

Flavor anomalies and their BSM interpretation

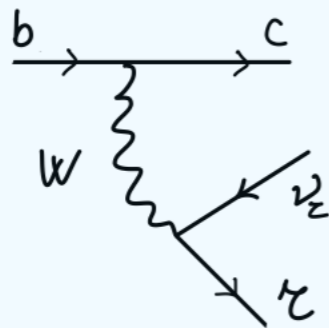
Claudia Cornella

Panic, 5-10.09.2021

The flavor anomalies

$$b \rightarrow cl\nu$$

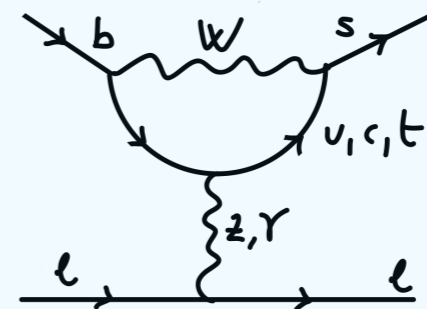
~15% deviation from $\tau/\mu, e$ universality



$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \nu)}{\mathcal{B}(B \rightarrow D^{(*)} \ell \nu)}$$

$$b \rightarrow sll$$

~15% deviation from μ/e universality



$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)} \mu^+ \mu^-)}{\mathcal{B}(B \rightarrow K^{(*)} e^+ e^-)}$$

What do we learn from B anomalies?

If they are NP signals, where else should we see something?

Data in $b \rightarrow sll$

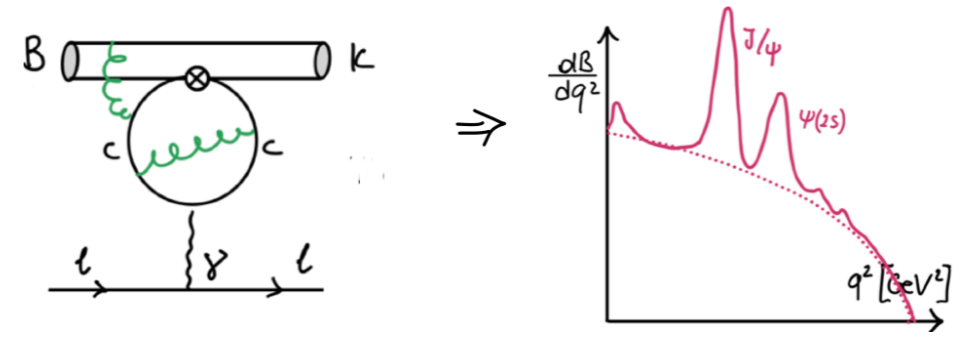
Several discrepancies from the SM in $b \rightarrow sll$:

- P'_5 ($B \rightarrow K^{(*)}\mu\mu$ angular distribution)
- deficit in $\mathcal{B}(B \rightarrow X_s\mu\mu)$ $X_s = K, K^*, \phi$

- μ/e LFUV in $B \rightarrow K^*ll, B \rightarrow Kll$
- deficit in $\mathcal{B}(B_s \rightarrow \mu\mu)$

“clean”

“not-so-clean”



$$R_{K^{(*)}}^{[1.1,6] \text{ GeV}^2} = 1.00 \pm 0.01$$

[Bordone et al, [1605.07633](#)]

$$\mathcal{B}(B_s \rightarrow \mu\mu)_{\text{SM}} = (3.66 \pm 0.14) \times 10^{-9}$$

[Beneke et al., [1908.07011](#)]

Recent LHCb updates:

2.3 σ deficit in $B_s \rightarrow \mu\mu$ combining ATLAS, CMS, LHCb

3.1 σ in $R_K^{[1.1,6]}$ (evidence of LFUV) [[LHCb, 2103.11769](#)]

EFT for $b \rightarrow sll$

$b \rightarrow sll$ data have a simple EFT solution:

Fit to **clean observables** only:

left-handed NP hypothesis ($\Delta C_9^\mu = -\Delta C_{10}^\mu$)
preferred by 4.6σ over SM

$\sim 4.8 \sigma$ adding all $b \rightarrow sll$ data
(treating ΔC_9^U as nuisance parameter)

$\gg 5\sigma$ with current best estimate of $\bar{c}c$ loop

[Altmannshofer, Stangl, [2103.13370](#); Algueró et al, [2104.08921](#); Hurth et al. [2104.10058...](#)]

These are **local significances** in NP space.

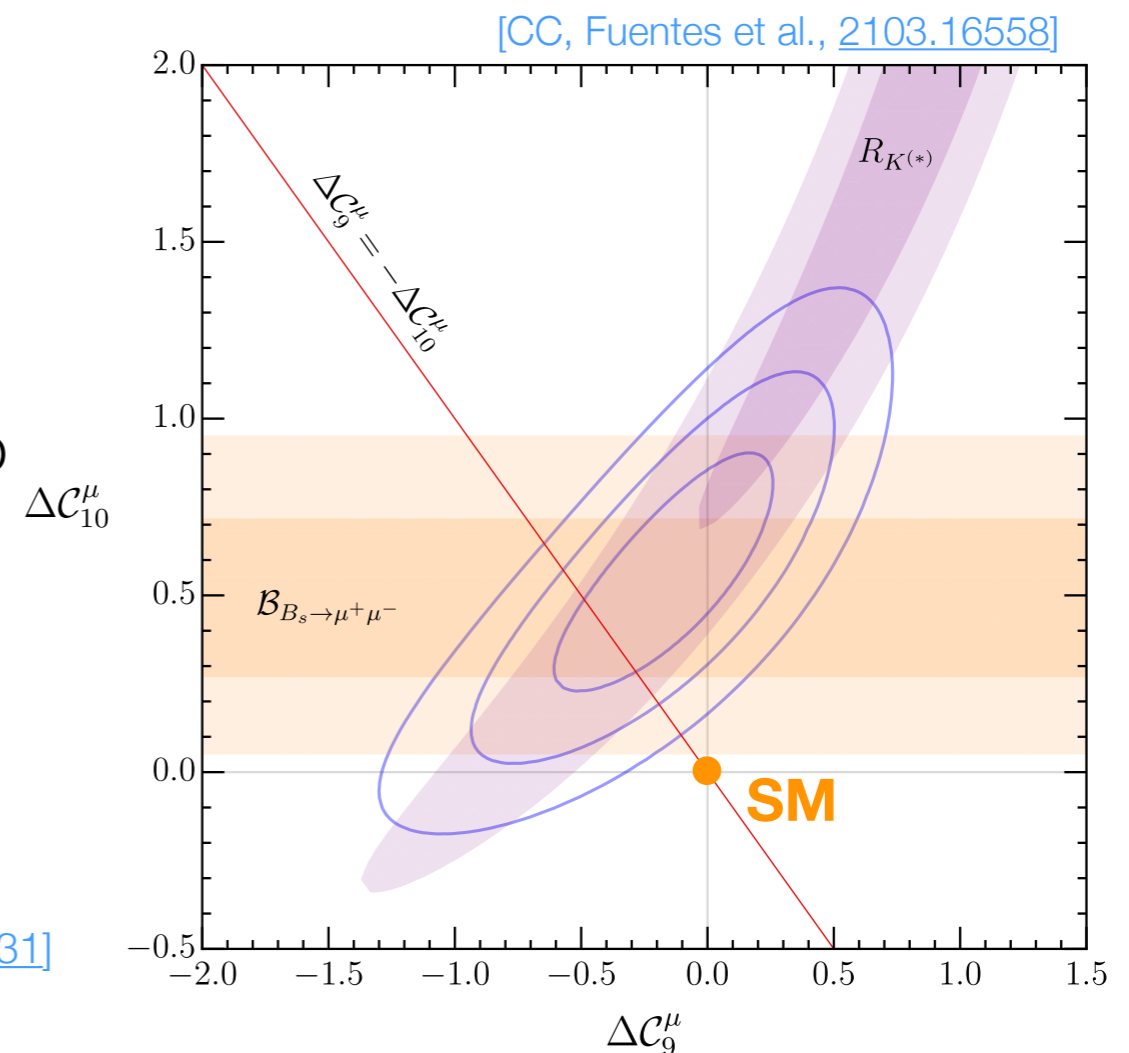
Global? LEE non negligible, still 3.9σ

[Isidori, Lancierini, Owen, Serra, [2104.05631](#)]

$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \frac{\alpha}{4\pi} \sum_i C_i \mathcal{O}_i$$

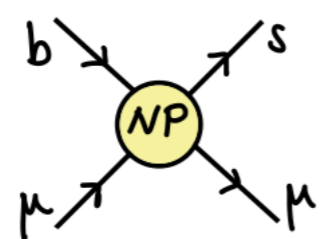
$$\mathcal{O}_9^\mu = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \mu)$$

$$\mathcal{O}_{10}^\mu = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \gamma_5 \mu)$$



Theory lessons from $b \rightarrow sll$

It's a **very weak** interaction:



A Feynman diagram showing a central yellow circle labeled 'NP'. An incoming line from the top-left is labeled 'b' with an arrow pointing towards the circle. An outgoing line to the top-right is labeled 's' with an arrow pointing away from the circle. An incoming line from the bottom-left is labeled with the Greek letter mu (μ) with an arrow pointing towards the circle. An outgoing line to the bottom-right is labeled with the Greek letter mu (μ) with an arrow pointing away from the circle.

$$\sim 3 \times 10^{-5} G_F \quad \Rightarrow \quad \frac{g_{\text{NP}}^2}{M_{\text{NP}}^2} \sim \frac{1}{(40 \text{ TeV})^2}$$

Direct production could be out of reach, but good chances for **indirect** discovery:

- build up evidence in $b \rightarrow sll$ @LHCb and Belle II
- possibly test LFUV in $b \rightarrow dll$

$$R_\pi[q_{\min}^2, q_{\max}^2] = \frac{\int_{q_{\min}^2}^{q_{\max}^2} dq^2 \frac{d\mathcal{B}}{dq^2} (B^+ \rightarrow \pi^+ \mu^+ \mu^-)}{\int_{q_{\min}^2}^{q_{\max}^2} dq^2 \frac{d\mathcal{B}}{dq^2} (B^+ \rightarrow \pi^+ e^+ e^-)} \stackrel{U(2)}{\approx} R_K^{(*)}$$

clean: pollution from long distance effects < 10% in large q^2 regions

[Bordone, CC, König, Isidori, [2101.11626](#)]

Which models? Z' , leptoquarks

Data in $b \rightarrow cl\nu$

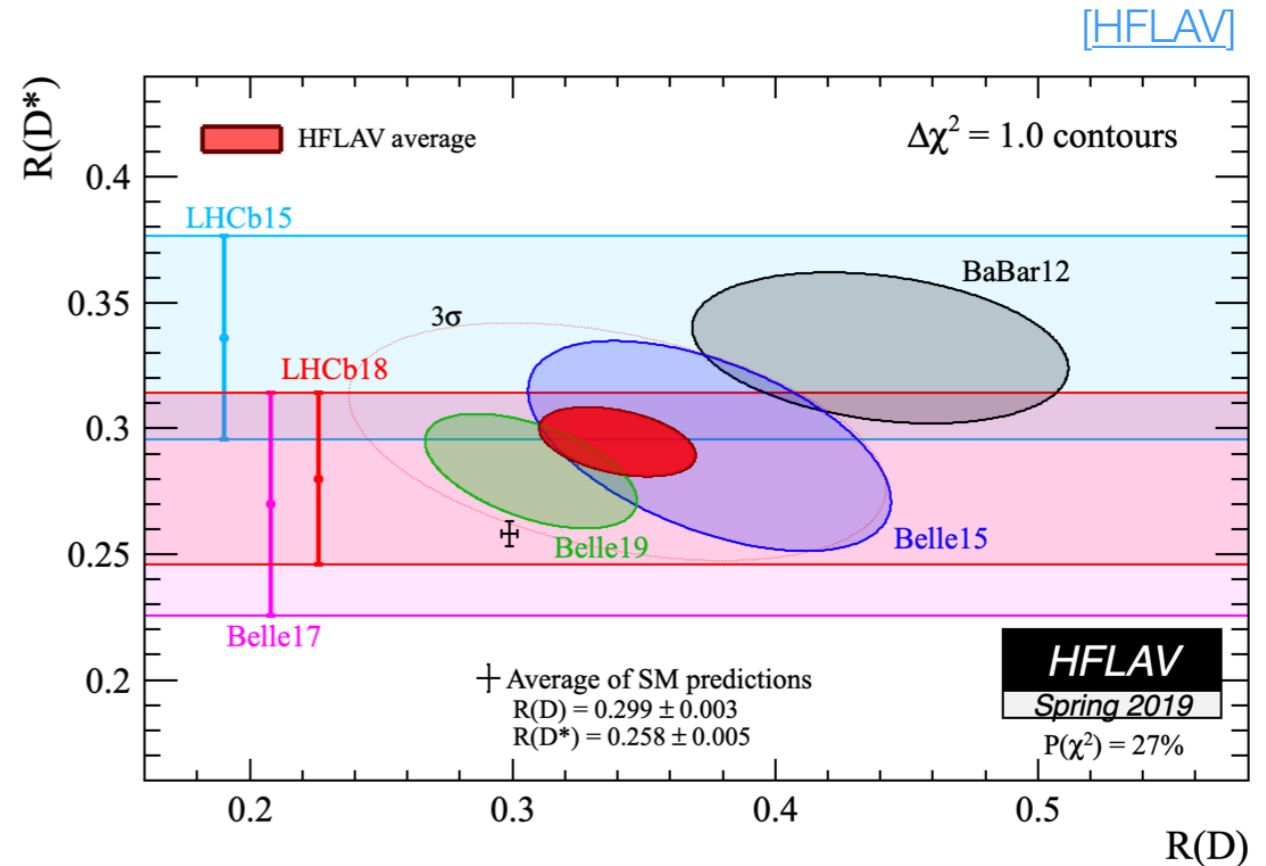
$$R_{D^{(*)}} \equiv \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)}\ell\bar{\nu})}$$

$\sim 10\%$ enhancement of $R_{D^{(*)}}$,
excess seen in **tau** mode.

Also **clean** SM prediction.

[new: $B \rightarrow D^*$ form factors from lattice [\[Bazavov et al., 2105.14019\]](#)]

3.1 σ tension combining BaBar, Belle & LHCb results for R_D, R_{D^*}

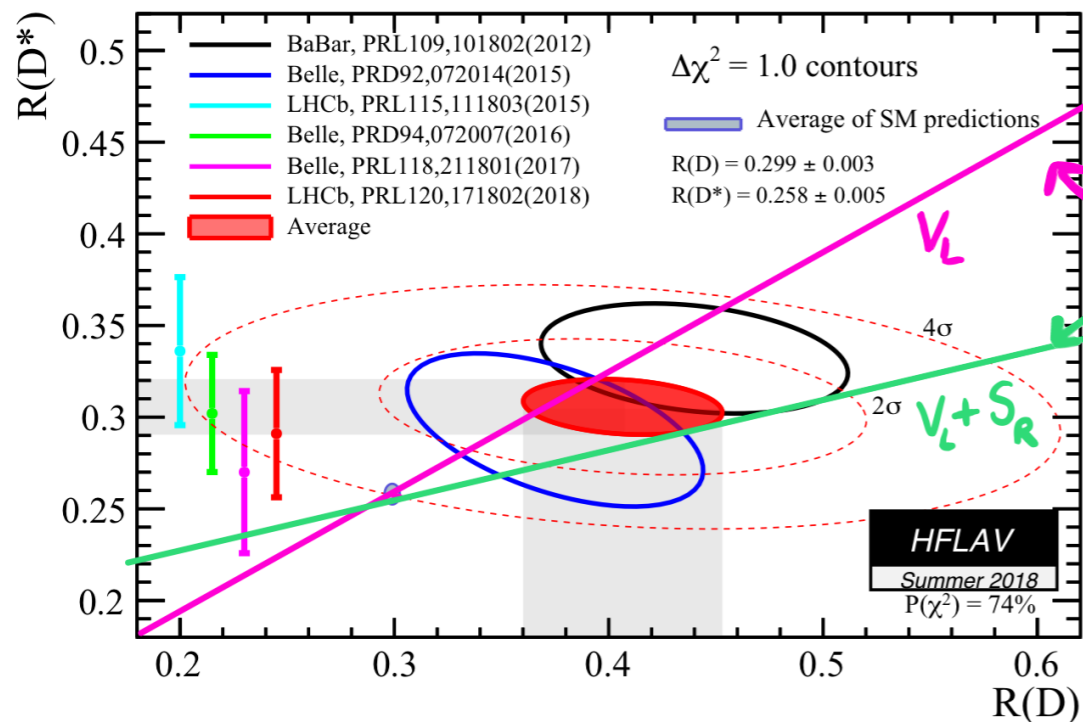


EFT for $b \rightarrow c\tau\nu$

$$\mathcal{L}_{\text{eff}} = -2\sqrt{2}G_F V_{cb} \left[(1+g_{V_L})(\bar{c}_L\gamma^\mu b_L)(\bar{\tau}_L\gamma_\mu\nu_L) + g_{V_R}(\bar{c}_R\gamma^\mu b_R)(\bar{\tau}_L\gamma_\mu\nu_L) + g_{S_R}(\bar{c}_L b_R)(\bar{\tau}_R\nu_L) \right. \\ \left. + g_{S_L}(\bar{c}_R b_L)(\bar{\tau}_R\nu_L) + g_T(\bar{c}_R\sigma^{\mu\nu} b_L)(\bar{\tau}_R\sigma_{\mu\nu}\nu_L) \right]$$

Several options: g_{V_L} [Fermi-like interaction], $g_{V_L} + g_{S_R}$, $g_{S_L} = \pm 4g_T$

Any update in $b \rightarrow c\tau\nu$ helps pin down the Lorentz structure:



- NP in LFU ratios:

$$V_L \text{ only} \quad \Delta R_D = \Delta R_D^* \quad (= \Delta R_{J/\Psi} = \Delta R_{\Lambda_c})$$

$$S, T \neq 0 \quad \Delta R_D \neq \Delta R_D^*$$

- Angular obs.: largely insensitive to V_L , powerful probe of S, T

Theory lessons from $b \rightarrow c\tau\nu$

It's a **large effect** (competes with *tree level* SM amplitude)

$$\sim 10^{-2} G_F \Rightarrow \frac{g_{\text{NP}}^2}{M_{\text{NP}}^2} \sim \frac{1}{(2 \text{ TeV})^2}$$

Constraints from low- and high-energy (LFU in τ decays, $pp \rightarrow \tau\tau$ tails) make model building for $R_{D^{(*)}}$ challenging

- charged Higgs excluded by τ_{B_c} ,
- W' in tension with $pp \rightarrow \tau\tau$

\Rightarrow **Leptoquarks** (scalar/vector) are the favored candidates.

Model	$R_{K^{(*)}}$	$R_{D^{(*)}}$	$R_{K^{(*)}} \& R_{D^{(*)}}$
$S_3 \ (\bar{\mathbf{3}}, \mathbf{3}, 1/3)$	✓	✗	✗
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[Sumensari et al., 2103.12504]

Towards a combined explanation

To account for **both** anomalies:

- $\Lambda_{\text{NP}} \sim \text{few TeV}$
- NP coupled dominantly to 3rd family, smaller couplings to 1st and 2nd
- NP in semileptonic processes only (no 4-lepton/4-quark)

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Minimal SMEFT setup: left-handed NP in semi-leptonic operators

$$\mathcal{L} = -\frac{1}{v^2} \left(C_{\ell q}^{(3)} (\bar{\ell}_L \gamma^\mu \tau^a \ell_L) (\bar{q}_L \gamma^\mu \tau^a q_L) C_{\ell q}^{(1)} (\bar{\ell}_L \gamma^\mu \ell_L) (\bar{q}_L \gamma^\mu q_L) \right)$$

works well: ✓ $\Delta C_9^\mu = -\Delta C_{10}^\mu$

✓ V_L solution to $b \rightarrow c\tau\nu$

$SU(2)_L \downarrow$

large $b_L \rightarrow s_L \tau_L \tau_L$

\downarrow

✓ ΔC_9^U by RGE mixing

[\[Crivellin et al., 1807.02068\]](#)

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[Crivellin et al., [1807.02068](#)]

! $b \rightarrow s\nu_{(\tau)}\bar{\nu}_{(\tau)}$ requires $C_{\ell q}^{(3)} \approx C_{\ell q}^{(1)}$

(automatically satisfied for U_1 ,
needs to be enforced otherwise)

Explaining both sets of anomalies

$$\mathcal{L} = -\frac{2}{v^2} C_{LL}^{ij\alpha\beta} (\bar{q}_L^i \gamma^\mu l_L^\alpha) (\bar{l}_L^\beta \gamma_\mu q_L^j)$$

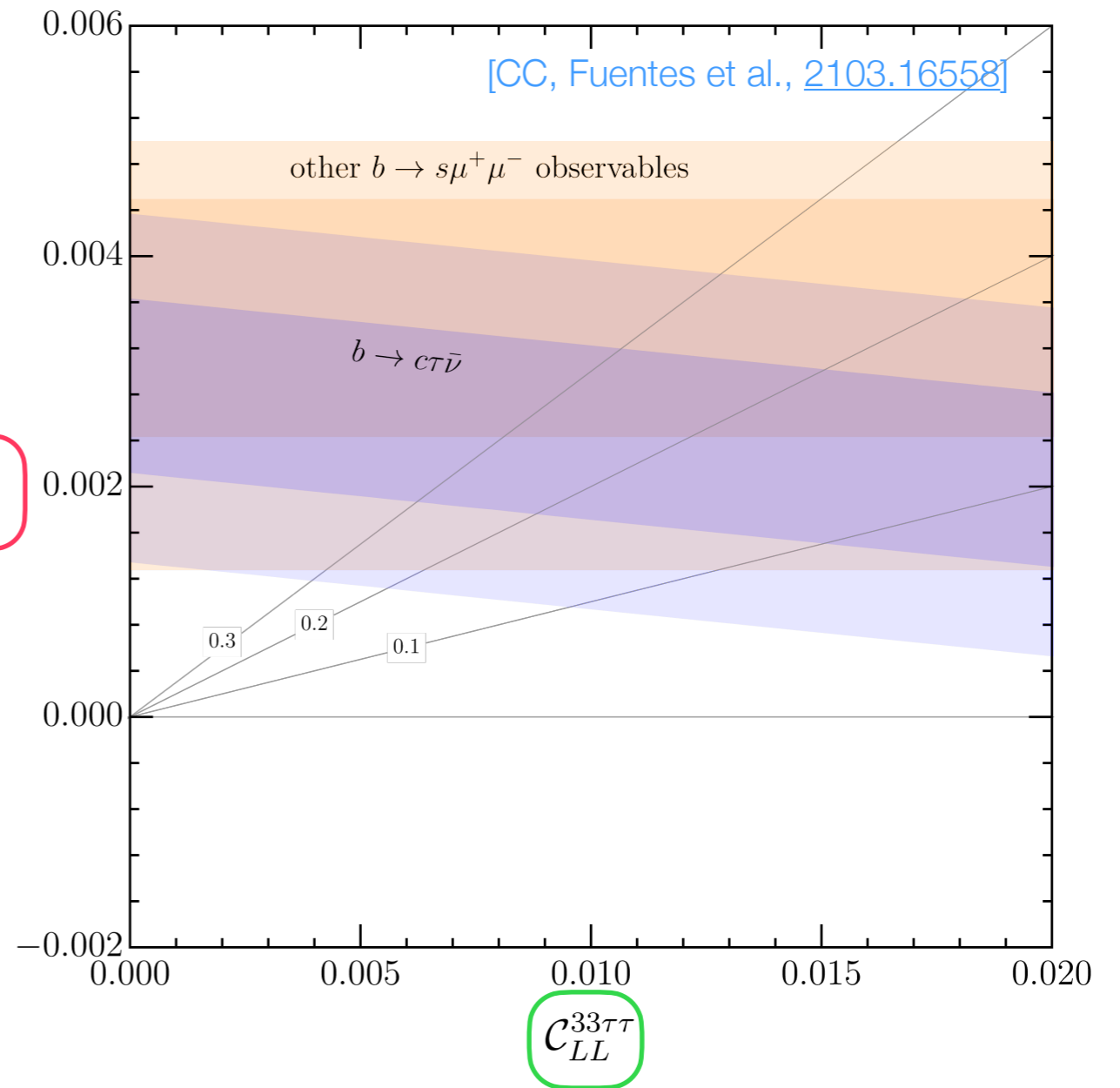
$$R_{D^{(*)}} = R_{D^{(*)}}^{\text{SM}} \left[1 + 2\text{Re} \left(C_{LL}^{33\tau\tau} + \frac{V_{cs}}{V_{cb}} C_{LL}^{23\tau\tau} \right) \right]$$

$$\Delta C_9^\mu = -\Delta C_{10}^\mu \approx -\frac{2\pi}{\alpha V_{ts}^* V_{tb}} C_{LL}^{23\mu\mu}$$

$$\Delta C_9^U \approx \frac{1}{V_{ts}^* V_{tb}} \frac{2}{3} C_{LL}^{23\tau\tau} \ln \frac{\Lambda^2}{m_b^2},$$

$$C_{LL}^{23\tau\tau}$$

Looks like plenty of space to explain $b \rightarrow sll$ [orange] and $b \rightarrow c\tau\nu$ [purple].

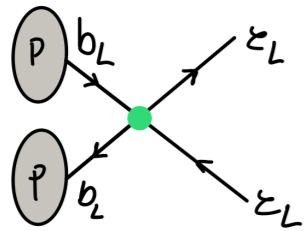


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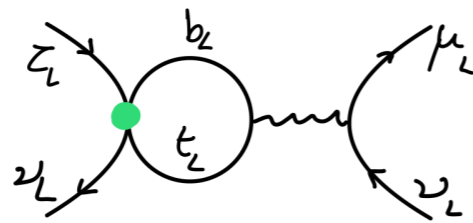
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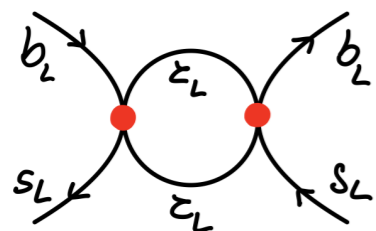
Many other constraints:



$pp \rightarrow \tau\tau$



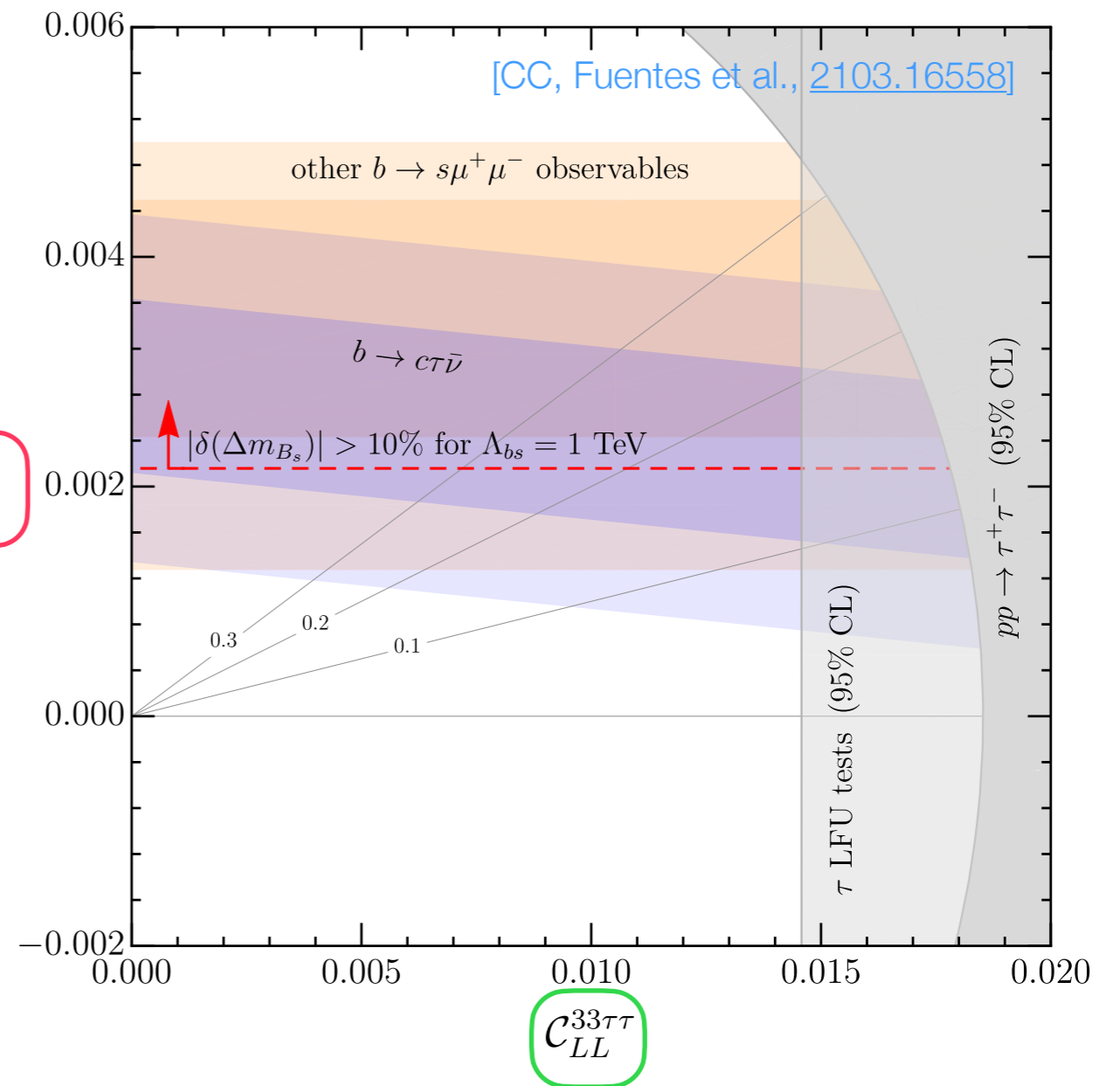
τ LFU tests



$B_s - \bar{B}_s$

...not much wiggle room!

$C_{LL}^{23\tau\tau}$

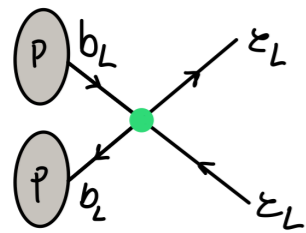


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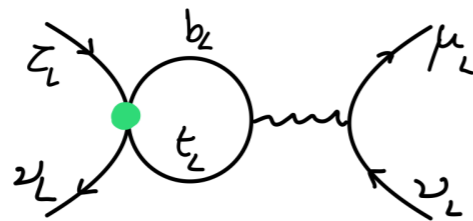
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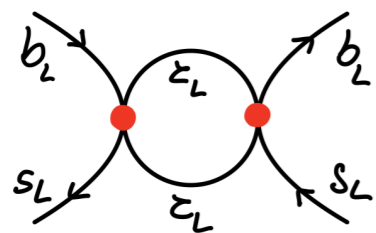
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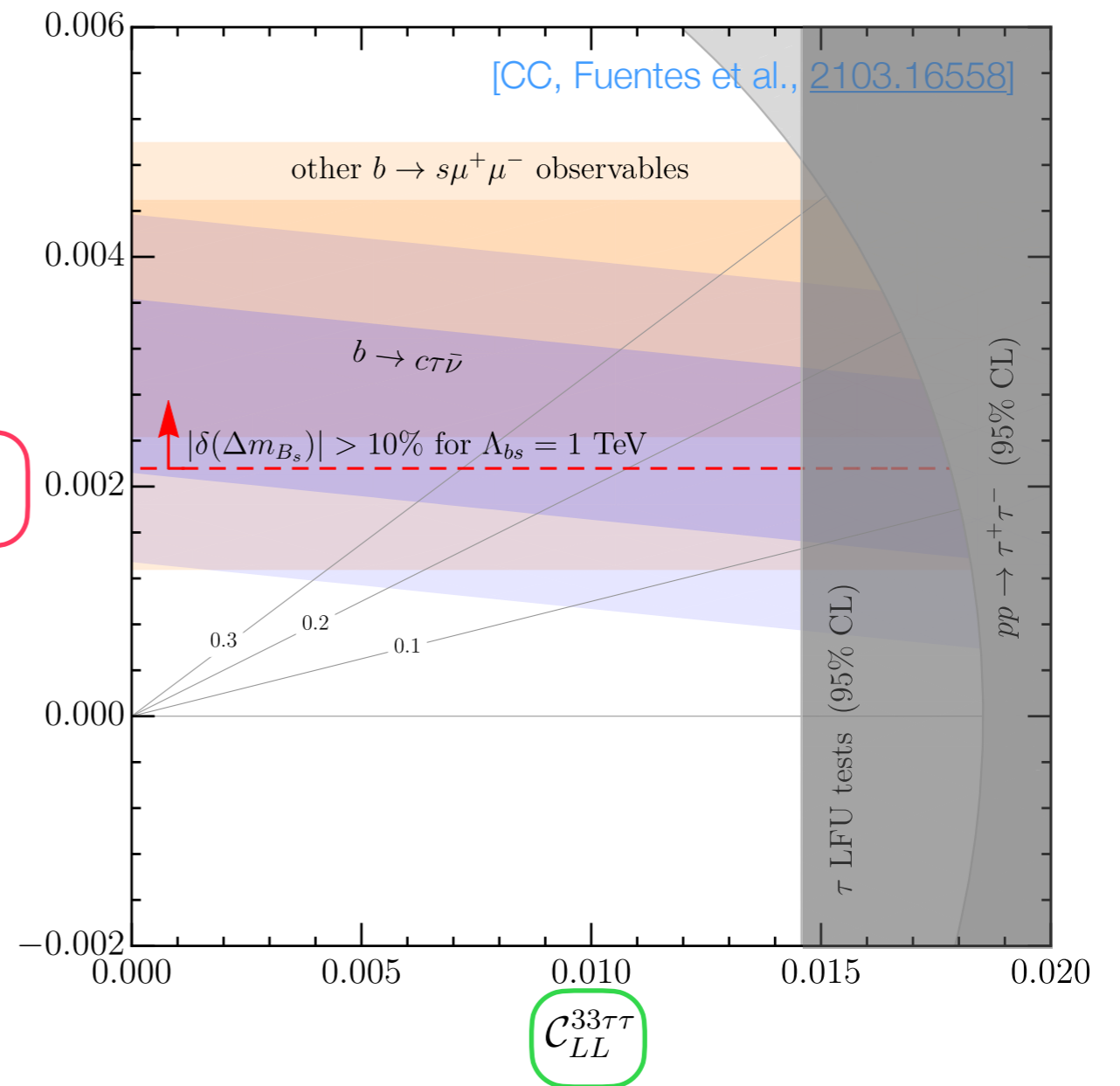
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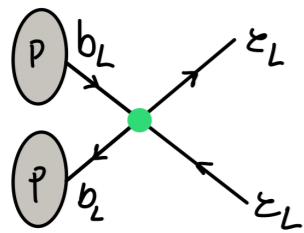


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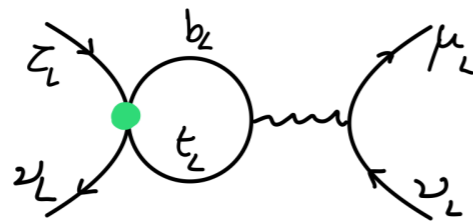
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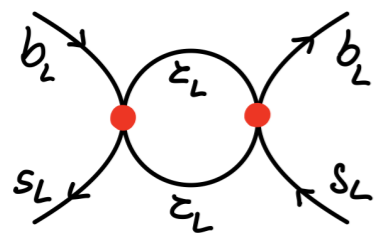
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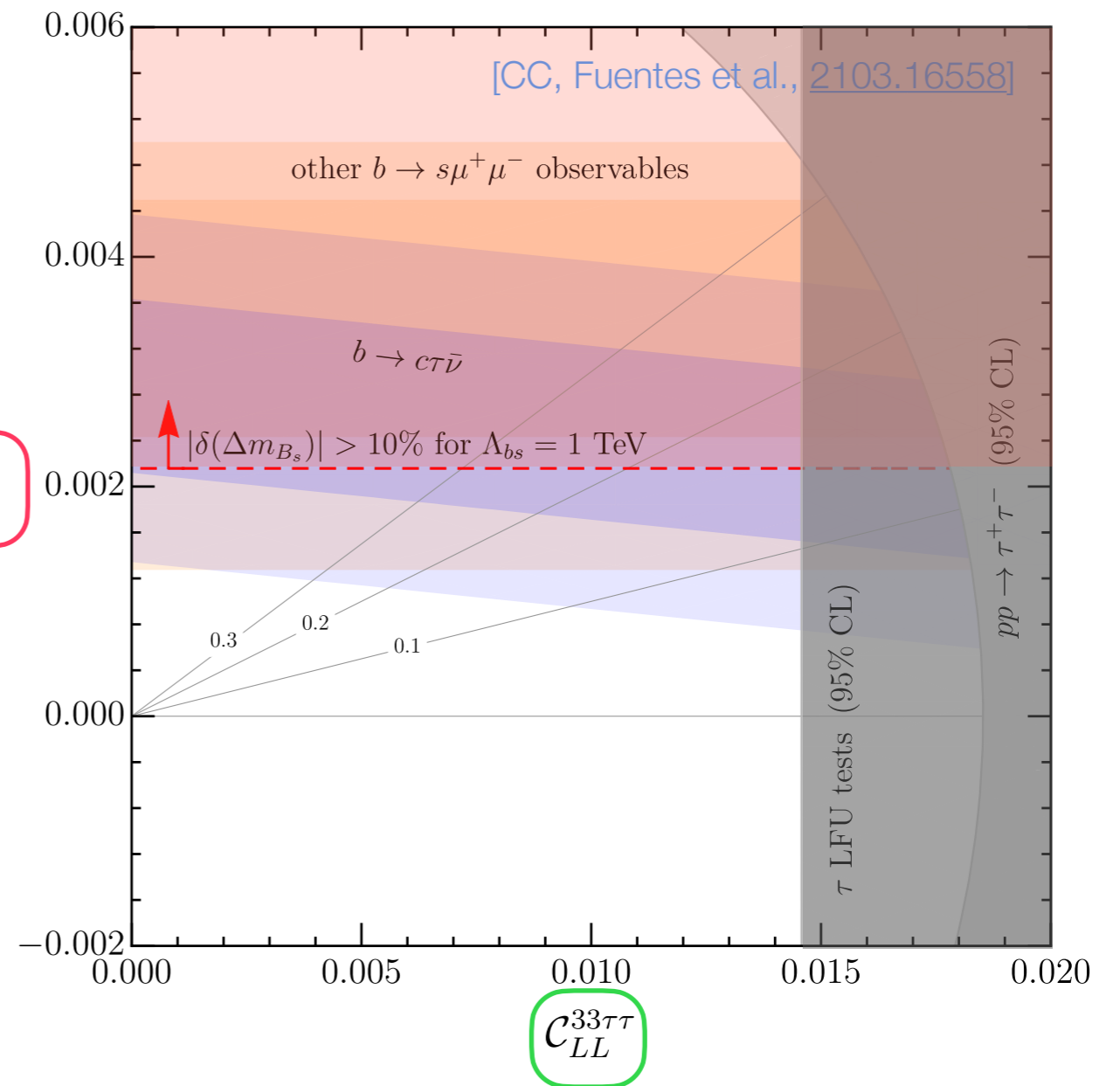
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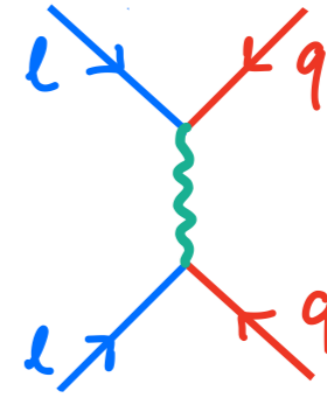
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Simplified models for B anomalies: leptoquarks.

The **only** viable mediators are **leptoquarks**:

- no 4-lepton and 4-quark at tree level,
- no resonant production at high- p_T



Which LQ explains what?

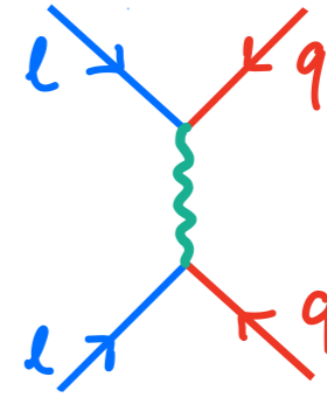
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[Sumensari et al., [2103.12504](#)]

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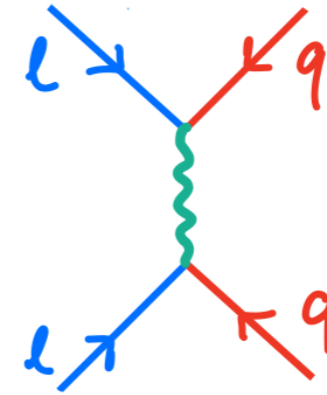
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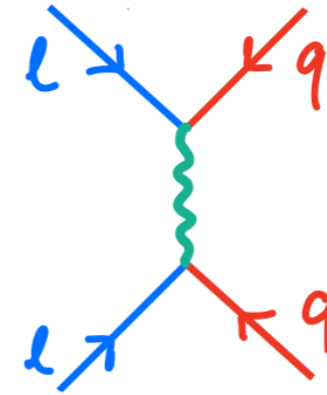
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[Crivellin et al. [1703.09226](#); Buttazzo et al. [1706.07808](#); Marzocca [1803.10972...](#)]

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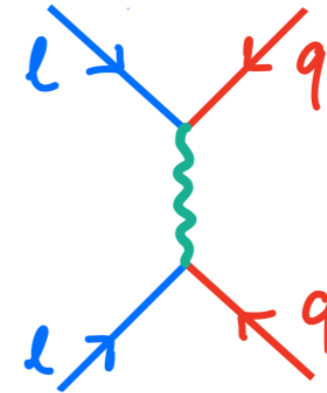
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$S_3 + R_2$ [[Bečirević et al., 1806.05689](#)]

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$S_3 + R_2$ [[Bečirević et al., 1806.05689](#)]

$U_1 \sim (\mathbf{3}, \mathbf{1})_{2/3}$ (+ UV completion)

[[di Luzio et al., 1708.08450](#); [Calibbi et al., 1709.00692](#); [Bordone, CC, et al. 1712.01368](#); [Barbieri, Tesi 1712.06844](#); [Heck, Teresi 1808.07492...](#)]

The U_1 simplified model

The **vector leptoquark** is the only single mediator solution:

- ✓ no tree-level contribution to $b \rightarrow s\nu\bar{\nu}$, protected from proton decay
- ! does not come alone: additional massive vectors (Z' , G'), vector-like fermions

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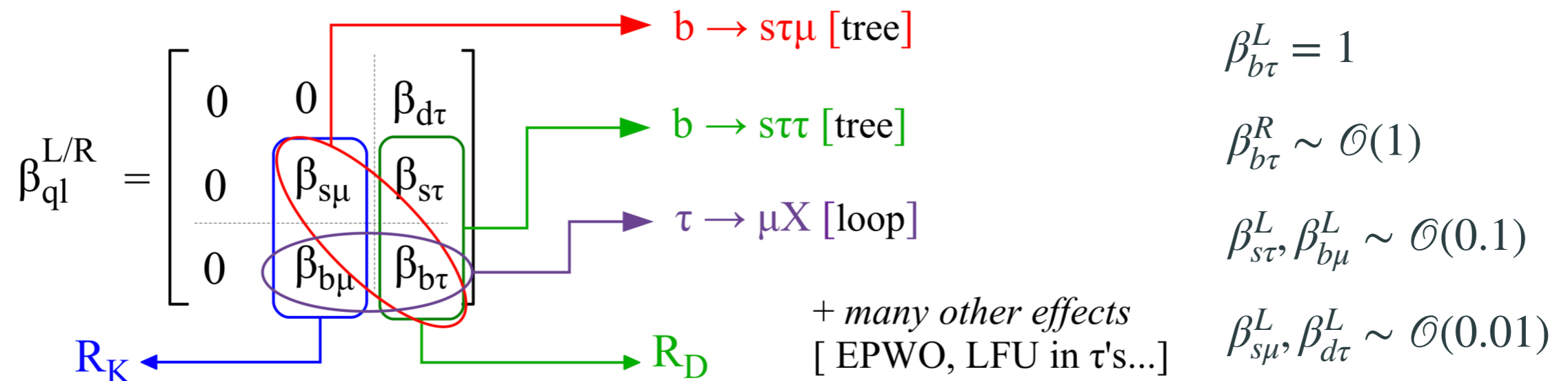
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Stick to simplified model:

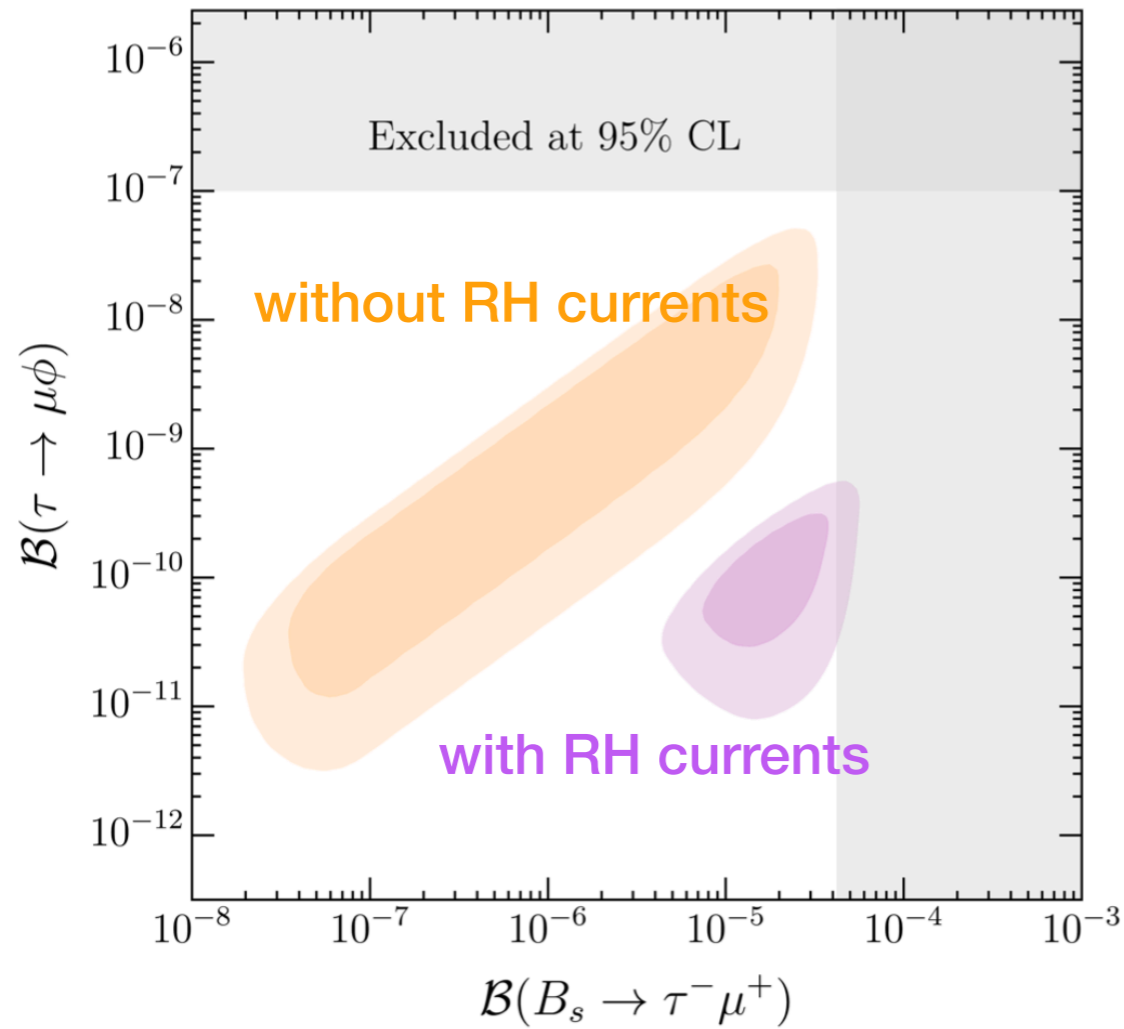
$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} U_1^\mu \left[\beta_L^{i\alpha} (\bar{q}_{L\mu}^i \gamma_\mu \ell_L^\alpha) + \beta_R^{i\alpha} (\bar{d}_{R\mu}^i \gamma_\mu e_R^\alpha) \right] + \text{h.c.} \quad U_1 \sim (\mathbf{3}, \mathbf{1}, 2/3)$$

Good description of all low-energy data with a “natural” flavor structure:



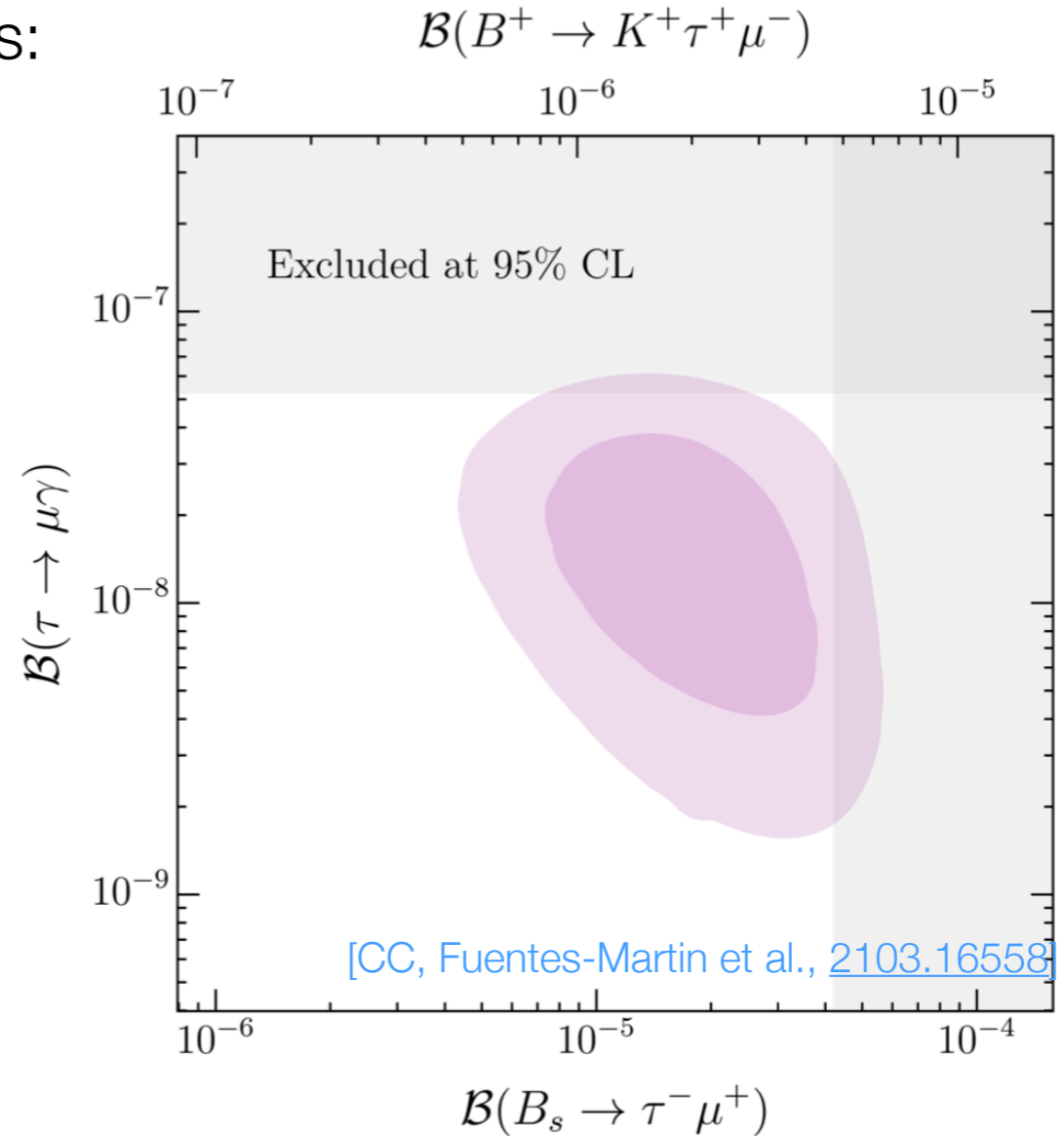
Low-energy predictions for the U_1

Large τ/μ LFV in $b \rightarrow s\tau\mu$ and τ decays:



$$\mathcal{B}(B_s \rightarrow \tau\mu) \approx \mathcal{B}(B \rightarrow K\tau\mu) \approx 10^{-7} - 10^{-6}$$

$$\mathcal{B}(\tau \rightarrow \mu\phi) \approx 10^{-10} - 10^{-8}$$



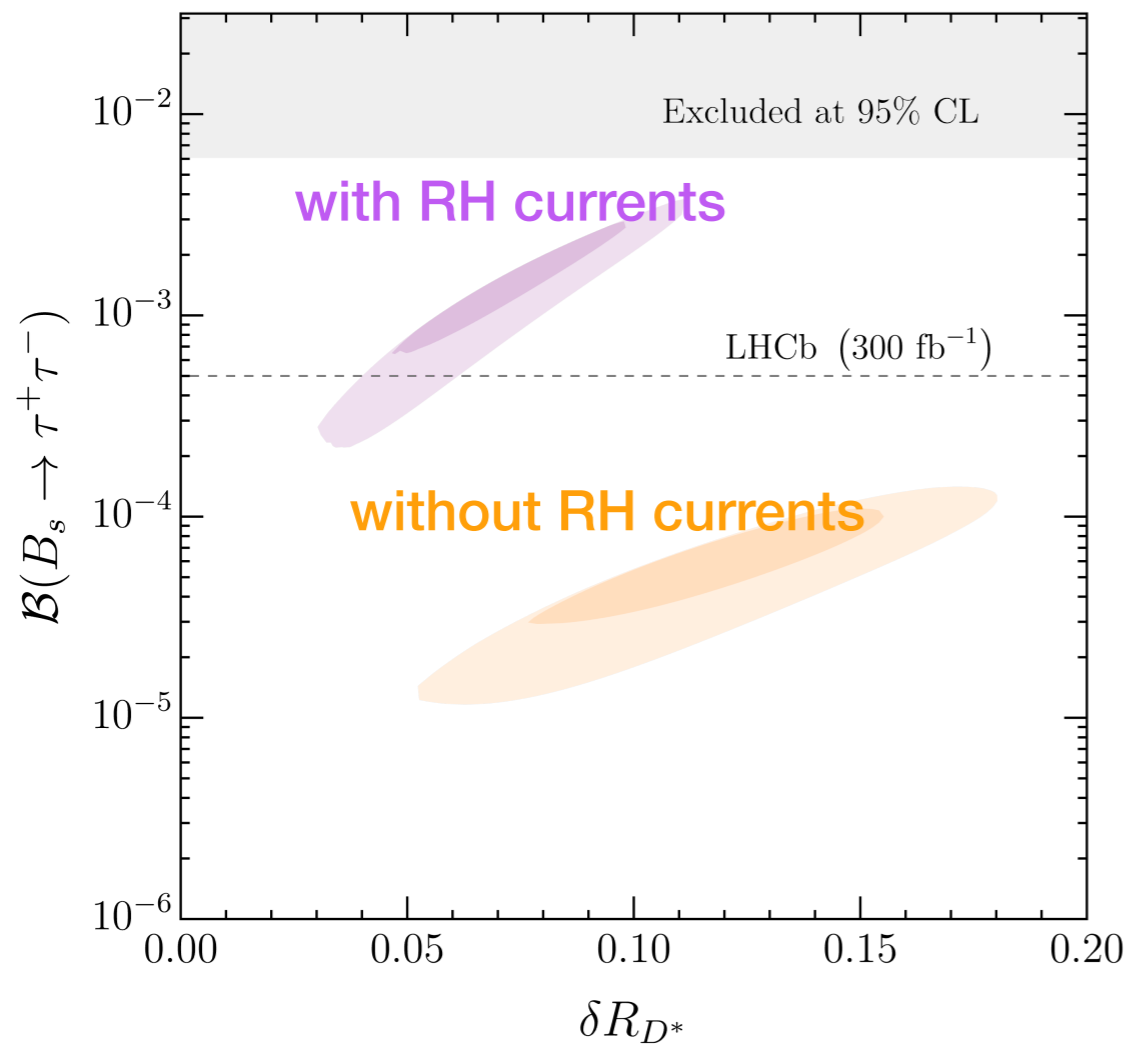
$$\mathcal{B}(B_s \rightarrow \tau\mu) \approx 1 \times 10^{-5}$$

$$\mathcal{B}(B \rightarrow K\tau\mu) \approx 1 \times 10^{-6}$$

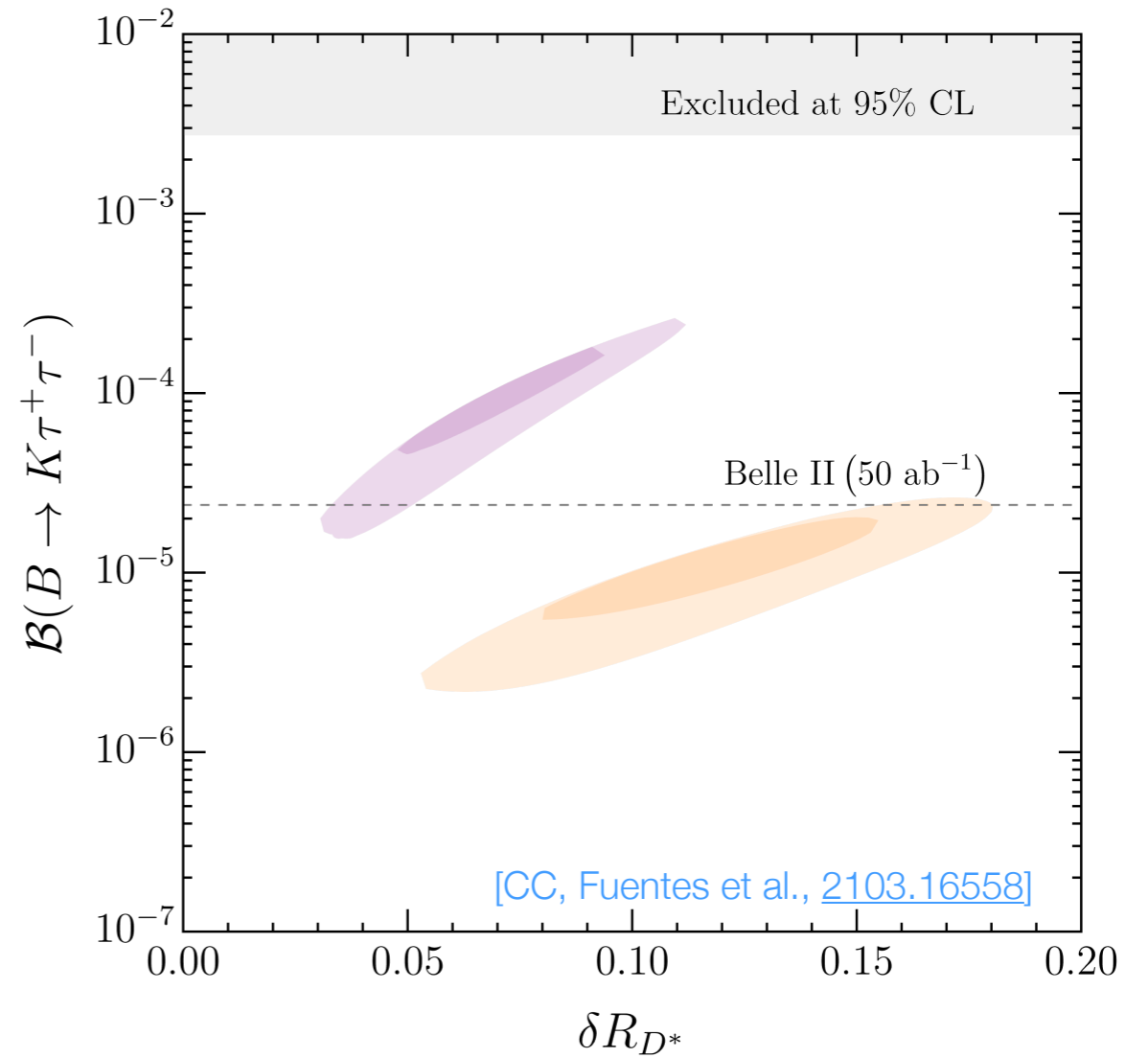
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Low-energy predictions for the U_1

Large $b \rightarrow s\tau\tau$:



$$\frac{\mathcal{B}(B_s \rightarrow \tau\tau)}{\mathcal{B}(B_s \rightarrow \tau\tau)_{\text{SM}}} \approx \frac{\mathcal{B}(B \rightarrow K\tau\tau)}{\mathcal{B}(B \rightarrow K\tau\tau)_{\text{SM}}} \approx 1 \times 10^2$$

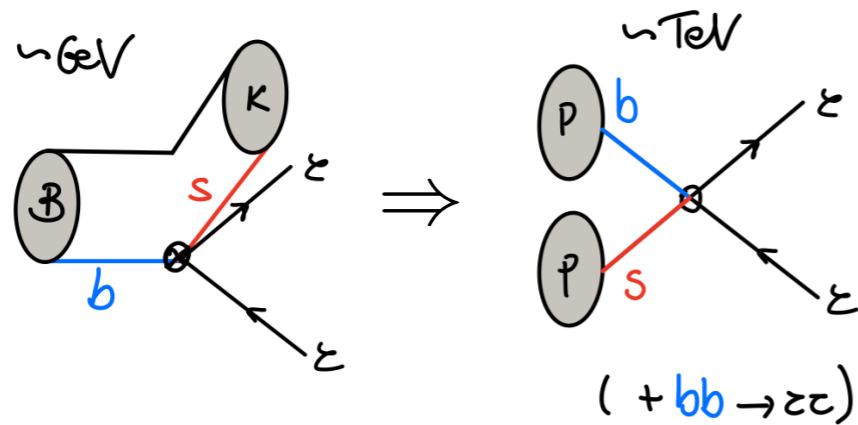


$$\frac{\mathcal{B}(B_s \rightarrow \tau\tau)}{\mathcal{B}(B_s \rightarrow \tau\tau)_{\text{SM}}} \approx \frac{\mathcal{B}(B \rightarrow K\tau\tau)}{\mathcal{B}(B \rightarrow K\tau\tau)_{\text{SM}}} \approx 1 - 2 \times 10^3$$

High-pT bounds for the U_1

The same interaction can be probed in **di-tau tails**.

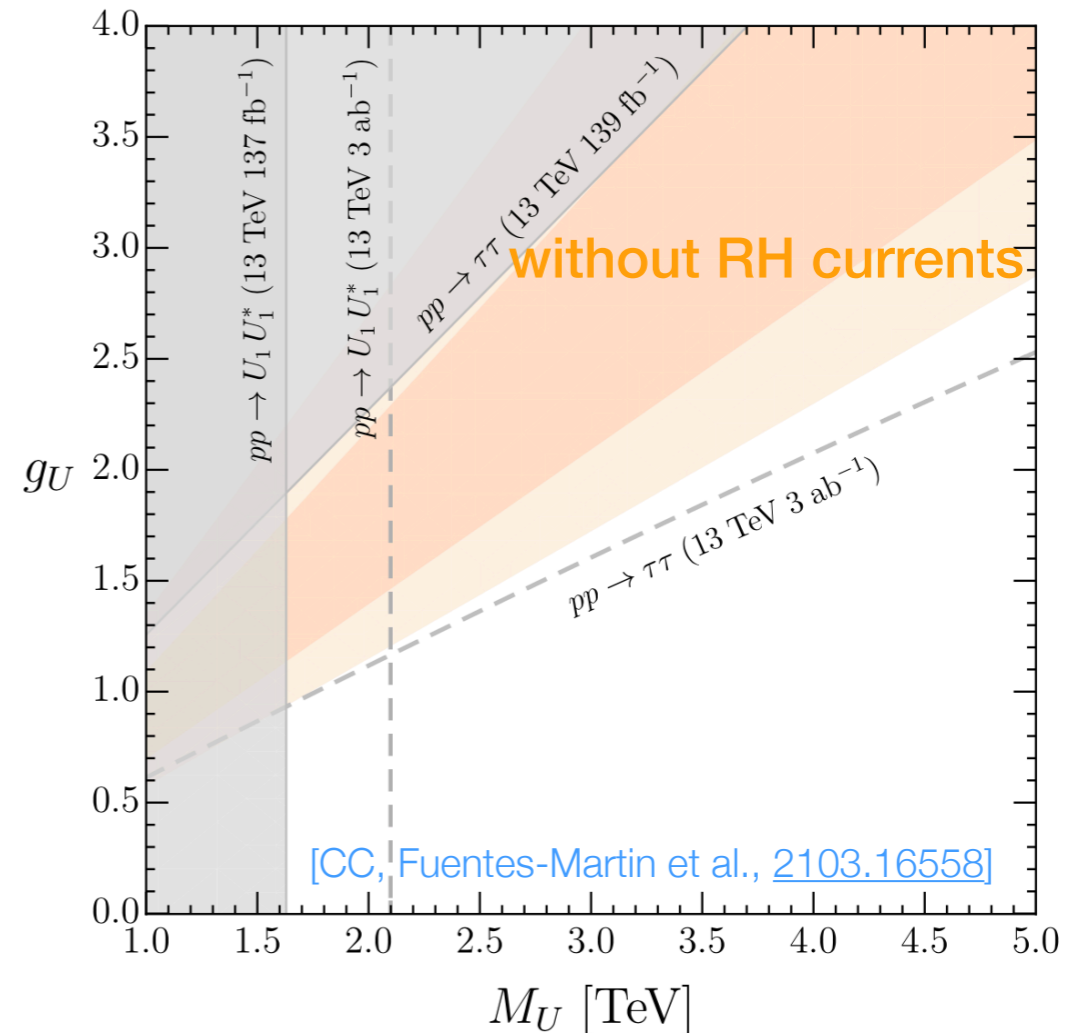
[Faroughy et al, [1609.07138](#);
Fuentes-Martin et al. [2003.12421..](#)]



U_1 solution completely **falsifiable at HL-LHC!**
(same for $R_2 + S_3$, still space left for $S_1 + S_3$)

Similar enhancements in all models for $R_D^{(*)}$.

Caveat: these conclusions are **very sensitive** to the central value of $R_D^{(*)}$!

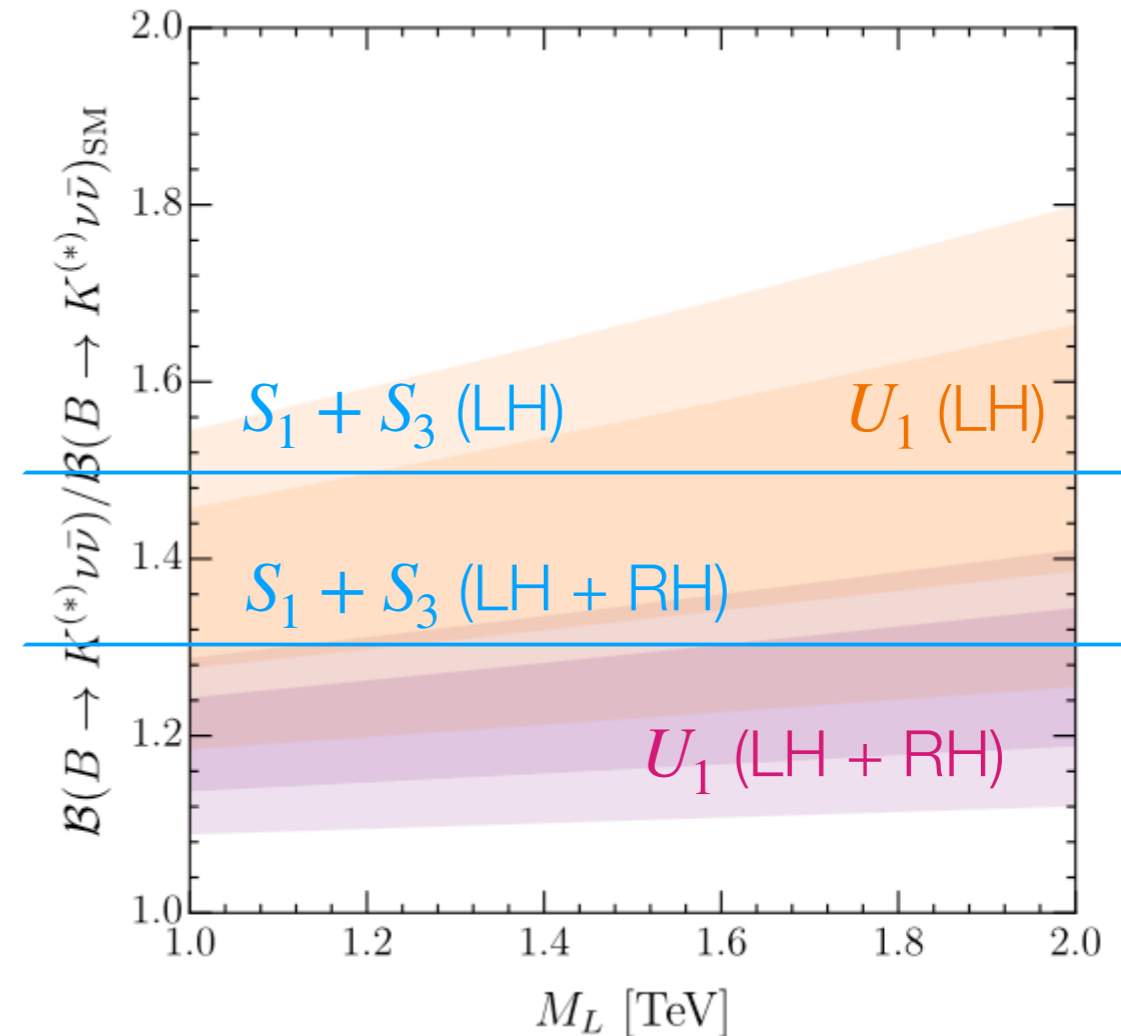


$B \rightarrow K\nu\bar{\nu}$ and B anomalies

One of the toughest constraints for model building; all combined solutions enhance it.

$$U_1, S_1 + S_3 \quad \frac{\mathcal{B}(B \rightarrow K\nu\nu)}{\mathcal{B}(B \rightarrow K\nu\nu)_{\text{SM}}} \approx 1.2 - 1.5$$

- for the U_1 absent at tree level, but any UV completion generates it at one loop
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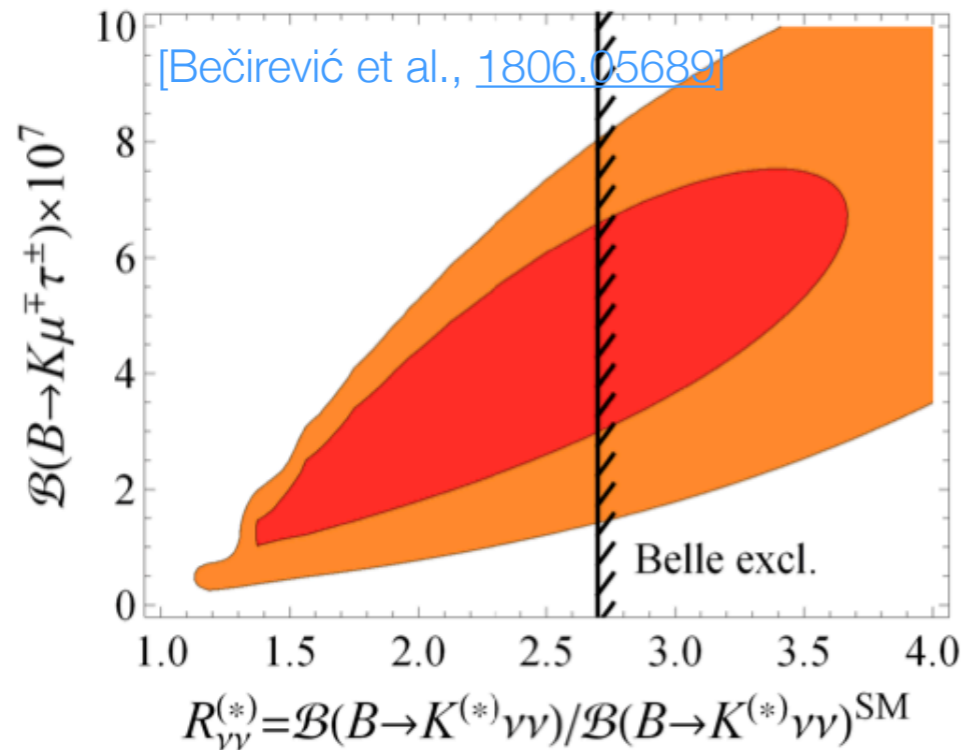


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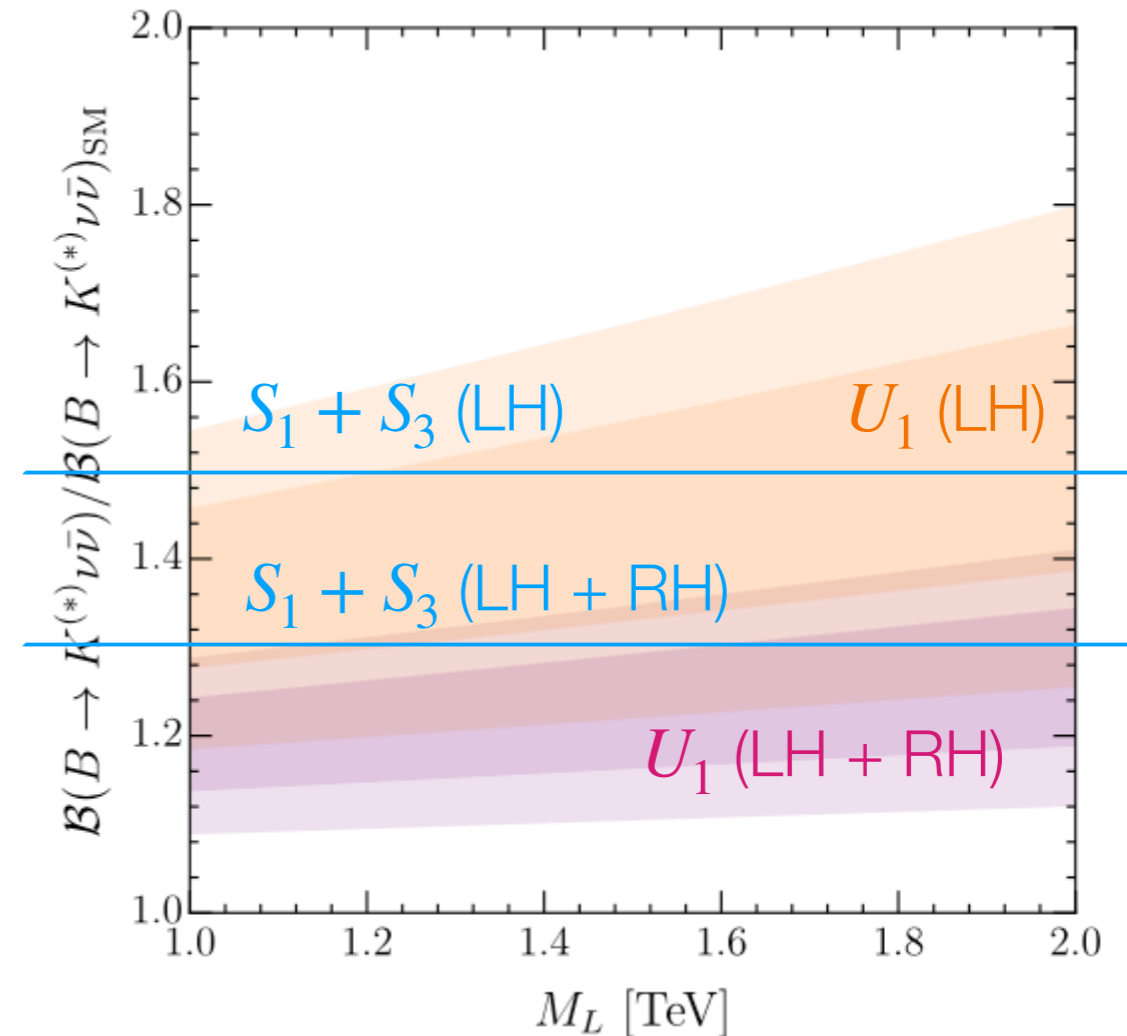
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$$R_2 + S_3$$

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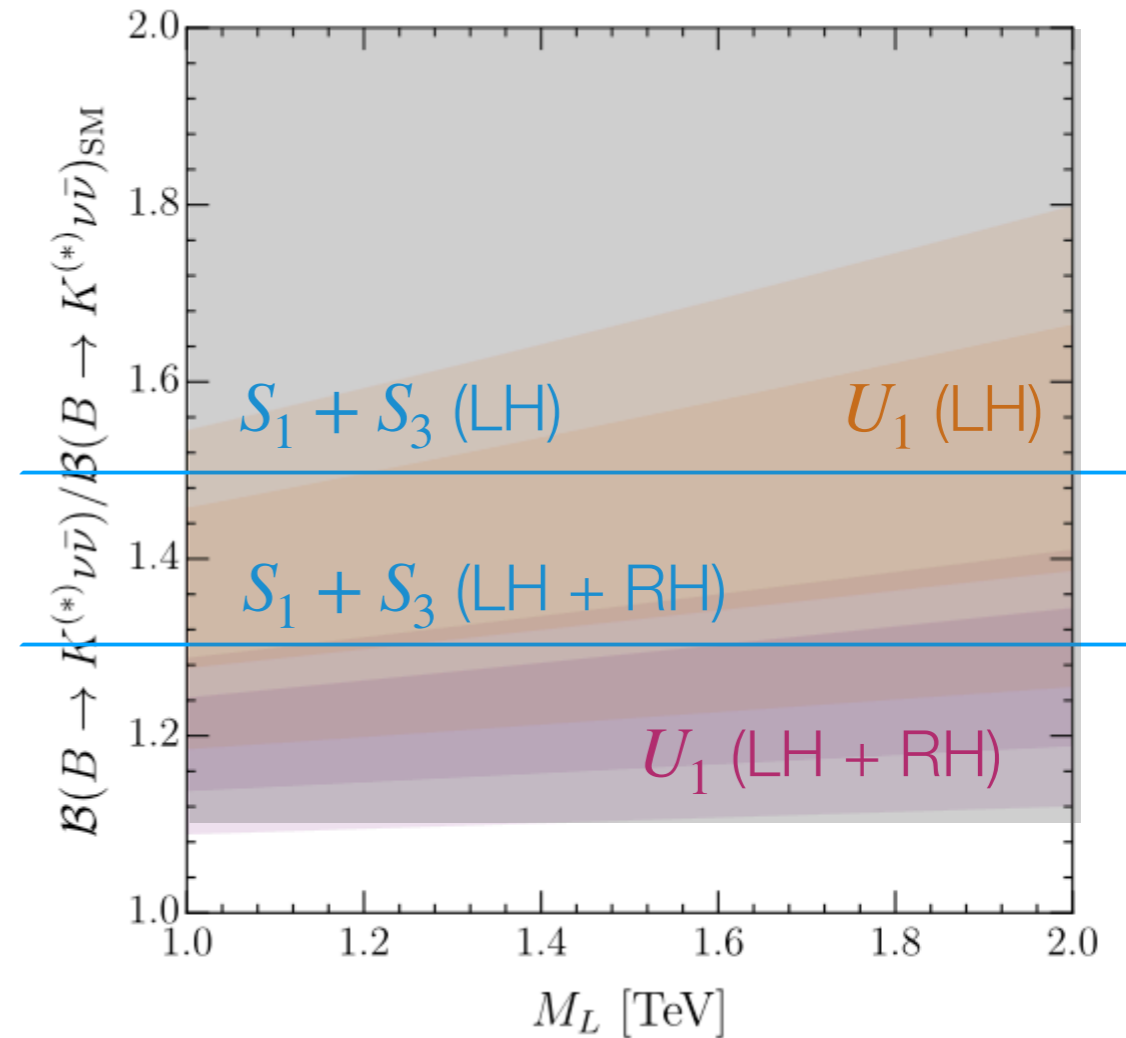
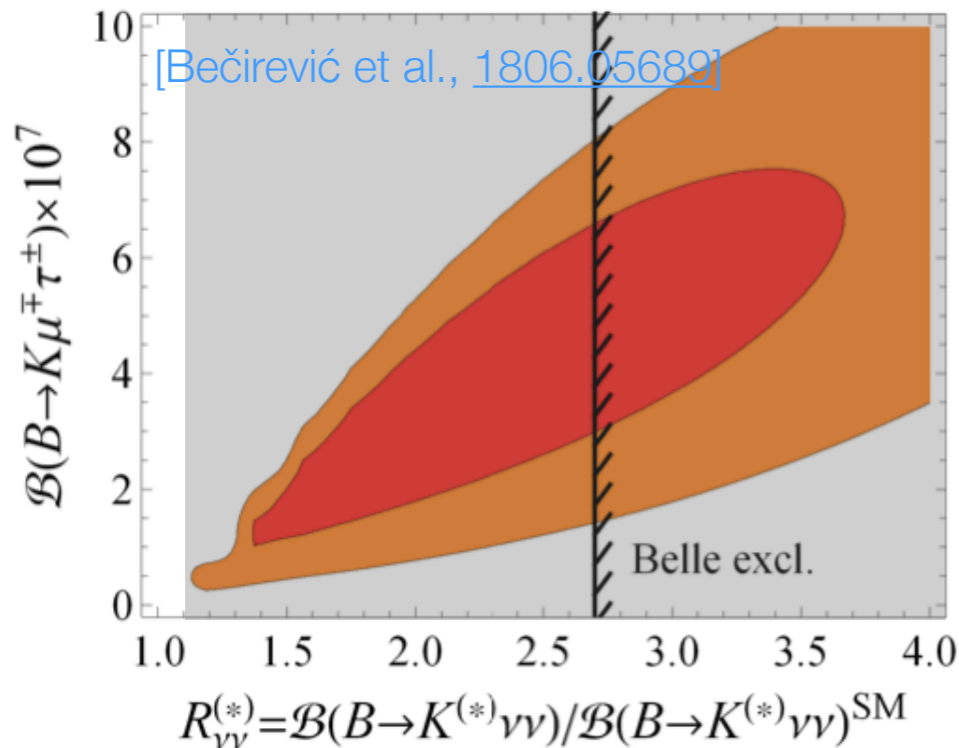


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Th: model building for $b \rightarrow sll$ is *very different* from that for $b \rightarrow sll$ and $b \rightarrow c\tau\nu$ without $R_{D^{(*)}}$ no obvious connection flavor anomalies - flavor hierarchies

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[Looking forward to the plenty upcoming measurements from both the energy and intensity frontiers!](#)