

High energy $\pi\eta^{(\prime)}$ production and the double Regge exchange

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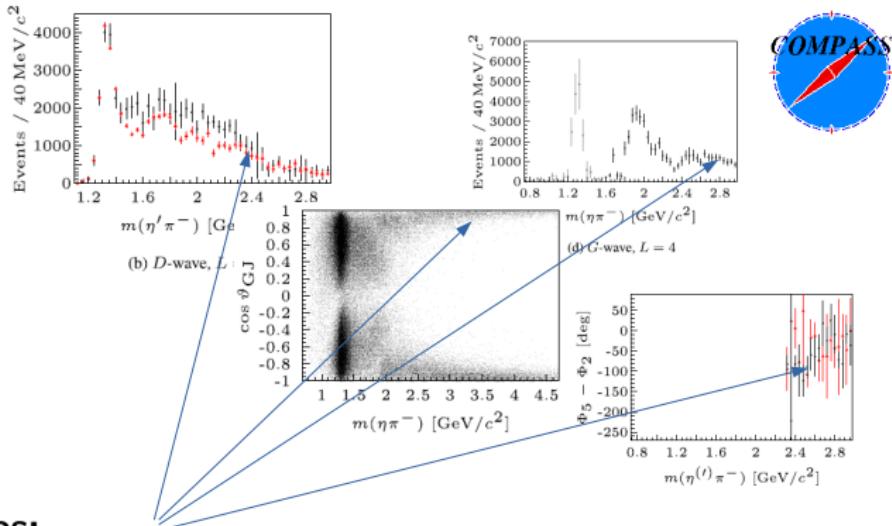


Motivation of the $\pi^- \eta^{(\prime)}$ channel study

- $\pi^- \eta^{(\prime)}$ pairs constitute a golden channel for the searches of hybrid exotics (odd partial waves are exotic)
- COMPASS experiment at CERN ([Adolph et al. PLB 740, 2015](#)) observed a strong forward-backward asymmetry in the $\pi^- \eta^{(\prime)}$ channels (stronger in the $\pi\eta'$ channel)
- the asymmetry is due to odd-even partial wave interference
- the strongest odd wave is the P -wave, which in the resonance region can be attributed to the π_1 hybrid
- relation between the high and low invariant mass region can be described in terms of Finite Energy Sum Rules (special kind of dispersion relations)



Experimental motivation



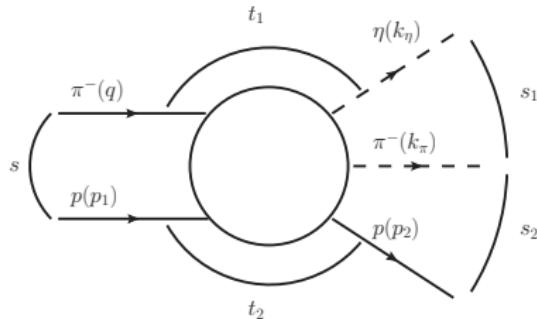
Objectives:

- Describe the $\pi^- \eta^{(\prime)}$ production above the resonance region.
- Identify dominant Regge exchanges and amplitude strengths.



Kinematics

- The $2 \rightarrow 3$ reaction (upolarized) is determined by 5 kinematical variables
- In Regge analysis it is customary to choose 5 Lorentz invariants, eg:
 $s = (q + p_1)^2$, $t_1 = (q - k_\eta)^2$, $t_2 = (p_1 - p_2)^2$, $s_1(k_\eta + k_\pi)^2$,
 $s_2 = (k_\pi + p_2)^2$.



- COMPASS analysis was performed in the Gottfried-Jackson frame, with θ and ϕ defining the direction of outgoing $\eta^{(\prime)}$
- In this frame $\cos \theta$ is related to t_1 and ϕ is related to s_2 (the other variables being s , s_1 and t_2).



Model vs. experimental data

- The COMPASS experiment operated with a fixed beam momentum of 191 GeV and t_2 was integrated in the region $t_2 \in [-1.0, -0.1]$ GeV 2 .
- The invariant mass ($m = \sqrt{s_1}$) and angular dependent intensity was parametrized as:

$$I(m, \Omega)_{COMPASS} = \sum_{\epsilon=\pm 1} \left| \sum_{L,M} f_{LM}^{\epsilon}(m) \Psi_{LM}^{\epsilon}(\Omega) \right|^2,$$

where ϵ is the reflectivity (\approx naturality in high energy limit), kept $\epsilon=+1$ and

$$\Psi_{LM}^{\epsilon=+1}(\theta, \phi) = \sqrt{2} Y_{LM}(\theta, 0) \sin(M\phi)$$

$$f_{LM}^{\epsilon=+1}(m) = \sqrt{I_{LM}} e^{i\phi_{LM}}$$

where the I_{LM} and ϕ_{LM} are the COMPASS experimental partial wave intensities and phaseshifts (determined relative to ϕ_{21}). 

Fitting experimental data at large invariant masses

- The easiest approach: fit the model partial waves to experimental ones

But:

- COMPASS analysis was based on partial wave expansion truncated at $L = 6$ and $M = 1$, even though it converges slowly for the $\pi\eta^{(\prime)}$ invariant masses above the resonance region,
- The intensity was normalized to the number of events in the mass bin.

Therefore:

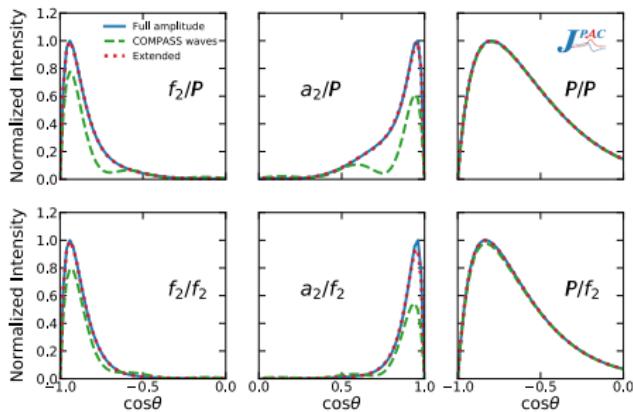
- The angular distribution constructed from the truncated partial wave expansion represents total yield.



Fitting experimental data at large invariant masses

On the model side

- Model intensities for full amplitudes and partial wave amplitudes truncated at $L = 6$ differ considerably.



Recipe:

- Fit the mass and angular dependent intensity.



Fitting experimental data at large invariant masses

- Practically we fit the extended log likelihood ([Barlow, Nucl.Instr.Meth. A 297, 1990](#))

$$\mathcal{L}_{\text{ext}} = \sum_{m_i} \int d\Omega [I_{JPAC}(m_i, \Omega) - I_{COMPASS}(m_i, \Omega)] \log I_{JPAC}(m_i, \Omega)$$

with model intensity defined as:

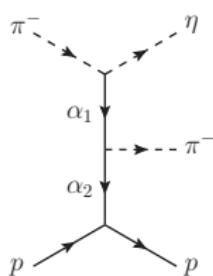
$$I_{JPAC}(m, z = \cos \theta, \phi) = q |T|^2 = q \left| \sum_{i=1}^6 a_i T^i(m_{\pi\eta}, z, \phi) \right|^2$$

where $q = \lambda^{1/2}(m^2, m_\pi^2, m_\eta^2)/m$.

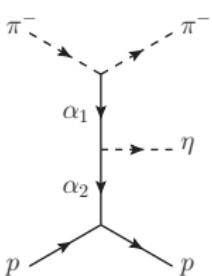


Double Regge exchange model

- Leading natural exchanges



Type I - fast η Type II - fast π



Type of diagram	α_1	α_2
<i>I</i> (fast η)	α_{a_2}	α_P
	α_{a_2}	α_{f_2}
<i>II</i> (fast π)	α_{f_2}	α_P
	α_{f_2}	α_{f_2}
	α_P	α_P
	α_P	α_{f_2}

- Individual diagram strengths depend on the 6 coupling triples:
 $G_{a_2\pi\eta} G_{a_2\pi P} G_{PNN}$, $G_{a_2\pi\eta} G_{a_2\pi f_2} G_{f_2NN}$, $G_{f_2\pi\pi} G_{f_2\eta P} G_{PNN}$,
 $G_{f_2\pi\pi} G_{f_2\eta f_2} G_{f_2NN}$, $G_{P\pi\pi} G_{\eta PP} G_{PNN}$, $G_{P\pi\pi} G_{f_2\eta P} G_{f_2NN}$, wherein at least one coupling is unknown.
- Therefore we treat diagram strengths as parameters to be fitted.



General form of the amplitude

- Double reggeon exchange (we follow Shimada *et al.* NPB 142, 1978)

$$T = -K\Gamma(1-\alpha_1)\Gamma(1-\alpha_2)$$

$$\left[(\alpha' s)^{\alpha_1-1} (\alpha' s_2)^{\alpha_2-\alpha_1} \xi_1 \xi_{21} \hat{V}_1 + (\alpha' s)^{\alpha_2-1} (\alpha' s_1)^{\alpha_1-\alpha_2} \xi_2 \xi_{12} \hat{V}_2 \right]$$

$\xi_i = \frac{1}{2}(\tau_i + e^{-i\pi\alpha_i})$ and $\xi_{ij} = \frac{1}{2}(\tau_i\tau_j + e^{-i\pi(\alpha_i-\alpha_j)})$ are signature factors.

$$\hat{V}_1(\alpha_1, \alpha_2, \eta) = \frac{\Gamma(\alpha_1 - \alpha_2)}{\Gamma(1 - \alpha_2)} {}_1F_1(1 - \alpha_1, 1 - \alpha_1 + \alpha_2, -\kappa)$$

$$\hat{V}_2 = \hat{V}_1(\alpha_1 \leftrightarrow \alpha_2) \quad \text{with} \quad \kappa^{-1} \equiv s/(\alpha' s_1 s_2)$$

and $K = -4\sqrt{s_1}|q||k_\eta||p_2|\sin\theta_2\sin\theta\sin\phi$

- Both α_1 and α_2 are of 2^{++} type so only positive signature and naturality.
- Regge trajectories:

$$\alpha_{f_2}(t) = 0.47 + 0.89t \quad \alpha_{a_2}(t) = 0.53 + 0.90t \quad \alpha_P(t) = 1.08 + 0.25t$$

Fitting procedure

- ① Full amplitude (a_i to be fitted):

$$T(s, s_1, t_2, \cos \theta, \phi; a_i) = K \sum_{i \in I, II} a_i \tilde{T}_i(\alpha_1(t_1), \alpha_2(t_2); s, s_1, t_2, \cos \theta, \phi)$$

- ② Model intensity:

$$I_{JPAC}(m, z = \cos \theta, \phi) = q |T|^2 = q \left| \sum_{i=1}^6 a_i T^i(m_{\pi\eta}, z, \phi) \right|^2$$

- ③ Extended log likelihood fit:

$$\mathcal{L}_{\text{ext}} = \sum_{m_i} \int d\Omega [I_{JPAC}(m_i, \Omega) - I_{COMPASS}(m_i, \Omega) \log I_{JPAC}(m_i, \Omega)]$$

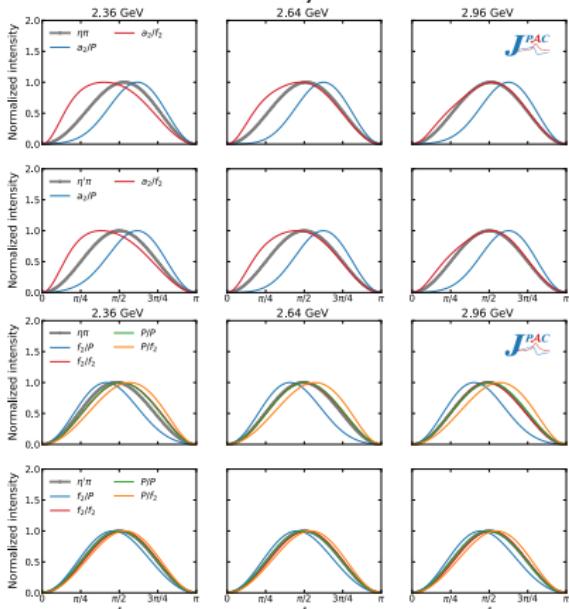
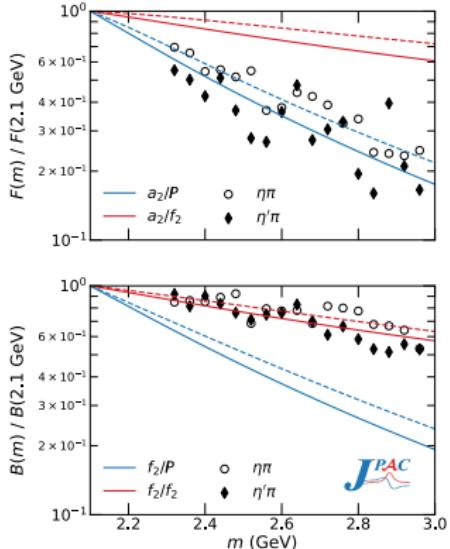


Fit results



Minimal set of amplitudes

Fitting the full set of 6 amplitudes is statistically inconclusive.



From the analysis of the forward and backward mass and ϕ distributions we infer that the minimal set of amplitudes should include a_2/P , a_2/f_2 and f_2/f_2 .

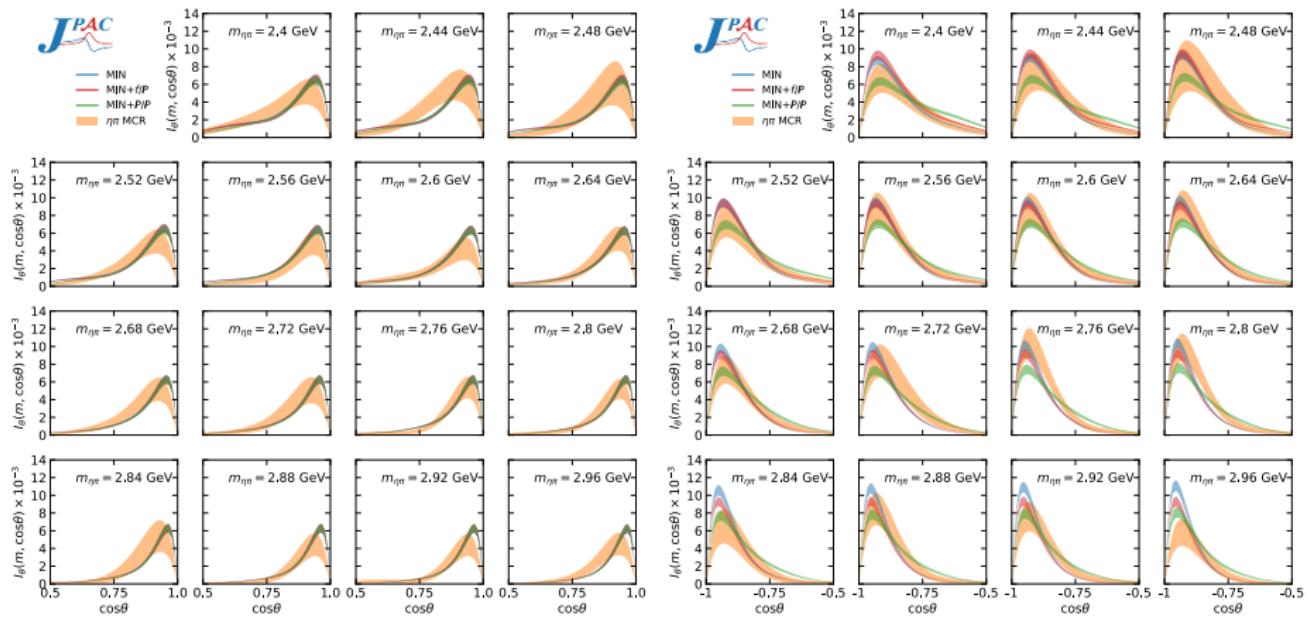
Parameter values from the fits

- To make the parameter evaluation stable we can fit at most 4 amplitudes.
- We exclude P/f_2 from all fits as it may disrupt the ϕ distribution.

Channel		MIN		MIN+ f/P		MIN+ P/P	
		MVR	MCR	MVR	MCR	MVR	MCR
$\eta\pi$	$\mathcal{L} \times 10^{-4}$	-22.8	-21.9 ± 0.9	-22.7	-22.0 ± 0.9	-22.8	-22.1 ± 0.8
	$c_{a2}P$	0.29	0.42 ± 0.03	0.28	0.40 ± 0.04	0.29	0.36 ± 0.04
	$c_{a2}f_2$	3.67	3.3 ± 0.4	3.70	3.4 ± 0.4	3.59	3.8 ± 0.4
	$c_{f_2}P$	—	—	-0.20	-0.30 ± 0.05	—	—
	$c_{f_2}f_2$	-11.82	-11.0 ± 0.3	-8.99	-6.6 ± 0.7	-10.86	-8.9 ± 0.4
	CP_P	—	—	—	—	0.0073	0.0135 ± 0.002
$\eta'\pi$	$\mathcal{L} \times 10^{-4}$	-11.7	-10.9 ± 1.0	-11.7	-11.0 ± 1.0	-11.8	-11.4 ± 1.0
	$c_{a2}P$	0.16	0.37 ± 0.07	0.16	0.34 ± 0.05	0.19	0.35 ± 0.05
	$c_{a2}f_2$	1.50	0.4 ± 0.6	1.51	0.7 ± 0.5	1.22	0.6 ± 0.5
	$c_{f_2}P$	—	—	-0.21	-0.29 ± 0.03	—	—
	$c_{f_2}f_2$	-11.42	-11.0 ± 0.5	-7.73	-5.5 ± 0.7	-9.01	-7.1 ± 0.6
	CP_P	—	—	—	—	0.012	0.018 ± 0.002

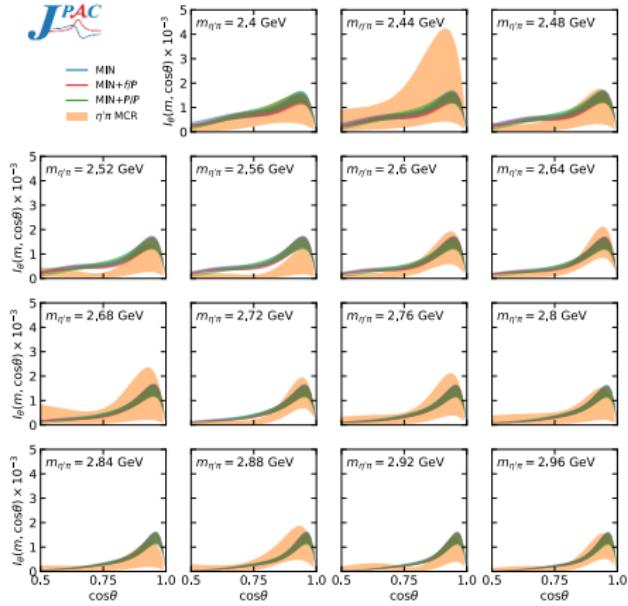


Forward and backward peaks for $\pi\eta$

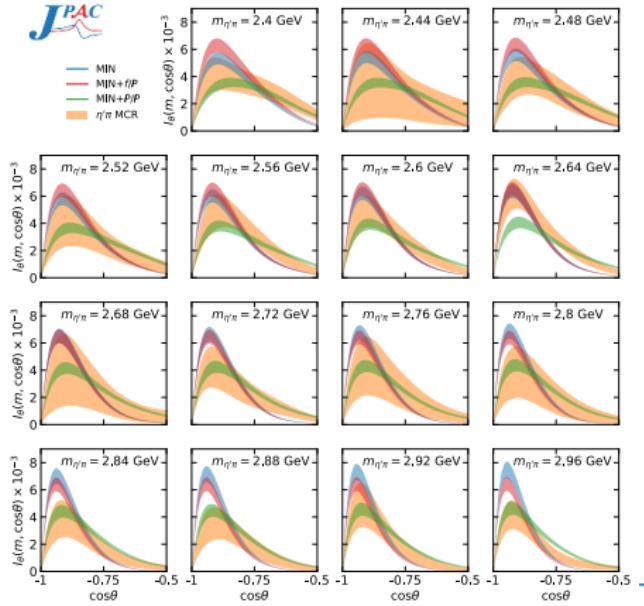


Forward and backward peaks for $\pi\eta'$

J_{PAC}



J_{PAC}

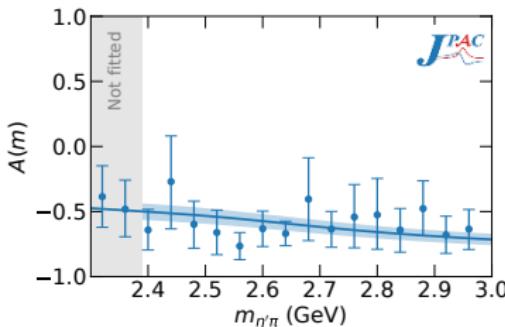
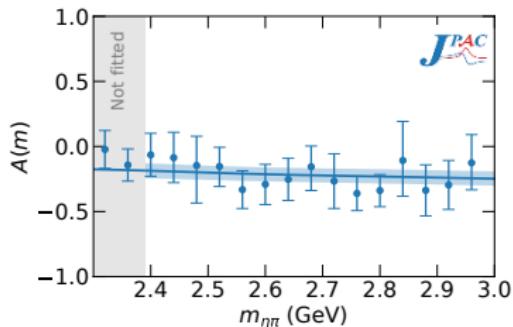


Forward-backward asymmetry

Define the forward-backward asymmetry as:

$$A(m) \equiv \frac{F(m) - B(m)}{F(m) + B(m)},$$

with $F(m) \equiv \int_0^1 d \cos \theta I_\theta(m, \cos \theta)$, $B(m) \equiv \int_{-1}^0 d \cos \theta I_\theta(m, \cos \theta)$



Summary

- We found that a_2/P , a_2/f_2 , f_2/f_2 and either f_2P or P/P (the data do not show clear preference for either exchange) are required to describe $\pi\eta$ intensity.
- To describe $\pi\eta'$ intensity the a_2/P , a_2/f_2 , f_2/f_2 and P/P exchanges are necessary. Glue rich exchange points toward π_1 hybrid production.
- Quite surprisingly, lower f_2 exchange is necessary to describe COMPASS data,



Thank you for your attention

