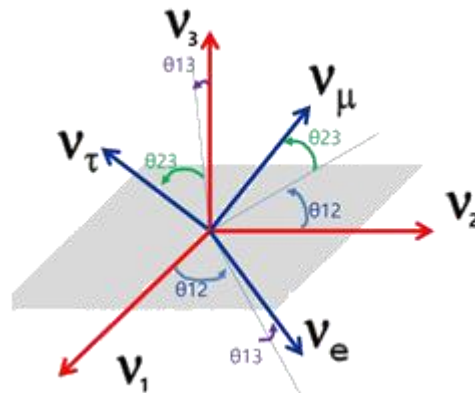
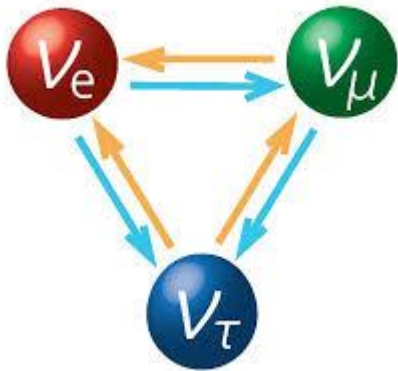


Neutrino mixing from flavour symmetry

Gui-Jun Ding

University of Science and Technology of China

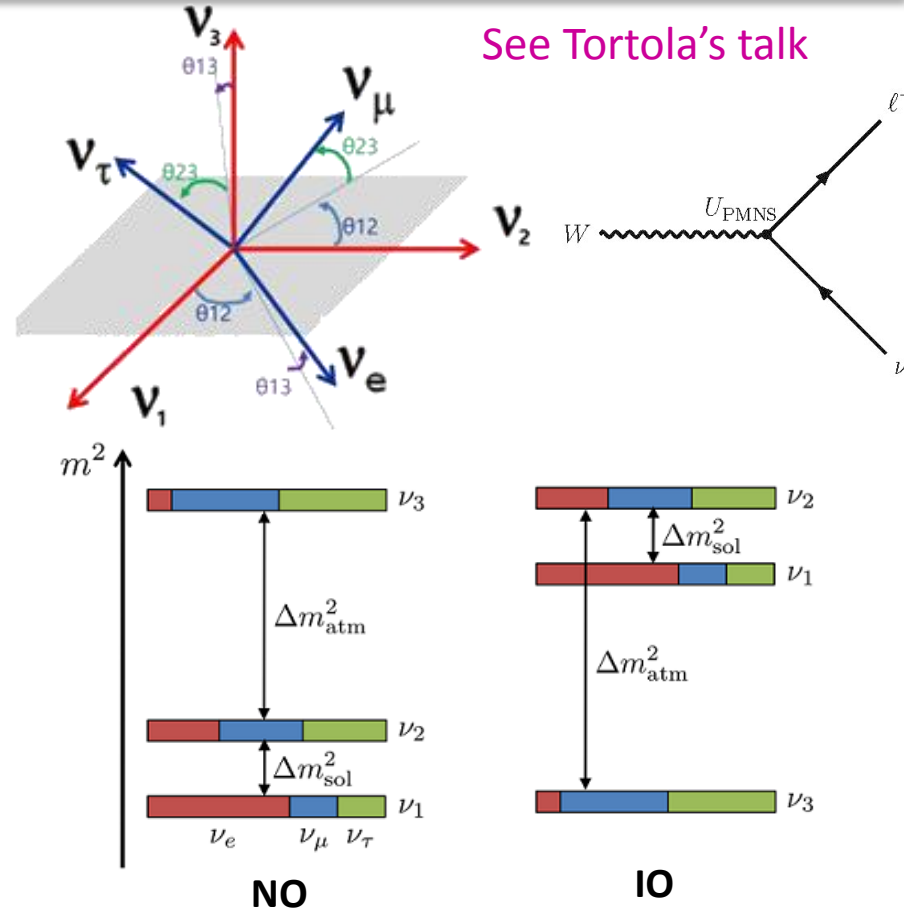
22nd Particles and Nuclei International Conference,
September 8th, 2021, Lisbon, Portugal



Current experimental results on lepton mixing

[P.F. de Salas et al, arXiv: 2006.11237]

parameter	best fit $\pm 1\sigma$	2σ range	3σ range
$\Delta m_{21}^2 [10^{-5} \text{eV}^2]$	$7.50^{+0.22}_{-0.20}$	7.12–7.93	6.94–8.14
$ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$ (NO)	$2.55^{+0.02}_{-0.03}$	2.49–2.60	2.47–2.63
$ \Delta m_{31}^2 [10^{-3} \text{eV}^2]$ (IO)	$2.45^{+0.02}_{-0.03}$	2.39–2.50	2.37–2.53
$\sin^2 \theta_{12} / 10^{-1}$	3.18 ± 0.16	2.86–3.52	2.71–3.69
$\theta_{12} / ^\circ$	34.3 ± 1.0	32.3–36.4	31.4–37.4
$\sin^2 \theta_{23} / 10^{-1}$ (NO)	5.74 ± 0.14	5.41–5.99	4.34–6.10
$\theta_{23} / ^\circ$ (NO)	49.26 ± 0.79	47.37–50.71	41.20–51.33
$\sin^2 \theta_{23} / 10^{-1}$ (IO)	$5.78^{+0.10}_{-0.17}$	5.41–5.98	4.33–6.08
$\theta_{23} / ^\circ$ (IO)	$49.46^{+0.60}_{-0.97}$	47.35–50.67	41.16–51.25
$\sin^2 \theta_{13} / 10^{-2}$ (NO)	$2.200^{+0.069}_{-0.062}$	2.069–2.337	2.000–2.405
$\theta_{13} / ^\circ$ (NO)	$8.53^{+0.13}_{-0.12}$	8.27–8.79	8.13–8.92
$\sin^2 \theta_{13} / 10^{-2}$ (IO)	$2.225^{+0.064}_{-0.070}$	2.086–2.356	2.018–2.424
$\theta_{13} / ^\circ$ (IO)	$8.58^{+0.12}_{-0.14}$	8.30–8.83	8.17–8.96
δ / π (NO)	$1.08^{+0.13}_{-0.12}$	0.84–1.42	0.71–1.99
$\delta / ^\circ$ (NO)	194^{+24}_{-22}	152–255	128–359
δ / π (IO)	$1.58^{+0.15}_{-0.16}$	1.26–1.85	1.11–1.96
$\delta / ^\circ$ (IO)	284^{+26}_{-28}	226–332	200–353



- Octant of θ_{23} : $>$ or $< 45^\circ$?
- What is the value of δ_{CP} ?

- Mass hierarchy: **NO** or **IO**?
- Absolute mass scale: $m_{\text{lightest}} = ?$
- **Majorana** or **Dirac** neutrinos?
- Why m_ν so small?
- Connections to other new physics?

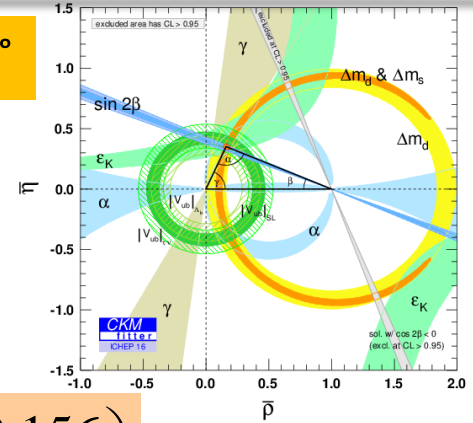
CKM vs PMNS

[Particle Data Group 2020]

$$\alpha = (88.8 \pm 2.3)^\circ$$

Quarks:

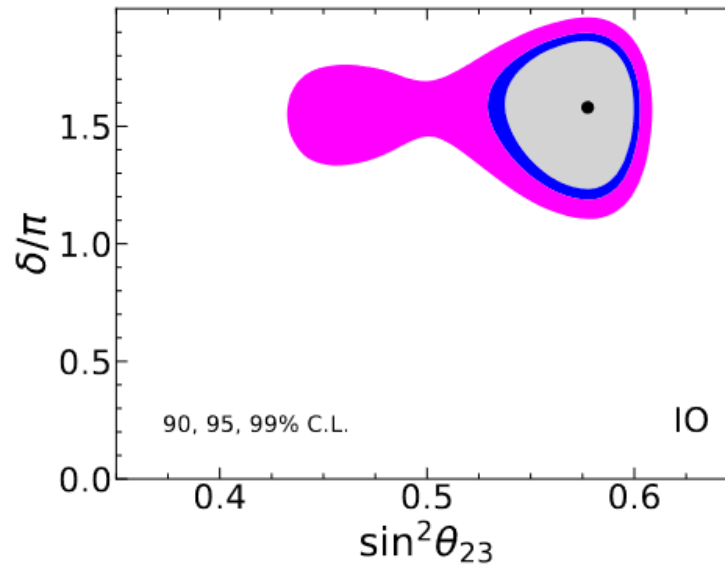
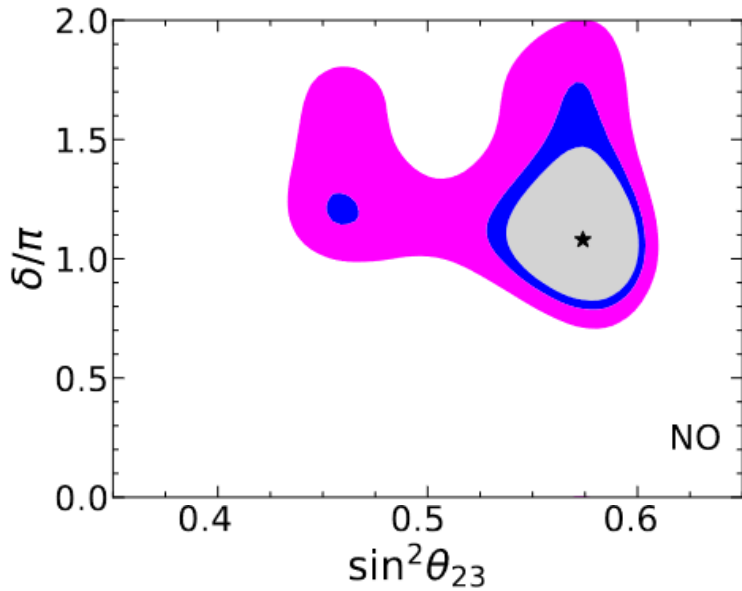
$$|V_{CKM}| \approx \begin{pmatrix} 0.97434 & 0.22506 & 0.00357 \\ 0.22492 & 0.97351 & 0.0414 \\ 0.00875 & 0.0403 & 0.99915 \end{pmatrix}$$



Leptons:

$$|U_{PMNS}| = \begin{pmatrix} 0.799 \sim 0.844 & 0.516 \sim 0.582 & 0.141 \sim 0.156 \\ 0.242 \sim 0.494 & 0.467 \sim 0.678 & 0.639 \sim 0.774 \\ 0.284 \sim 0.521 & 0.490 \sim 0.695 & 0.615 \sim 0.754 \end{pmatrix}$$

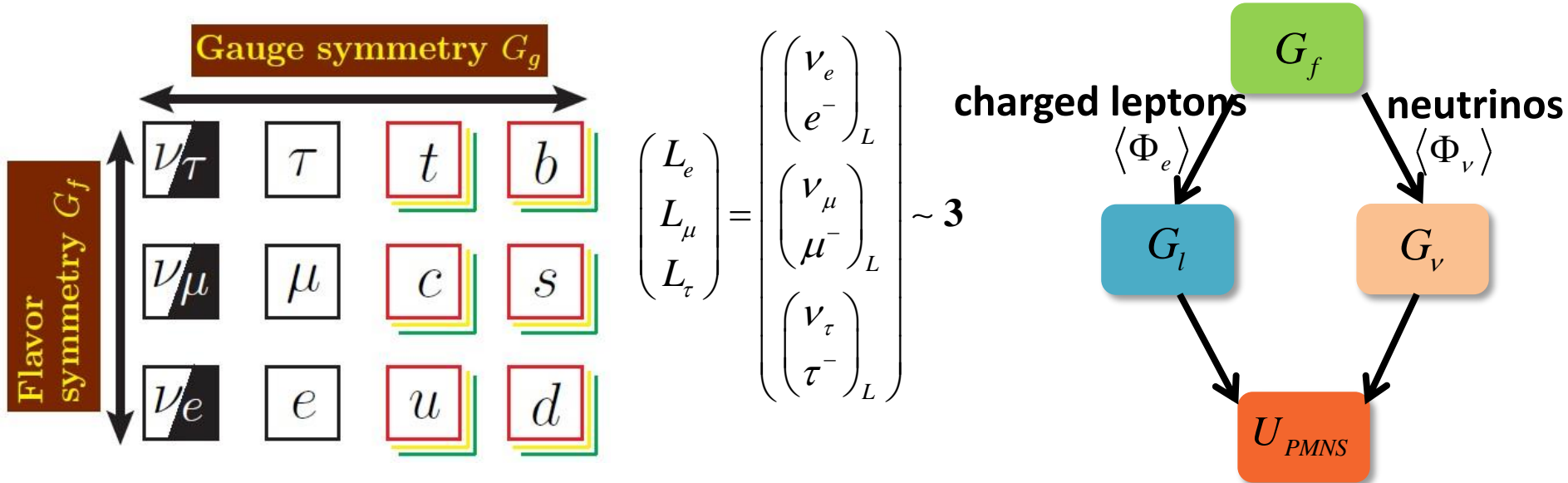
[P.F. de Salas et al, arXiv: 2006.11237]



Why quark and lepton mixings are so different?

Symmetry as a guiding principle to flavor puzzle

Relate lepton mixing to how are G_f broken, lepton mixing matrix arises from mismatch of the different residual subgroups G_l and G_ν



$$\mathcal{L}_m = -Y_{ij}^e(\langle \Phi_e \rangle) \bar{L}_i H e_{Rj} - \frac{1}{2} Y_{ij}^\nu(\langle \Phi_\nu \rangle) \bar{L}_i^c H H^T L_j + \dots$$

	Continuous	Discrete
Abelian	U(1)	Z_n
Non-Abelian	U(2), SU(3), SO(3) ...	$A_4, S_4, A_5, \Delta(6n^2), \dots$

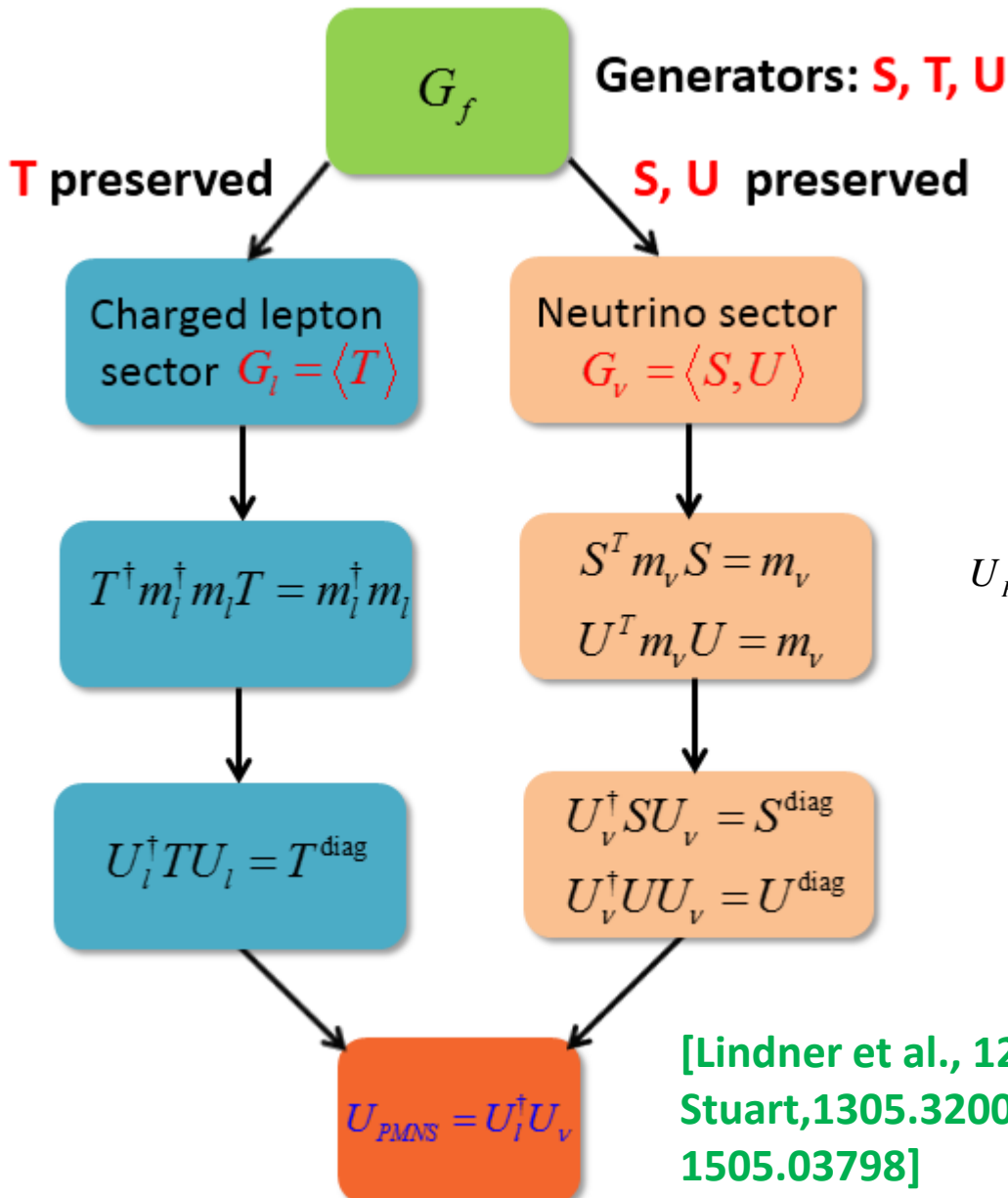
high dimensional operators:
 $\bar{\Phi}_e, \bar{\Phi}_\nu$ cross terms

[Altarelli, Feruglio, 1002.0211; Tanimoto et al., 1003.3552; King and Luhn, 1301.1340; Feruglio, Romanino, arXiv:1912.06028]

Direct approach of flavor symmetry

[King and Luhn, 1301.1340]

For Majorana neutrinos



- Klein symmetry **S, U** and **T** are **each** identified as subgroups of some flavor symmetry G_f
- Only **trimaximal** mixing can be compatible with data

$$U_{PMNS} = \frac{1}{\sqrt{3}} \begin{pmatrix} \sqrt{2} \cos \theta & 1 & -\sqrt{2} \sin \theta \\ -\sqrt{2} \cos(\theta - \pi/3) & 1 & \sqrt{2} \sin(\theta - \pi/3) \\ -\sqrt{2} \cos(\theta + \pi/3) & 1 & \sqrt{2} \sin(\theta + \pi/3) \end{pmatrix}$$

θ depends on groups G_f, G_l, G_ν

- Dirac CP phase δ_{CP} is predicted to be **conserved**.

[Lindner et al., 1212.2411; King, Neder, Stuart, 1305.3200; Fonseca and Grimus, 1405.3678; Yao, Ding, 1505.03798]

➤ If the lepton mixing matrix is **partially** determined by the flavor symmetry G_f , G_l and G_ν , e.g. $G_\nu = Z_2$

[Ge, Dicus, Repko, 1104.0602;
Hernandez and Smirnov, 1204.0445]

For example, two deformations of TBM

TM₁

$$U = U_{TBM} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta & \sin \theta e^{-i\delta} \\ 0 & -\sin \theta e^{i\delta} & \cos \theta \end{pmatrix},$$

TM₂

$$U = U_{TBM} \begin{pmatrix} \cos \theta & 0 & \sin \theta e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin \theta e^{i\delta} & 0 & \cos \theta \end{pmatrix}$$

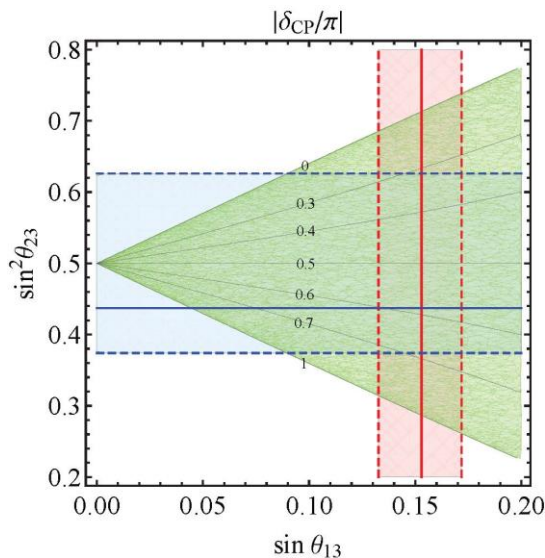
Two predictions in terms of sum rules

$$3 \cos^2 \theta_{12} \cos^2 \theta_{13} = 2$$

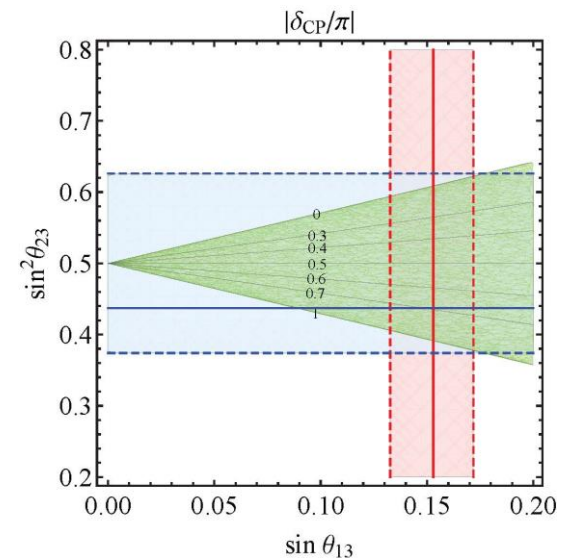
$$\sin^2 \theta_{23} \approx \frac{1}{2} - \sqrt{2} \sin \theta_{13} \cos \delta_{CP}$$

$$3 \sin^2 \theta_{12} \cos^2 \theta_{13} = 1$$

$$\sin^2 \theta_{23} \approx \frac{1}{2} + \frac{1}{\sqrt{2}} \sin \theta_{13} \cos \delta_{CP}$$

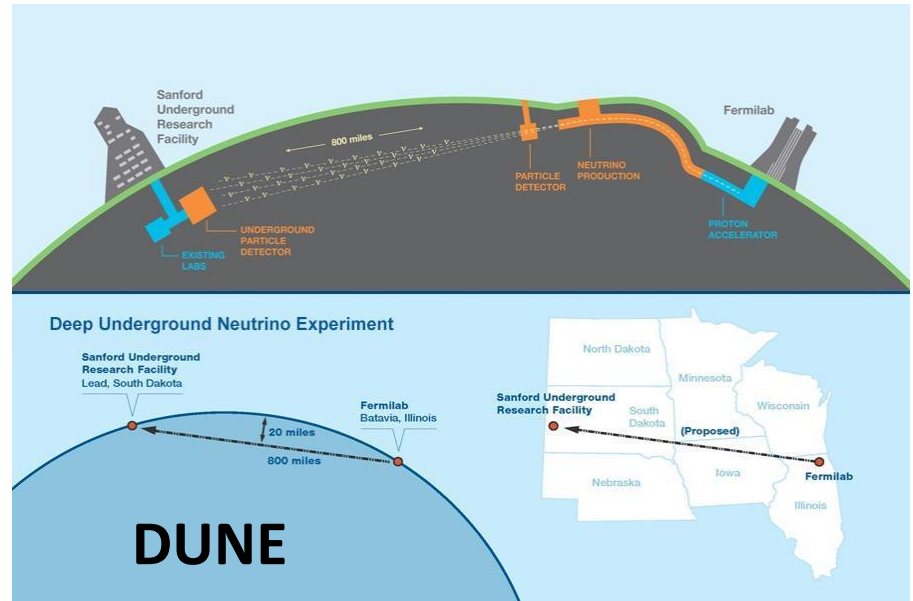
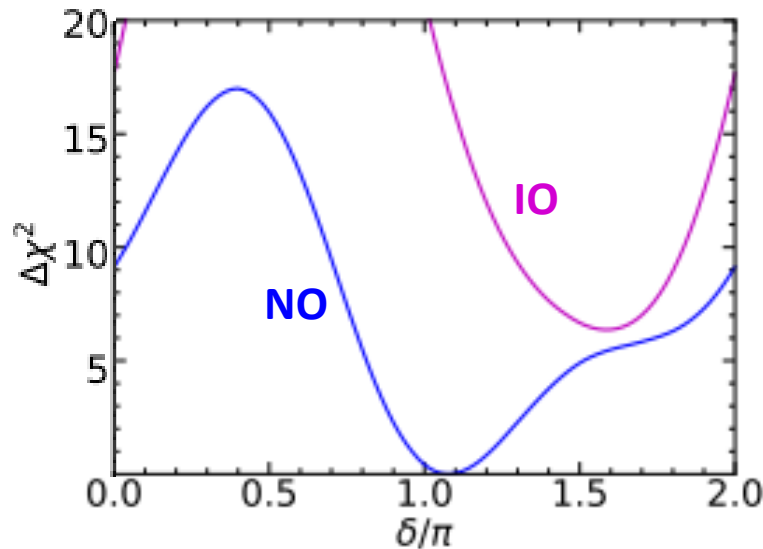


He, Zee, 07 and 11; Grimus, Lavoura, 08; Albright, Rodejohann, 09; King, Luhn 11; Chen, Vale et., 1812.04663, 1905.11997...

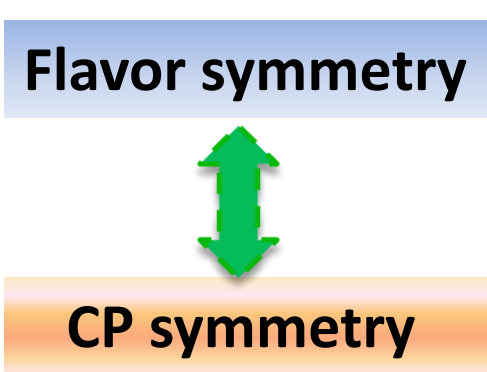


Leptonic CP violation

[P.F. de Salas et al, arXiv: 2006.11237]

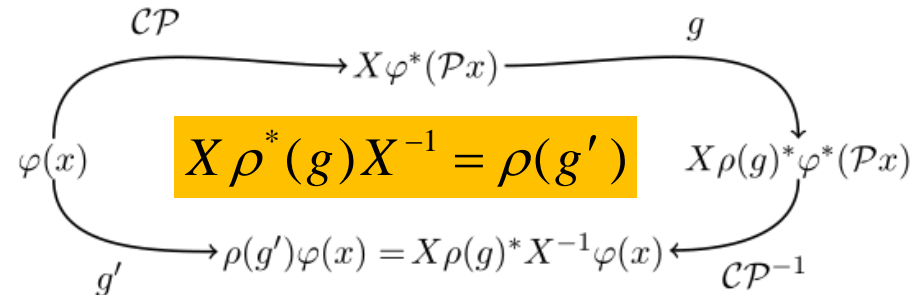


➤ Theoretical idea: **flavor symmetry** → **flavor+CP symmetries**



Mixing angles
& CP phases

"closure" relations have to hold!



[Grimus, Rebelo, hep-ph/9506272; Feruglio et al., 1211.5560; Lindner et al., 1211.6953]

a simple predictive CP symmetry: $\mu\tau$ reflection

$\mu\tau$ reflection = $\mu\tau$ exchange + canonical CP

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \xrightarrow{\text{CP}} \begin{pmatrix} \nu_e^c \\ \nu_\tau^c \\ \nu_\mu^c \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} \nu_e^c \\ \nu_\mu^c \\ \nu_\tau^c \end{pmatrix}$$

This CP transformation is **not** a unit matrix.

If the neutrino mass matrix is **invariant** under the $\mu\tau$ reflection

$$m_\nu = \begin{pmatrix} a & b & b^* \\ b & c & d \\ b^* & d & c^* \end{pmatrix} \xrightarrow{\text{CP}} |U_{PMNS}| = \begin{bmatrix} \text{red} & \text{green} & \text{purple} \\ \text{blue} & \text{green} & \text{orange} \\ \text{blue} & \text{green} & \text{orange} \end{bmatrix} \xrightarrow{\text{CP}} \theta_{23} = \frac{\pi}{4}, \quad \delta_{CP} = \pm \frac{\pi}{2}$$

$\nu_e \quad \nu_\mu \leftrightarrow \nu_\tau^c$

The last two rows have equal magnitudes

[Harrison, Scott, hep-ph/0210197; Grimus, Lavoura, hep-ph/0305309]

Generalized $\mu\tau$ reflection

the generalized $\mu\tau$ reflection of neutrino

[Chen, Ding, Valle et al,
1512.01551, 1802.04275]

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} \rightarrow \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \Theta & i \sin \Theta \\ 0 & i \sin \Theta & \cos \Theta \end{pmatrix} \begin{pmatrix} \nu_e^c \\ \nu_\mu^c \\ \nu_\tau^c \end{pmatrix},$$

$\mu\tau$ reflection: $\Theta = \pi / 2$

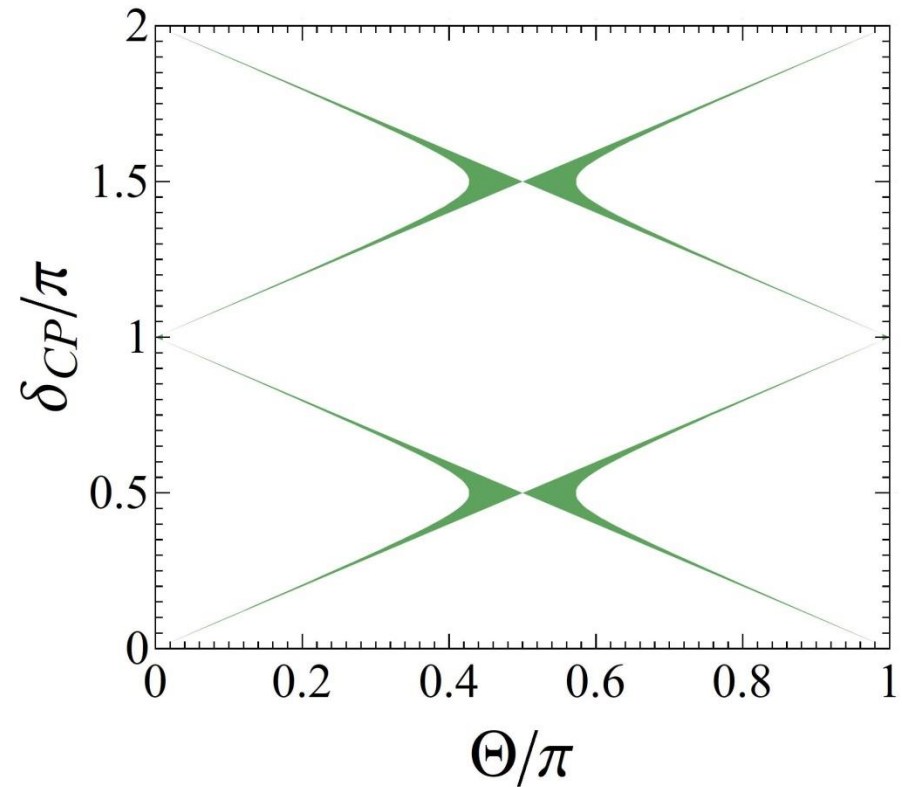
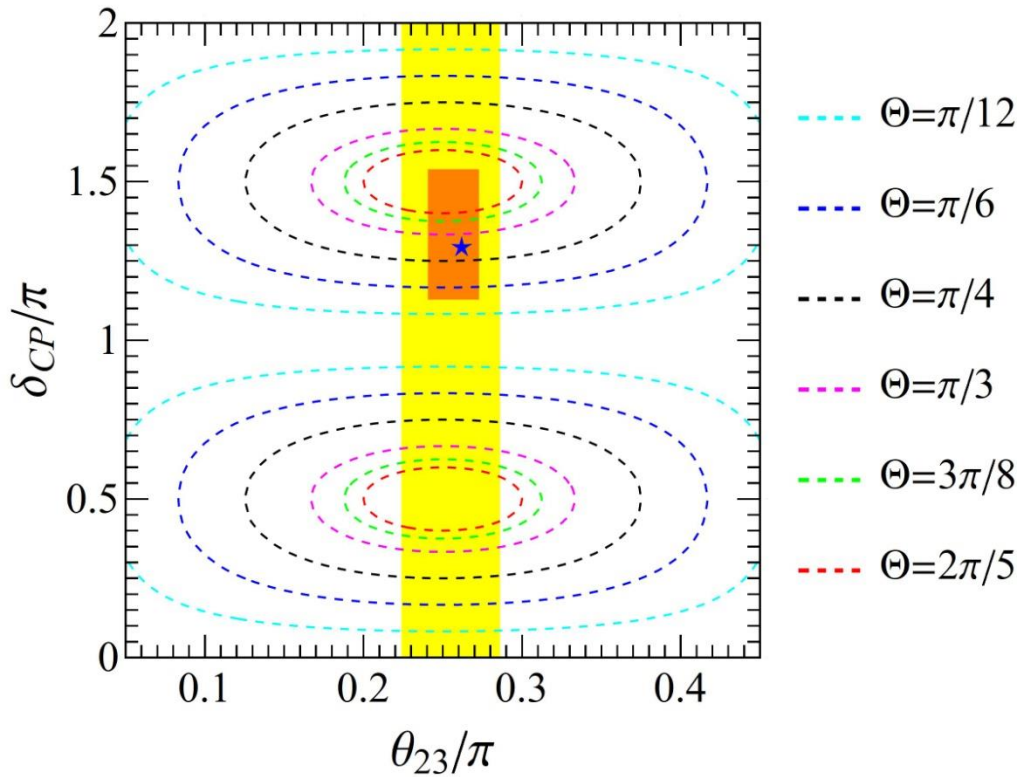
In the charged lepton diagonal basis, the lepton mixing matrix is

$$U = \Sigma_{\mu\tau} O_{3\times 3} \text{diag}(1, i^{k_1}, i^{k_2}), \quad \Sigma_{\mu\tau} \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \frac{\Theta}{2} & i \sin \frac{\Theta}{2} \\ 0 & i \sin \frac{\Theta}{2} & \cos \frac{\Theta}{2} \end{pmatrix}$$

Lepton mixing matrix is determined up to a real orthogonal matrix $O_{3\times 3}$

Predictions for lepton mixing parameters:

$$\sin^2 \delta_{CP} \sin^2 2\theta_{23} = \sin^2 \Theta, \quad \sin \alpha_{21} = \sin \alpha'_{31} = 0$$



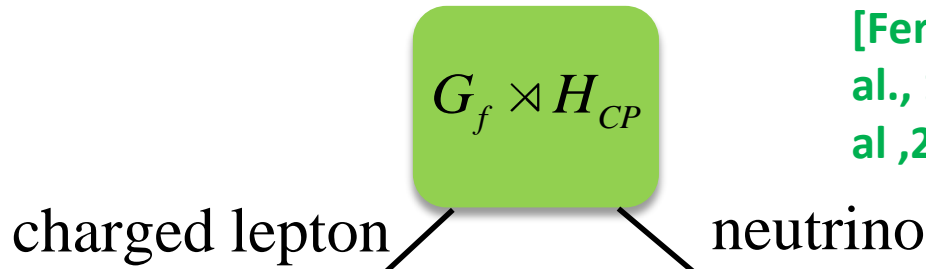
Numerical example:

$$\Theta = \frac{\pi}{3}, \quad \theta_1 = 0.274\pi, \quad \theta_2 = 0.0475\pi, \quad \theta_3 = 0.813\pi,$$

$$\sin^2 \theta_{12} = 0.307, \quad \sin^2 \theta_{13} = 0.0221, \quad \sin^2 \theta_{23} = 0.538, \quad \delta_{CP} = 1.335\pi$$

Semi-direct approach to lepton mixing

[Feruglio et al., 1211.5560; Ding et al., 1303.6180,...., Hagedorn et al., 2107.07537]



$$Z_n^{g_l}, n \geq 3$$

$$Z_2^{g_\nu} \times X_\nu$$

$$X_\nu = \Sigma_\nu \Sigma_\nu^T,$$

$$\Sigma_\nu^\dagger \rho_3(g_\nu) \Sigma_\nu = \pm \text{diag}(1, -1, -1)$$

$$U_l^\dagger \rho_3(g_l) U_l = \rho_3(g_l)_{\text{diag}}$$

$$U_l$$

$$U_\nu = \Sigma_\nu R_{23}(\theta) Q_\nu$$

$$U_l^\dagger m_l^\dagger m_l U_l \equiv \text{diag}(m_e^2, m_\mu^2, m_\tau^2)$$

$$U_\nu^T m_\nu U_\nu \equiv \text{diag}(m_1, m_2, m_3)$$

$$U_{PMNS} = U_l^\dagger \Sigma_\nu R_{23}(\theta) Q_\nu$$

➤ The mixing angles and CP violating phases are predicted in terms of a **single real** parameter $0 \leq \theta < \pi$. One column is fixed.

➤ quark mixing can not be explained in semi-direct approach.

Benchmark examples

For popular flavor symmetries A_4, S_4, A_5 : [Li et al, 1503.03711, Iura et al., 1503.04140, Ballett et al., 1503.07543]

$$\delta_{CP} = \pm \pi / 2, \theta_{23} = \pi / 4 \quad \text{or} \quad \delta_{CP} = 0, \pi, \theta_{23} \neq \pi / 4$$

➤ Flavor group $G_f = \Delta(96)$, residual symmetry $G_e = Z_3^{ac^3}, G_\nu = Z_2^{c^2}, X_\nu = c^2 d$

$$U_{PMNS} = \frac{1}{\sqrt{3}} \begin{pmatrix} \omega^2 e^{i\pi/4} \cos \theta - \omega \sin \theta & 1 & \omega^2 e^{i\pi/4} \sin \theta + \omega \cos \theta \\ e^{i\pi/4} \cos \theta - \sin \theta & 1 & e^{i\pi/4} \sin \theta + \cos \theta \\ \omega e^{i\pi/4} \cos \theta - \omega^2 \sin \theta & 1 & \omega e^{i\pi/4} \sin \theta + \omega^2 \cos \theta \end{pmatrix}, \quad \omega \equiv e^{2\pi i/3}$$

[Ding et al., 1409.8005; Hagedorn et al., 1408.7118]

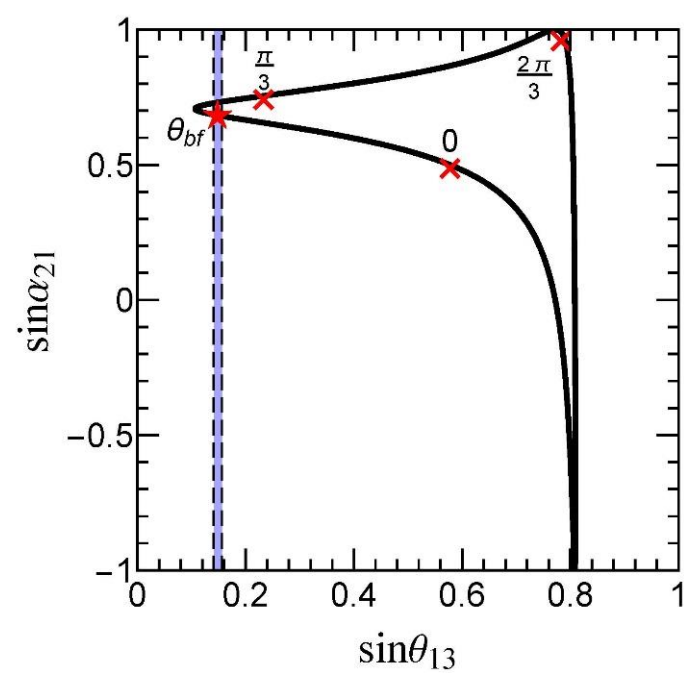
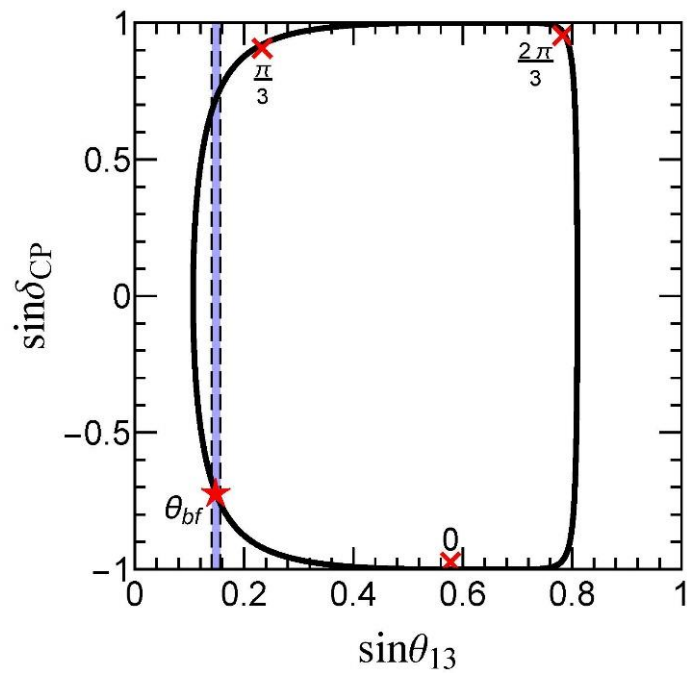
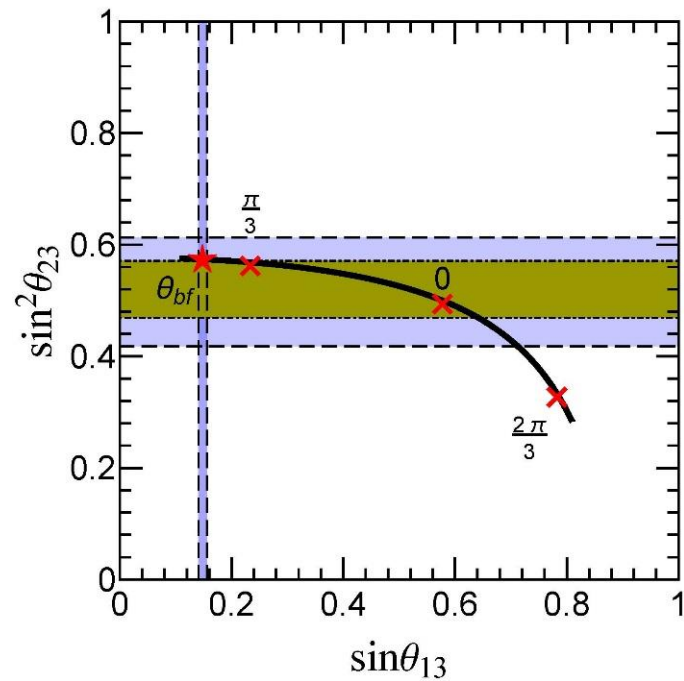
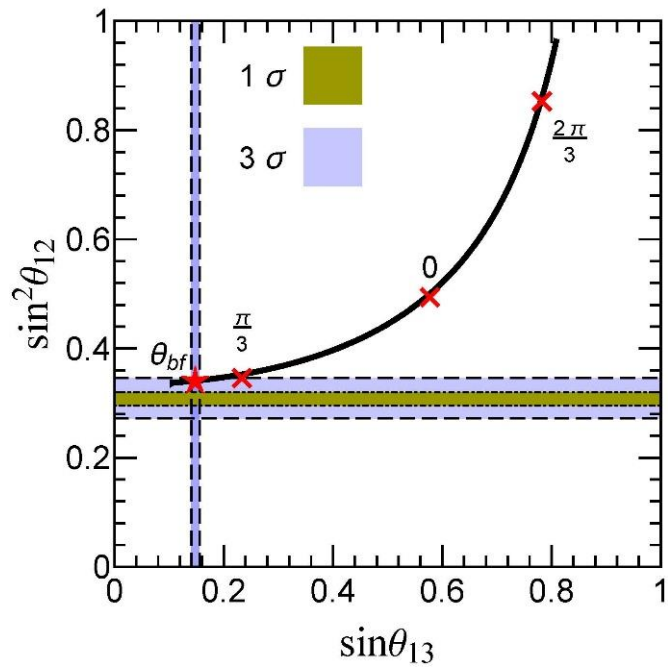
$$\Rightarrow \begin{cases} \sin^2 \theta_{13} = \frac{1}{3} - \frac{\sqrt{3}+1}{6\sqrt{2}} \sin 2\theta, & \sin^2 \theta_{12} = \frac{2\sqrt{2}}{4\sqrt{2} + (\sqrt{3}+1)\sin 2\theta}, & J_{CP} = -\frac{\cos 2\theta}{6\sqrt{3}}, \\ \sin^2 \theta_{23} = \frac{1}{2} + \frac{(3-\sqrt{3})\sin 2\theta}{8\sqrt{2} + 2(\sqrt{3}+1)\sin 2\theta}, & I_1 = \frac{1}{18} \left(1 + (\sqrt{3}-1)\sin^2 \theta + \sqrt{2} \sin 2\theta \right), & I_2 = \frac{\cos 2\theta}{18} \end{cases}$$

Numerical results:

non-regular values of CP phases

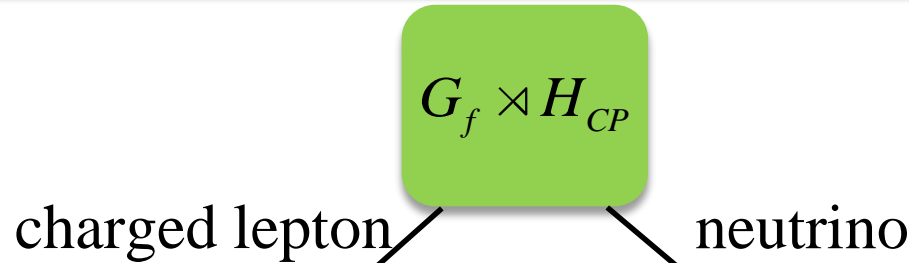
$$\theta_{bf} = 0.209\pi, \quad \chi_{min}^2 = 7.960, \quad \sin^2 \theta_{12} = 0.341, \quad \sin^2 \theta_{13} = 0.0220,$$

$$\sin^2 \theta_{23} = 0.574, \quad \sin \delta_{CP} = -0.722, \quad \sin \alpha_{21} = 0.683, \quad \sin \alpha_{31} = -0.091.$$



Universal flavor symmetry for quark and lepton mixing

[Lu, Ding, 1610.05682,
1806.02301; Li, Lu, Ding,
1706.04576]



$$Z_2^{g_l} \times X_l$$

$$Z_2^{g_\nu} \times X_\nu$$

$$X_l = \Sigma_l \Sigma_l^T,$$

$$X_\nu = \Sigma_\nu \Sigma_\nu^T,$$

$$\Sigma_l^\dagger \rho_3(g_l) \Sigma_l = \pm \text{diag}(1, -1, -1)$$

$$\Sigma_\nu^\dagger \rho_3(g_\nu) \Sigma_\nu = \pm \text{diag}(1, -1, -1)$$

$$U_l = \Sigma_l R_{23}(\theta_l) P_l \quad U_\nu = \Sigma_\nu R_{23}(\theta_\nu) P_\nu Q_\nu$$

$$U_l^\dagger m_l^\dagger m_l U_l \equiv \text{diag}(m_e^2, m_\mu^2, m_\tau^2)$$

$$U_\nu^T m_\nu U_\nu \equiv \text{diag}(m_1, m_2, m_3)$$

$$U_{PMNS}(\theta_l, \theta_\nu) = R_{23}^T(\theta_l) \Sigma_l^\dagger \Sigma_\nu R_{23}(\theta_\nu) P_\nu Q_\nu$$

- All mixing angles and CP phases are expressed in terms of **two free angles** $\theta_{l,\nu} \in [0, \pi)$
- This scheme can be extended to quark sector, and the CKM mixing matrix is of similar form

Quark and lepton mixing from Dihedral group D_n and CP

Quark sector:

[Lu, Ding, 1901.07414]

Assignment: 2+1

$$\begin{pmatrix} \begin{pmatrix} u_L \\ d_L \\ c_L \\ s_L \end{pmatrix} \end{pmatrix} \sim \mathbf{2}_1, \quad \begin{pmatrix} t_L \\ b_L \end{pmatrix} \sim \mathbf{1}_1$$

Residual symmetry: $Z_2^{g_u} = Z_2^{SR^x}$, $X_u = SR^{n/2}$, $Z_2^{g_d} = Z_2^{SR^y}$, $X_d = S$, $x, y = 0, \dots, n-1$

➤ The CKM matrix is determined to be $c_{u,d} \equiv \cos \theta_{u,d}$, $s_{u,d} \equiv \sin \theta_{u,d}$

$$V_{CKM} = \begin{pmatrix} -c_d \sin \varphi_1 & \boxed{\cos \varphi_1} & s_d \sin \varphi_1 \\ c_u c_d \cos \varphi_1 + i s_u s_d & c_u \sin \varphi_1 & -c_u s_d \cos \varphi_1 + i c_d s_u \\ -c_d s_u \cos \varphi_1 + i c_u s_d & -s_u \sin \varphi_1 & s_u s_d \cos \varphi_1 + i c_u c_d \end{pmatrix}$$

with $\varphi_1 = \frac{y-x}{n} \pi$ ← fixed by residual symmetry

➤ Sum rules

$$\begin{cases} \cos^2 \theta_{13}^q \sin^2 \theta_{12}^q = \cos^2 \varphi_1, \\ J_{CP}^q \approx \frac{1}{2} \sin 2\varphi_1 \sin \varphi_2 \sin \theta_{13}^q \sin \theta_{23}^q \end{cases}$$

➤ Benchmark value: $\varphi_1 = 3\pi/7$ which can be achieved in D_{14} group

	θ_1^{bf}/π	θ_2^{bf}/π	$\sin \theta_{12}^q$	$\sin \theta_{23}^q$	$\sin \theta_{13}^q$	J_{CP}^q
Our	0.01326	0.00117	0.22252	0.04166	0.00357	3.224×10^{-5}
Data	—	—	0.22500 ± 0.00100	0.04200 ± 0.00059	0.003675 ± 0.000095	$(3.120 \pm 0.090) \times 10^{-5}$

The measured quark mixing angles and CP violation phases can be accommodated.

➤ Lepton sector : $\varphi_1 = 2\pi/7$

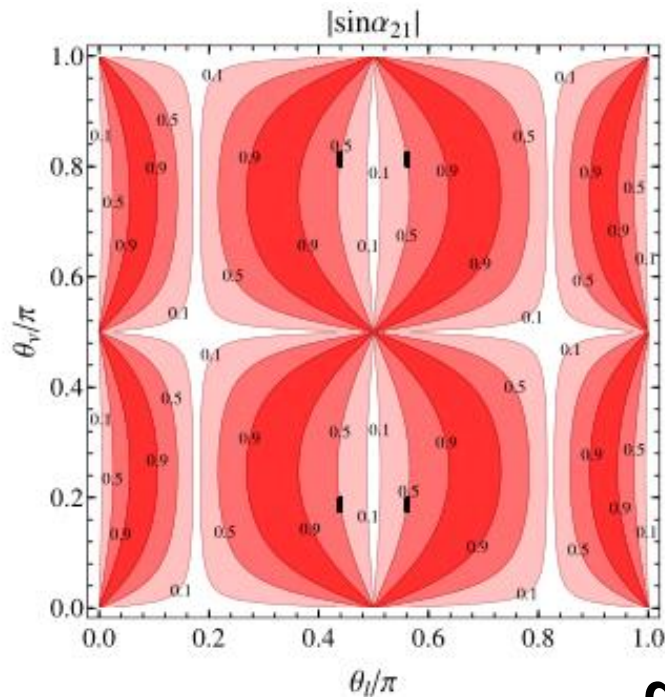
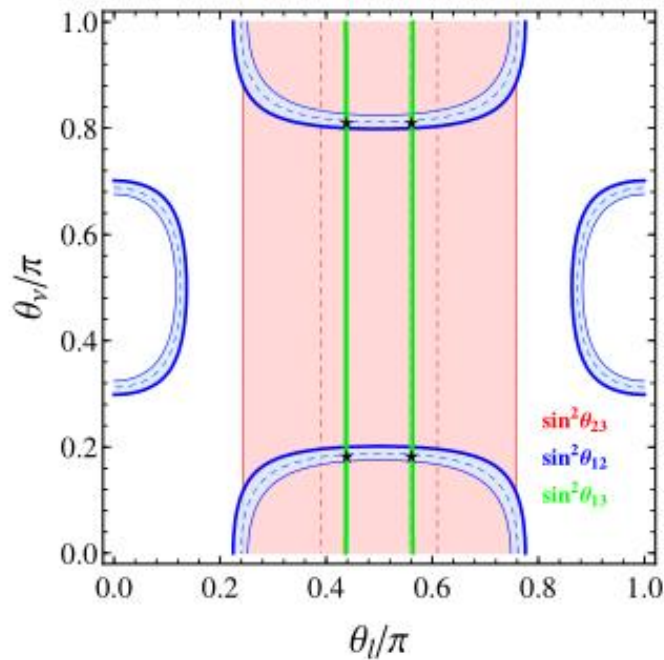
$$U_{PMNS} = R_{12}(\theta_l) \begin{pmatrix} 0 & \cos \frac{2\pi}{7} & \sin \frac{2\pi}{7} \\ i & 0 & 0 \\ 0 & -\sin \frac{2\pi}{7} & \cos \frac{2\pi}{7} \end{pmatrix} R_{13}(\theta_\nu) Q_\nu \longrightarrow \cos^2 \theta_{23} \cos^2 \theta_{13} = \cos^2 \frac{2\pi}{7}$$

Numerical benchmark:

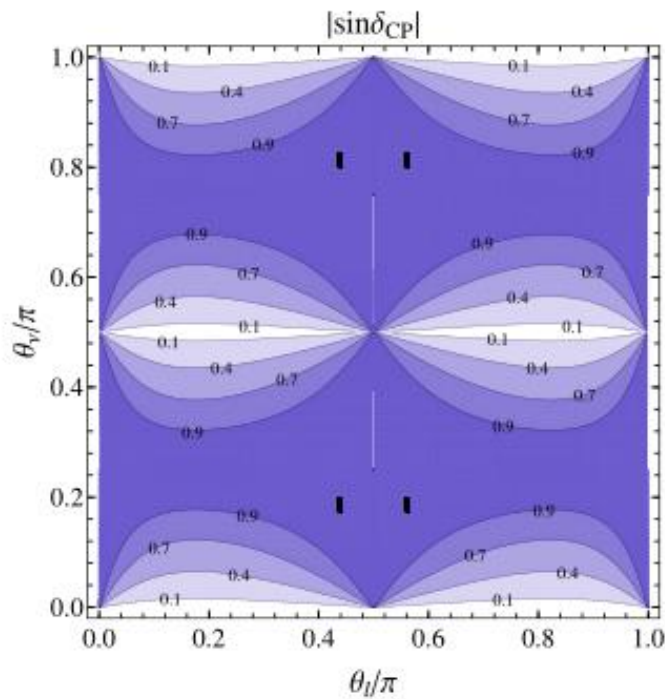
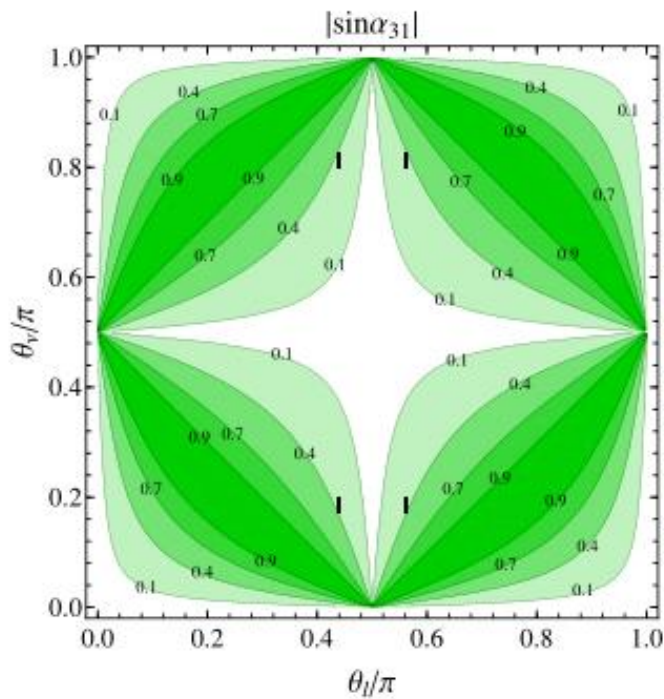
$$\theta_l = 0.562\pi, \quad \theta_\nu = 0.186\pi, \quad \chi_{\min}^2 = 1.824,$$

$$\sin^2 \theta_{13} = 0.0226, \quad \sin^2 \theta_{12} = 0.310, \quad \sin^2 \theta_{23} = 0.602,$$

$$\delta_{CP} = 1.467\pi, \quad |\sin \alpha_{21}| = 0.503, \quad |\sin \alpha_{31}| = 0.358$$

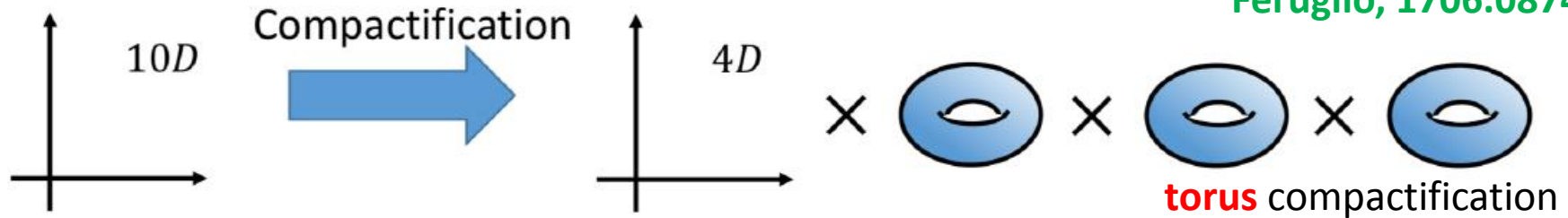


quite predictive!



Recent progress: modular symmetry

[Ferrara et al, 1989;
Feruglio, 1706.08749]



integrating over 6D: $S = \int d^4x d^6y \mathcal{L}_{10D} \Rightarrow \int d^4x \mathcal{L}_{\text{eff}}(\varphi, \tau_i)$

The shape of the torus is parametrized by the modulus τ , and it is invariant under the modular transformation

$$\tau \rightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d} \Leftrightarrow \begin{pmatrix} a & b \\ c & d \end{pmatrix} \quad \begin{matrix} ad - bc = 1 \\ a, b, c, d \text{ integers} \end{matrix} \quad \xrightarrow{\text{Modular group}} \quad \Gamma \cong SL(2, \mathbb{Z})$$

➤ Generators S and T : $S^4 = (ST)^3 = 1, S^2T = TS^2$

$$S = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad T = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad \tau \xrightarrow{S} -\frac{1}{\tau}, \quad \tau \xrightarrow{T} \tau + 1$$

➤ Finite modular groups: the quotient over the principal congruence subgroups $\Gamma(N)$

Γ_2	Γ_3	Γ_4	Γ_5
S_3	A_4	S_4	A_5
Γ'_2	Γ'_3	Γ'_4	Γ'_5
S_3	T'	S'_4	A'_5

$\Gamma_N \equiv SL(2, \mathbb{Z}) / \pm\Gamma(N), \quad \Gamma'_N \equiv SL(2, \mathbb{Z}) / \Gamma(N)$ Γ'_N is the “double covering” of Γ_N

Modular invariance as flavor symmetry

- Modular transformation of chiral superfield in MSSM

$$\gamma\tau = \frac{a\tau + b}{c\tau + d}, \quad \chi^{(I)} \rightarrow (c\tau + d)^{-k_I} \rho^{(I)}(\gamma) \chi^{(I)}$$

Modular weight

unitary representation of finite modular groups Γ'_N or Γ_N

- Modular invariant superpotential

$$W = \sum_n Y_{I_1 \dots I_n}(\tau) \chi^{(I_1)} \dots \chi^{(I_n)}$$

Yukawa couplings only depend on τ

Modular invariance requires a holomorphic $Y_{I_1 I_2 \dots I_n}(\tau)$ satisfying

$$Y_{I_1 \dots I_n}(\gamma\tau) = (c\tau + d)^{k_Y(n)} \rho^{(Y)}(\gamma) Y_{I_1 \dots I_n}(\tau)$$

$Y_{I_1 I_2 \dots I_n}(\tau)$ are modular forms of level N and weight k_Y

obeying

$$\begin{cases} k_Y(n) + k_{I_1} + \dots + k_{I_n} = 0 \\ \rho^{(Y)} \otimes \rho^{(I_1)} \otimes \dots \otimes \rho^{(I_n)} \supset 1 \end{cases}$$

- Lowest weight **2** modular forms of level 3 transforms as a **triplet of A_4**

$$Y_3^{(2)} = \begin{pmatrix} \varepsilon^2(\tau) \\ \sqrt{2} \mathcal{G}(\tau) \varepsilon(\tau) \\ -\mathcal{G}^2(\tau) \end{pmatrix} = \begin{pmatrix} 1 + 12q + 36q^2 + \dots \\ -6q^{1/3} (1 + 7q + 8q^2 + \dots) \\ -18q^{2/3} (1 + 2q + 5q^2 + \dots) \end{pmatrix} \quad \begin{aligned} \mathcal{G}(\tau) &= 3\sqrt{2} \frac{\eta^3(3\tau)}{\eta(\tau)}, \quad q = e^{2\pi i\tau}, \\ \varepsilon(\tau) &= -\frac{3\eta^3(3\tau) + \eta^3(\tau/3)}{\eta(\tau)} \end{aligned}$$

tensor product generate high weight modular forms

Example: a minimal model based on $\Gamma_3 = A_4$

➤ Field content


Including gCP symmetry to constrain all couplings real

	L	(e^c, μ^c, τ^c)	N^c	H_u	H_d
$SU(2)_L \times U(1)_Y$	$(2, -1/2)$	$(1, 1)$	$(1, 0)$	$(2, 1/2)$	$(2, -1/2)$
A_4	3	$(\mathbf{1}, \mathbf{1}, \mathbf{1}')$	3	1	1
k_I	-1	$(3, 5, 9)$	1	0	0

Charged lepton mass terms

[Yao, Lu, Ding, 2012.13390]


$$W_e = \alpha e^c (LY_3^{(2)})_1 H_d + \beta \mu^c (LY_3^{(4)})_1 H_d + \gamma \tau^c (LY_{3I}^{(8)})_{1'} H_d + \gamma' \tau^c (LY_{3II}^{(8)})_{1''} H_d$$



$$M_e = \begin{pmatrix} \alpha Y_{3,1}^{(2)} & \alpha Y_{3,3}^{(2)} & \alpha Y_{3,2}^{(2)} \\ \beta Y_{3,1}^{(4)} & \beta Y_{3,3}^{(4)} & \beta Y_{3,2}^{(4)} \\ \gamma Y_{3I,3}^{(8)} + \gamma' Y_{3II,3}^{(8)} & \gamma Y_{3I,2}^{(8)} + \gamma' Y_{3II,2}^{(8)} & \gamma Y_{3I,1}^{(8)} + \gamma' Y_{3II,1}^{(8)} \end{pmatrix} \nu_d$$

NO flavons

Neutrino mass terms: $W_\nu = g(N^c L)_1 H_u + \Lambda(N^c N^c Y_3^{(2)})_1$



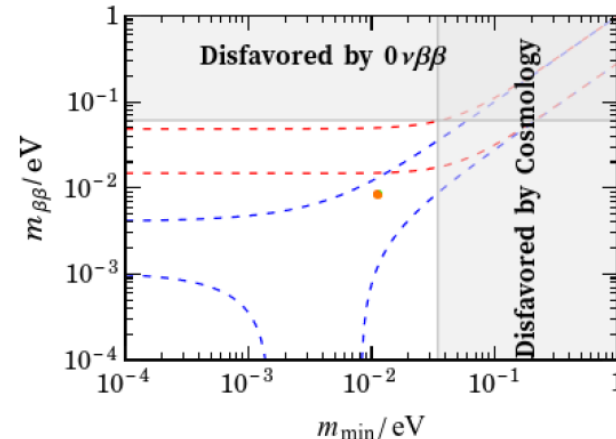
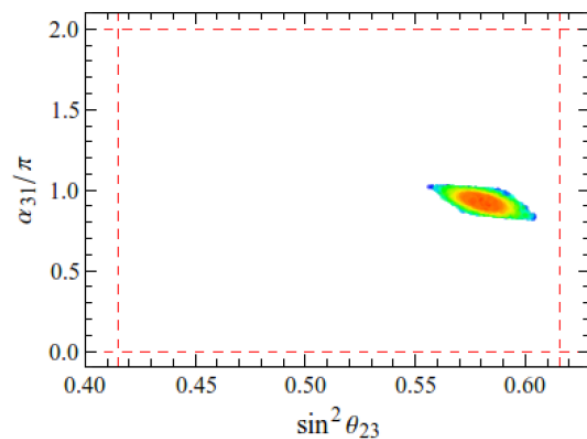
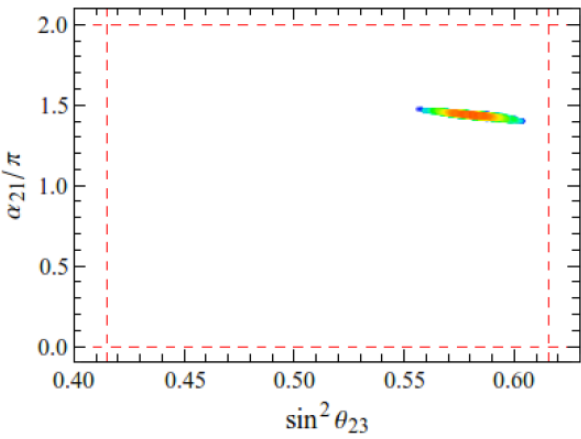
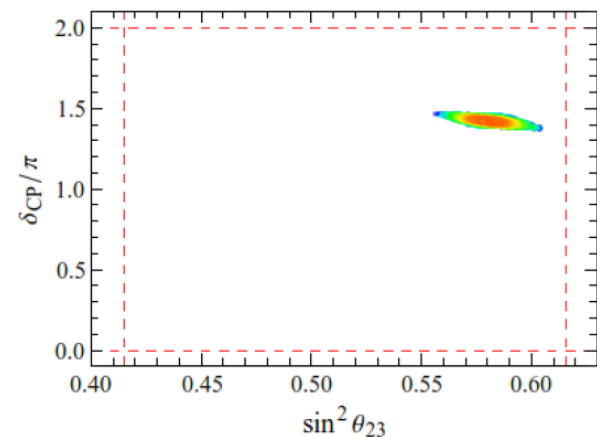
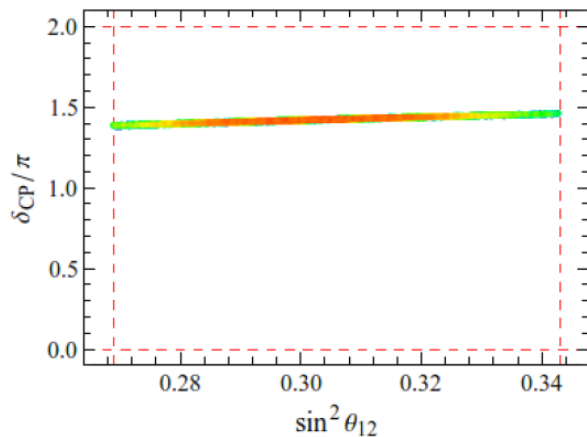
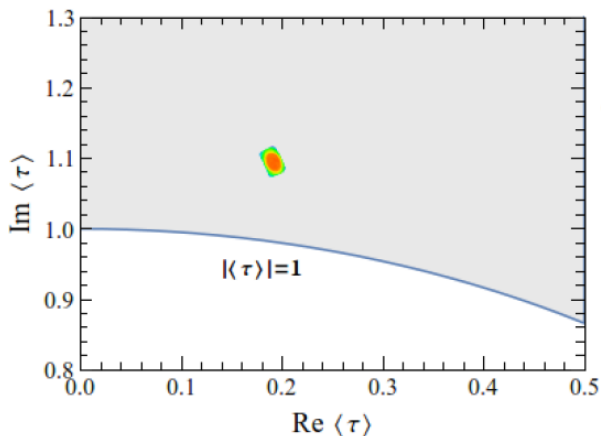
$$M_D = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0_3 \end{pmatrix} g \nu_u, \quad M_N = \begin{pmatrix} 2Y_{3,1}^{(2)} & -Y_{3,3}^{(2)} & -Y_{3,2}^{(2)} \\ -Y_{3,3}^{(2)} & 2Y_{3,2}^{(2)} & -Y_{3,1}^{(2)} \\ -Y_{3,2}^{(2)} & -Y_{3,1}^{(2)} & 2Y_{3,3}^{(2)} \end{pmatrix} \Lambda$$

The light neutrino mass matrix only depends on τ up to an overall scale

7 real input parameters describe 12 observables: $m_{e,\mu,\tau}, m_{1,2,3}, \theta_{12}, \theta_{13}, \theta_{23}, \delta, \alpha_{21}, \alpha_{31}$

✓ The complex modulus τ is the only source of both CP violation and flavor symmetry breaking

✓ The best fit value $\tau \simeq 0.19 + 1.09i$ is close to the self-dual point $\tau = i$



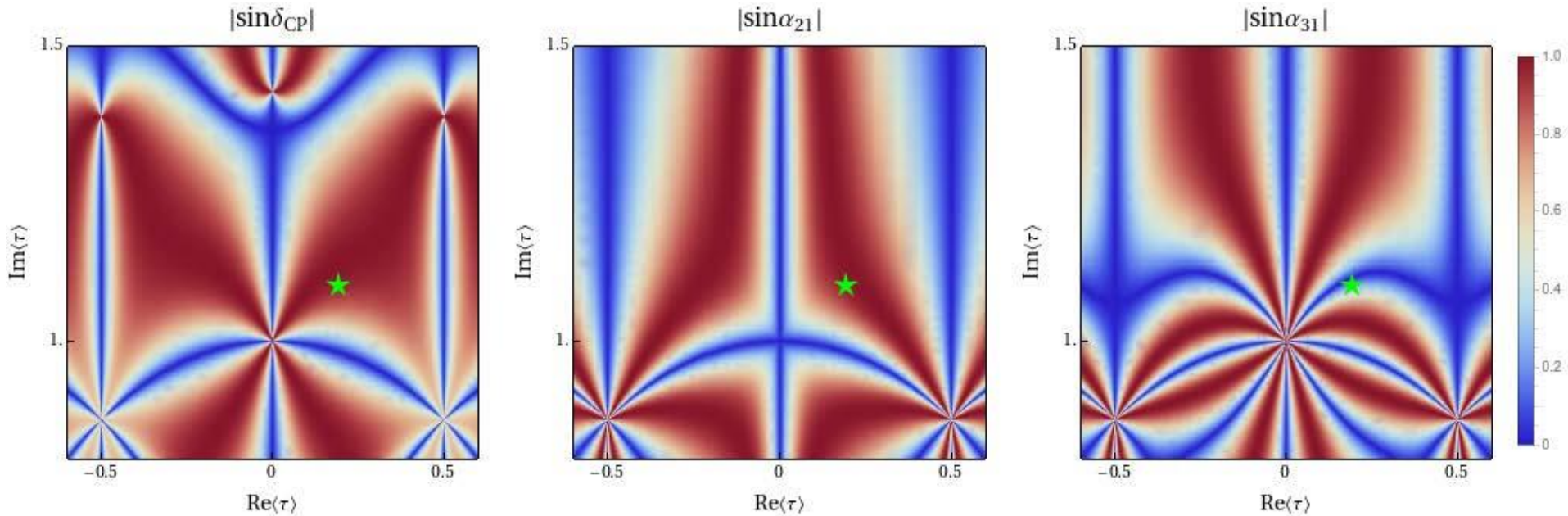
$\sin^2 \theta_{23} \sim 0.58$, $\delta \sim 1.42\pi$, $\alpha_{21} \sim 1.43\pi$, $\alpha_{31} \sim 0.92\pi$
 $\sum_i m_i \sim 76.18 \text{ meV}$, $m_{\beta\beta} \sim 8.57 \text{ meV}$, $\chi^2 \sim 4.87$

χ^2



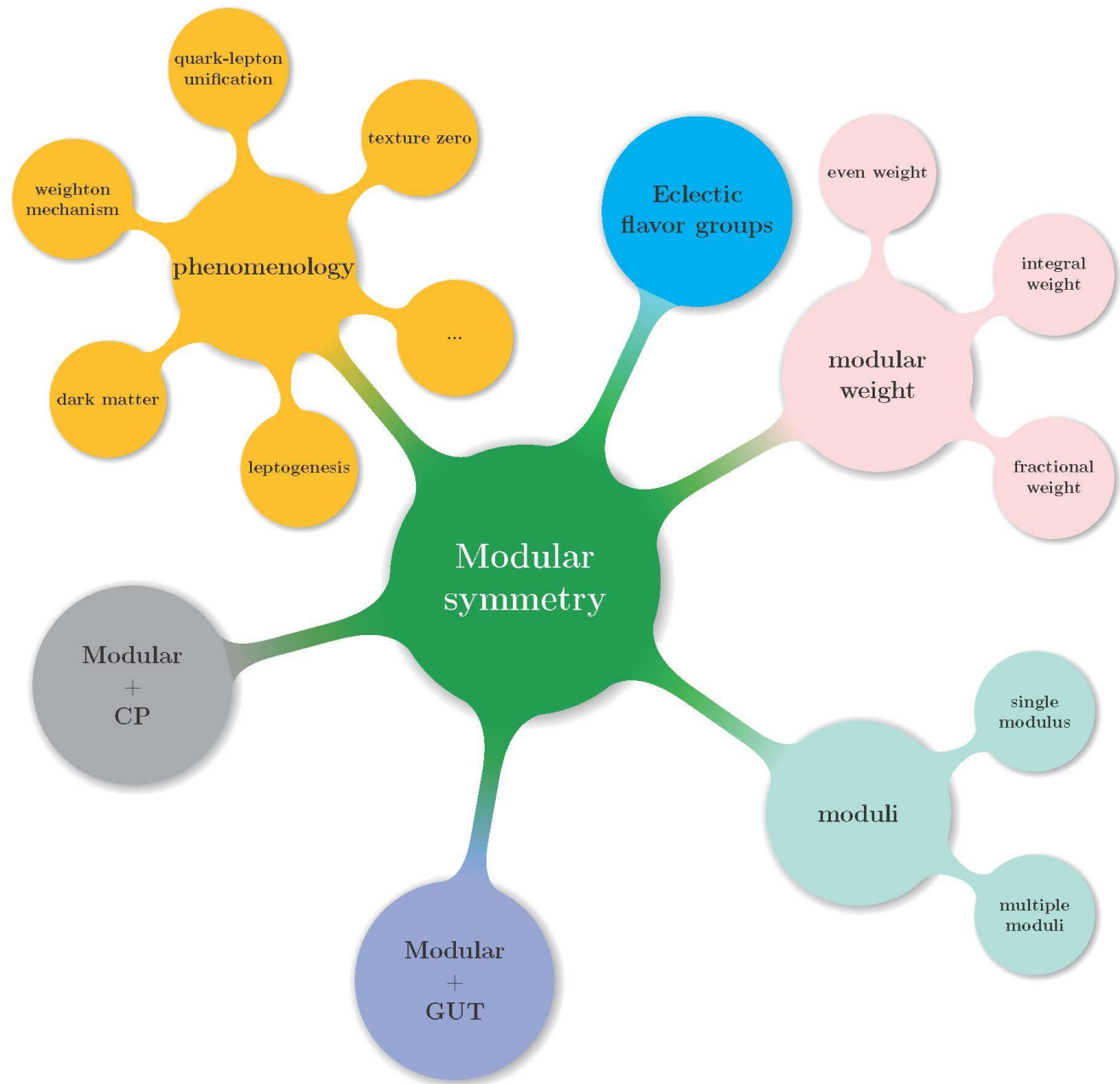
[Yao, Lu, Ding, 2012.13390]

Numerical results for CP violation phases:



CP is conserved for modulus τ imaginary or at the border of the fundamental domain, a small deviation can generate large CP violation.

Further development of modular symmetry



Baur, Chen, Criado, Ding, Feruglio, King, Kobayashi, Li, Liu, Lu, Nilles, Nomura, Novichkov, Okada, Penedo, Petcov, Ramos-Sanchez, Ratz, Y. Shimizu, Tanimoto, Tatsuishi, Titov, Trautner, Uemura, Varzielas, Vaudrevange, Wang, Yao, Yu, Zhang, Zhou...

Summary

- flavour symmetries are a useful tool to understand the origin flavor mixing and CP violation, but no compelling and unique picture have emerged so far.
- More precise neutrino data calls for convincing model of neutrino masses and mixings, with testable and confirmed predictions. Modular symmetry is a new promising approach to the flavor puzzles. Modular symmetry is still at the early stage of its development, many aspects still need to be understood.
- Future experimental data on θ_{23} , δ_{CP} , the effective Majorana mass $m_{\beta\beta}$ can exclude many models, will provide important hints for the underlying principle.