

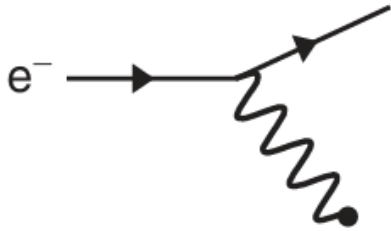
High Energy Physics

Hadron Structure and DIS

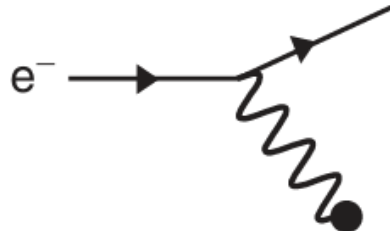
Luís Correia 2019230540

Scattering of an electron off a proton

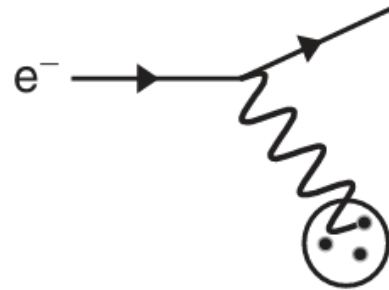
$$e^- p \rightarrow e^- p$$



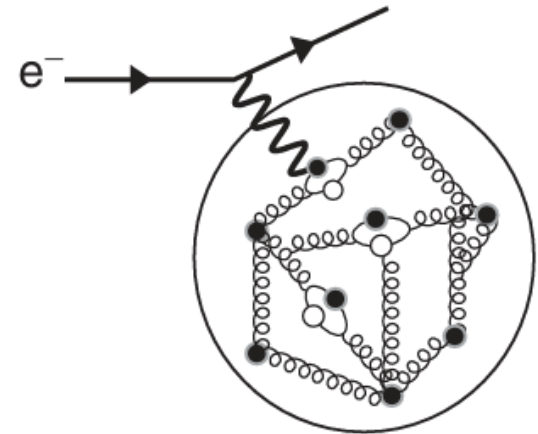
$$\lambda \gg r_p$$



$$\lambda \sim r_p$$

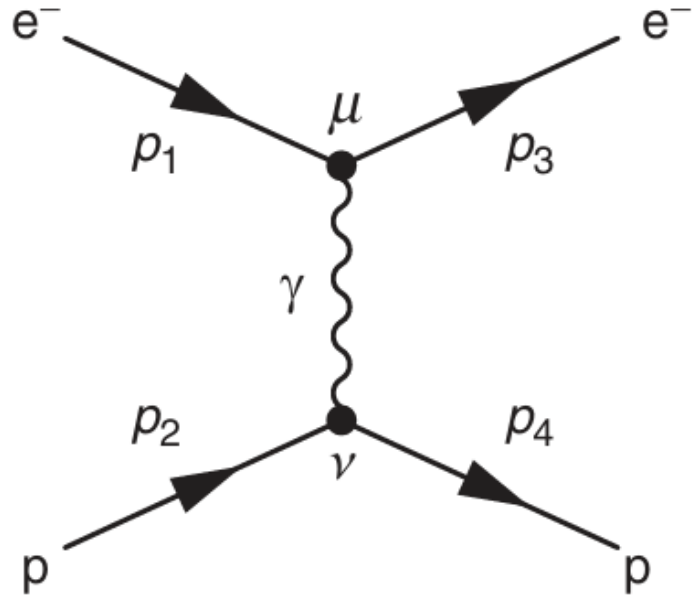


$$\lambda < r_p$$



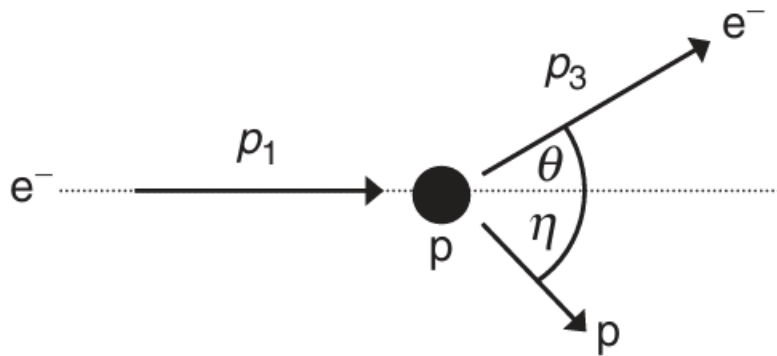
$$\lambda \ll r_p$$

Rutherford and Mott scattering



$$\mathcal{M}_{fi} = \frac{Q_q e^2}{q^2} [\bar{u}(p_3) \gamma^\mu u(p_1)] g_{\mu\nu} [\bar{u}(p_4) \gamma^\nu u(p_2)]$$

$$= \frac{e^2}{q^2} j_e \cdot j_p$$



Rutherford and Mott scattering

electron current

$$u_{\uparrow}(p_1) = N_e \begin{pmatrix} 1 \\ 0 \\ \kappa \\ 0 \end{pmatrix}, \quad u_{\downarrow}(p_1) = N_e \begin{pmatrix} 0 \\ 1 \\ 0 \\ -\kappa \end{pmatrix} \quad \text{and} \quad u_{\uparrow}(p_3) = N_e \begin{pmatrix} c \\ s \\ \kappa c \\ \kappa s \end{pmatrix}, \quad u_{\downarrow}(p_3) = N_e \begin{pmatrix} -s \\ c \\ \kappa s \\ -\kappa c \end{pmatrix}.$$

$$j_{e\uparrow\uparrow} = \bar{u}_{\uparrow}(p_3) \gamma^{\mu} u_{\uparrow}(p_1) = (E + m_e) \left[(\kappa^2 + 1)c, 2\kappa s, +2i\kappa s, 2\kappa c \right],$$

$$j_{e\downarrow\downarrow} = \bar{u}_{\downarrow}(p_3) \gamma^{\mu} u_{\downarrow}(p_1) = (E + m_e) \left[(\kappa^2 + 1)c, 2\kappa s, -2i\kappa s, 2\kappa c \right],$$

$$j_{e\downarrow\uparrow} = \bar{u}_{\uparrow}(p_3) \gamma^{\mu} u_{\downarrow}(p_1) = (E + m_e) \left[(1 - \kappa^2)s, 0, 0, 0 \right],$$

$$j_{e\uparrow\downarrow} = \bar{u}_{\downarrow}(p_3) \gamma^{\mu} u_{\uparrow}(p_1) = (E + m_e) \left[(\kappa^2 - 1)s, 0, 0, 0 \right].$$

$$N_e = \sqrt{E + m_e}$$

$$s = \sin(\theta/2)$$

$$c = \cos(\theta/2)$$

$$\kappa = \frac{p}{E + m_e} \equiv \frac{\beta_e \gamma_e}{\gamma_e + 1}$$

Rutherford and Mott scattering

proton current

$$u_{\uparrow}(p_2) = \sqrt{2m_p} \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad u_{\downarrow}(p_2) = \sqrt{2m_p} \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \quad u_{\uparrow}(p_4) \approx \sqrt{2m_p} \begin{pmatrix} c_{\eta} \\ -s_{\eta} \\ 0 \\ 0 \end{pmatrix} \quad u_{\downarrow}(p_4) \approx \sqrt{2m_p} \begin{pmatrix} -s_{\eta} \\ -c_{\eta} \\ 0 \\ 0 \end{pmatrix}$$

$$j_{p\uparrow\uparrow} = -j_{p\downarrow\downarrow} = 2m_p [c_{\eta}, 0, 0, 0]$$

$$c_{\eta} = \cos(\eta/2)$$

$$s_{\eta} = \sin(\eta/2)$$

$$j_{p\uparrow\downarrow} = j_{p\downarrow\uparrow} = -2m_p [s_{\eta}, 0, 0, 0]$$

Rutherford and Mott scattering

matrix element

$$\begin{aligned}\langle |\mathcal{M}_{fi}^2| \rangle &= \frac{1}{4} \sum |\mathcal{M}_{fi}^2| \\ &= \frac{1}{4} \frac{e^4}{q^4} \times 4m_p^2 (E + m_e)^2 \cdot [c_\eta^2 + s_\eta^2] \cdot [4(1 + \kappa^2)^2 c^2 + 4(1 - \kappa^2)^2 s^2] \\ &= \frac{4m_p^2 m_e^2 e^4 (\gamma_e + 1)^2}{q^4} [(1 - \kappa^2)^2 + 4\kappa^2 c^2] \\ &= \frac{16m_p^2 m_e^2 e^4}{q^4} \left[1 + \beta_e^2 \gamma_e^2 \cos^2 \frac{\theta}{2} \right]\end{aligned}$$

$$\langle |\mathcal{M}_{fi}^2| \rangle = \frac{m_p^2 m_e^2 e^4}{p^4 \sin^4(\theta/2)} \left[1 + \beta_e^2 \gamma_e^2 \cos^2 \frac{\theta}{2} \right]$$

Rutherford and Mott scattering

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2} \left(\frac{1}{m_p + E_1 - E_1 \cos \theta} \right)^2 \langle |\mathcal{M}_{fi}|^2 \rangle$$

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Rutherford}} = \frac{\alpha^2}{16E_K^2 \sin^4(\theta/2)} \quad \beta_e \gamma_e \ll 1 \quad E_1 \sim m_e \ll m_p$$

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} = \frac{\alpha^2}{4E^2 \sin^4(\theta/2)} \cos^2 \frac{\theta}{2} \quad \beta_e \gamma_e \gg 1$$

Form Factors

$$V(\mathbf{r}) = \int \frac{Q\rho(\mathbf{r}')}{4\pi|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}'$$

$$\begin{aligned}\mathcal{M}_{fi} &= \langle \psi_f | V(\mathbf{r}) | \psi_i \rangle = \int e^{-i\mathbf{p}_3 \cdot \mathbf{r}} V(\mathbf{r}) e^{i\mathbf{p}_1 \cdot \mathbf{r}} d^3\mathbf{r} = \int \int e^{i\mathbf{q} \cdot \mathbf{r}} \frac{Q\rho(\mathbf{r}')}{4\pi|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}' d^3\mathbf{r} \\ &= \int \int e^{i\mathbf{q} \cdot (\mathbf{r} - \mathbf{r}')} e^{i\mathbf{q} \cdot \mathbf{r}'} \frac{Q\rho(\mathbf{r}')}{4\pi|\mathbf{r} - \mathbf{r}'|} d^3\mathbf{r}' d^3\mathbf{r} = \int e^{i\mathbf{q} \cdot \mathbf{R}} \frac{Q}{4\pi|\mathbf{R}|} d^3\mathbf{R} \int \rho(\mathbf{r}') e^{i\mathbf{q} \cdot \mathbf{r}'} d^3\mathbf{r}'\end{aligned}$$

$$F(\mathbf{q}^2) = \int \rho(\mathbf{r}) e^{i\mathbf{q} \cdot \mathbf{r}} d^3\mathbf{r}$$

$$\left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \rightarrow \frac{\alpha^2}{4E^2 \sin^4(\theta/2)} \cos^2 \left(\frac{\theta}{2} \right) |F(\mathbf{q}^2)|^2$$

Rosenbluth formula

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1^2 \sin^4(\theta/2)} \frac{E_3}{E_1} \left(\frac{G_E^2 + \tau G_M^2}{(1 + \tau)} \cos^2 \frac{\theta}{2} + 2\tau G_M^2 \sin^2 \frac{\theta}{2} \right)$$

$$Q^2 \left(1 + \frac{Q^2}{4m_p^2} \right) = \mathbf{q}^2 \qquad \tau = \frac{Q^2}{4m_p^2}$$

In the limit: $Q^2 \ll 4m_p^2$

$$G_E(Q^2) \approx G_E(\mathbf{q}^2) = \int e^{i\mathbf{q}\cdot\mathbf{r}} \rho(\mathbf{r}) d^3\mathbf{r}$$

$$G_M(Q^2) \approx G_M(\mathbf{q}^2) = \int e^{i\mathbf{q}\cdot\mathbf{r}} \mu(\mathbf{r}) d^3\mathbf{r}$$

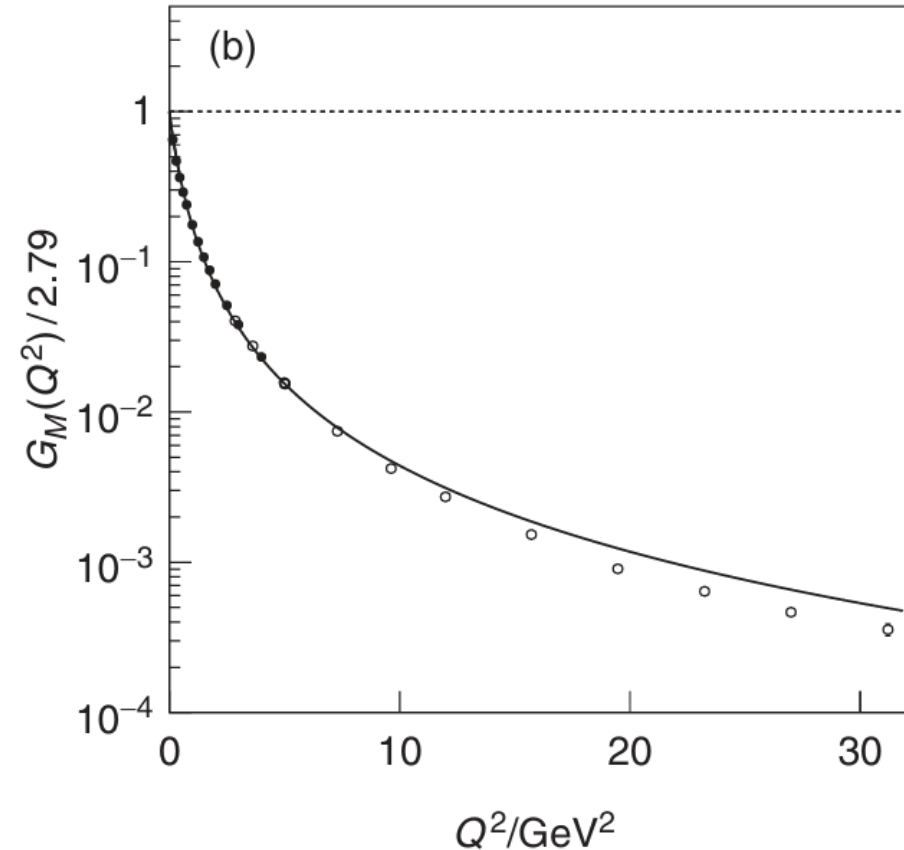
Rosenbluth formula

High Q

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{elastic}} \sim \frac{\alpha^2}{4E_1^2 \sin^4(\theta/2)} \frac{E_3}{E_1} \left[\frac{Q^2}{2m_p^2} G_M^2 \sin^2 \frac{\theta}{2} \right]$$

$$G_M(q^2) \propto q^{-4}$$

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{elastic}} \propto \frac{1}{Q^6} \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}}$$



Deep Inelastic Scattering

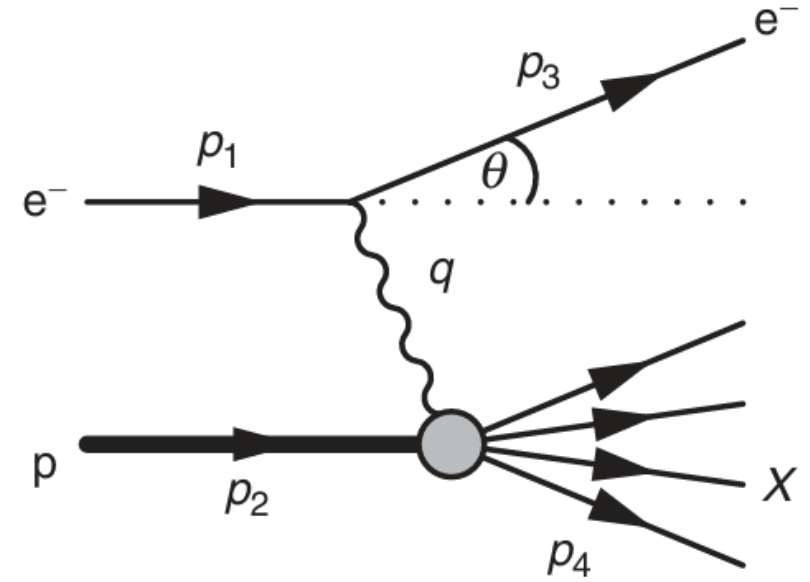
Kinematic variables: x, y, ν, Q^2

$$Q^2 \equiv -q^2$$

$$x \equiv \frac{Q^2}{2p_2 \cdot q} \quad 0 < x < 1$$

$$y \equiv \frac{p_2 \cdot q}{p_2 \cdot p_1} \quad 0 < y < 1$$

$$\nu \equiv \frac{p_2 \cdot q}{M}$$



$$W^2 = p_4^2 = (p_2 + q)^2$$

$$y = \left(\frac{2m_p}{s - m_p^2} \right) \nu$$

$$Q^2 = (s - m_p^2)xy$$

Deep Inelastic Scattering

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4E_1^2 \sin^4(\theta/2)} \frac{E_3}{E_1} \left(\frac{G_E^2 + \tau G_M^2}{(1 + \tau)} \cos^2 \frac{\theta}{2} + 2\tau G_M^2 \sin^2 \frac{\theta}{2} \right)$$

$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[\frac{G_E^2 + \tau G_M^2}{(1 + \tau)} \left(1 - y - \frac{m_p^2 y^2}{Q^2} \right) + \frac{1}{2} y^2 G_M^2 \right]$$

$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[\left(1 - y - \frac{m_p^2 y^2}{Q^2} \right) f_2(Q^2) + \frac{1}{2} y^2 f_1(Q^2) \right]$$

$$\frac{d^2\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[\left(1 - y - \frac{m_p^2 y^2}{Q^2} \right) \frac{F_2(x, Q^2)}{x} + y^2 F_1(x, Q^2) \right]$$

Deep Inelastic Scattering

Bjorken scaling and the Callan-Gross relation

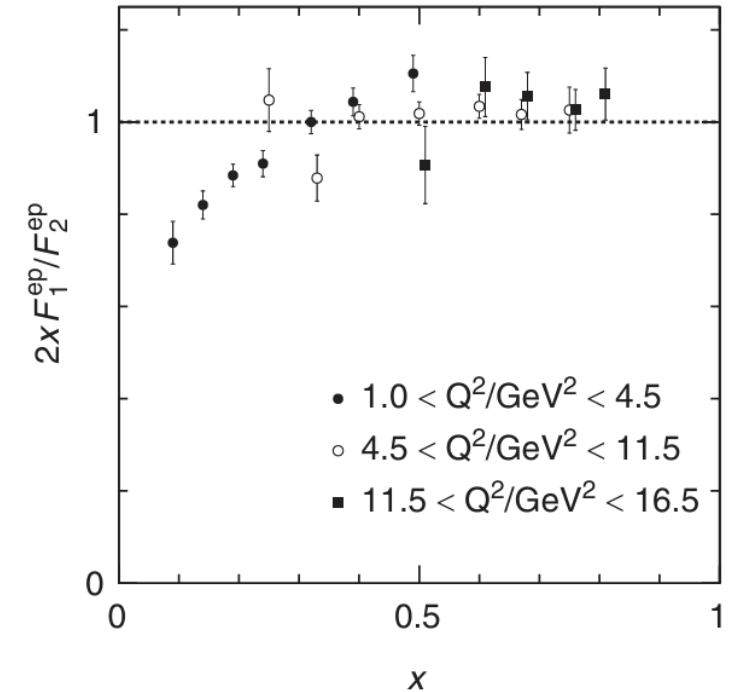
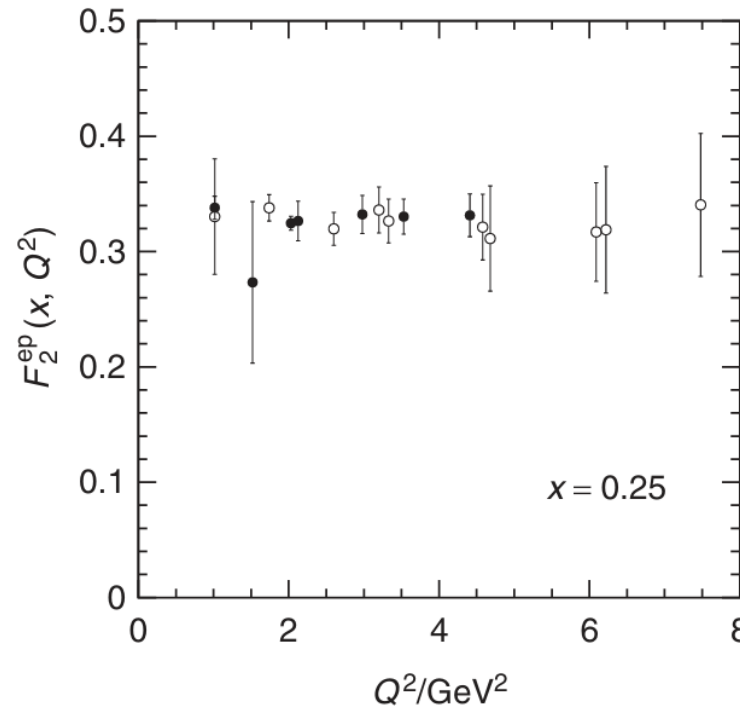
For deep inelastic scattering: $Q^2 \gg m_p^2 y^2$

$$\frac{d^2\sigma}{dx dQ^2} \approx \frac{4\pi\alpha^2}{Q^4} \left[(1-y) \frac{F_2(x, Q^2)}{x} + y^2 F_1(x, Q^2) \right]$$

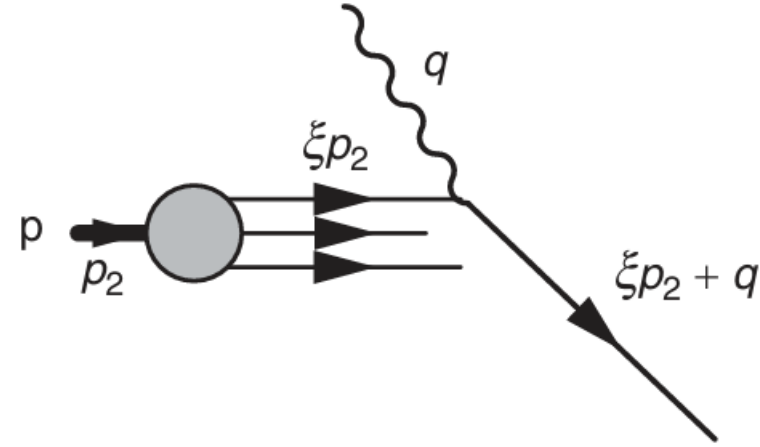
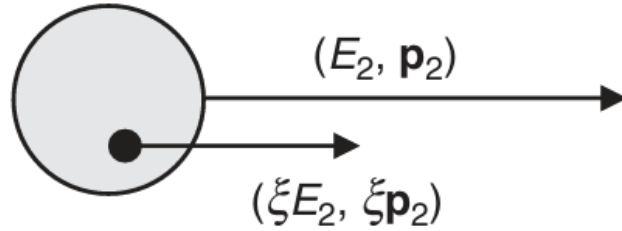
$$F_1(x, Q^2) \rightarrow F_1(x)$$

$$F_2(x, Q^2) \rightarrow F_2(x)$$

$$F_2(x) = 2xF_1(x)$$



Parton model



$$(\xi p_2 + q)^2 = \xi^2 p_2^2 + 2\xi p_2 \cdot q + q^2 = m_q^2$$

$$\xi = \frac{-q^2}{2p_2 \cdot q} = \frac{Q^2}{2p_2 \cdot q} \equiv x$$

In terms of the proton momentum

$$s = (p_1 + p_2)^2 \approx 2p_1 \cdot p_2$$

$$y = \frac{p_2 \cdot q}{p_2 \cdot p_1}$$

$$x = \frac{Q^2}{2p_2 \cdot q}$$

For the quark

$$s_q = x s, \quad y_q = y \quad \text{and} \quad x_q = 1$$

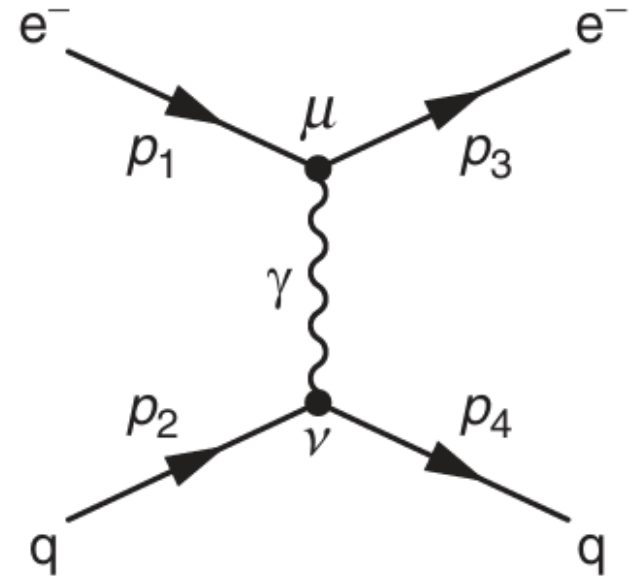
Parton model

$$\frac{d\sigma}{dq^2} = \frac{2\pi\alpha^2 Q_q^2}{q^4} \left[1 + \left(1 + \frac{q^2}{s} \right)^2 \right]$$

$$q^2 = -Q^2 = -(s_q - m_q^2)x_q y_q$$

$$\frac{q^2}{s_q} = -x_q y_q = -y$$

$$\frac{d\sigma}{dQ^2} = \frac{4\pi\alpha^2 Q_q^2}{Q^4} \left[(1 - y) + \frac{y^2}{2} \right]$$



Parton model

For a specific flavour we have

$$\frac{d^2\sigma}{dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[(1-y) + \frac{y^2}{2} \right] \times Q_i^2 q_i^p(x) \delta x$$

Summing over all quark flavours

$$\frac{d^2\sigma^{\text{ep}}}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4} \left[(1-y) + \frac{y^2}{2} \right] \sum_i Q_i^2 q_i^p(x)$$

$$F_2^{\text{ep}}(x, Q^2) = 2xF_1^{\text{ep}}(x, Q^2) = x \sum_i Q_i^2 q_i^p(x)$$