

# Probing the CP nature of the Higgs coupling in $t\bar{t}h$ events at the LHC

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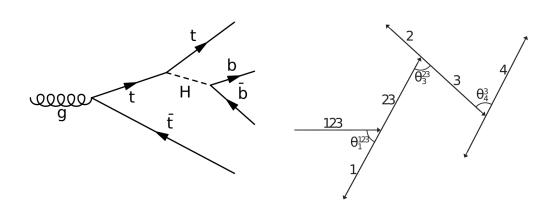
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#### Introduction

- New sources of CP violation are required to account for the observed baryon asymmetry in the Universe [1]
- Some BSM theories, namely 2HDMs [2], predict Higgs bosons which are not CP eigenstates, leading to CP-violating Higgs interactions
- The top quark is the most massive fermion and  $t\bar{t}h$  production will allow the first direct measurement of its Yukawa coupling
- Can angular observables in  $t\bar{t}h$  (dilepton) events at the LHC be used to measure the CP nature of the top quark Yukawa coupling?

#### Angular observables

With motivation from processes such as that on the left diagram, a generic chain of two-body decays (right diagram) was considered:
 (123) \( \to 1 + (23), (23) \( \to 2 + (3) \) and \( (3) \( \to 4 + 5 \)



- $\theta^X_Y$  is the angle between the direction of Y, in the rest frame of X, and the direction of X, in the rest frame of its parent system
- Three families of observables:  $f(\theta_1^{123})g(\theta_4^3)$ ,  $f(\theta_1^{123})g(\theta_3^2)$  and  $f(\theta_3^{23})g(\theta_4^3)$ , with  $f,g=\{\sin,\cos\}$
- A set of observables was obtained by replacing 1, 2 and 3 with all permutations of  $(t,\bar t,h)$  and replacing 4 with all the decay products of  $t,\bar t$  or h
- Observables proposed in previous works were also studied. Two of the most compelling ones were  $b_4$ , from [11], and  $\beta_{b\bar{b}}\Delta\theta^{\ell h}(\ell^+,\ell^-)$ , from [12]

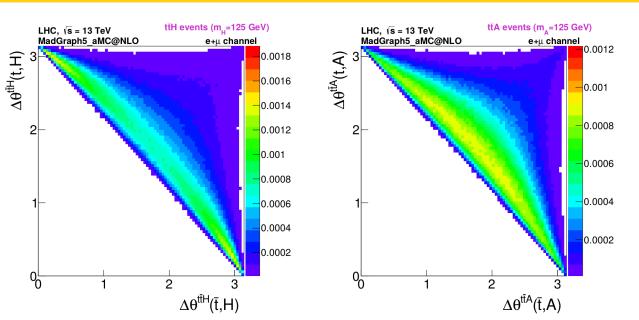
#### **Event Generation and Simulation**

- MADGRAPH5\_AMC@NLO [3] with NNPDF2.3 PDF sets [4]
- UFO Model for signal: HC\_NLO\_XO [5]
- Top Yukawa interaction was parametrised as

$$\mathcal{L} = \kappa y_t \, \bar{t} \left(\cos \alpha + i \gamma_5 \sin \alpha\right) t \, h$$

- $ullet \ tar th \ ext{signal:} \ m_h = 125 \ ext{GeV}, \coslpha \ ext{from -1 to 1}$
- $\cos \alpha = \pm 1 \Rightarrow$  SM, pure CP-even Higgs boson (h = H)
- $\cos \alpha = 0 \Rightarrow$  pure CP-odd Higgs boson (h = A)
- pp collisions with  $\sqrt{s} = 13 \text{ TeV}$
- $t\bar{t}h$  and  $t\bar{t}b\bar{b}$  were generated at NLO in QCD.  $t\bar{t}$  + jets (c- and light) and other SM backgrounds generated at LO with additional jets
- MADSPIN [6] used to decay h and the  $t\bar{t}$  system (dileptonically)
- Parton shower (PS) and hadronization: PYTHIA6 [7]
- Fast detector simulation with DELPHES [10]. The *b*-tagging efficiency goes up to 50% (20%) for true *b*-jets (*c*-jets).

# Angular distributions at generator+PS level

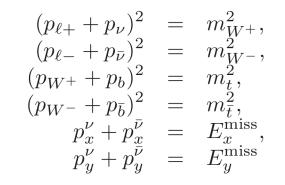


2D distributions of  $t\bar{t}h$  events, at generator+PS level, over the angle between the t and h directions (y-axis) and the angle between the  $\bar{t}$  and h directions (x-axis), all evaluated in the  $t\bar{t}h$  rest frame. The plot on the left shows  $t\bar{t}H$  events and the plot on the right shows  $t\bar{t}A$  events.

- CP nature of the Yukawa interaction significantly affects the kinematics of production
- In  $t\bar{t}H$ , more events are produced in which t (or  $\bar{t}$ ) recoils against the  $\bar{t}h$  (th) system. In  $t\bar{t}A$ , events are more evenly spread

#### Event selection and kinematic reconstruction

- Event selection:
- $\geq$ 4 jets and 2 opposite-sign leptons  $(e \text{ or } \mu)$ , all with  $p_T \geq 20 \text{ GeV}$  and  $|\eta| \leq 2.5$
- $|m_{\ell\ell} m_Z| > 10 \text{ GeV}$
- At least 3 b-tagged jets
- 16% (17%) of  $t\bar{t}H$  ( $t\bar{t}A$ ) signal events are accepted
- For the reconstruction of t,  $\bar{t}$  and h, a jet assignment is picked using a BDT:
- Training on  $t\bar{t}H$  events
- Object-pair variables used as input:  $\Delta R$ ,  $\Delta \Phi$ , labframe angles  $\Delta \theta$  and invariant masses
- Neutrinos are reconstructed by imposing W and top mass constraints and that the  $\vec{p_T}$  of  $\nu$  and  $\bar{\nu}$  add up to the  $\vec{E_T}^{\rm miss}$  in the detector
- If multiple solutions are found, a likelihood is maximised to choose the most adequate one:



$$L_{t\bar{t}h} \sim \frac{1}{p_{T\nu}p_{T\bar{\nu}}}P(p_{T\nu})P(p_{T\bar{\nu}})P(p_{T\bar{t}})P(p_{T\bar{t}})P(p_{Tt\bar{t}})P(m_t,m_{\bar{t}})P(m_h),$$

where P(X) is the probability density function (p.d.f.) at generator+PS level of the observable X, evaluated at the reconstructed value of X

#### **Asymmetries**

• Asymmetries associated to each of the observables  $x_Y$  under study were defined according to (see [13])

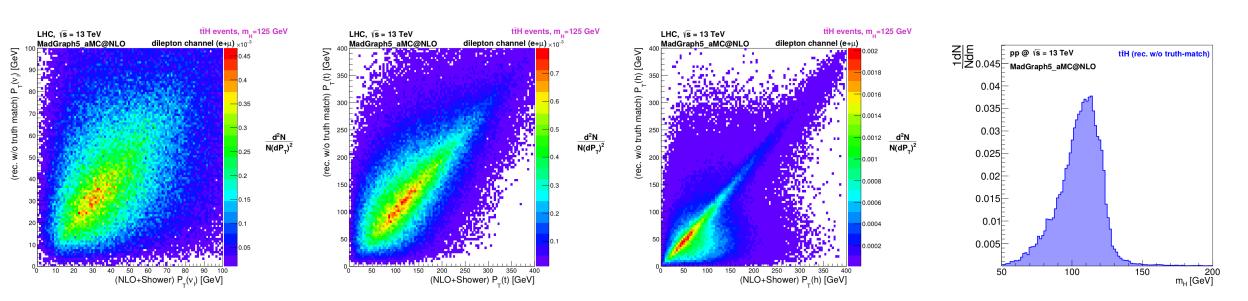
$$A^{Y} = \frac{N(x_{Y} > 0) - N(x_{Y} < 0)}{N(x_{Y} > 0) + N(x_{Y} < 0)}$$

Observable $x_Y$	$A^Y$ at gen.+PS level before cuts $t ar{t} H/t ar{t} A$		$A^Y$ after cuts and reconstruction $tar t H/tar t A \qquad tar t bar b$	
$\cos\left(\theta_{h}^{\bar{t}h}\right)\cos\left(\theta_{\ell-}^{h}\right)$	+0.37/+0.41	+0.17	+0.42/+0.39	+0.24
$\sin{( heta_h^{tar{t}h})}\sin{( heta_{ar{b}_{ar{ au}}}^{ar{t}})}$ (seq.)	+0.28/+0.33	-0.17	+0.25/+0.28	+0.03
$\sin{( heta_h^{tar{t}h})}\cos{( heta_{b_h}^{ar{t}^t})}$ (seq.)	-0.65/-0.77	-0.62	-0.78/-0.83	-0.76
$\sin{( heta_t^{tar{t}h})}\sin{( heta_{W+}^{h})}$ ( seq.)	-0.03/-0.46	-0.60	+0.17/-0.06	-0.04
$\sin{( heta_{ar{t}}^{tar{t}h})}\sin{( heta_{b_h}^h)}$ (seq.)	+0.25/-0.08	+0.07	+0.37/+0.16	+0.23
$\sin{( heta_h^{tar{t}h})}\sin{( heta_{ar{t}}^{tar{t}})}$	+0.16/+0.37	-0.21	+0.23/+0.31	+0.01
$b_4 = (p_t^z . p_{\bar{t}}^z) / ( \vec{p}_t  .  \vec{p}_{\bar{t}} )$ [11]	+0.35/-0.10	+0.33	+0.16/-0.17	+0.12

 $A^Y$  values in  $t\bar{t}H$ ,  $t\bar{t}A$  and  $t\bar{t}b\bar{b}$  events for a selection of observables with significant signal/background or  $t\bar{t}H/t\bar{t}A$  separation, at generator+PS level (before selection cuts) and after full kinematic reconstruction (including cuts).

# Reconstruction performance on $t\bar{t}H$ events

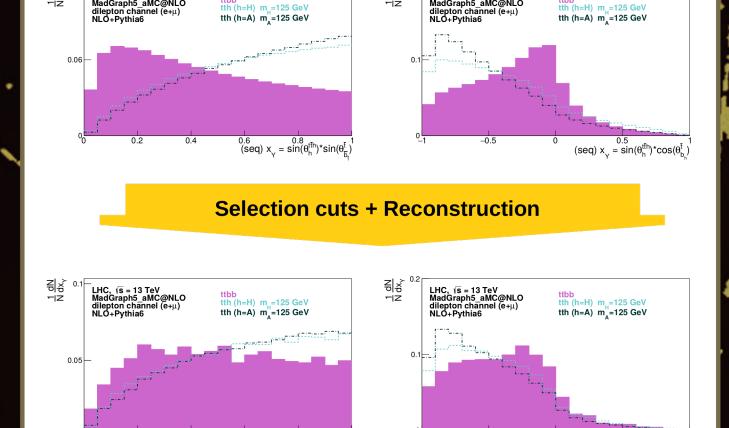
Solution is found in 62% (61%) of selected  $t\bar{t}H$  ( $t\bar{t}A$ ) signal events. In 31% (34%) of the  $t\bar{t}H$  ( $t\bar{t}A$ ) signal events, the reconstruction results in the same jet assignment as a  $\Delta R$  truth-matching



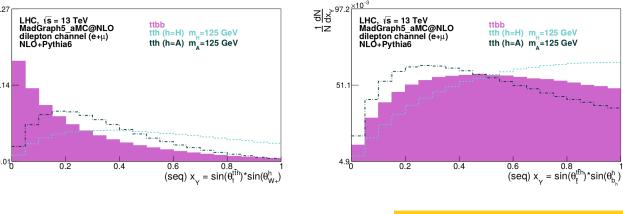
 $p_T$  distributions of  $\nu$  (left), t (middle-left) and H (middle-right) after  $t\bar{t}H$  reconstruction. On the x-axis is the generator+PS level  $p_T$  and on the y-axis is the reco-level  $p_T$ . On the right, reconstructed  $m_H$ .

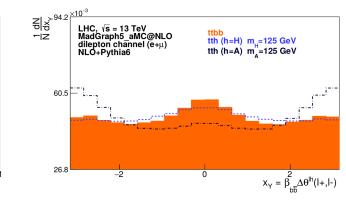
• Excess of high- $p_T$  neutrinos o reconstruction uses their free momenta to compensate for radiation losses

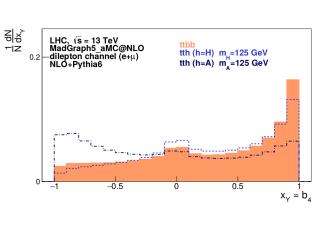
# Signal/Background



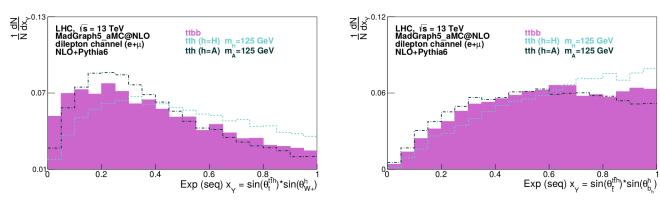
# CP-even/CP-odd

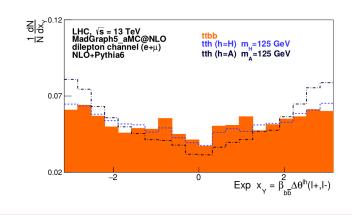


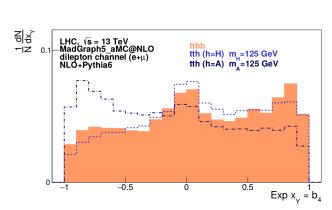




#### Selection cuts + Reconstruction



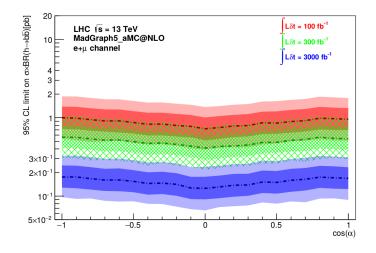


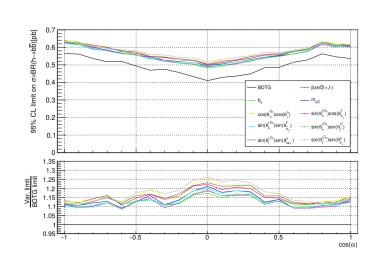


# Expected limits at 95% CL

Exp (seq)  $x_v = \sin(\theta_h^{t\bar{t}h}) * \cos(\theta_b^{\bar{t}})$ 

- A BDT was trained for signal/background classification for each value of  $\cos \alpha$ . Input variables include the observables studied and variables used in the ATLAS 8 TeV search [14].
- **Expected limits at 95% CL were obtained using the BDT as a discriminant.** Limits were also extracted using instead each single input variable





Limits on  $\sigma \times BR(h \to b\bar{b})$  obtained with the BDT discriminant for integrated luminosites of 100, 300 and 3000 fb<sup>-1</sup> (left). Comparison, for 300 fb<sup>-1</sup>, of limits obtained using each angular observable as discriminant (right).

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### Conclusions

- Angular observables in  $t\bar{t}h$  events keep information about the CP nature of the top-Higgs coupling, even after cuts, detector simulation and kinematic reconstruction
- Most powerful observables require  $tar{t}h$  reconstruction
- Search is not too sensitive to the CP-phase  $\alpha$

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