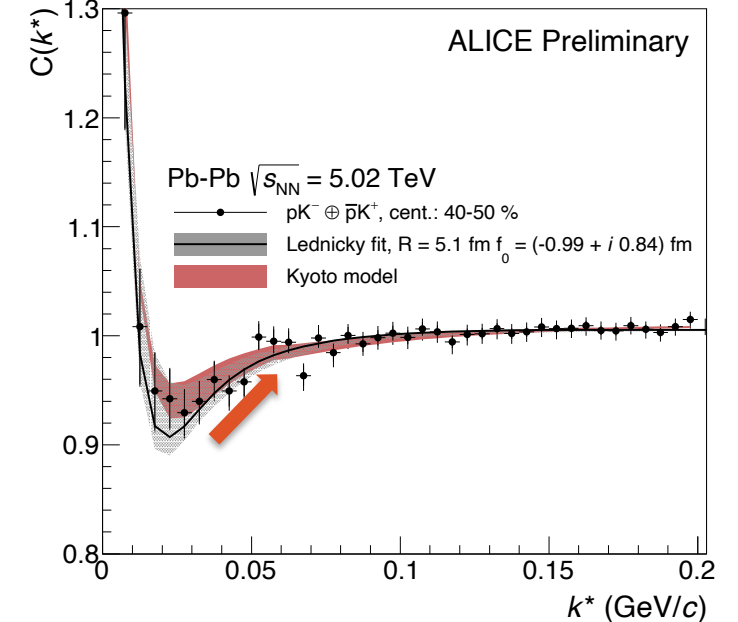
ALICE Coll. Phys.Rev.Lett. 124 (2020) 9, 092301



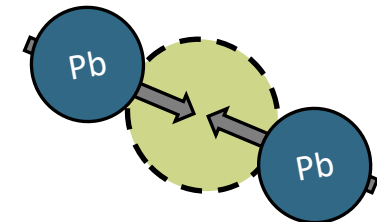
Paper in preparation



ALICE Coll. arXiv: 2105.05683

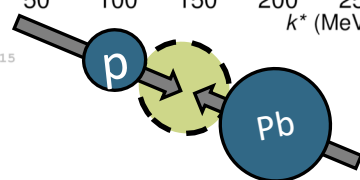


ALI-PREL-354647



assuming $w_j = 1$

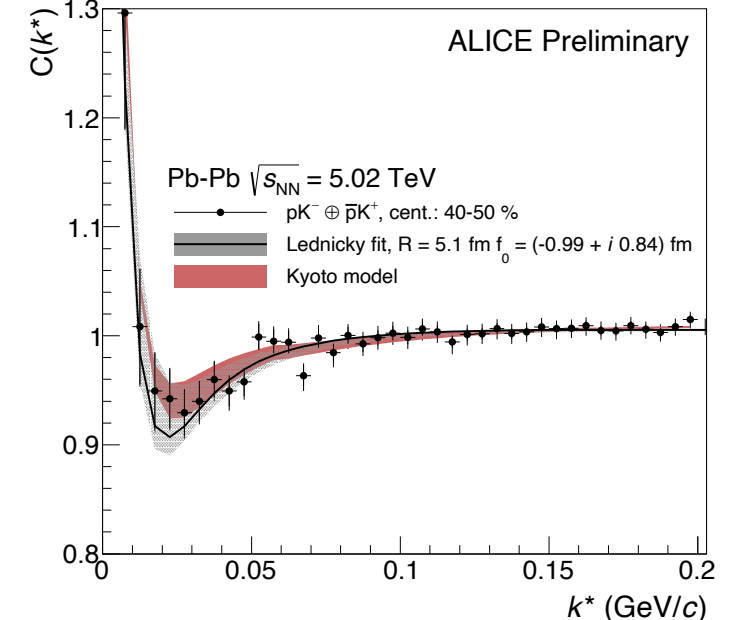
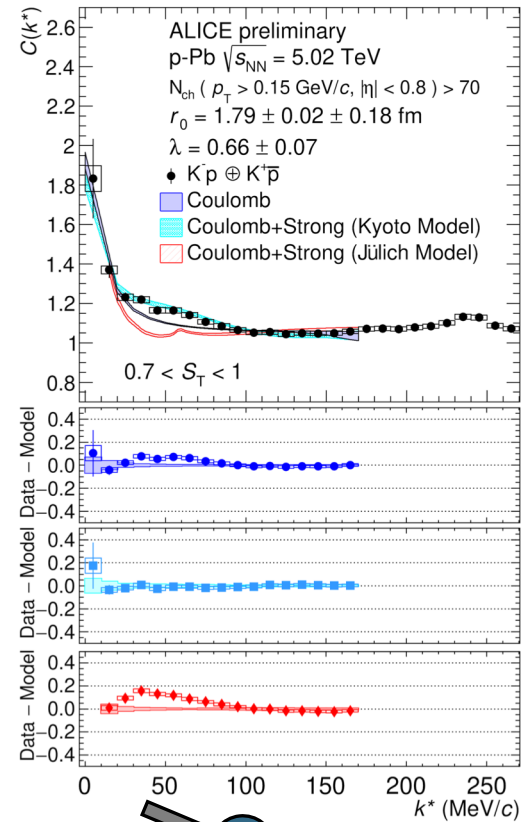
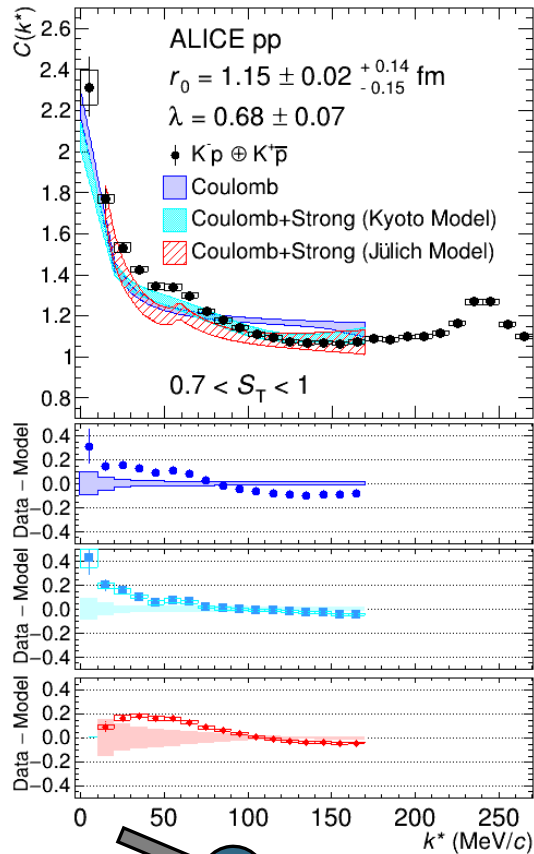
ALI-PREL-316315



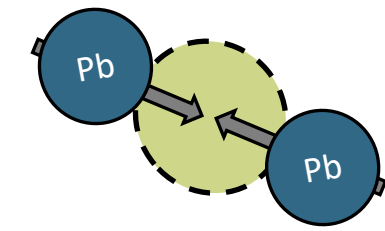
K⁻-p femtoscopy in different colliding systems

- Disappearance of \bar{K}^0 -n cusp at $k^* \sim 60$ MeV/c as the source increases
- χ EFT Kyoto potential describes the data in all three systems
- Pb-Pb data well described by Lednicky single channel \rightarrow negligible contribution from inelastic channels
- Pb-Pb results in agreement with scattering parameters from kaonic hydrogen

*Ikeda et al. Nucl.Phys.A 881 (2012),
PLB706 (2011)
Kamiya et al. PRL 124 (2020)
Mihayara et al. PRC95 (2017)*

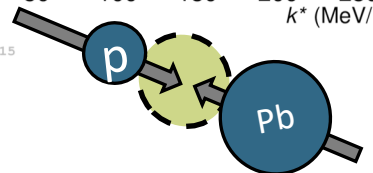


ALI-PREL-354647

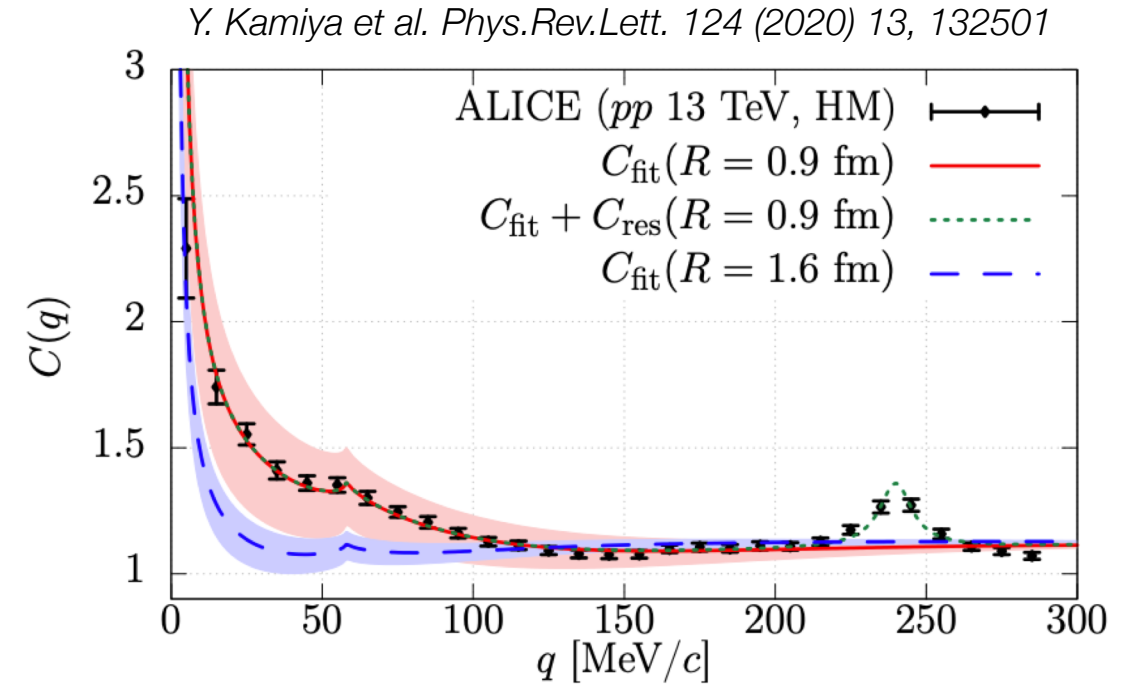


assuming $w_j = 1$

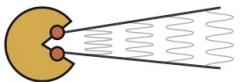
ALI-PREL-316315



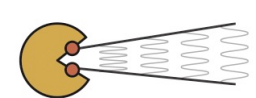
- Modification of the conversion weights leads to a better agreement with the data
 - see *Phys.Rev.Lett.* 124 (2020) 13, 132501
- Study on conversion weights in pp, p–Pb and Pb–Pb
 - still agreement data/model?
 - provide constraints for the \bar{K}^0n coupling
- Extend the femtoscopic studies to K-d pairs in all colliding systems



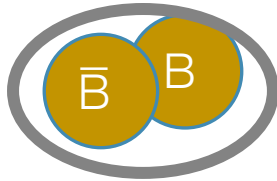
Work in progress with
 Dr. Y. Kamiya
 Prof. A. Ohnishi
 Prof. T. Hyodo



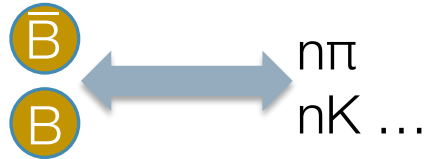
Baryon-antibaryon femtoscopy



Predictions for Baryonia??



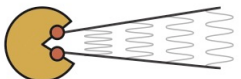
?



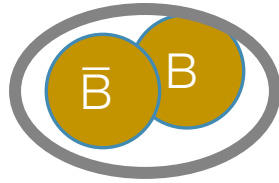
Annihilation dynamics
 $n\pi, nK, \pi K, \dots$



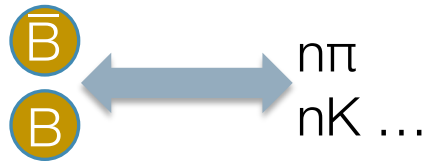
E. Klempt et al. Phys. Rept. 413 (2005)
E. Klempt et al. Phys. Rept. 368 (2002)
D.Zhou and R.G. E. Timmermans PRC86 (2012)
J. Haidenbauer et al. JHEP 1707 (2017)



Predictions for Baryonia??



?

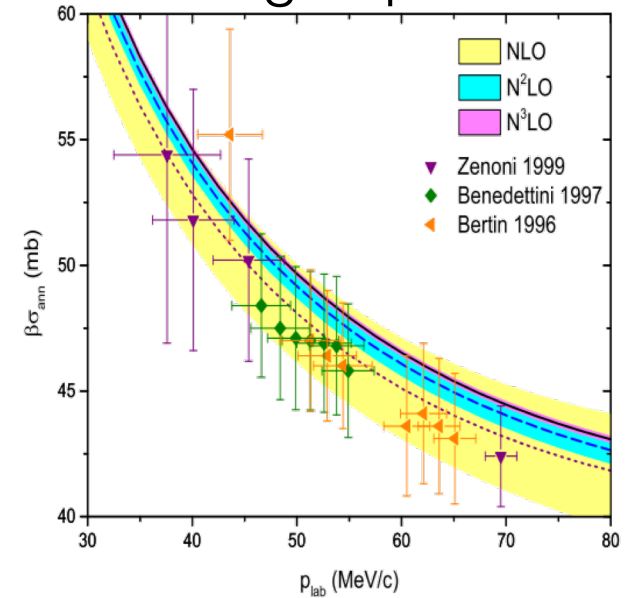


Annihilation dynamics
 $n\pi, nK, \pi K, \dots$

antiproton atoms

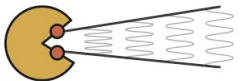
$p-\bar{p}$

Scattering Experiments

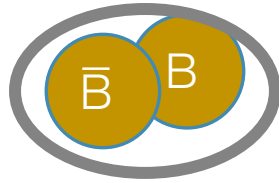


Energy

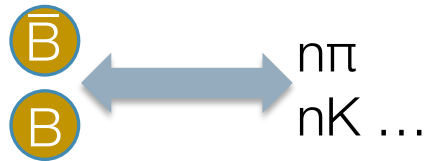
E. Klempt et al. Phys. Rept. 413 (2005)
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J. Haidenbauer et al. JHEP 1707 (2017)



Predictions for Baryonia??



?



Annihilation dynamics
 $n\pi, nK, \pi K, \dots$

antiproton atoms

$p-\bar{p}$

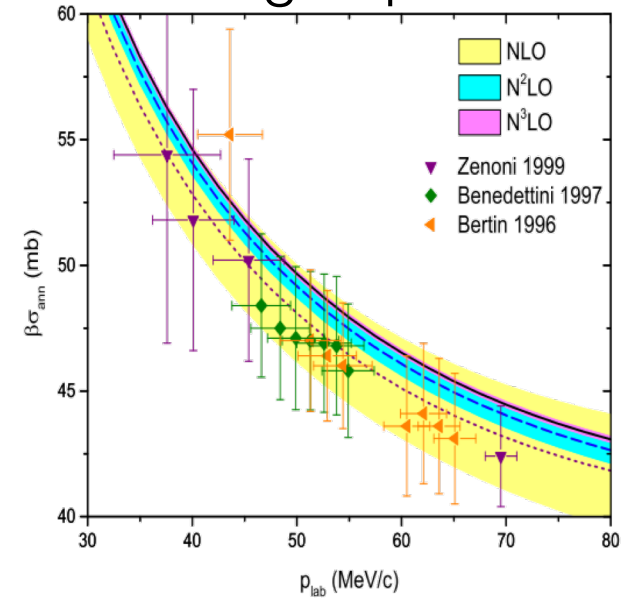
$p-\bar{\Lambda}$

$\Lambda-\bar{\Lambda}$

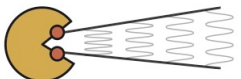
177 MeV

Energy

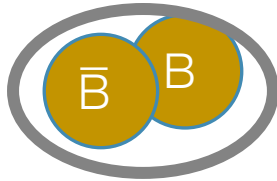
Scattering Experiments



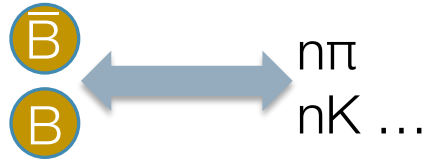
E. Klempt et al. Phys. Rept. 413 (2005)
E. Klempt et al. Phys. Rept. 368 (2002)
D.Zhou and R.G. E. Timmermans PRC86 (2012)
J. Haidenbauer et al. JHEP 1707 (2017)



Predictions for Baryonia??

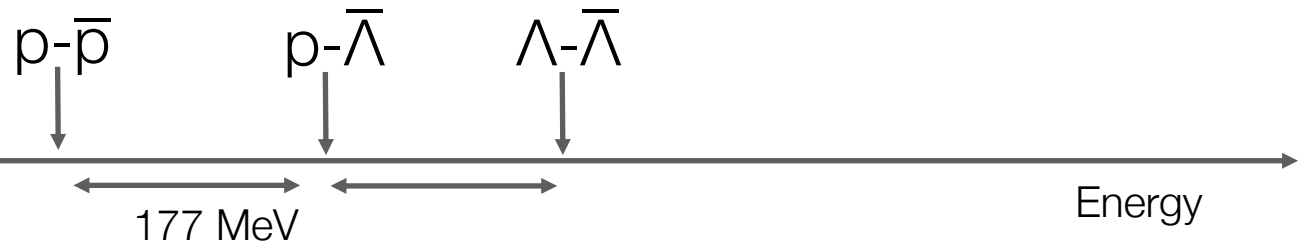
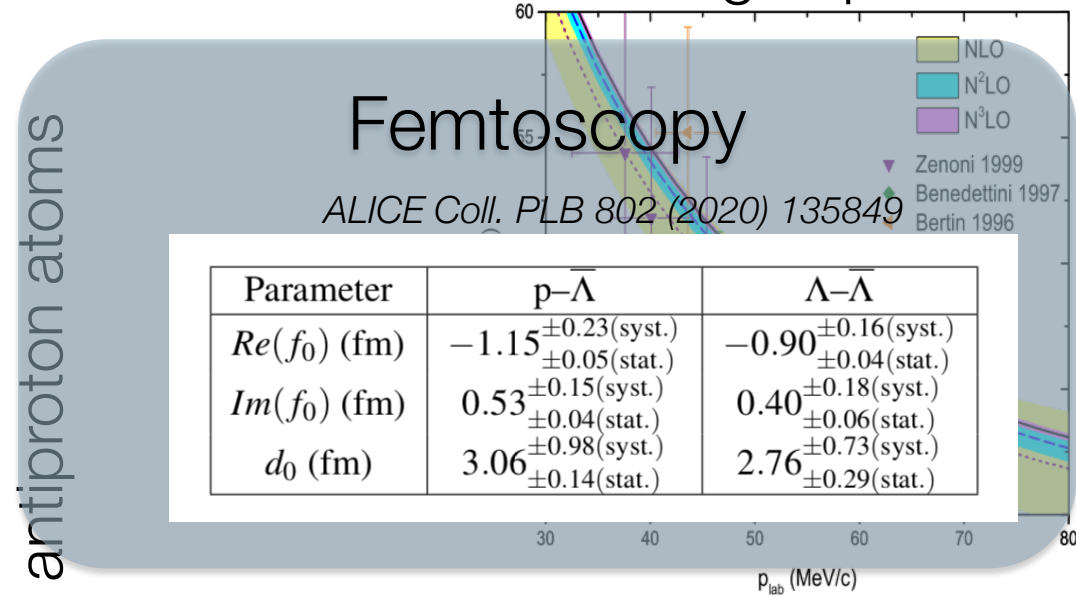


?

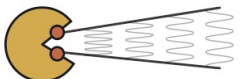


Annihilation dynamics
nπ, nK, πK, ...

Scattering Experiments



E. Klempt et al. Phys. Rept. 413 (2005)
E. Klempt et al. Phys. Rept. 368 (2002)
D.Zhou and R.G. E. Timmermans PRC86 (2012)
J. Haidenbauer et al. JHEP 1707 (2017)



$$C_{p\text{-}\bar{p}}(k^*) = \int S(r) |\psi_{p\bar{p} \rightarrow p\bar{p}}|^2 d^3 r + \int S(r) |\psi_{n\bar{n} \rightarrow p\bar{p}}|^2 d^3 r$$

elastic
 $n\text{-}\bar{n} \rightarrow p\text{-}\bar{p}$

- Chiral Effective Field Theory at N³LO (with $n\text{-}\bar{n}$ coupled-channel) wavefunctions with Coulomb
 - S and P waves, tuned to scattering data and protonium

J. Haidenbauer et al. JHEP 1707 (2017)
D.Mihailov et al. Eur.Phys.J.C 78 (2018) 5, 394
Haidenbauer et al.,PRD 92 (2015) 5, 054032

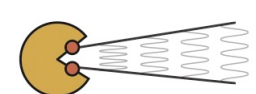
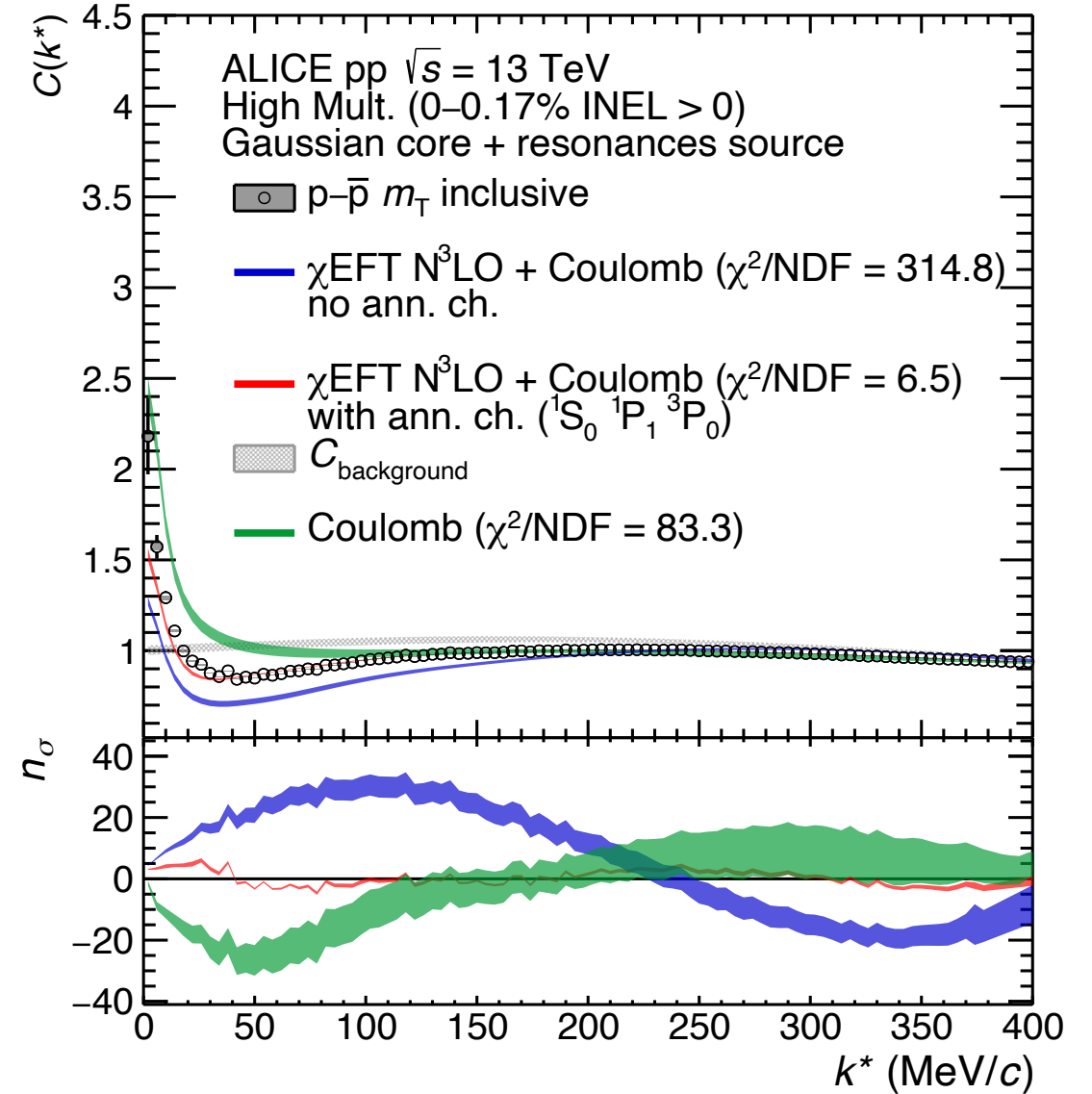
$$C_{p\text{-}\bar{p}}(k^*) = \int S(r) |\psi_{p\bar{p} \rightarrow p\bar{p}}|^2 d^3 r + \int S(r) |\psi_{n\bar{n} \rightarrow p\bar{p}}|^2 d^3 r + \sum_{PW} \rho_{PW} \omega_{PW} \int S(r) |\psi_{p\bar{p} \rightarrow p\bar{p}}^{PW}|^2 d^3 r$$

elastic
 $n\text{-}\bar{n} \rightarrow p\text{-}\bar{p}$
multi-meson annihilation channels

- **Chiral Effective Field Theory at N³LO** (with $n\text{-}\bar{n}$ coupled-channel) wavefunctions with Coulomb
 - S and P waves, tuned to scattering data and protonium
- **Approximate inclusion of annihilation channels** ($X \rightarrow p\text{-}\bar{p}$) using the Migdal-Watson approximation
 - elastic WF rescaled by a coupling weight ω_{PW} to be fitted to data
 - Investigation on the shape of each PWs to reduce number of parameters
 - 1S_0 for S states
 - 3P_0 and 1P_1 for P states
- Calculations performed with CATS framework

J. Haidenbauer et al. JHEP 1707 (2017)
D.Mihailov et al. Eur.Phys.J.C 78 (2018) 5, 394
Haidenbauer et al.,PRD 92 (2015) 5, 054032

- No cusp of $n\text{-}\bar{n}$ opening at $k^* \sim 50 \text{ MeV}/c \rightarrow$ in agreement with charge-exchange cross-sections
- rise of CF at low k^*
 - no agreement with Coulomb only
 - χEFT calculations with no explicit CC terms do not reproduce the data at low k^*
 - evidence of annihilation channels feeding into $p\text{-}\bar{p}$ pairs

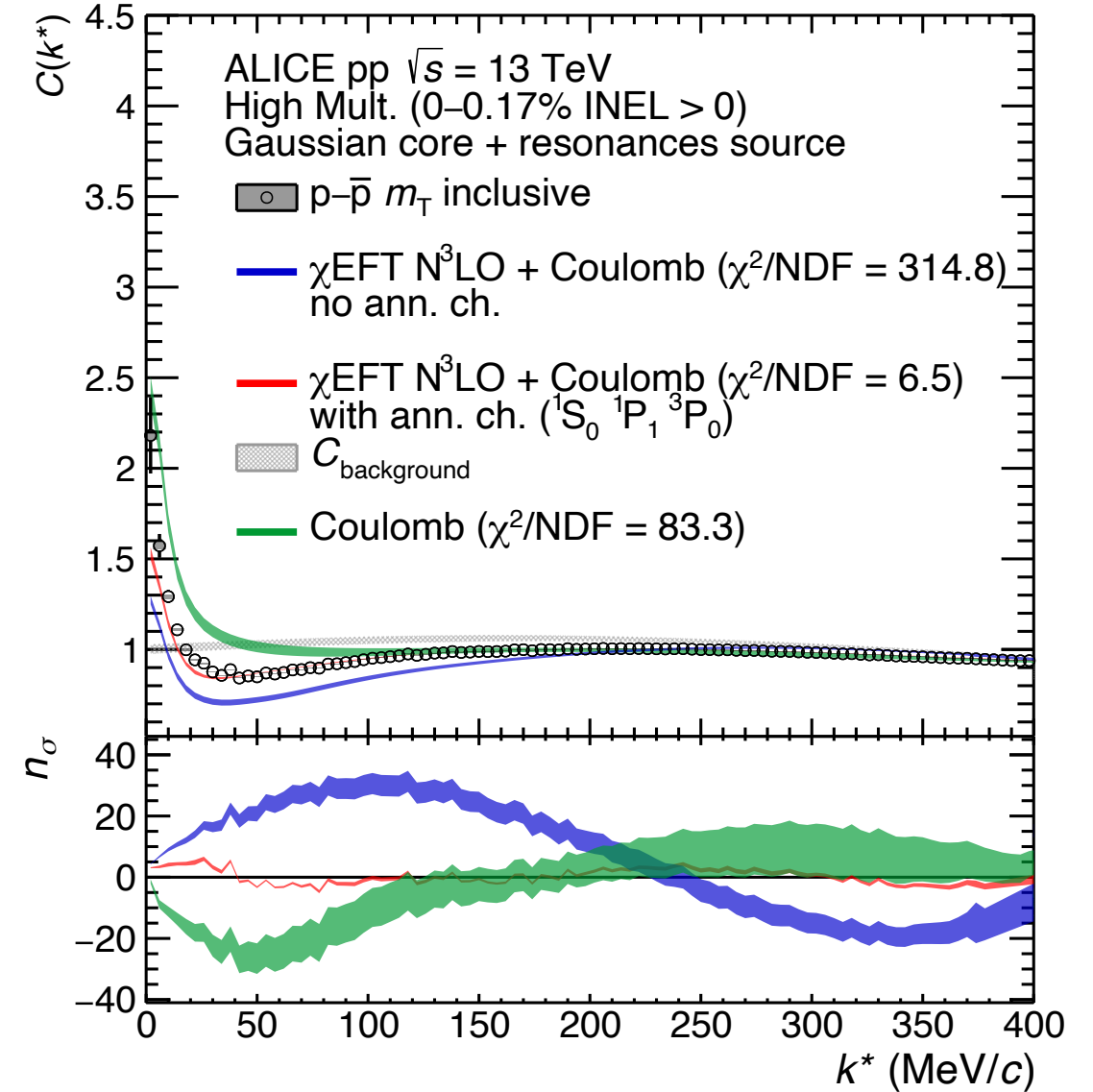


TUM Correlation function for $p\text{-}\bar{p}$

- No cusp of $n\text{-}\bar{n}$ opening at $k^* \sim 50 \text{ MeV}/c \rightarrow$ in agreement with charge-exchange cross-sections
- rise of CF at low k^*
 - no agreement with Coulomb only
 - χEFT calculations with no explicit CC terms do not reproduce the data at low k^*
 - evidence of annihilation channels feeding into $p\text{-}\bar{p}$ pairs
- **Annihilation channels $X \rightarrow p\text{-}\bar{p}$ included**
 - better agreement with the data is obtained
 - Dominant coupling weights in 3P_0 and 1S_0

$$\omega_{^3P_0} = 40.04 \pm 4.06 \text{ (stat)} \pm 4.24 \text{ (syst)}$$

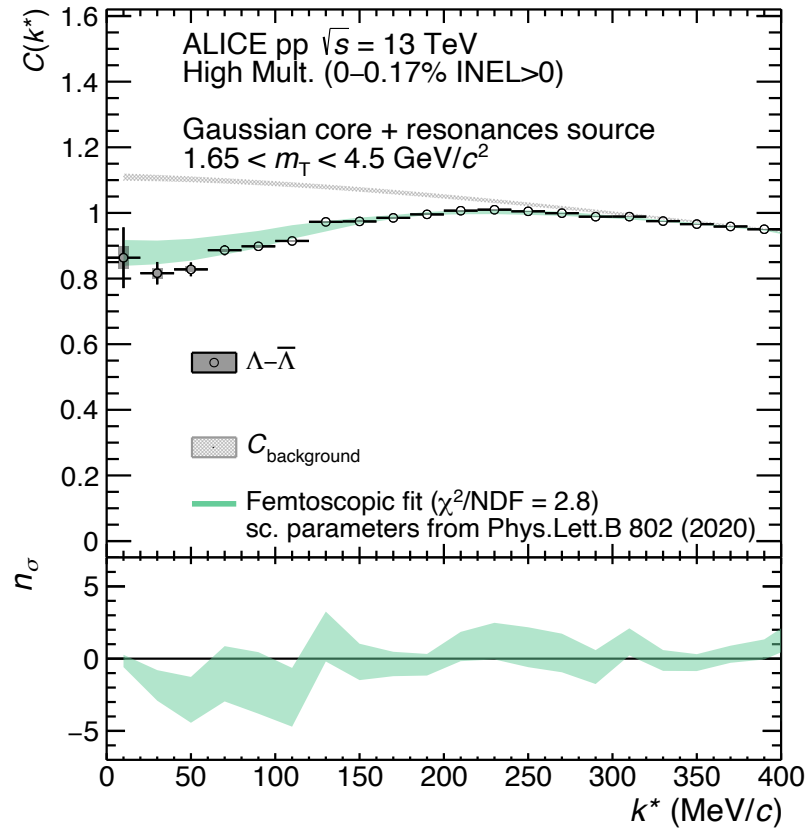
$$\omega_{^1S_0} = 1.19 \pm 0.10 \text{ (stat)} \pm 0.19 \text{ (syst)}$$



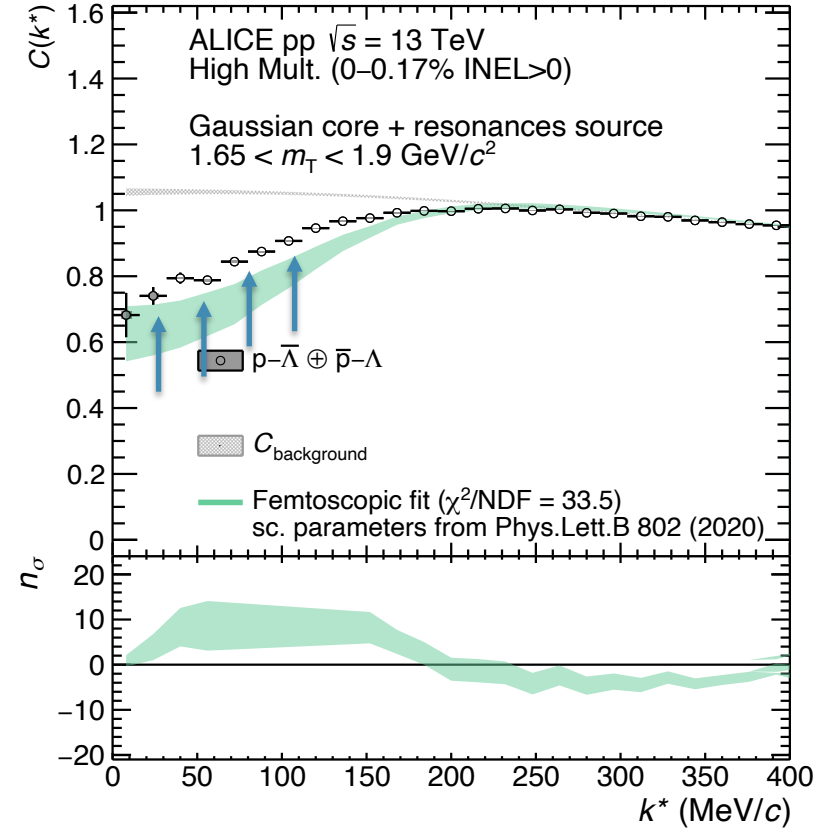
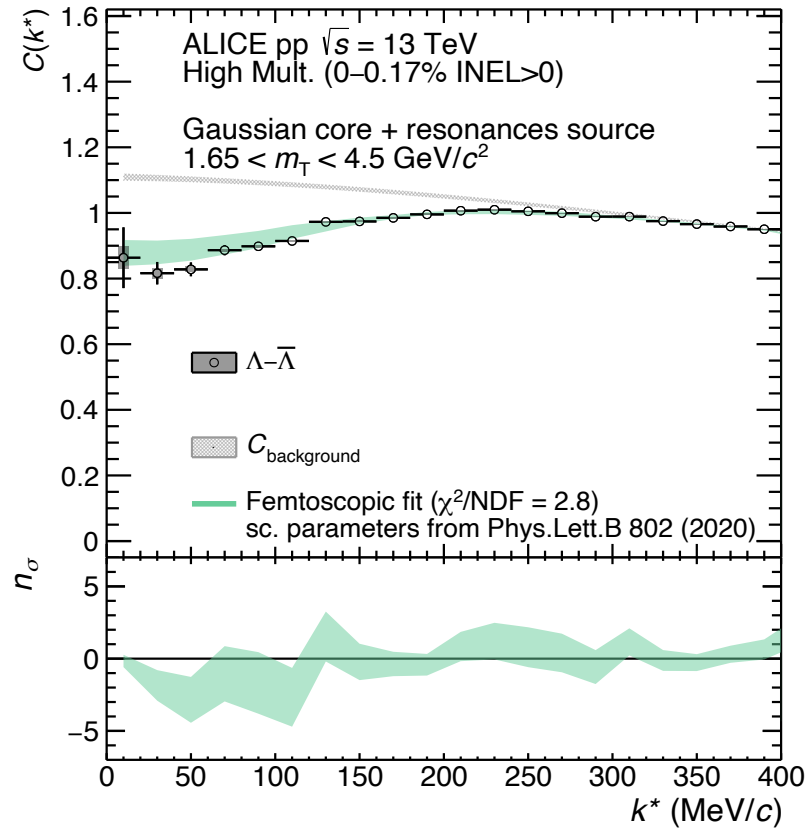
PWA: D.Zhou and R.G. E. Timmermans PRC86 (2012)

ALICE Coll. arXiv: 2105.05190

- No exact wavefunctions available \rightarrow single-channel Lednický-Lyuboshits formula
- Assuming the scattering parameters obtained in Pb-Pb
 - nice agreement with Λ - $\bar{\Lambda}$ data \rightarrow inelastic part present but not dominant

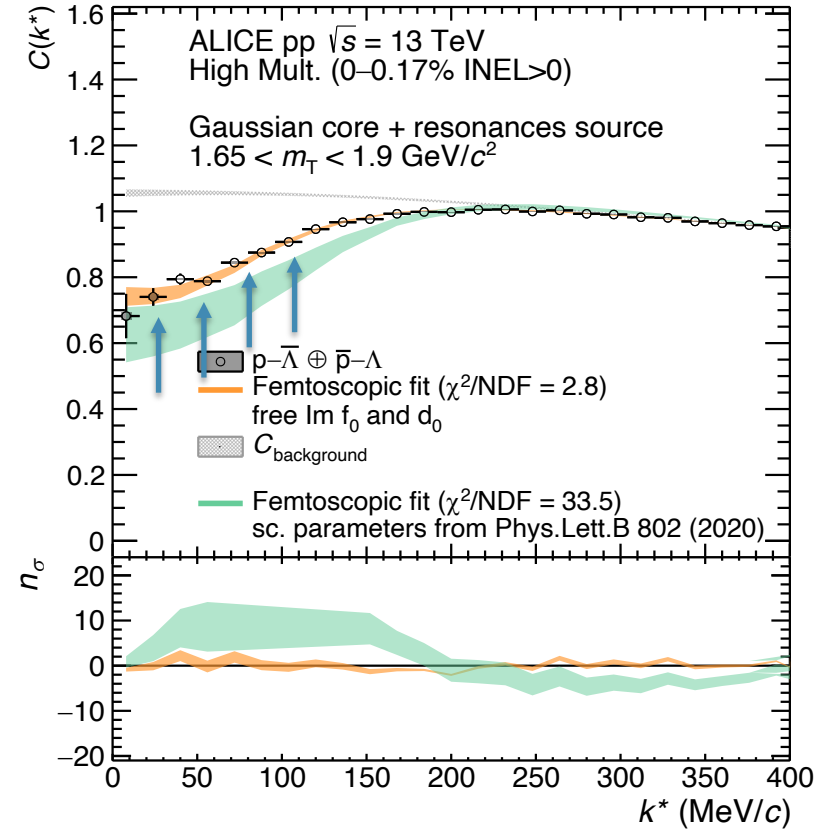
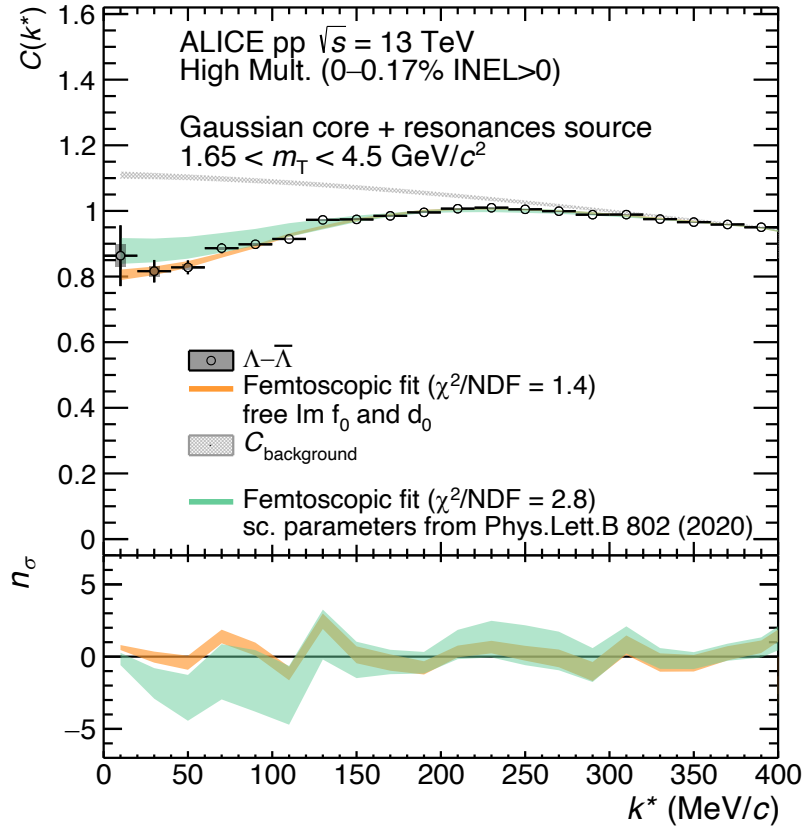


- No exact wavefunctions available \rightarrow single-channel Lednický-Lyuboshits formula
- Assuming the scattering parameters obtained in Pb-Pb
 - nice agreement with Λ - $\bar{\Lambda}$ data \rightarrow inelastic part present but not dominant
 - underestimate of p - $\bar{\Lambda}$ data \rightarrow large coupling to multi-meson annihilation channels



ALICE Coll. arXiv: 2105.05190

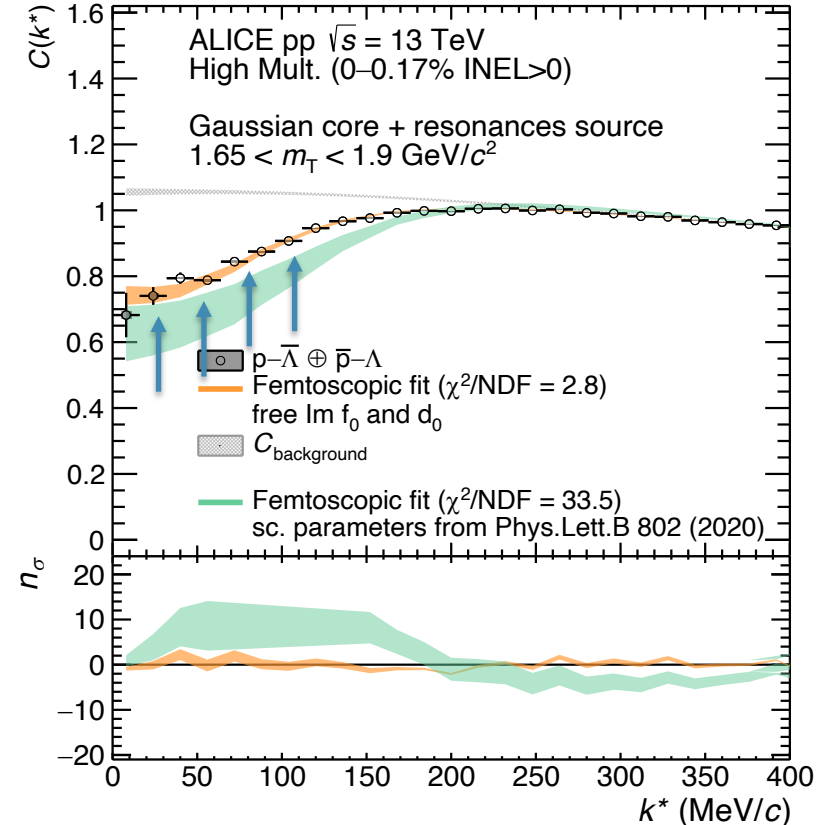
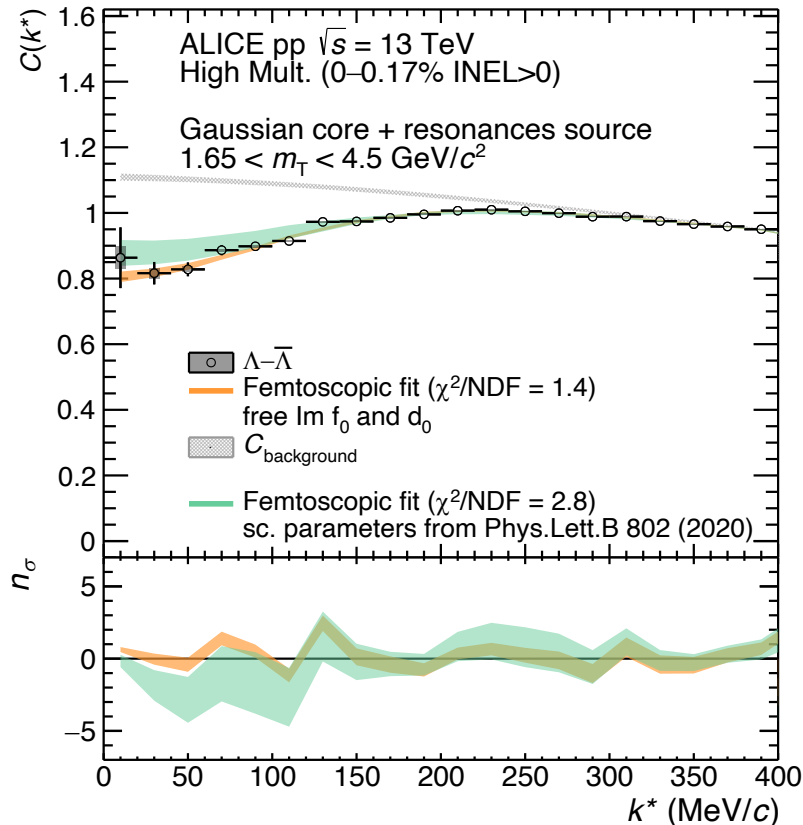
- Elastic part $\text{Re}(f_0)$ fixed from Pb-Pb data, free inelastic $\text{Im}(f_0)$ and d_0
 - extracted values for Λ - $\bar{\Lambda}$ are compatible with Pb-Pb scatt. pars
 - to reproduce p - $\bar{\Lambda}$ data $\text{Im}f_0$ has to be increased by a factor ~ 5.3
- Larger presence of multi-meson annihilation channels in p - $\bar{\Lambda} \rightarrow$ no bound-states?



ALICE Coll. arXiv: 2105.05190

- Estimates based on kinematics (EPOS) and SU(3) flavour symmetry for 2-meson channels ($\pi\pi, \pi K$)
 - similar amount of $p\text{-}\bar{\Lambda}$ and $\Lambda\text{-}\bar{\Lambda}$ pairs at low k^* ($\sim 6.4\%$)
 - coupling strength from meson-baryon SU(3) lagrangian provides largest coupling for $p\text{-}\bar{\Lambda}$

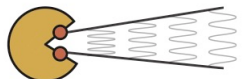
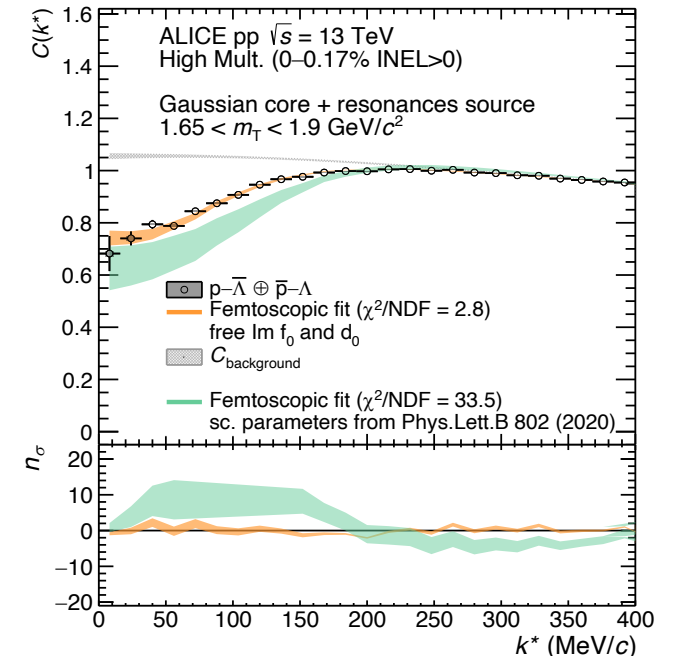
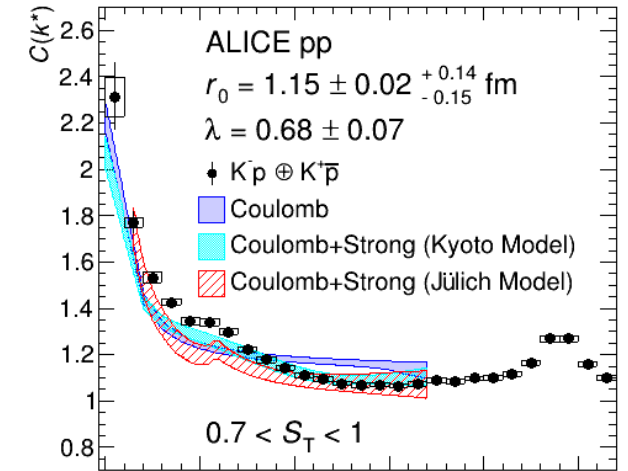
$$\frac{g_{2M \rightarrow p\text{-}\bar{\Lambda}} \times N_{2M \rightarrow p\text{-}\bar{\Lambda}}}{g_{2M \rightarrow \Lambda\text{-}\bar{\Lambda}} \times N_{2M \rightarrow \Lambda\text{-}\bar{\Lambda}}} \approx 6.3$$



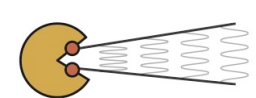
ALICE Coll. arXiv: 2105.05190

TUM Conclusions

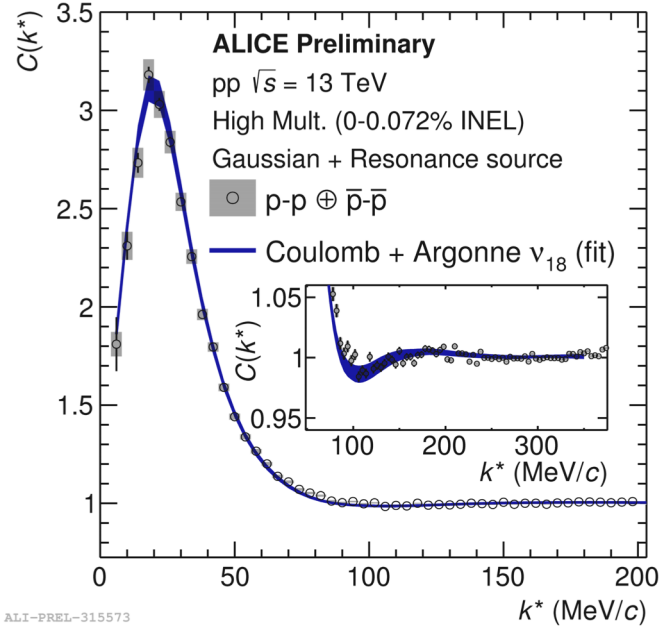
- Femtoscopy in small colliding systems
 - high-precision data at low momenta
 - sensitive to inelastic channels
- K^- - p in pp, p-Pb and Pb-Pb collisions
 - first experimental evidence of \bar{K}^0 - n channel
 - disappearance of CC contributions in Pb-Pb data
 - extend the analysis with detailed study on conversion weights
- Baryon-antibaryon in pp collisions
 - Λ - $\bar{\Lambda}$ results annihilation not dominant and room for baryonia
 - p - \bar{p} and p - $\bar{\Lambda}$ large presence of annihilation channels \rightarrow no formation of bound states?
 - need for theoretical input on p - $\bar{\Lambda}$ and Λ - $\bar{\Lambda}$ interactions



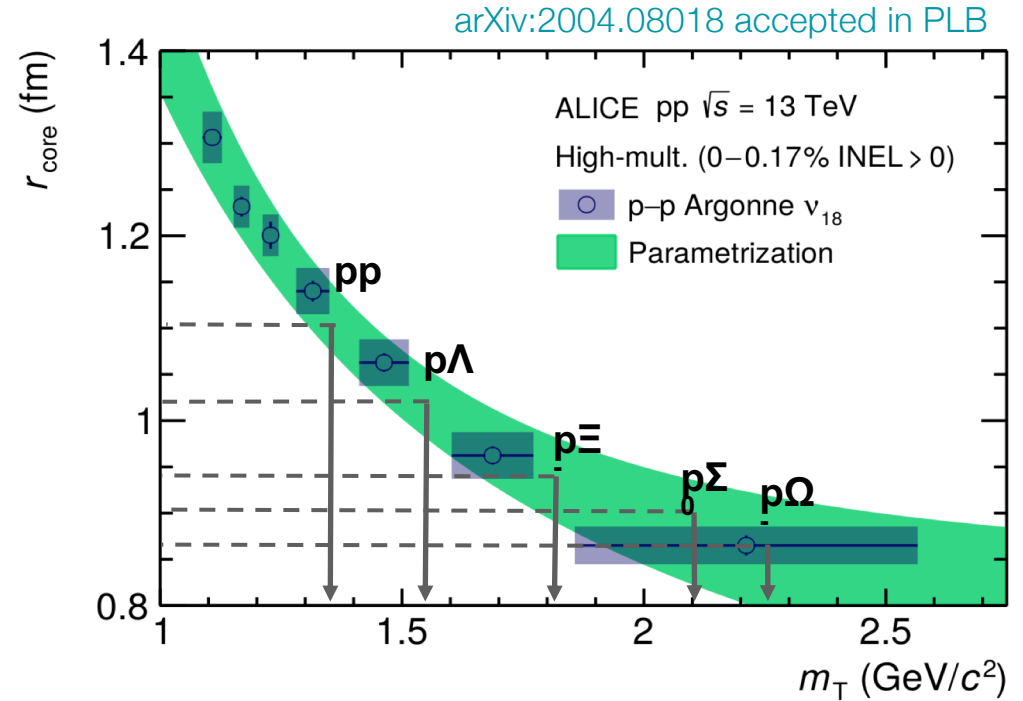
Additional slides



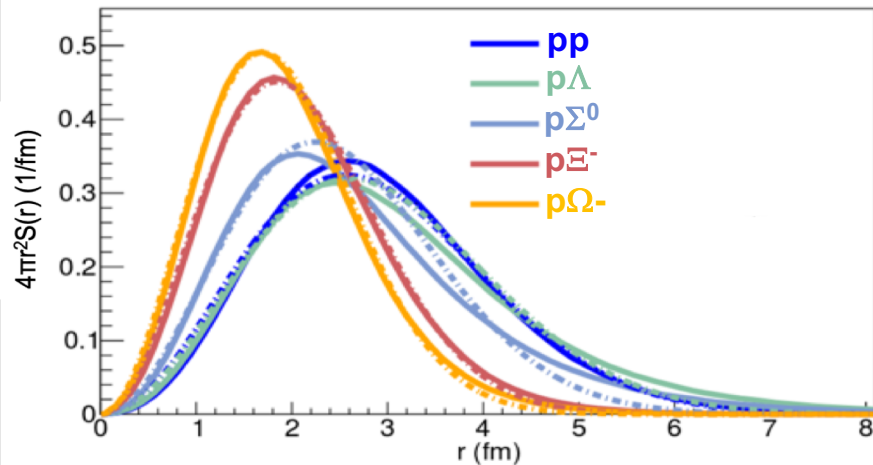
Fit with a 'core' Gaussian + Resonances



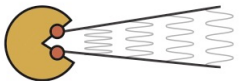
Core Radius



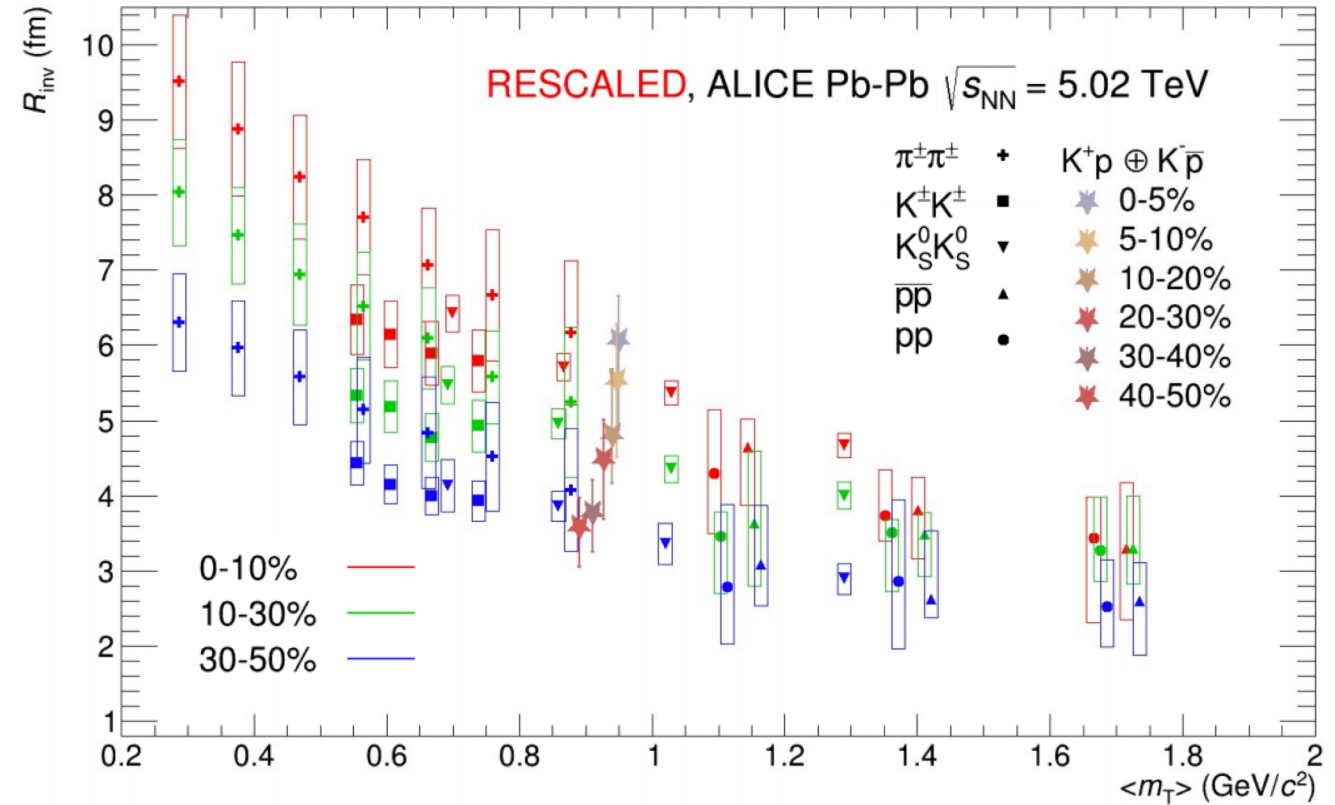
Effective Radius



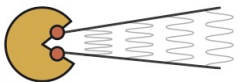
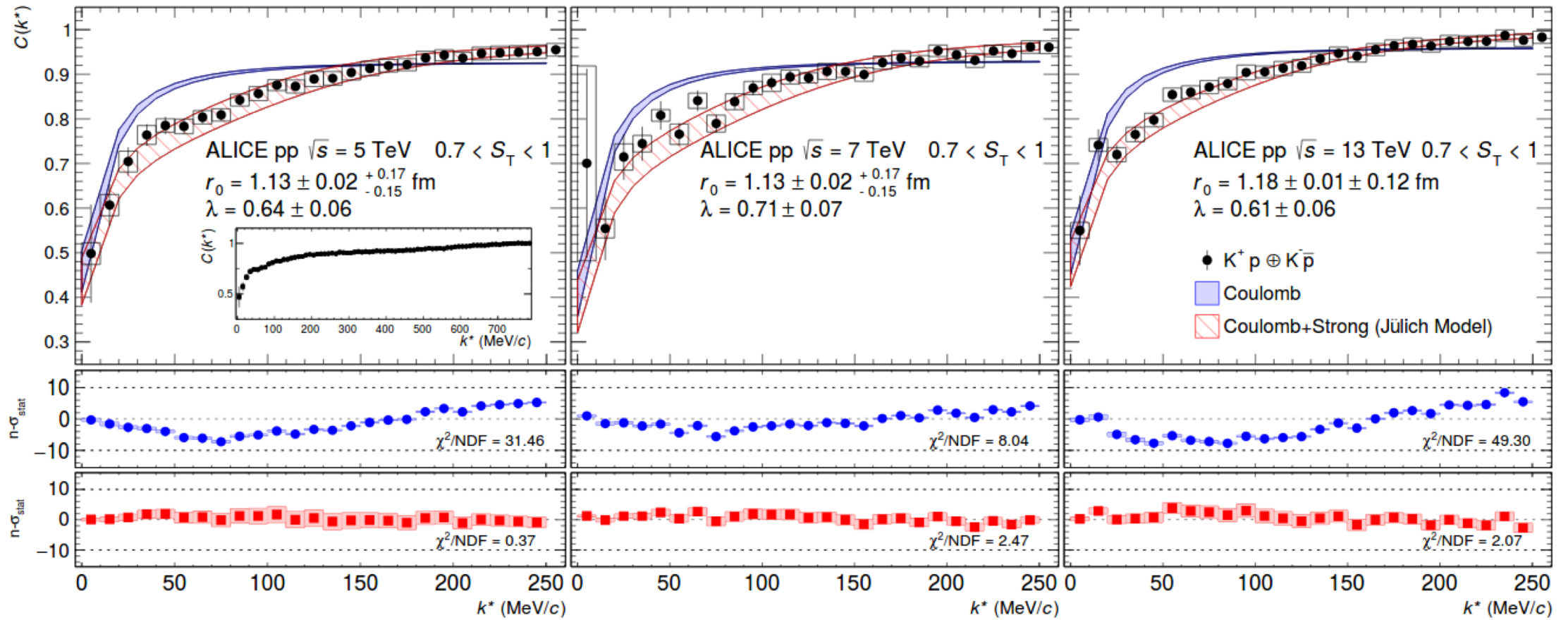
Pair	r_{Core} [fm]	r_{Eff} [fm]
pp	0.96	1.28
p Λ	0.88	1.3
p Σ^0	0.75	1.12
p Ξ^-	0.8	0.92
p Ω^-	0.73	0.85

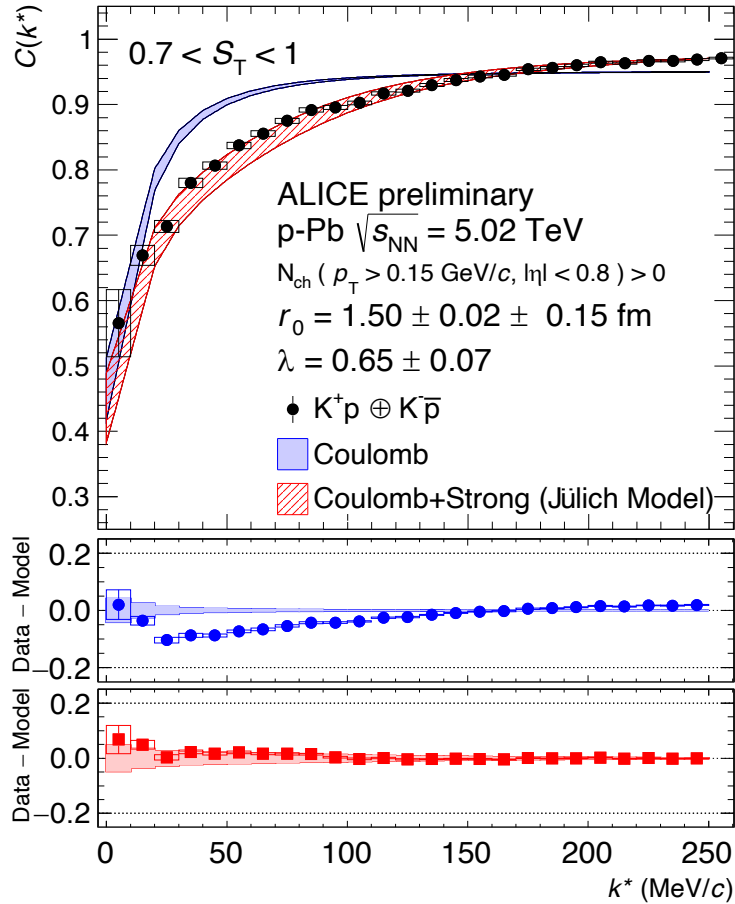


- Evolution of the system in heavy-ion is affected by freeze-out and rescattering → scaling with multiplicity and m_T
- For Pb–Pb studies at 5.02 TeV:
 - radius for $B-\bar{B}$ pairs same as for B-B
 - radius for K-p fixed from K+p
 - agreement with rescaled results on pion, protons and kaon pairs at 2.76 TeV

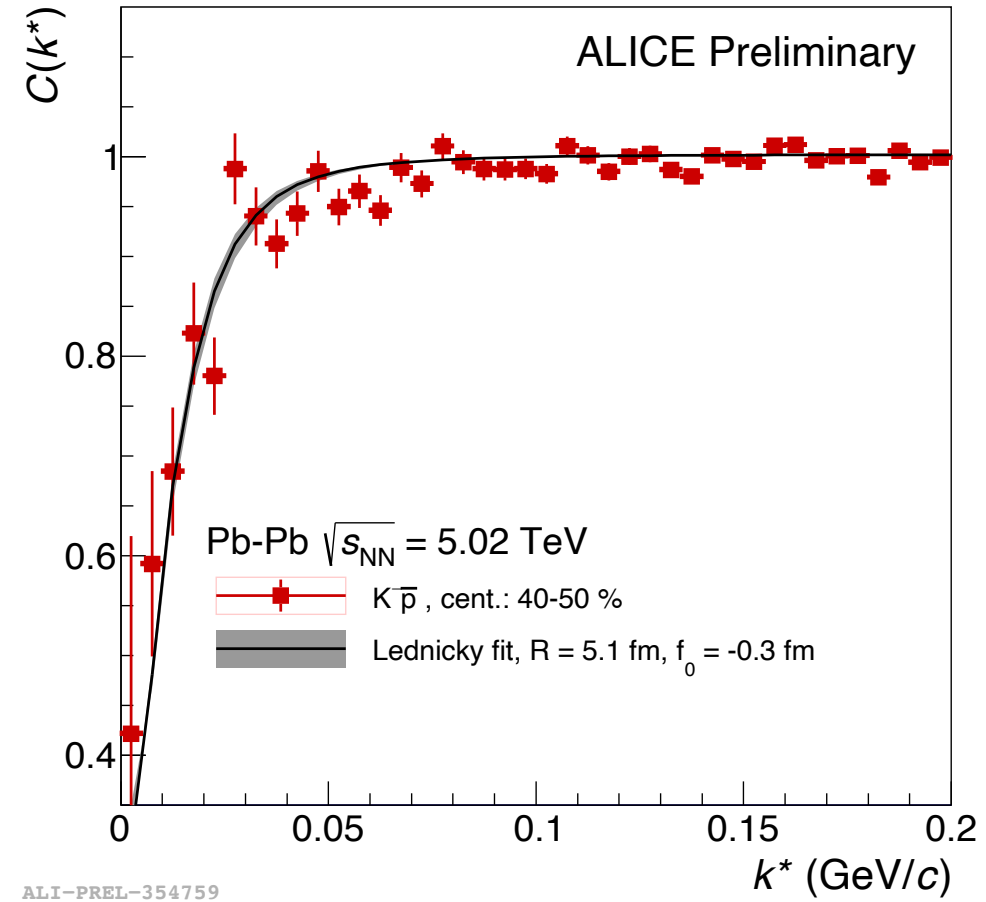


ALICE Coll. Phys.Rev.Lett. 124 (2020) 9, 092301

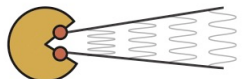




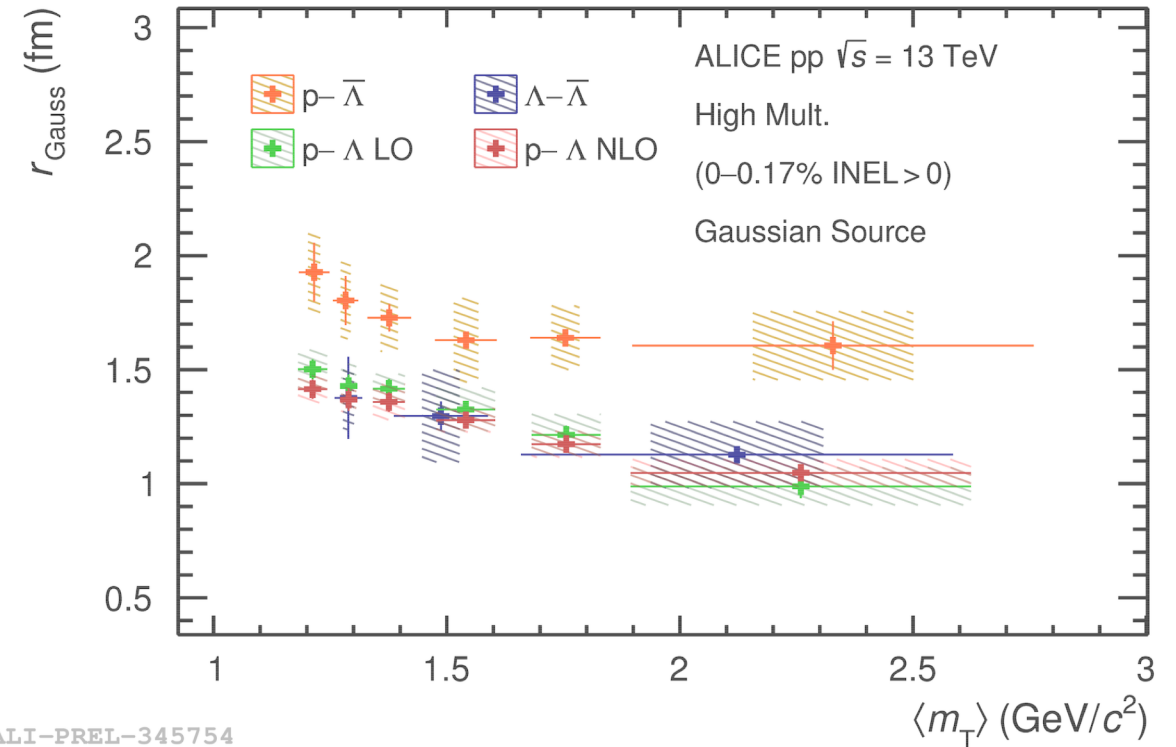
ALI-PREL-316319



ALI-PREL-354759



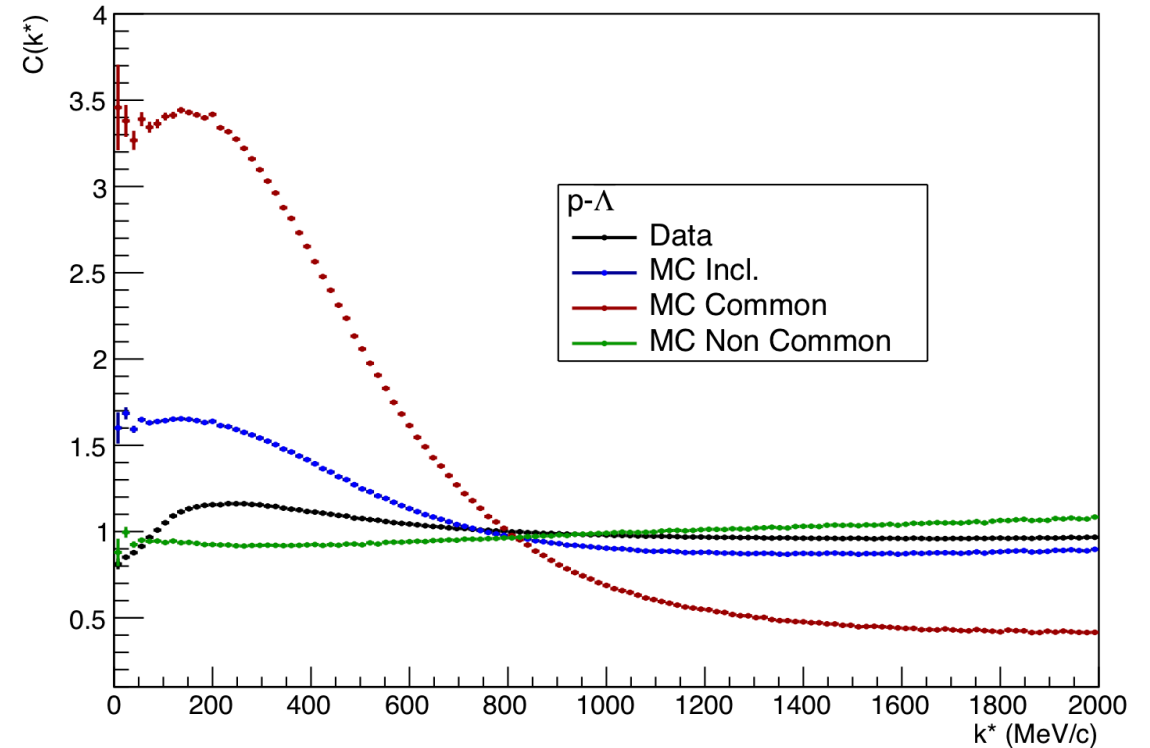
- Λ -anti Λ correlation function: well reproduced by Lednicky-Lyuboshits (Pb-Pb), extracted radius in agreement with B-B scaling
- p -anti Λ correlation function: well reproduced by Lednicky-Lyuboshits, BUT extracted radius larger wrt to B-B scaling
- Physical interpretation: p -anti Λ is more coupled-channels contributions (supported by $SU(3)_f$ considerations) \Rightarrow larger amount of inelastic channels translates into larger radius



- Large background that cannot be removed by applying S_T / S_0 cuts \Rightarrow results are obtained WITHOUT any S_T/S_0 cuts
- Use of Pythia to model the background:
 - jetty contributions \Rightarrow particles from Common ancestors
 - non-jetty part \Rightarrow particles from Non Common Ancestors

$$C_{mini-jet}(k^*) = [w_C C_C(k^*) + (1 - w_C) C_{NC}(k^*)]$$

- Weights w_C to be fitted to data since they cannot be anchored to $\Delta\eta\Delta\phi$ correlations (more than a rescaling is needed)

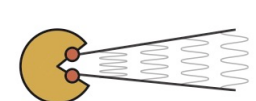


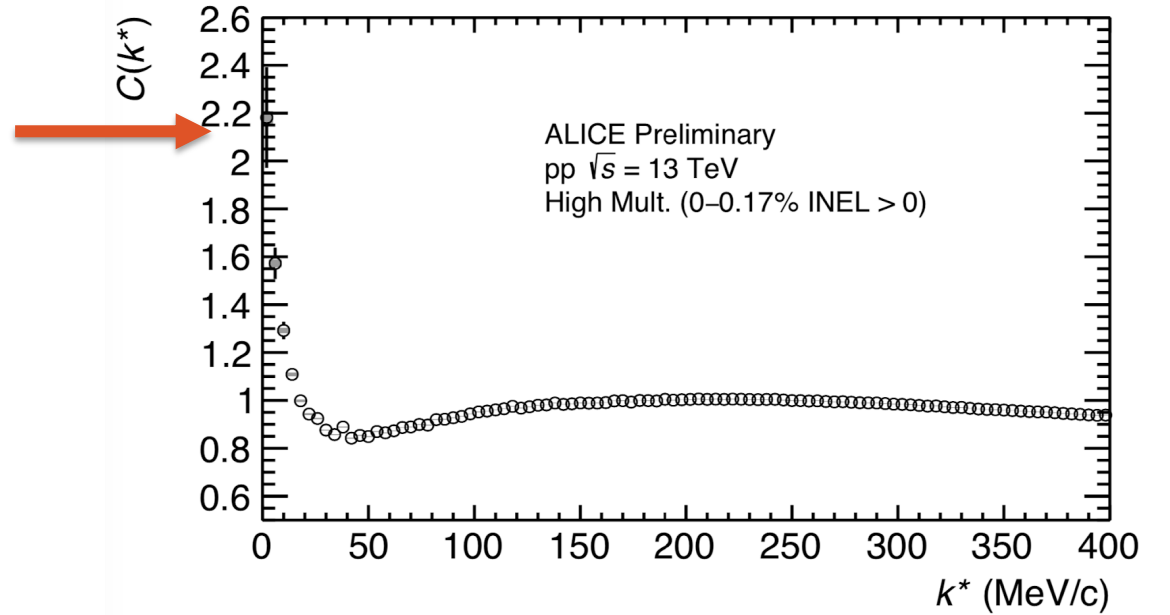
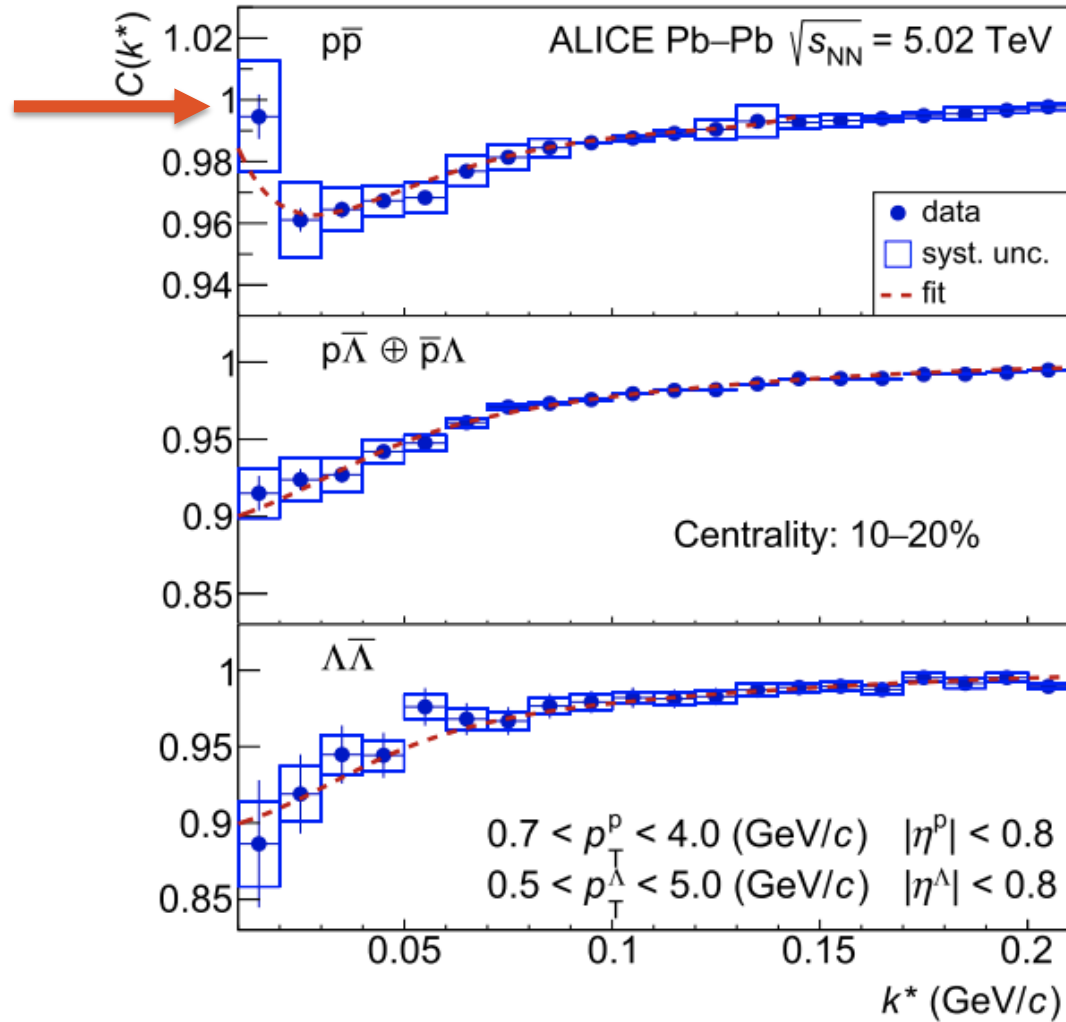
- Residual contributions included through λ parameters (DCA/CPA Template Fits)
- Non-femtoscopic background:
 - Mini-jet background \Rightarrow Shape fixed by Ancestors Template
 - Large k^* kinematics effects \Rightarrow Pol1/Pol2 (prefit in k^* [400-2500] MeV/c and kept fixed in the final fit)

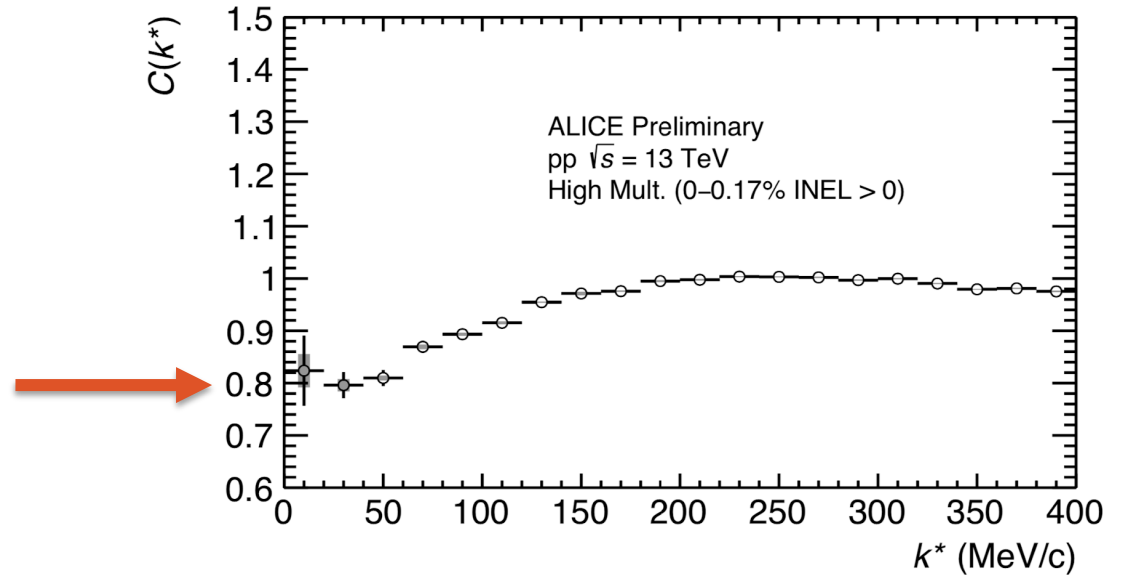
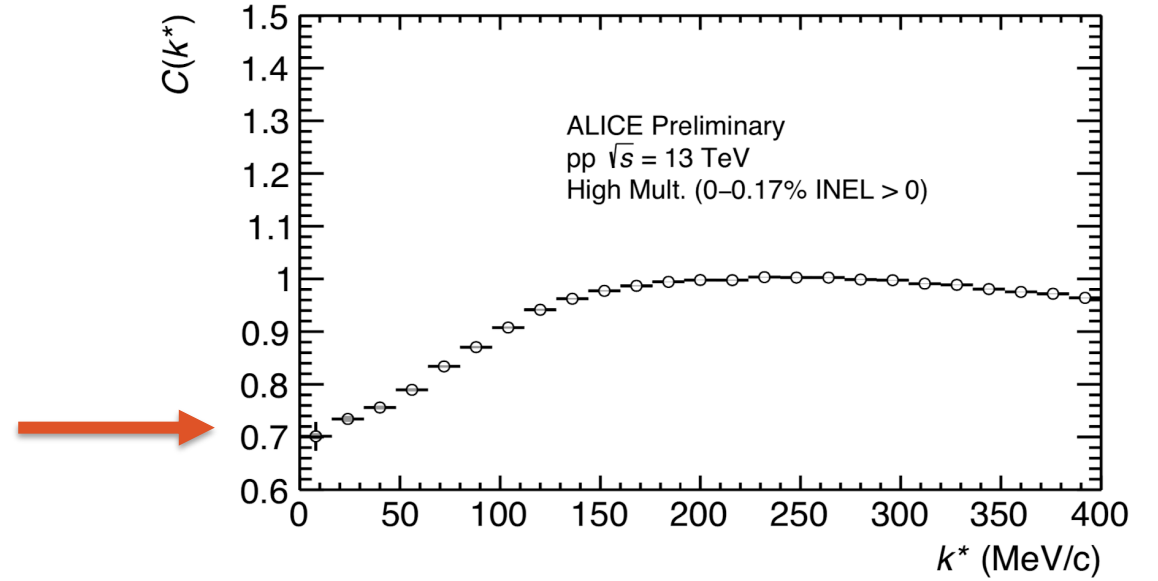
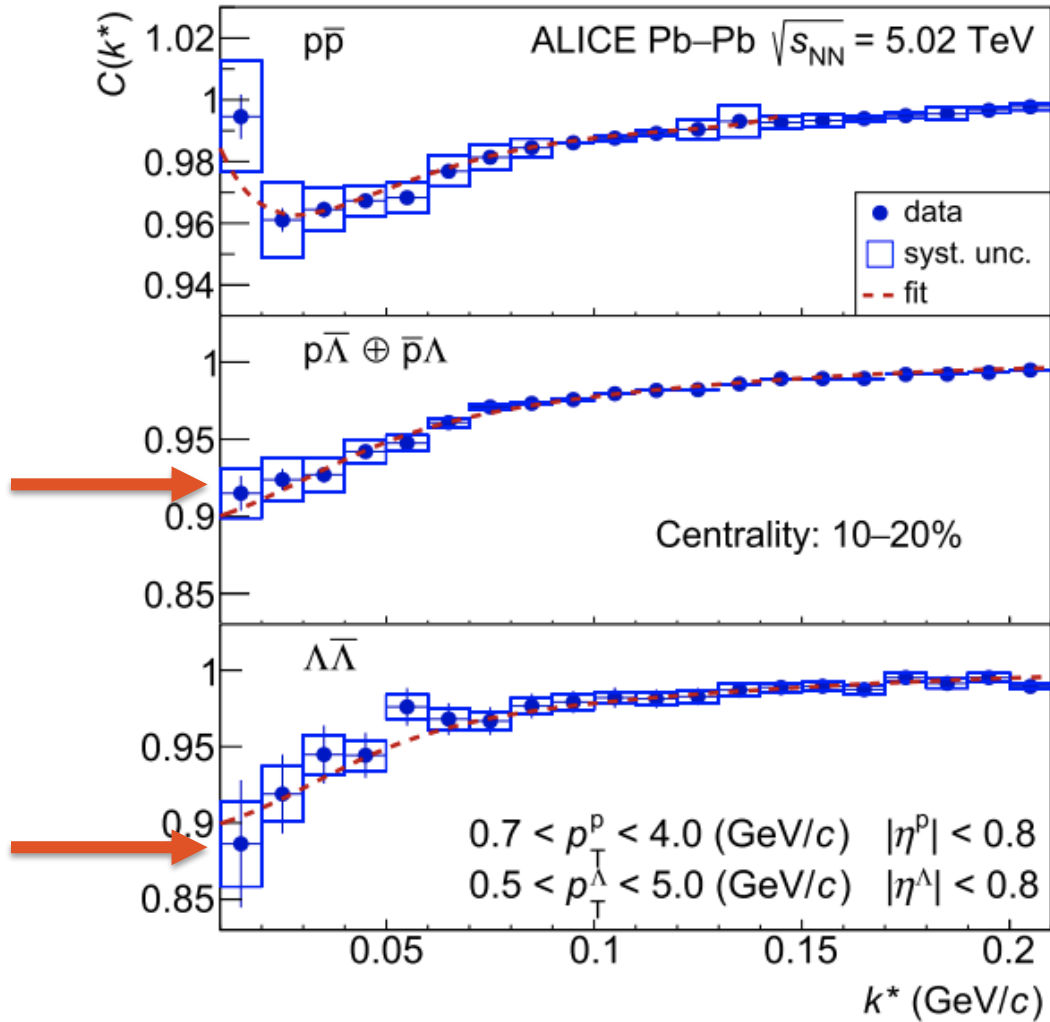
$$C_{BCKG}(k^*) = [w_C C_C(k^*) + (1 - w_C) C_{NC}(k^*) + (a + bk^* + c(k^*)^2)]$$

- Total correlation function:
 - Free parameters: weights w_C , Norm N_D
 - Coupled-channel modeling affects ONLY C_{model}

$$C_{tot}(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{BCKG}(k^*),$$







- Pairs close in mass with the same quantum numbers: e.g. $p-\Xi^-$ and $\Lambda-\Lambda$
- Schrödinger equation of one pair \rightarrow Equation system of all 1, ... , N pairs

$$\hat{\mathcal{H}} \cdot \psi = E \cdot \psi \mapsto \begin{pmatrix} \hat{\mathcal{H}}_{11} & \cdots & \hat{\mathcal{H}}_{1N} \\ \vdots & \ddots & \vdots \\ \hat{\mathcal{H}}_{N1} & \cdots & \hat{\mathcal{H}}_{NN} \end{pmatrix} \cdot \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_N \end{pmatrix} = E \cdot \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_N \end{pmatrix}$$

•

1) Coupled channels influence the elastic channels of the two-particle wave function ψ_j

