

# Quantum materials, lattice field theory, chiral anomaly

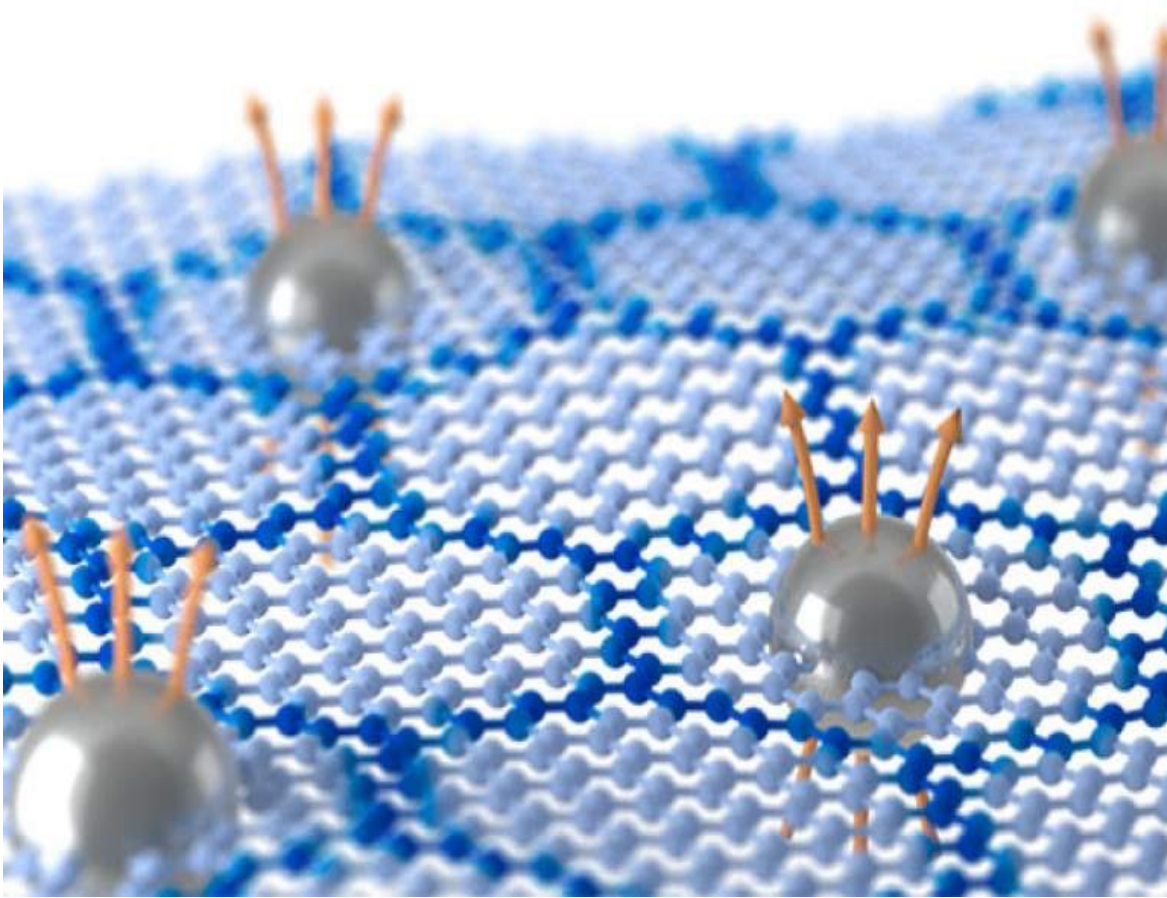
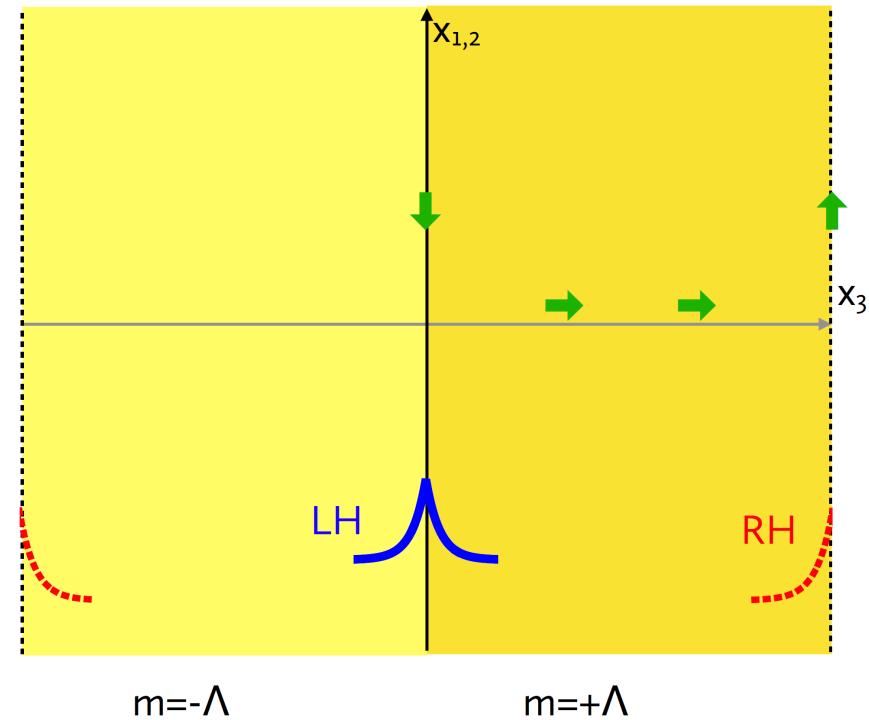


Figure courtesy: Nationalmaglab, Phys.org, David Kaplan

*Phys.Rev.Lett.* 132 (2024) 14, 141603,

*Phys.Rev.Lett.* 132 (2024) 14, 141604

arxiv: [2412.02024](https://arxiv.org/abs/2412.02024)



Srimoyee Sen,  
Iowa State University

work with David Kaplan, University of Washington

“Weyl and Dirac Semimetals as a Laboratory for High-Energy Physics” 2025,  
University of Minho, Braga,

# Plan of the talk

- Thinking about ideas at the intersection of condensed matter physics (topological phases, weyl and Dirac semimetals), nuclear and particle physics profitable for all the fields involved.
- Will discuss a striking example from the past connecting lattice field theory, chiral anomalies and quantum Hall effect
- Discuss problems where a similar line of thinking may help make progress.

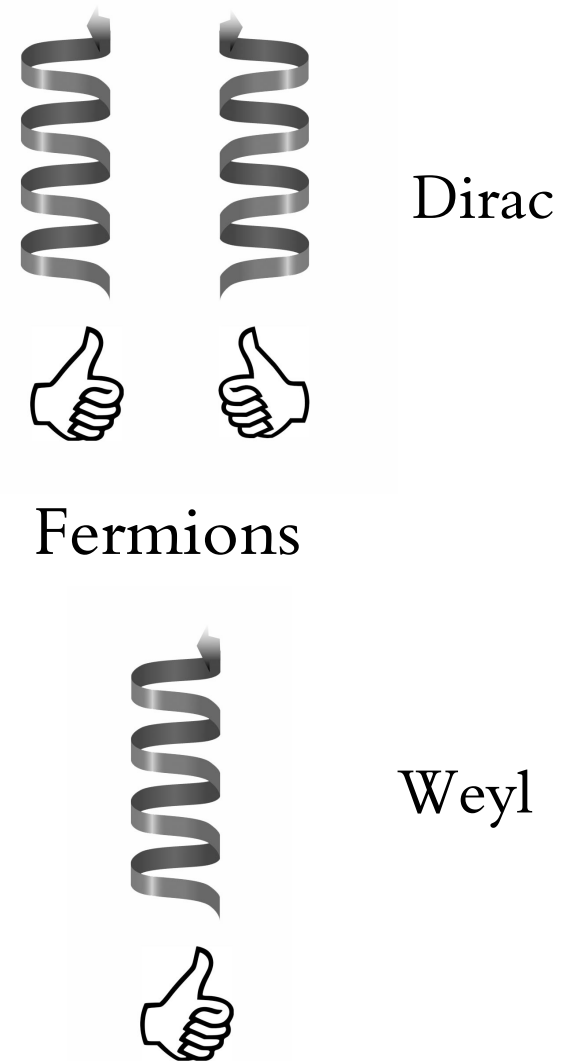
# Field Theory ingredients: fundamental forces and fermions (massless)

**Standard Model of Elementary Particles**

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
QUARKS	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> higgs
	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
LEPTONS	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.433 \text{ GeV}/c^2$	
	0	0	0	$\pm 1$	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	

**SCALAR BOSONS**

**GAUGE BOSONS VECTOR BOSONS**



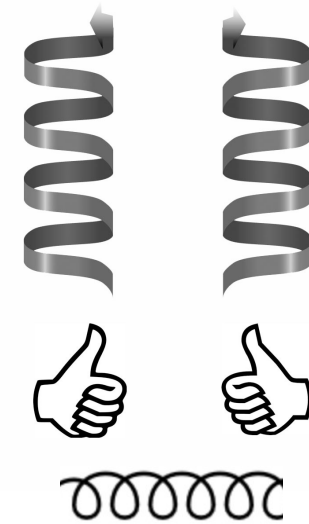
# Field Theory ingredients: fundamental forces and fermions (massless)

**Standard Model of Elementary Particles**

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
<b>QUARKS</b>	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> higgs
	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
<b>LEPTONS</b>	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$	$\approx 91.19 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.433 \text{ GeV}/c^2$	
	0	0	0	$\pm 1$	
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	

**SCALAR BOSONS**

**GAUGE BOSONS  
VECTOR BOSONS**



Dirac

Quantum chromodynamics:

Gluons talk to fermions of both chirality **equally**

**Vector gauge theory**

# Field Theory ingredients: fundamental forces and fermions (massless/gapless)

**Standard Model of Elementary Particles**

	three generations of matter (fermions)			interactions / force carriers (bosons)	
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> higgs
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon	
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson	
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson	

**QUARKS** (vertical label on the left of the quark section)

**LEPTONS** (vertical label on the left of the lepton section)

**GAUGE BOSONS VECTOR BOSONS** (vertical label on the right of the gauge boson section)

**SCALAR BOSONS** (vertical label on the right of the scalar boson section)



Weyl

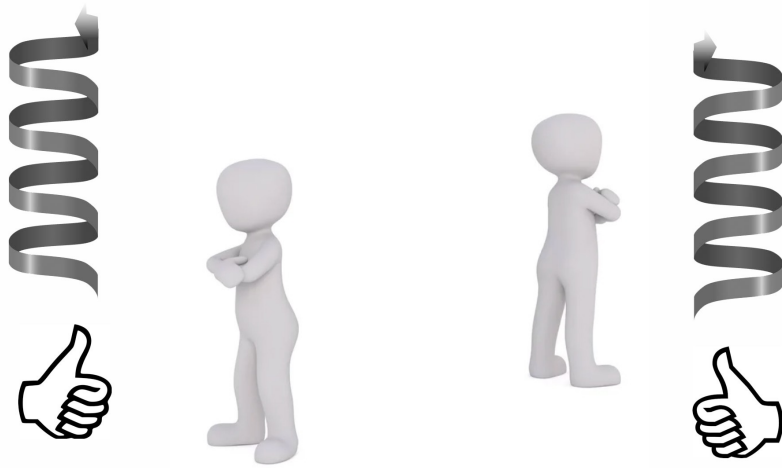
Standard model, weak interactions:

Talks only to one chirality

Chiral gauge theory

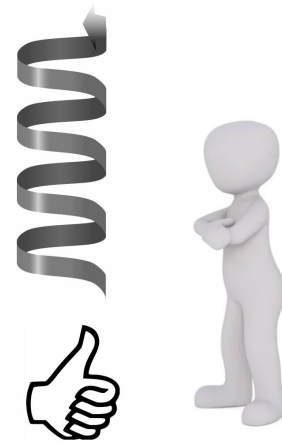
# QCD vs weak interactions

$m = 0 \Rightarrow$  chiral symmetry



QCD: good to have chiral symmetry, i.e. exact masslessness, can't allow the chiralities to talk.

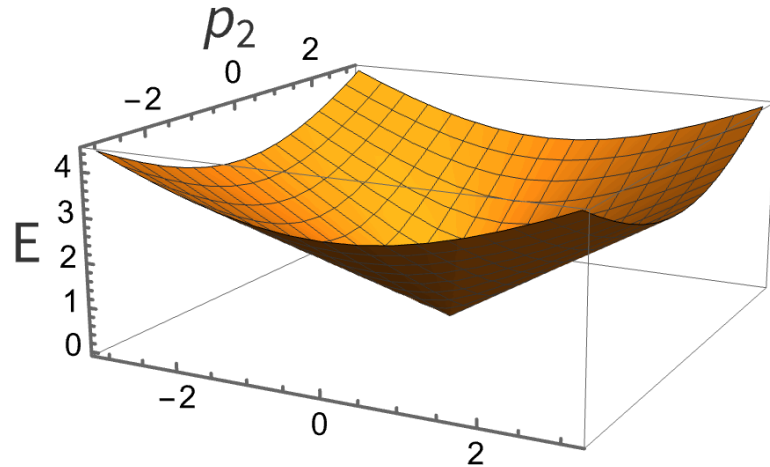
Solved on the lattice.



SM: Chiral symmetry absolutely essential, unpair the chirality

“Chiral gauge theories”, yet to be formulated on the lattice.

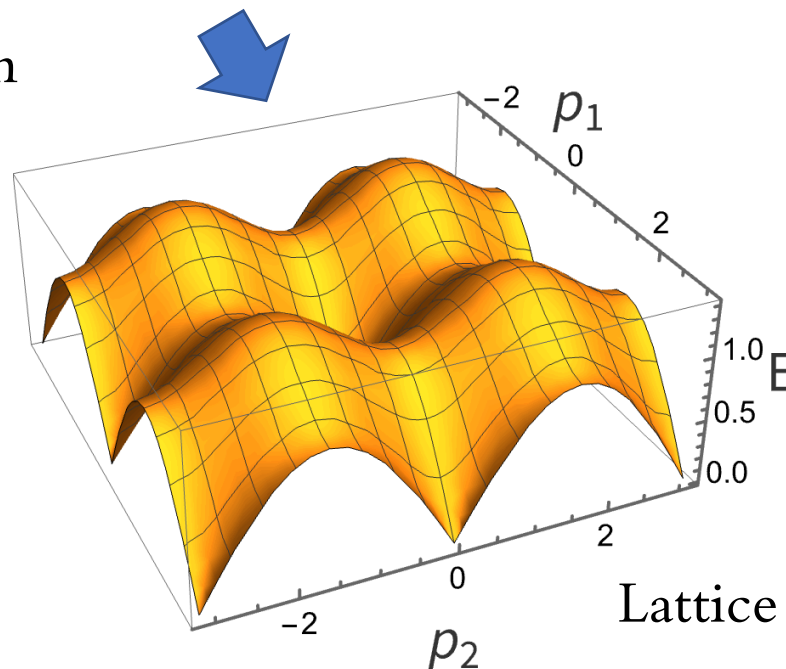
# On the lattice (massless Dirac fermion problem for QCD)



Continuum dispersion

We want to formulate fermion theories on the lattice.

We want to keep the fermions practically gapless (to simulate light quarks).

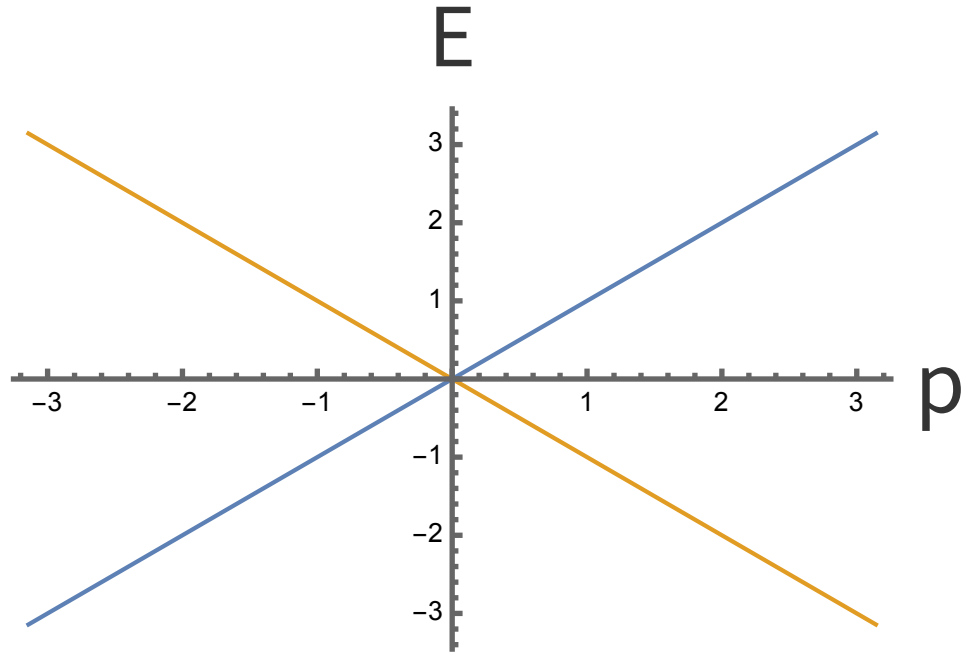


Lattice Brillouin zone

This combination i.e. “discrete space-time + gapless fermion”: extremely difficult to engineer.

Nielsen-Ninomiya No-Go

# Why chiral symmetry is hard: Dispersion (1 spatial dimension)



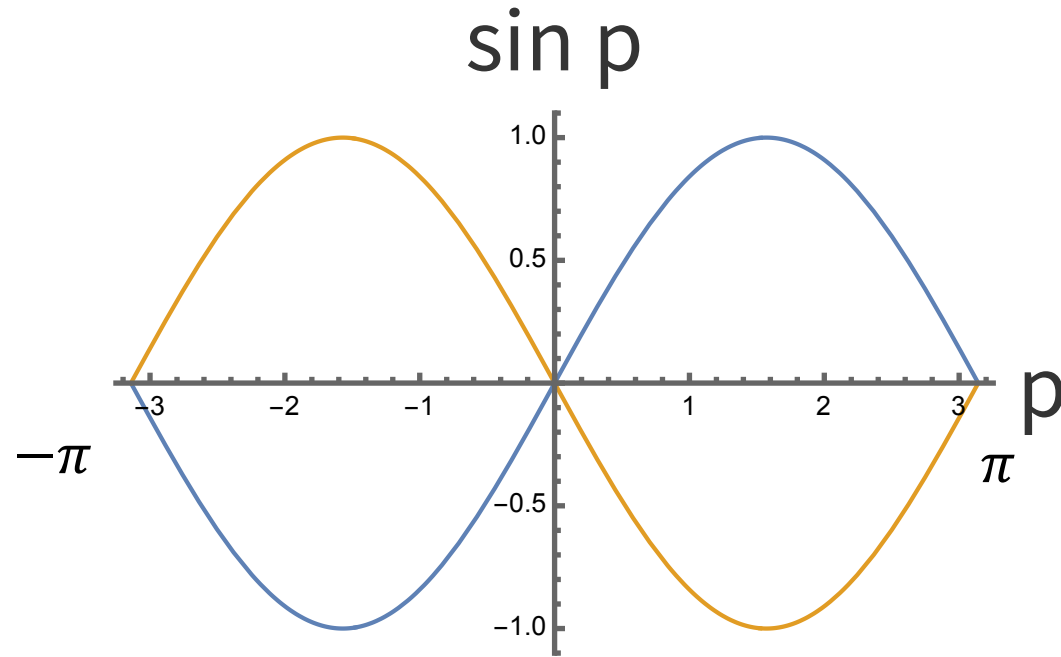
Continuum dispersion for a  
Dirac Hamiltonian

The no-go is better visualized  
using dispersion relation in  
Minkowski space-time (time  
continuous).

Hamiltonian formulation.

$$E = \pm p$$

# Brilluoin zones (Dirac)



Two Dirac fermions

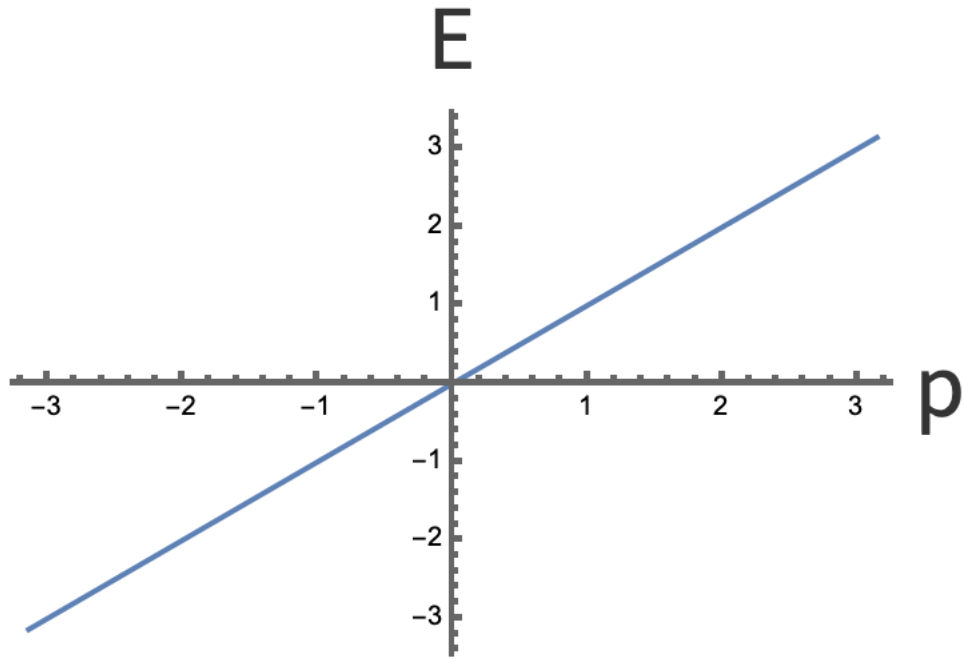
Lattice in space.

Time not discretized.

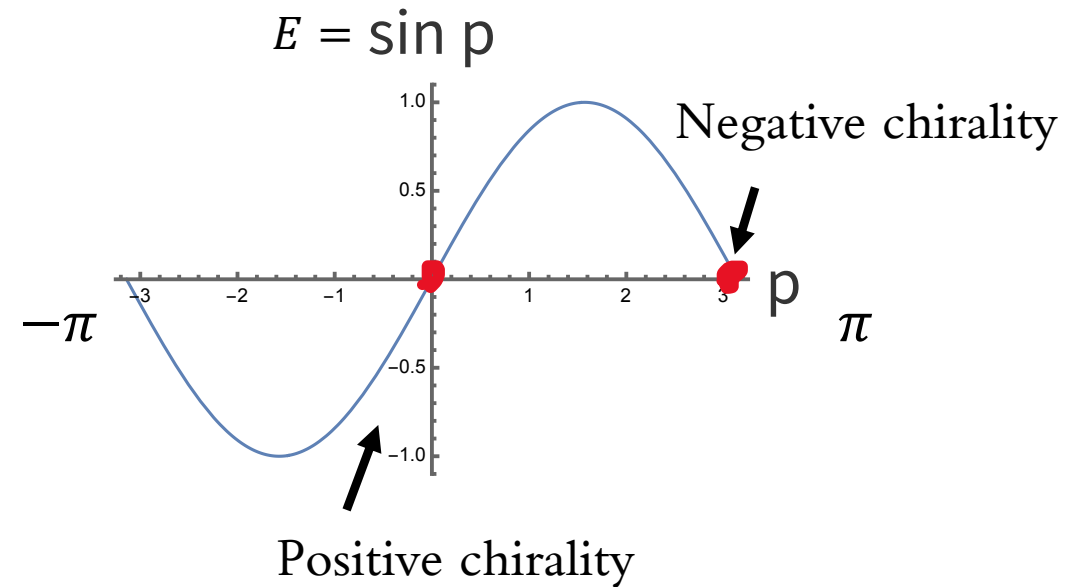
Solving the naively discretized  
Dirac Hamiltonian with  
eigenvalues  $\pm \sin p$

$$E = \pm \sin p$$

# Weyl fermion



Continuum



Lattice

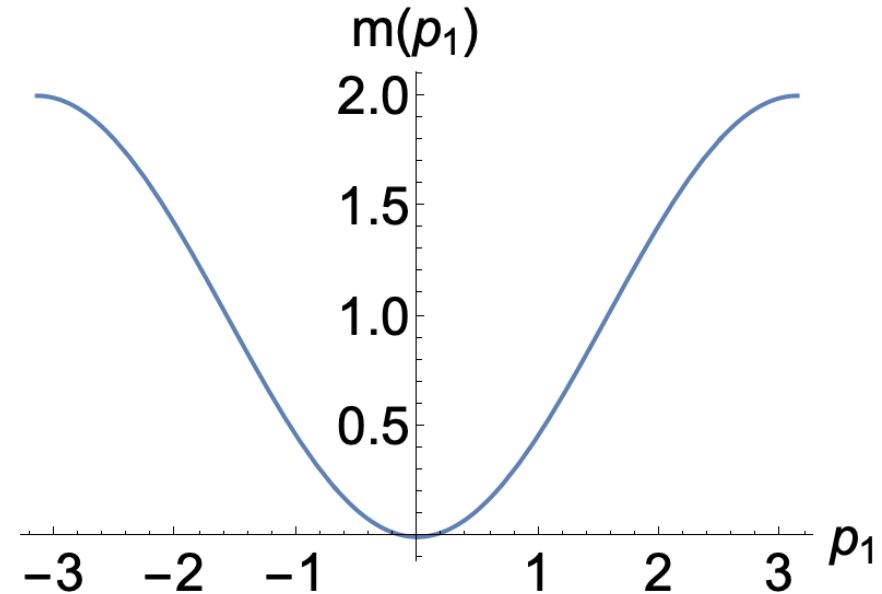
Even number of zero crossing of periodic functions

# Quick fix for Dirac (leads to other problems)

Add a momentum dependent mass:

$$(p_0)^2 - (\sin p_1)^2 - m(p_1)^2 = 0$$

$$(p_0 \sim \pm p_1 \sim 0) , \cancel{(p_0 \sim \pi \pm p_1 \sim 0)}$$



Extra species gone. But this causes other problems.

Chiral symmetry is killed.

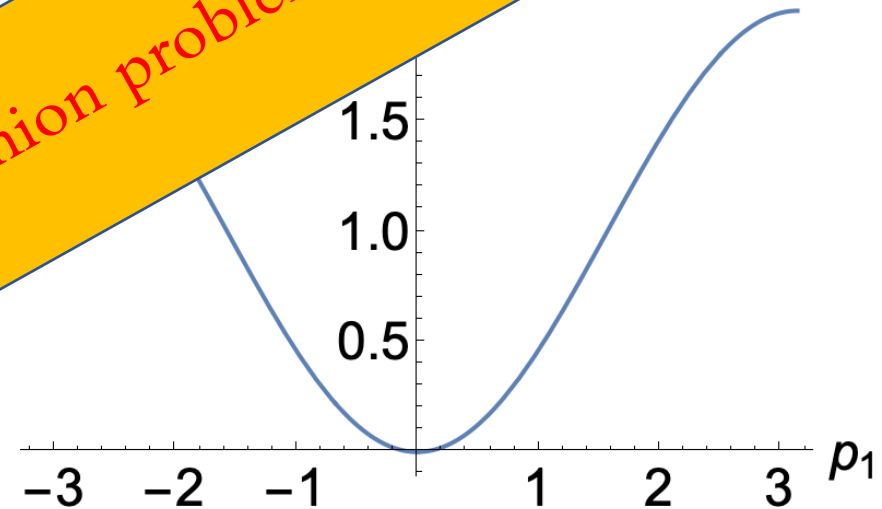
# Quick fix for Dirac (leads to other problems)

Add a momentum dependent mass:

$$(p_0)^2 - (\sin p_1)^2 - m(p_1)^2 = 0$$

$$(p_0 \sim \pm p_1 \sim 0)$$

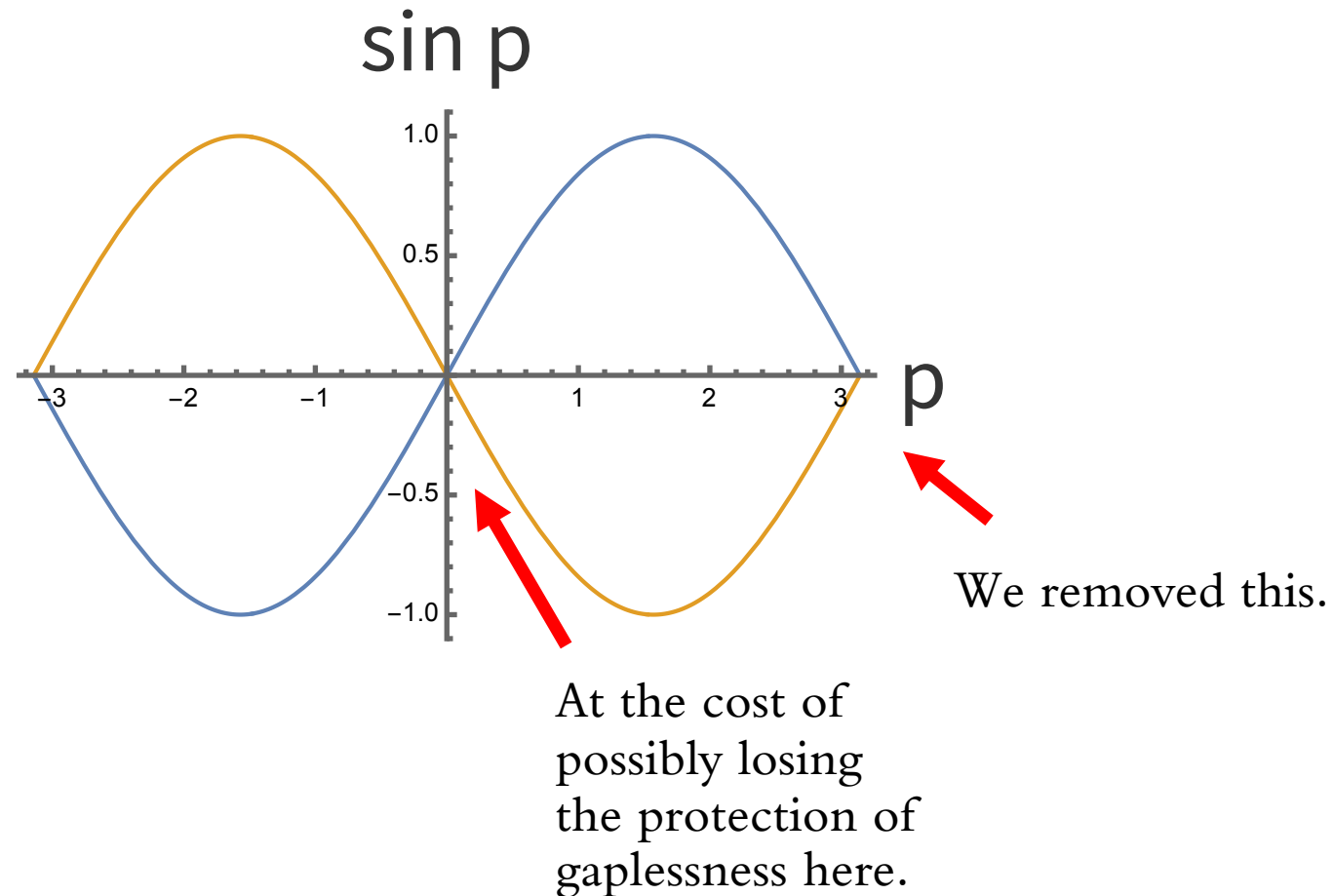
This is the chiral fermion problem



Extra species appear. But this causes other problems.

Chiral symmetry is killed.

# Let's take another look at momentum dependent mass



# What if you could physically separate the two chiralities?

By going to one higher dimension?

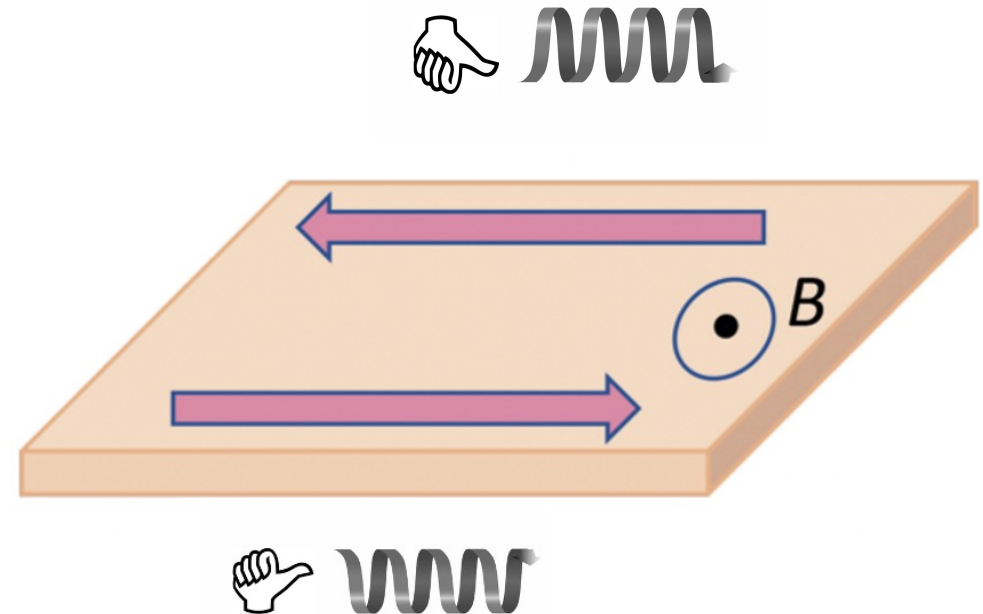
Where does this happen?

# What if you could physically separate the two chiralities?

Quantum Hall Effect

Relativistic Analog: Dirac fermion  
in  $2 + 1$  dimensions

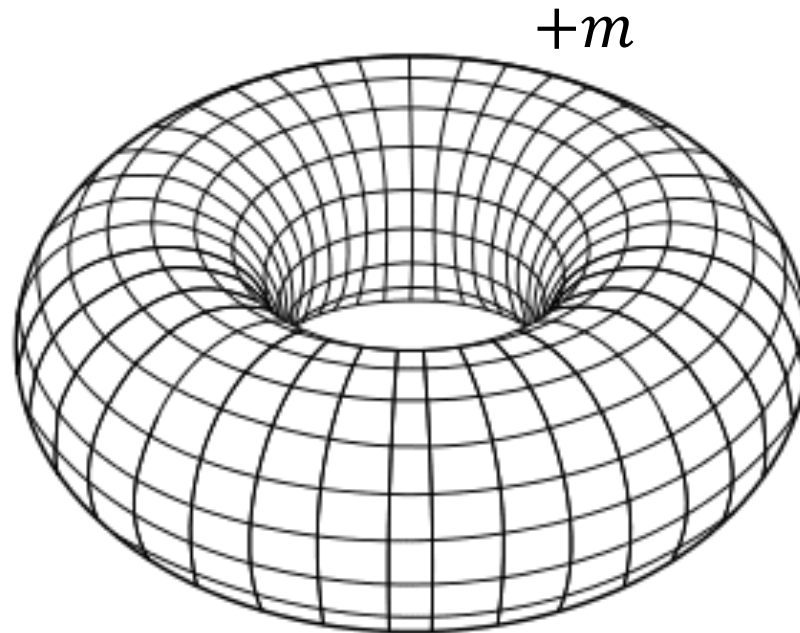
Anomalous Hall effect.



# How do you achieve this in a relativistic theory?

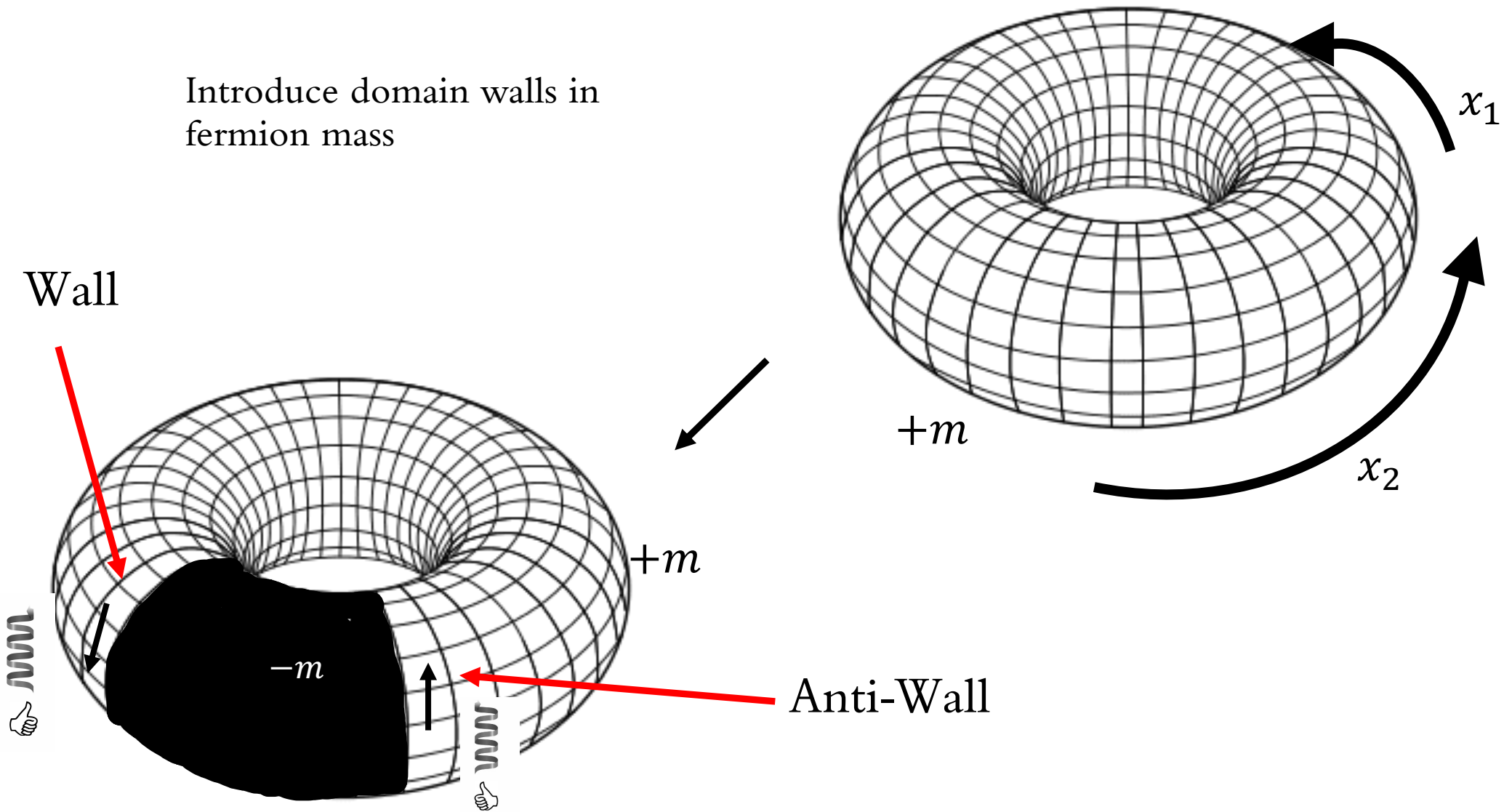
Go to higher dimension. Consider 2+1 D Dirac fermion. Let's go to finite volume.

Periodic boundary condition in all directions with uniform Dirac mass.

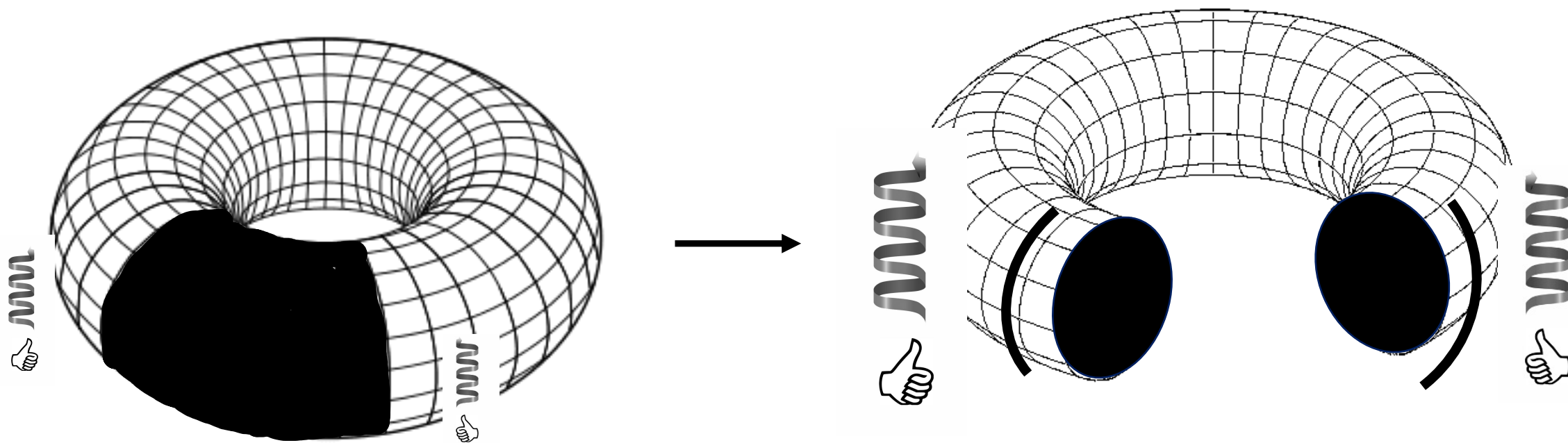


# How do you separate the two chirality?

Introduce domain walls in  
fermion mass



# Mass defect to cylinder (open boundary condition)

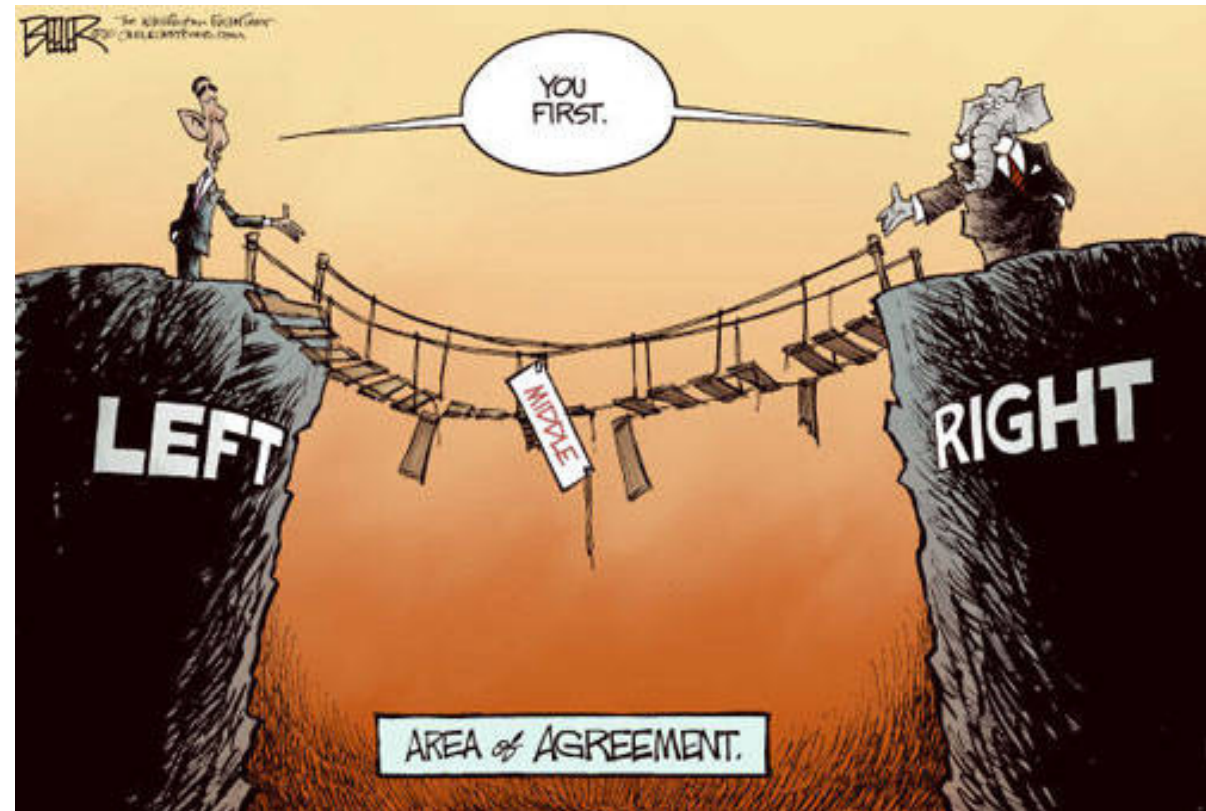


# So, how do you simulate QCD?

- Physically separate the right and the left: **right** lives on the **wall** and **left** on the **anti-wall** (Kaplan 1992).

Quantum corrections suppressed in the separation.

- Allow the gluons to talk to both walls.
- And the physics of the two walls together will give you QCD.



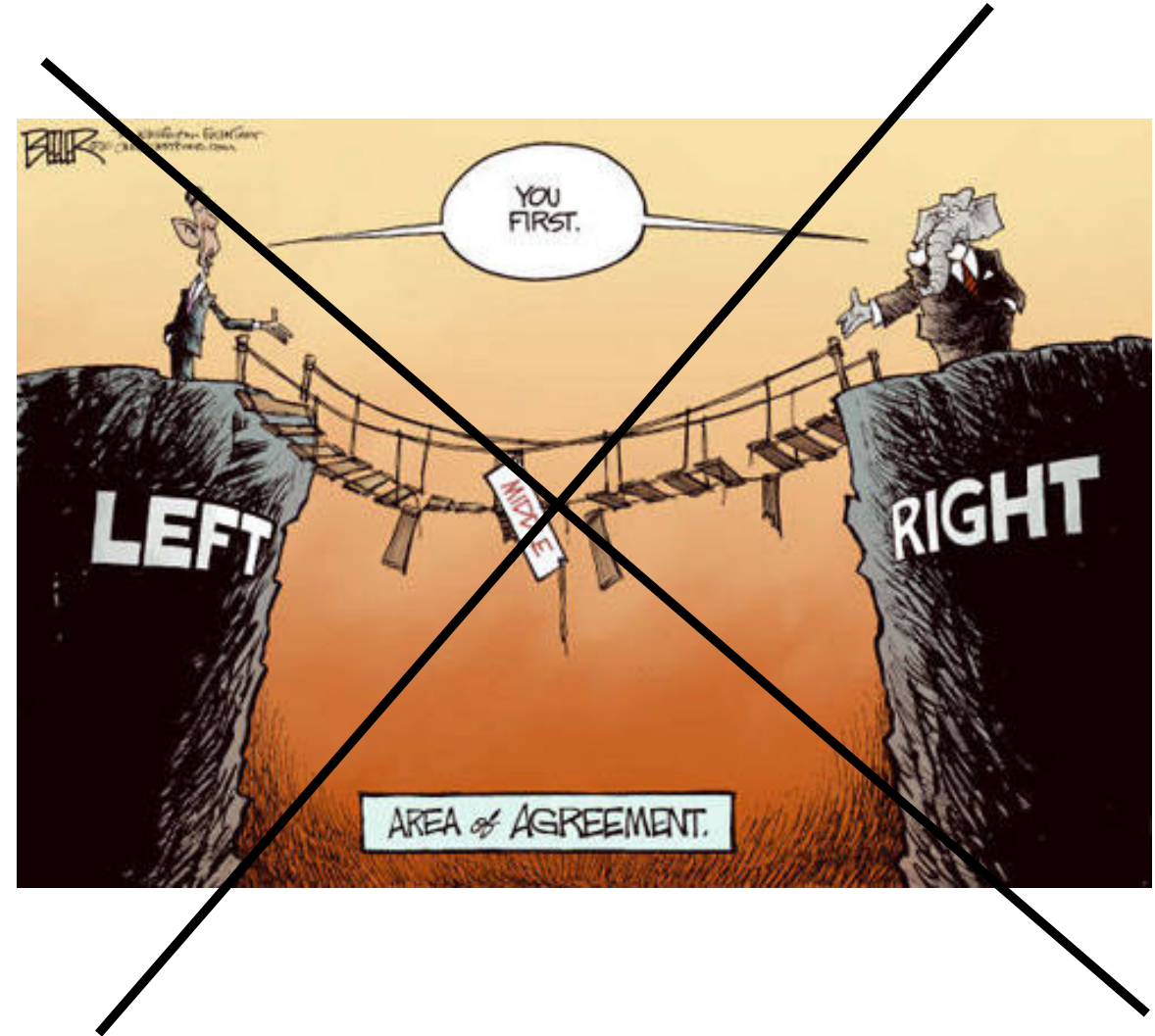
Quantum Hall effect

# Doesn't work for chiral gauge theories

The idea does not work for **chiral gauge theories** though.

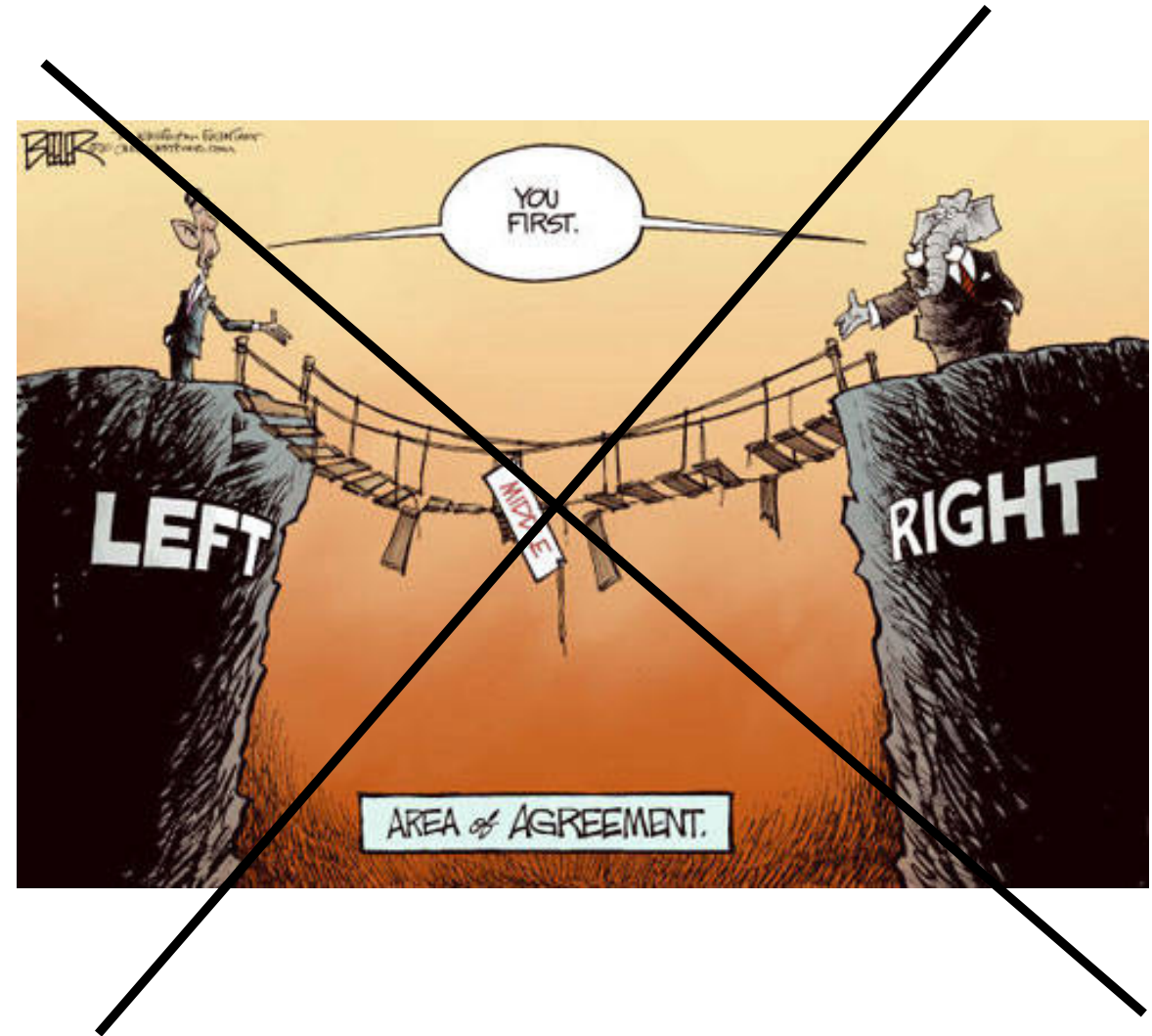
The construction in finite volume necessarily has two defects.

Two defects lead to opposite chiralities producing non-chiral theory.

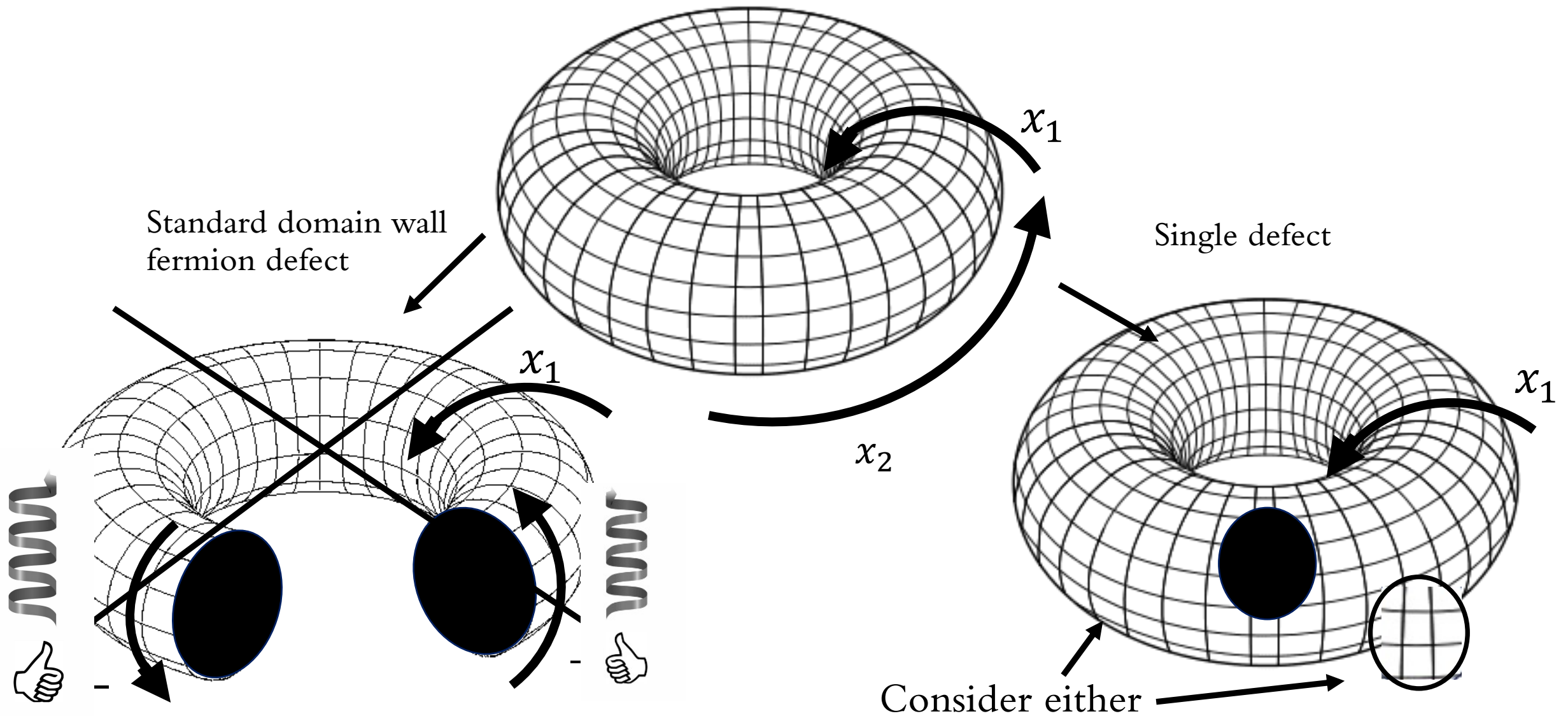


# Doesn't work for chiral gauge theories

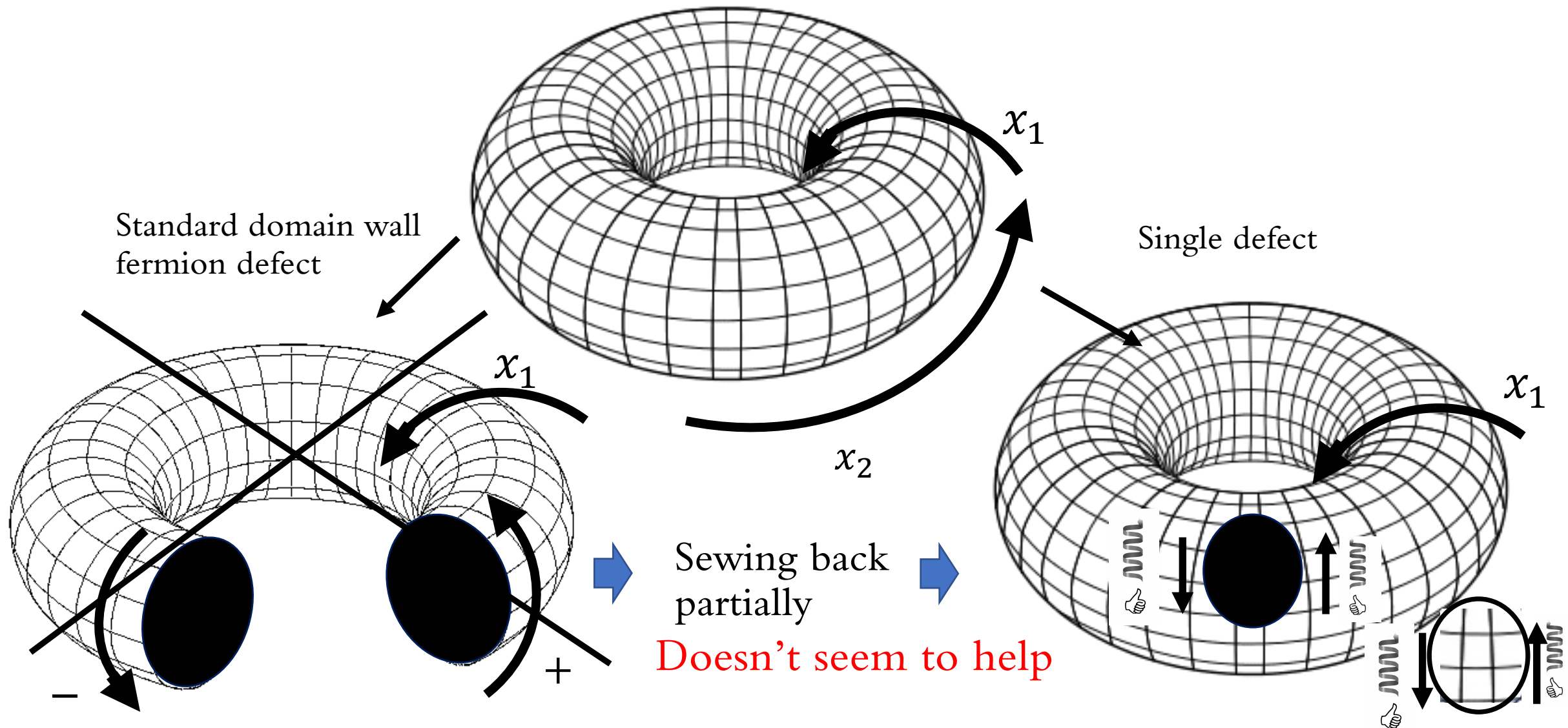
We need to isolate a single Weyl fermion of a particular chirality --- impossible with the standard domain wall setup.



# What if we had a single defect?

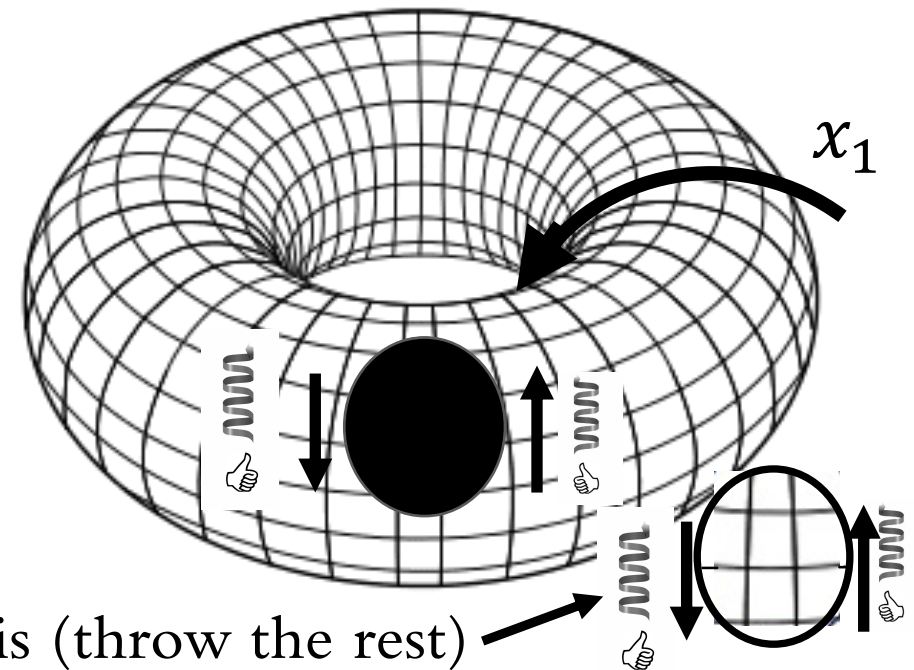


# Opposite chirality on the two sides..



# Opposite chirality on the two sides..

Maybe the problem is that we are keeping the definition of chirality position independent.



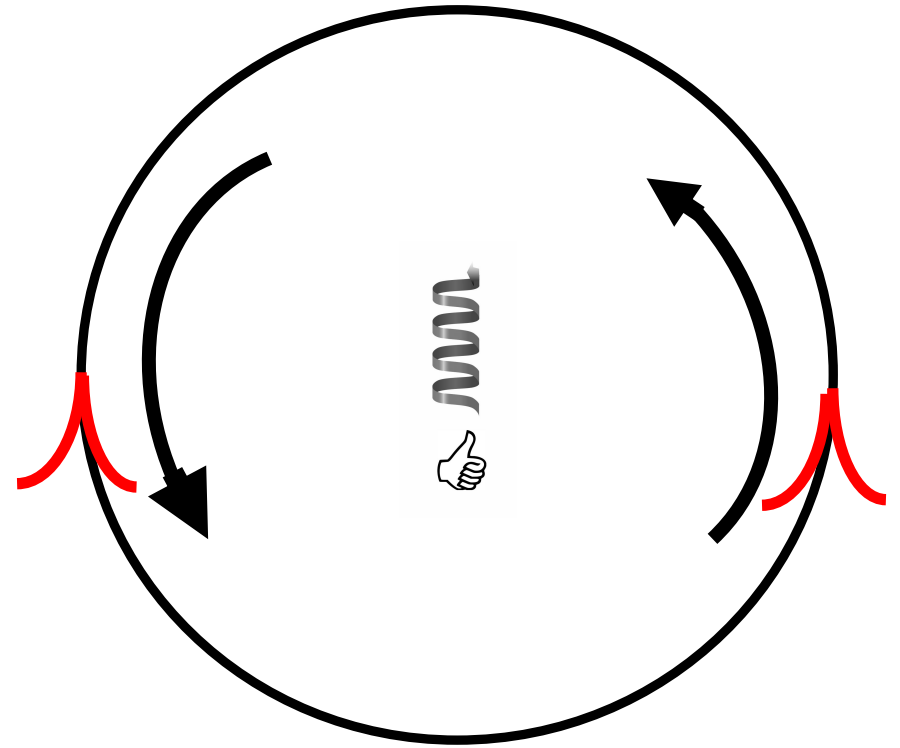
# Define chirality in a position dependent manner

Define chirality as clockwise travel vs anticlockwise travel:

counter-clockwise



clockwise



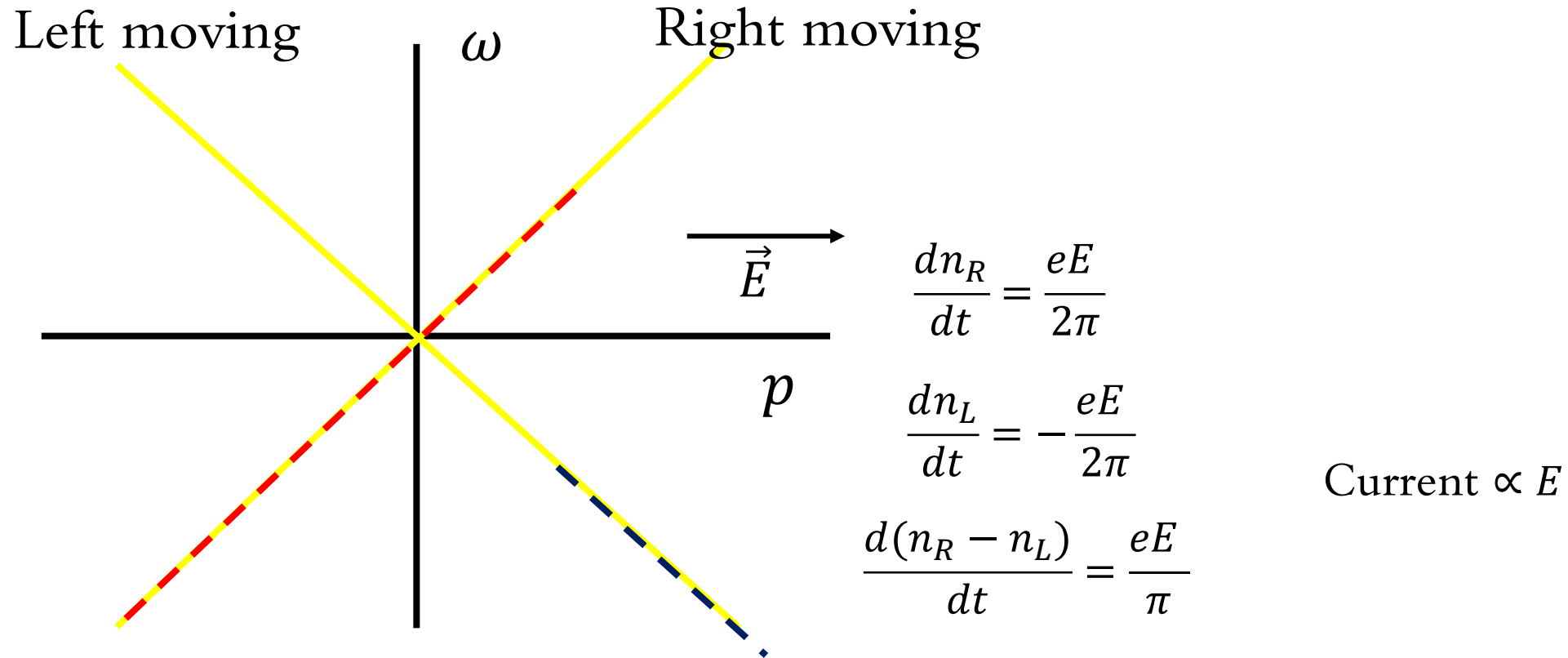
# Gauging

Can engineer any number of Weyl fermions on the boundary.

We can gauge any subgroup of the available global symmetry of the free fermion theory. ---- **Must make sense only if the theory is anomaly free.**

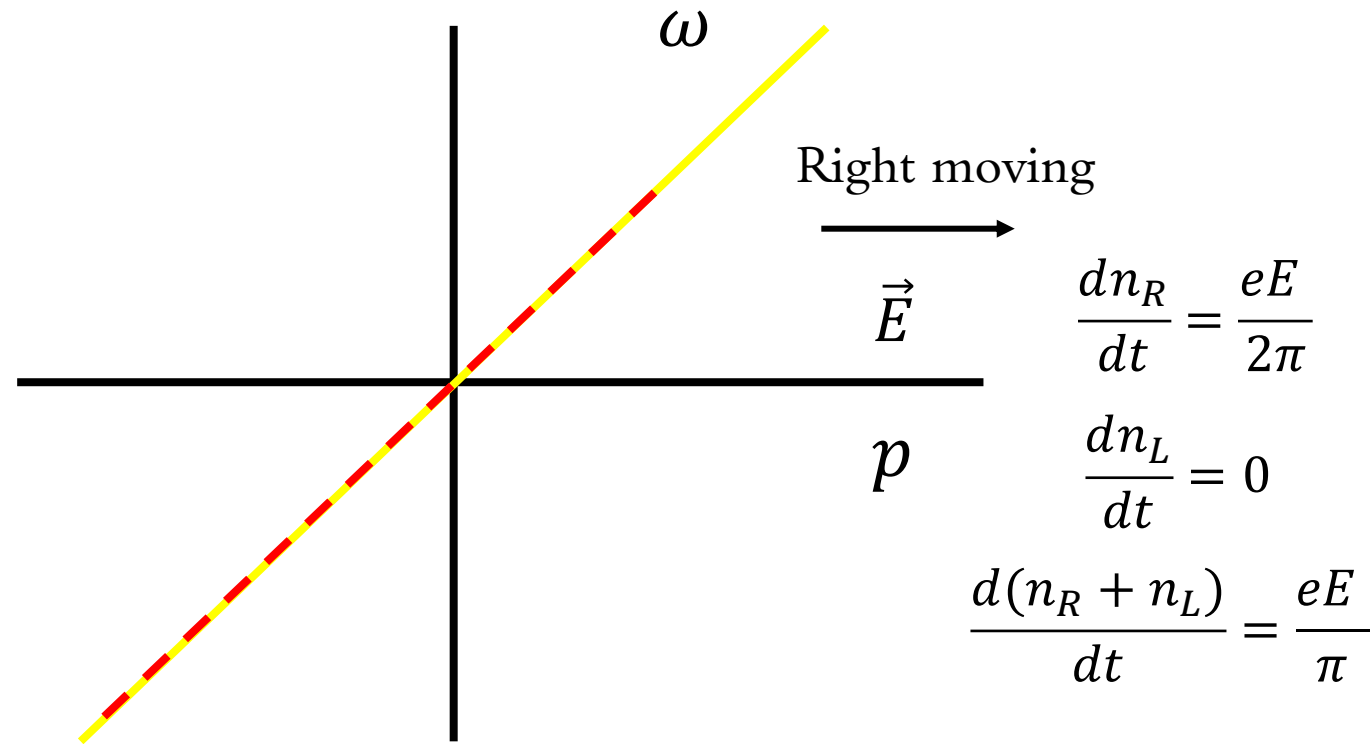
**How do we see this?**

# 1+1 D massless Dirac fermion spectrum



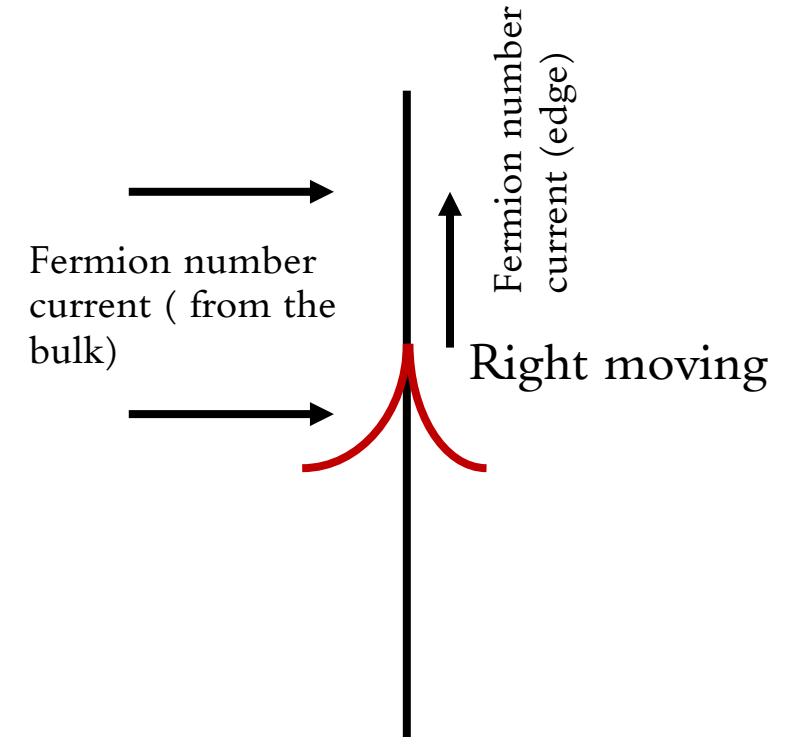
Vector current or charge  $n_R + n_L$   
conserved, axial not so.

# Edge world: Anomaly, Weyl fermion



Current  $\propto E$

Vector current not conserved, by itself is sick in an electric field.



Can exist on the boundary of a higher dimensional theory

# Domain wall + anomaly

Left moving

$$m < 0$$

$$\Rightarrow j_2 = \partial_0 A_1 - \partial_1 A_0 = E_1$$

Fermion number  
current (bulk)

$x_2$

zoom in

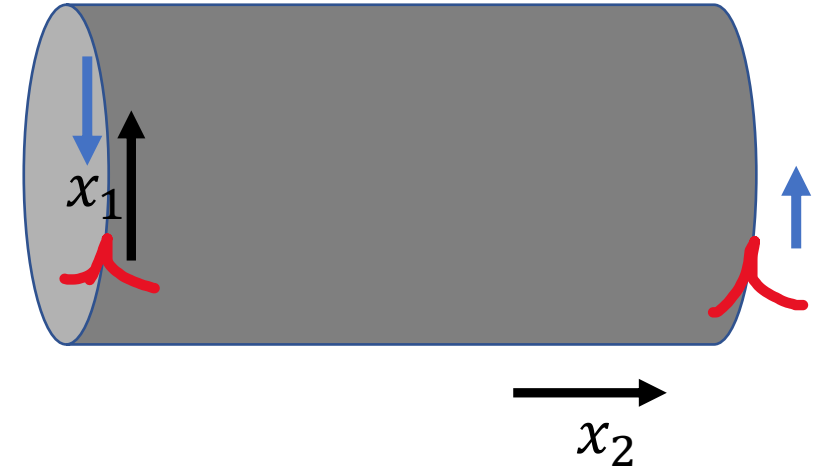
$x_1$

$$E_1 = \partial_0 A_1 - \partial_1 A_0$$

Right moving

$$m > 0$$

$$m < 0$$

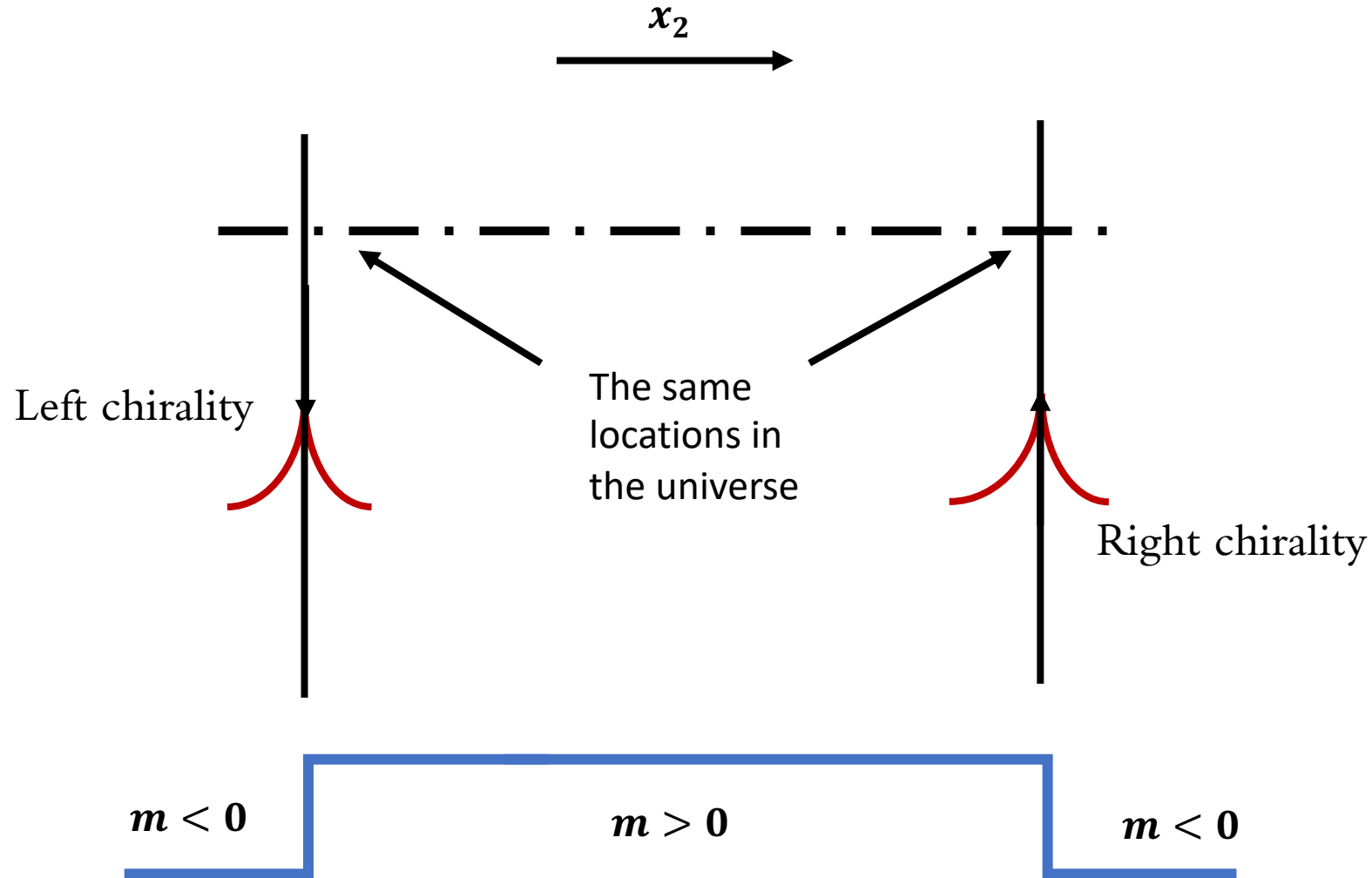


2 + 1 D Dirac fermion  
setup with domain walls

Gauge the Dirac  
fermion

Bulk talks to the  
boundary, great for  
QCD

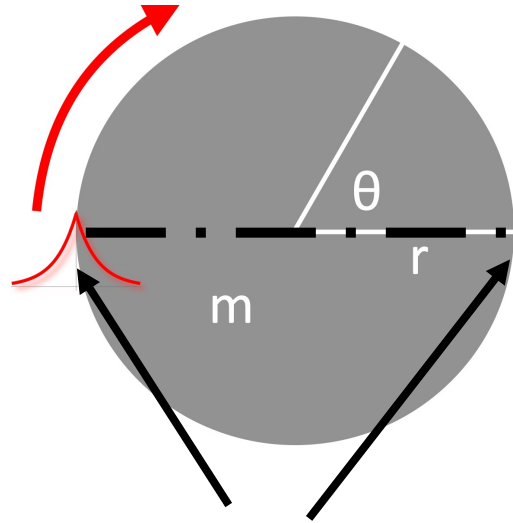
# Domain wall + anomaly + QCD



Bulk talks to the boundary, great for QCD

The two walls talk to each other

# disk + anomaly



Different locations in the universe, shouldn't communicate across the defect

But they will if the boundary theory has gauge anomaly: **nonlocal!!**

Thankfully, the standard model is anomaly free.

So, the disk construction makes sense, and the boundary theory is **local**.

# Summary and outlook

The plan is to engineer the fermion content of the CGT (e.g. the standard model) on the boundary and then turn on the desirable gauge interactions.

Several outstanding questions remain relating to the gauging of the theory.

The ideas discussed here rely on separating the chiralities in extra dimension. Demonstrate ties between quantum Hall physics and lattice field theory.

In Dirac and Weyl semimetals, one may be able to separate different Weyl nodes in momentum space.

Could they be used in lattice formulation of QCD or CGT in real time?