

# Detector and physics simulation

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#### (LIP/IST)

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# **Physics Simulation**



#### Liliana Apolinário (LIP/IST)

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Not easy to evaluate through analytical calculations...



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(skeleton process has been dressed up and is no longer directly visible)

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Event Generators to the rescue!

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Problem 1	



Problem 1	
Problem 2	







Problem 1		
Problem 2		



#### MC Event Generator



Factorization into simpler (and reasonably accurate) components



MC Event Generator





Same average behaviour and fluctuations as real data



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Detector performance (propagation, magnetic field, shower calorimeter, ....)

Detector Simulation GEANT

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Same format as the real data recorded by the detector



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Detector Simulation GEANT

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  - Beam Remnants and Multi-particle
    Interactions (MPI) (rest of the collision)

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#### Hard Process



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  - SM: Hard QCD, Soft QCD, Heavy-Flavour, DIS, W/Z, Higgs Production...
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  - SM: Hard QCD, Soft QCD, Heavy-Flavour, DIS, W/Z, Higgs Production...
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- Given the topology and kinematics, one can evaluate the cross-section,  $\sigma$ .

## **Parton Distributions**

• Initial topology and kinematics is not fixed, but rather sampled from the parton distribution of the two incoming protons...
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Elementary cross-section:  $\sigma_{ij \to k_1 k_2}(x_1, x_2) = \sigma_{ij \to k_1 k_2}(z_1, y_2)$ 

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Elementary cross-section:  $\sigma_{ij \rightarrow k_1 k_2}(x_1, x_2) + \sigma_{ij \rightarrow k_1 k_2}(z_1, y_2)$ 

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- Cross-section for a process ij  $\rightarrow k$ :  $\sigma_{ij \rightarrow k} = \int dx_1 \int dx_2 f_i^1(x_1) f_j^2(x_2) \hat{\sigma}_{ij \rightarrow k}$



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Parton Distribution Functions (PDFs)

Elementary cross-section (hard process)

Probability to find a parton 'i' inside beam particle '1' carrying a fraction  $x_1$  of the total momentum

(dependent on the hard process scale, Q<sup>2</sup>)



- Derivation from first principles does not yet exist. But its evolution, in Q<sup>2</sup>, can be described analytically.
  - Rely on parameterisations:
    - conjunction of experimental data and evolution equations
  - Once established,
     (proton, Pb, Au, ...)
     they are universal.



















$$\sigma_{ij\to k} = \int \mathrm{d}x_1 \int \mathrm{d}x_2 f_i^1(x_1) f_j^2(x_2) \,\hat{\sigma}_{ij\to k}$$





$$\frac{d^2 \sigma^h}{dy d^2 p_T} = \int dx_a dx_b f_a(x_a, \mu_f) f_b(x_b, \mu_f)$$
$$\frac{d\sigma_{ab \to c}(x_a p_a, x_b p_b, \mu, \mu_f, \mu'_f, p_T/z)}{d\hat{t}} D^h_c(z, \mu'_f)$$





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  - Matrix elements (few particle corrections but higher order)



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. Fortunately, these collinear contributions can be resummed by renormalization group

As most of the collinear emissions are well separated in scale from the probe,  $q_0$ , these **Two approaches to calculate additional radiation to the hard scattering:** a can be reinterpreted as modifying the hadron structure as opposed to corrections to  $C_a$ .

Its in a *Q*-dependence of the PDF that evolves the probe from the hard scale to lower Matrix elements (few particle corrections but higher order) Im scales (indicated by the red sub-diagram in Fig 2.9). For the change  $Q \rightarrow Q + \Delta Q$  the

al probability of an emission with energy fraction z and transverse momentum  $Q < p_{\perp} < P_{\perp} < P_{\perp}$  Parton shower (more particle corrections but LO and NLO only) is given by

• Evolution equation  $\frac{\alpha}{b} \frac{dp_{\perp}^2}{dp} P_{a} = \frac{\alpha}{b} \frac{\Delta Q}{dp} =$ 

 $_{\leftarrow b}(z)$  is the splitting function for parton of b splitting into type a, and can be computed diagrams  $\frac{\partial D_a^h(x, Q^2)}{\partial Q^2} + \frac{\alpha_s(Q^2)}{2\pi} \int_x^1 \frac{dz}{z} \sum d\hat{B}_z d\hat{B}_z tribult d\hat{B}_d (f_z^h a Q^2) a$  at momentum fraction ne from splittings of other partons at  $x' \stackrel{b}{=} x/z$ , and can be written as

$$\Delta f_a(x,Q) = \sum_b^{-\frac{\alpha_s(Q^2)}{2f_0^1}} \int_0^1 dz \sum_{\substack{dz \neq b \\ \pi}} \hat{P}_{Q \leftarrow b}(z) D_a^h(x,Q^2).$$

$$= \Delta \ln Q \sum_b^{-\frac{\alpha}{\pi}} \int_x^1 \frac{dz}{z} f_b(\frac{x}{z},Q^2) P_{a \leftarrow b}(z).$$

$$W'$$

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(-b)(z) is the splitting function for parton of b splitting into type a, and can be computed diagrams shown in Fig 2. 10: Changes in z be distributed in  $\mathcal{D}_{b}^{h}$  of pactor a at momentum fraction  $\partial Q^{2}$  for  $dz = 2\pi \int_{x}^{a} dz = 2\pi \int_{x}^{a} d$ 

$$\Delta f_a(x,Q) = \sum_b \int_0^{2\pi} dx' \int_0^0 dz' \frac{\Delta \varphi}{\pi} \frac{\Delta \varphi}{Q} P_{a\leftarrow b}(z) f_b(x',Q) \delta(x-z)$$
$$= \Delta \ln Q \sum \frac{\alpha}{\pi} \int_x^1 \frac{dz}{z} f_b(\frac{x}{z},Q^2) P_{a\leftarrow b}(z).$$

 $\dot{q}_0$ 

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$$\Delta f_{a}(x,Q) = \sum_{b}^{-\frac{\alpha_{s}(Q^{-})}{2}} \int_{0}^{1} dx \int_{0}^{0} \int_{0}^{q} dz \sum_{a \leftarrow b} \hat{P}_{a}(z) D_{a}^{h}(x,Q^{2}).$$

Splitting Function  $(SE)_{z} f_{b}(\frac{x}{z}, Q^{2})P_{a \leftarrow b}(z)$ . Probability of parton 'b' splits into parton 'a' with a fraction of energy z  $\dot{q}_0$ 

 $\dot{q}_{2}$ 

- Quantum mechanics = amplitudes (concept of randomness)
- Event generators = Monte Carlo techniques
  - Selection from a probability distribution function
  - Veto algorithm
  - ...

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Given a random number, R, what is  $t_1$ ? At  $t_1$ , it decays.

- Results into spray of partons/particles that will form jets;
  - Resulting pattern will contribute to the event structure (2, 3,... jet event)



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- Dominant  $2 \rightarrow 2$  QCD cross-sections are divergent for  $p_T \rightarrow 0$  but drop rapidly for large  $p_T$ .
  - Probability of multiple parton interactions is not negligible for ep, pp or AA collisions



#### Hadronization



- Mechanism that confines back quarks and gluons into hadrons;
- QCD perturbation theory, formulated in terms of quarks and gluons, is valid at short distances only
- At long distances, in the confinement regime, coloured pardons are transformed into hadrons, a process called hadronization (or fragmentation)
  - Fragmentation process not understood from first principles (rely on phenomenological models)
    - All of them rely on the color flow between the constituents

#### Hadronization

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- hard scattering
- (QED) initial/final state radiation
- partonic decays, e.g.  $t \rightarrow bW$
- parton shower evolution
- nonperturbative gluon splitting
- colour singlets
- colourless clusters
- cluster fission
- $\bullet \ cluster \rightarrow hadrons$
- hadronic decays

### Summary

- Result of an Event Generator:
  - 'Real' event as if could be observed by a perfect detector.
  - Output can be used now to interface to the detector simulation



## **More MC Event Generators**

• Typical hadronic event generator (PYTHIA) contains the subprocesses mentioned so far:

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Problem 2	

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Hard Scattering		
IS Shower	FS Shower	
PDFs	FFs	
Beam Remnants/M	PI Hadro	

## **More MC Event Generators**

 Typical hadronic event generator (PYTHIA) contains the subprocesses mentioned so far:

• Other type of event generators include:

Hard Scattering			
IS Shower		FS Shower	
PDFs	FFs		
Beam Remnants/MPI		Hadro	

- Cosmic Rays (for Extensive Air Showers)
- Heavy-ions (+ Nuclear initial-state, High multiplicity, soft processes, inmedium energy loss, Collective behavior of the medium)
- Multi-purpose parton event generators (BSM physics)

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# **Detector Simulation**



#### Patrícia Gonçalves (LIP/IST)