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ICSC Centro Nazionale di Ricerca in HPC, Big Data and Quantum Computing

Sum rules for Chiral, Conformal and Gravitational anomaly form factors

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Based on works 2502.03182 and 2504.01904

UNIVERSITÀ DEL SALENTO L'Ateneo tra i due mari
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ICSC Centro Nazionale di Ricerca in HPC, Big Data and Quantum Computing

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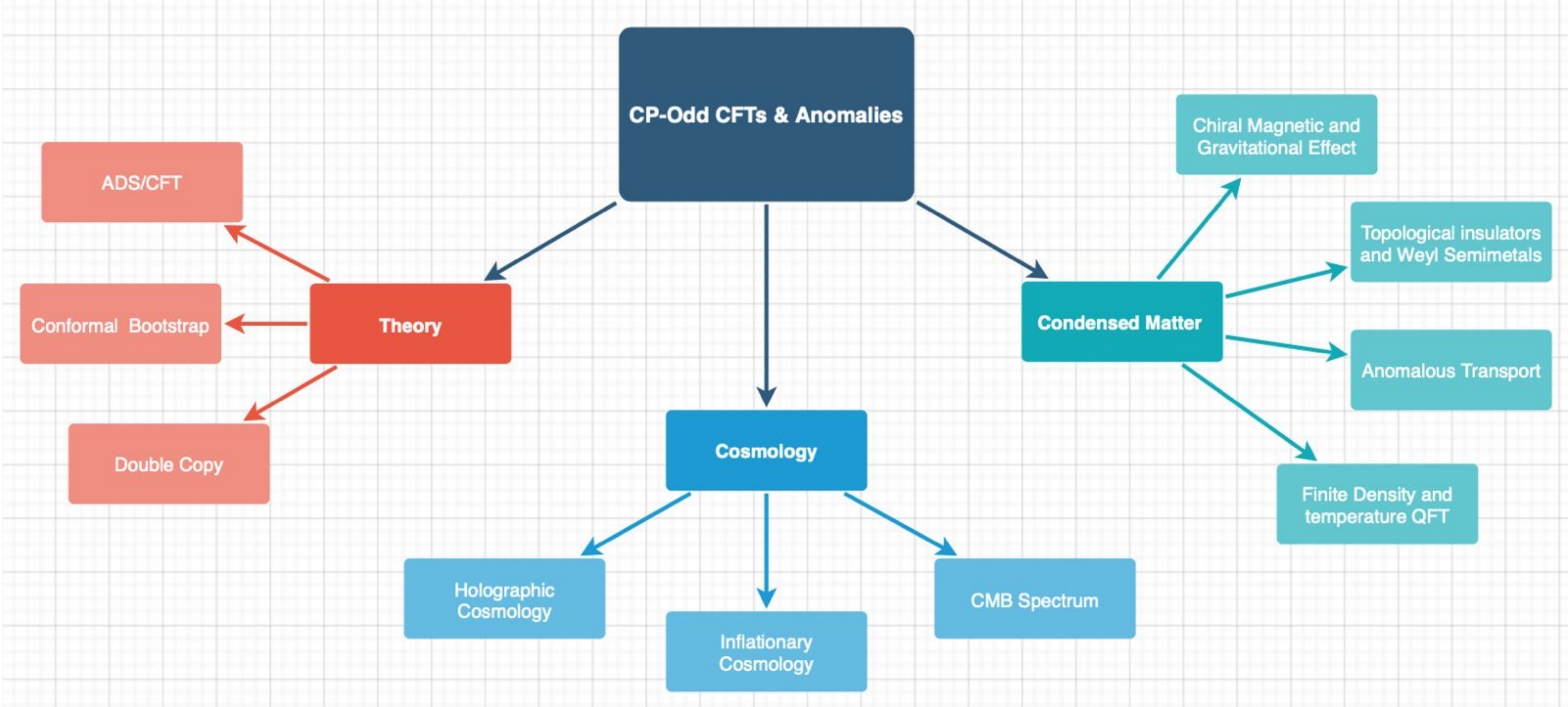
Outline

- Some motivation
- Chiral sum rules and the anomaly pole of the <AVV>
- Sum rules in the <ATT>
- Conformal sum rules and the <TJJ> extra pole
- Future Developments


CP-Odd CFTs & Anomalies
Theory
ADS/CFT
Conformal Bootstrap
Double Copy
Cosmology
Holographic Cosmology
Inflationary Cosmology
CMB Spectrum
Condensed Matter
Chiral Magnetic and Gravitational Effect
Topological insulators and Weyl Semimetals
Anomalous Transport
Finite Density and temperature QFT

Some motivation

- **Weyl semimetals** exhibit **chiral anomalies** through observable effects:
    - **Chiral magnetic effect**
    - **Anomalous Hall conductivity**

- **Luttinger's relation** relates **thermal gradients** to **gravitational gradients**

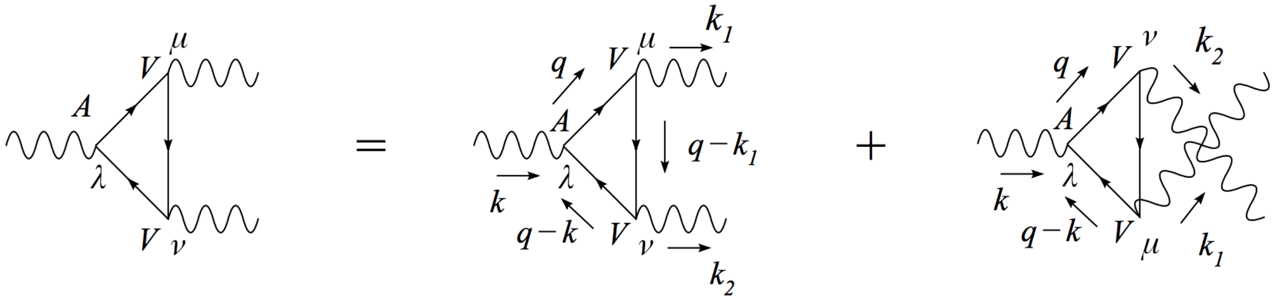
- **Gravitational anomalies** emerge in **thermal and energy transport**

- **Conformal anomalies** appear in scale-invariant regimes or surface states of **Dirac semimetals**

- **Correlator-based methods** allow:
    - Systematic derivation of anomaly-induced transport via **Kubo formulas**,
    - Unified treatment of electric, thermal, and mixed responses,
    - Clear signatures to connect **topological band structure** to **field-theoretic anomalies**.

AVV diagram

$$\Delta_{\alpha\mu\nu}(p_1,p_2) = \int \frac{d^4k}{(2\pi)^4}\mathrm{Tr}\left(\frac{1}{\not k - \not p_1 - m}\gamma_\mu \frac{1}{\not k - m}\gamma_\nu \frac{1}{\not k + \not p_2 - m}\gamma_\alpha\gamma_5\right) + [(p_1,\mu)\leftrightarrow(p_2,\nu)]$$



Schoutens used in this parametrization

$$p_2^{\mu_1}\epsilon^{\mu_2\mu_3p_1p_2} = p_2^{\mu_2}\epsilon^{\mu_1\mu_3p_1p_2} - p_2^{\mu_3}\epsilon^{\mu_1\mu_2p_1p_2} - p_2^2\epsilon^{\mu_1\mu_2\mu_3p_1} + (p_1\cdot p_2)\,\epsilon^{\mu_1\mu_2\mu_3p_2}$$
$$p_1^{\mu_2}\epsilon^{\mu_1\mu_3p_1p_2} = p_1^{\mu_1}\epsilon^{\mu_2\mu_3p_1p_2} + p_1^{\mu_3}\epsilon^{\mu_1\mu_2p_1p_2} - p_1^2\epsilon^{\mu_1\mu_2\mu_3p_2} + (p_1\cdot p_2)\,\epsilon^{\mu_1\mu_2\mu_3p_1}$$

$$\langle J^{\mu_1}(p_1)J^{\mu_2}(p_2)J_5^{\mu_3}(q)\rangle = F_1\left(\epsilon^{\mu_1\mu_2p_1p_2}p_1^{\mu_3} + \epsilon^{\mu_1\mu_2p_1p_2}p_2^{\mu_3}\right)$$
$$+F_2\left(\epsilon^{\mu_2\mu_3p_1p_2}p_1^{\mu_1} - \epsilon^{\mu_1\mu_2\mu_3p_2}s_1\right) + F_3\left(\epsilon^{\mu_1\mu_3p_1p_2}p_2^{\mu_2} - \epsilon^{\mu_1\mu_2\mu_3p_1}s_2\right)$$

$$\nabla_\mu\langle J_5^\mu\rangle = a_1\varepsilon^{\mu\nu\rho\sigma}F_{\mu\nu}F_{\rho\sigma} + a_2\varepsilon^{\mu\nu\rho\sigma}R^{\alpha\beta}{}_{\mu\nu}R_{\alpha\beta\rho\sigma},$$

$$
\begin{aligned}
F_1 =&\frac{1}{64\pi^2\lambda(s,s_1,s_2)^2}\Big(2\left(s^2-2(s_1+s_2)s+(s_1-s_2)^2\right)(s-s_1-s_2)\\
&+2\left((s_1+s_2)s^2-2\left(s_1^2-4s_2s_1+s_2^2\right)s+(s_1-s_2)^2(s_1+s_2)\right)\mathrm{B}_0\left(s,m^2\right)\\
&-2s_1\left((s-s_1)^2-5s_2^2+4(s+s_1)s_2\right)\mathrm{B}_0\left(s_1,m^2\right)\\
&-2s_2\left(s^2+4s_1s-2s_2s-5s_1^2+s_2^2+4s_1s_2\right)\mathrm{B}_0\left(s_2,m^2\right)\\
&+4\big(m^2(s-s_1-s_2)\left(s^2-2(s_1+s_2)s+(s_1-s_2)^2\right)\\
&-s_1s_2\left(-2s^2+(s_1+s_2)s+(s_1-s_2)^2\right)\big)\mathrm{C}_0\left(s,s_1,s_2,m^2\right)\Big)
\end{aligned}
$$

Off-Shell

$$
\lim_{q^2\to 0} q^2\left\langle J^{\mu_1}(p_1)J^{\mu_2}(p_2)J^{\mu_3}(q)\right\rangle = 0.
$$

$$
\begin{aligned}
F_2 =&\frac{1}{2\pi^2\lambda(s,s_1,s_2)^2}\Big(\left(s^2-2(s_1+s_2)s+(s_1-s_2)^2\right)(s-s_1+s_2)\\
&+s\left((s-s_1)^2-5s_2^2+4(s+s_1)s_2\right)\mathrm{B}_0\left(s,m^2\right)\\
&-\left(s^3-(2s_1+s_2)s^2+\left(s_1^2+8s_2s_1-s_2^2\right)s+(s_1-s_2)^2s_2\right)\mathrm{B}_0\left(s_1,m^2\right)\\
&+\left(-5s^2+4(s_1+s_2)s+(s_1-s_2)^2\right)s_2\mathrm{B}_0\left(s_2,m^2\right)\\
&+2\big((s-s_1+s_2)\left((s-s_1)^2+s_2^2-2(s+s_1)s_2\right)m^2\\
&+ss_2\left(s^2+s_1s-2s_2s-2s_1^2+s_2^2+s_1s_2\right)\big)\mathrm{C}_0\left(s,s_1,s_2,m^2\right)\Big)
\end{aligned}
$$

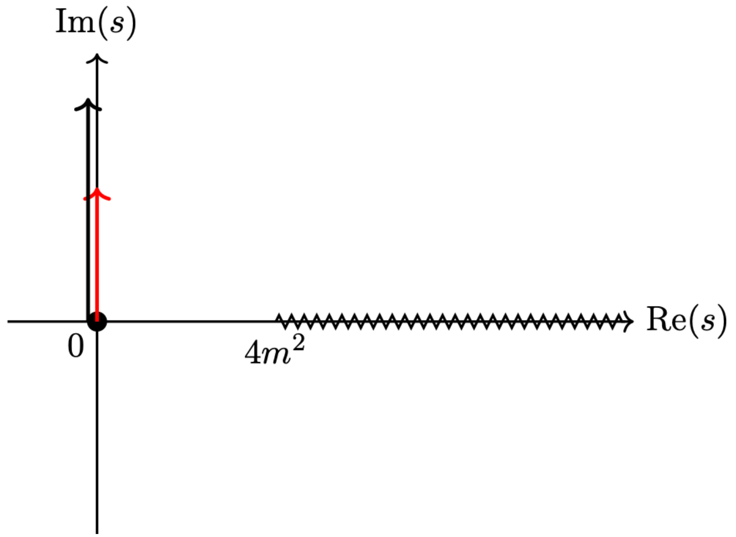
On-Shell m=0

$$
\lim_{q^2\to 0} \quad q^2\left\langle J^{\mu_1}(p_1)J^{\mu_2}(p_2)J^{\mu_3}(q)\right\rangle = \frac{1}{2\pi^2}q^{\mu_3}\epsilon^{\mu_1\mu_2p_1p_2}
$$

$$
\begin{aligned}
F_3 =&-\frac{1}{2\pi^2\lambda(s,s_1,s_2)^2}\Big(\left(s^2-2(s_1+s_2)s+(s_1-s_2)^2\right)(s+s_1-s_2)\\
&+s\left(s^2+4s_1s-2s_2s-5s_1^2+s_2^2+4s_1s_2\right)\mathrm{B}_0\left(s,m^2\right)\\
&+s_1\left(-5s^2+4(s_1+s_2)s+(s_1-s_2)^2\right)\mathrm{B}_0\left(s_1,m^2\right)\\
&-\left((s+s_1)(s-s_1)^2+(s+s_1)s_2^2-2\left(s^2-4s_1s+s_1^2\right)s_2\right)\mathrm{B}_0\left(s_2,m^2\right)\\
&+2\big((s+s_1-s_2)\left((s-s_1)^2+s_2^2-2(s+s_1)s_2\right)m^2\\
&+ss_1\left((s-s_1)^2-2s_2^2+(s+s_1)s_2\right)\big)\mathrm{C}_0\left(s,s_1,s_2,m^2\right)\Big).
\end{aligned}
$$

No particle pole

$$p_1^\mu \Delta_{\alpha\mu\nu}(p_1,p_2)=0, \quad p_2^\nu \Delta_{\alpha\mu\nu}(p_1,p_2)=0.$$



The Longitudinal part is:

$$\Phi_0(p_1^2,p_2^2,p_3^2,m^2) = \frac{g^2 m^2}{2\pi^2}\frac{1}{q^2}\mathrm{C}_0\left(q^2,p_1^2,p_2^2,m^2\right)+\frac{g^2}{4\pi^2}\frac{1}{q^2}.$$

$$\Phi_0(q^2,s_1,s_2,m^2) = \frac{1}{2\pi i}\oint_C \frac{\Phi_0(s,s_1,s_2,m^2)}{s-q^2}\,ds,$$

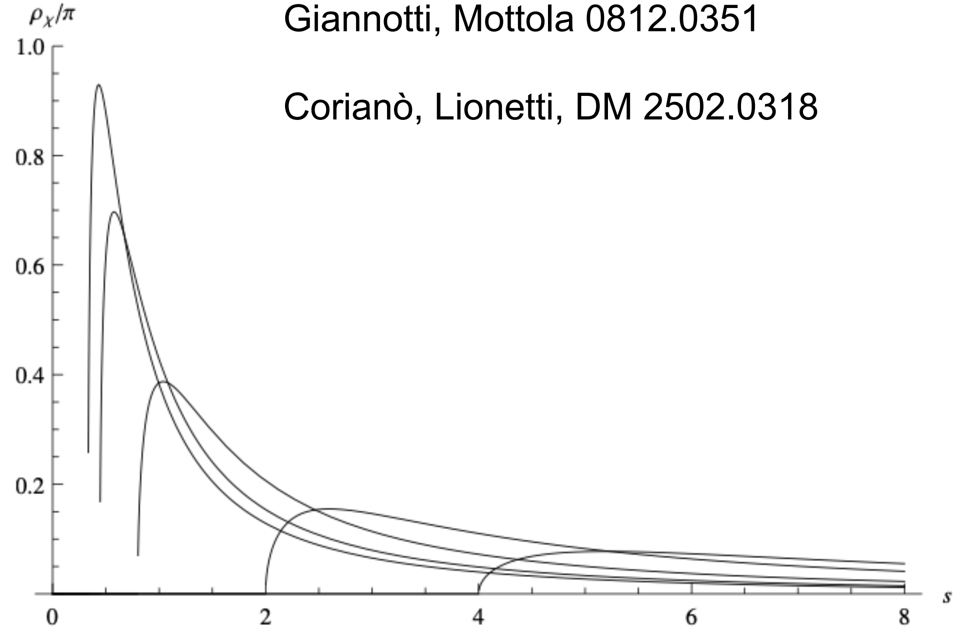
$$\mathrm{disc}\,\Phi_0(q^2,p_1^2,p_2^2,m^2) = -4i\pi m^2\, C_0(q^2,p_1^2,p_2^2,m^2)\,\delta(q^2)-2\pi i\,\delta(q^2)+\frac{2m^2}{q^2}\mathrm{disc}\,C_0(q^2,p_1^2,p_2^2,m^2).$$

$$\begin{aligned}\rho(s,s_1,s_2,m^2) &\equiv \frac{2m^2}{q^2}\mathrm{disc}\,C_0(q^2,p_1^2,p_2^2,m^2)\theta(q^2-4m^2)\\ &= \frac{2m^2\pi\log\left(\frac{m^2-s(w-\bar z)(\bar w-z)}{m^2-s(w-z)(\bar w-\bar z)}\right)}{s^2(z-\bar z)},\end{aligned}$$

$$\frac{1}{\pi}\int_{4m^2}^{\infty}\Delta\Phi_0(s,p_1^2,p_2^2,m^2)ds = \frac{g^2}{2\pi^2},$$

Giannotti, Mottola 0812.0351

Corianò, Lionetti, DM 2502.0318



(a)

$$\Phi_0(q^2) = \frac{1}{\pi} \int_0^\infty \frac{\Delta\Phi_0(s)}{s - q^2}\, ds$$

$$\mathcal{L}^{-1}\left\{\int_0^\infty \frac{\Delta\Phi_0}{s - q^2}\, ds\right\}(t) = \frac{1}{\pi} \int_0^\infty \Delta\Phi_0(s) e^{q^2 t}\, ds$$

$$f(t) = \mathcal{L}^{-1}\{F(s)\} = \frac{1}{2\pi i} \int_{\gamma - i\infty}^{\gamma + i\infty} F(s) e^{st}\, ds,$$

$$\lim_{t \to 0} \mathcal{L}^{-1}\{\Phi_0\}(t) = \frac{1}{\pi} \int_0^\infty \Delta\Phi_0\, ds$$

$$\mathcal{L}^{-1}\{\Phi_0\}(t) = -\int_0^1 dx \int_0^{1-x} dy\, \frac{g^2 \left(m^2 \exp\left(\frac{t\left(-m^2 + p_1^2 xy - p_2^2 x^2 - p_2^2 xy + p_2^2 x\right)}{xy + y^2 - y}\right) - 2m^2 + p_1^2 xy - p_2^2 x^2 - p_2^2 xy + p_2^2 x\right)}{2\pi^2 \left(m^2 - p_1^2 xy + p_2^2 x^2 + p_2^2 xy - p_2^2 x\right)}$$

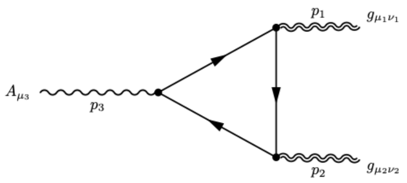
$$\frac{1}{\pi} \int_{4m^2}^\infty \Delta\Phi_0(s, p_1^2, p_2^2, m^2) ds = \frac{g^2}{2\pi^2},$$

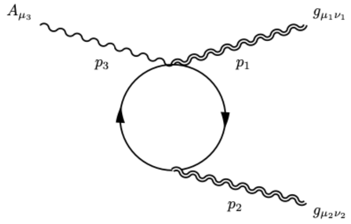
ATT correlator

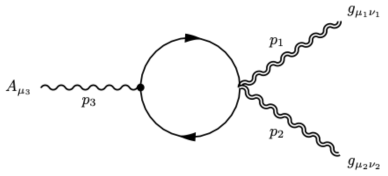
$$\langle T^{\mu_1\nu_1}T^{\mu_2\nu_2}J_A^{\mu_3}\rangle = \langle t^{\mu_1\nu_1}t^{\mu_2\nu_2}j_A^{\mu_3}\rangle + \langle t^{\mu_1\nu_1}t^{\mu_2\nu_2}j_{A\,loc}^{\mu_3}\rangle + \langle t_{loc}^{\mu_1\nu_1}t^{\mu_2\nu_2}j_A^{\mu_3}\rangle + \langle t^{\mu_1\nu_1}t_{loc}^{\mu_2\nu_2}j_A^{\mu_3}\rangle$$

$$\langle t^{\mu_1\nu_1}t^{\mu_2\nu_2}j_{A\,loc}^{\mu_3}\rangle = \frac{p_3^{\mu_3}}{p_3^2}\Pi^{\mu_1\nu_1}_{\alpha_1\beta_1}(p_1)\Pi^{\mu_2\nu_2}_{\alpha_2\beta_2}(p_2)\epsilon^{\alpha_1\alpha_2 p_1 p_2}\left(\bar{F}_1 g^{\beta_1\beta_2}(p_1\cdot p_2) + \bar{F}_2 p_1^{\beta_2}p_2^{\beta_1}\right)$$

$$\begin{aligned}\bar{F}_1 &= \frac{1}{24\pi^2} + \frac{m^2}{2\pi^2\lambda\,(s-s_1-s_2)}\Big\{2\left[\lambda m^2 + s s_1 s_2\right]C_0\left(s,s_1,s_2,m^2\right) + Q(s)\left[s(s_1+s_2)-(s_1-s_2)^2\right]\\ &\quad - s_2 Q(s_2)\left[s+s_1-s_2\right] - s_1 Q(s_1)\left[s-s_1+s_2\right] + \lambda\Big\},\\ \bar{F}_2 &= -\frac{1}{24\pi^2} + \frac{m^2}{2\pi^2\lambda^2}\Big\{2\left[\lambda\left(m^2(-s+s_1+s_2)-2s_1s_2\right) + 3s_1s_2\left((s_1-s_2)^2 - s(s_1+s_2)\right)\right]C_0\left(s,s_1,s_2,m^2\right)\\ &\quad + s_1 Q(s_1)\left[\lambda + 6s_2(s+s_1-s_2)\right] + s_2 Q(s_2)\left[\lambda + 6s_1(s-s_1+s_2)\right] - Q(s)\left[12 s s_1 s_2 + \lambda(s_1+s_2)\right]\\ &\quad + \lambda(-s+s_1+s_2)\Big\}.\end{aligned}$$



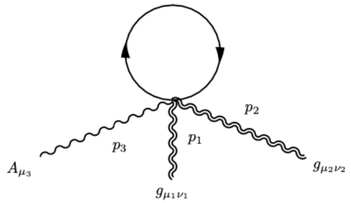




$$\langle t^{\mu_1\nu_1}t^{\mu_2\nu_2}j_{A\,loc}^{\mu_3}\rangle = (\bar{F}_1 - \bar{F}_2)\frac{p_3^{\mu_3}}{4\,p_3^2}\left\{\left[\varepsilon^{\nu_1\nu_2 p_1 p_2}\left(p_1\cdot p_2\, g^{\mu_1\mu_2} - p_1^{\mu_2}p_2^{\mu_1}\right) + (\mu_1\leftrightarrow\nu_1)\right] + (\mu_2\leftrightarrow\nu_2)\right\}$$

$$+(\bar{F}_1 + \bar{F}_2)\ \frac{p_3^{\mu_3}}{4\,p_3^2}\left\{\left[\varepsilon^{\nu_1\nu_2 p_1 p_2}\left(p_1\cdot p_2\, g^{\mu_1\mu_2} + p_1^{\mu_2}p_2^{\mu_1}\right) + (\mu_1\leftrightarrow\nu_1)\right] + (\mu_2\leftrightarrow\nu_2)\right\}$$

$$+(\bar{F}_1 + \bar{F}_2)\frac{(p_1\cdot p_2)}{p_1^2 p_2^2}\frac{p_3^{\mu_3}}{4\,p_3^2}\left\{\left[\varepsilon^{\nu_1\nu_2 p_1 p_2}\left((p_1\cdot p_2)\,p_1^{\mu_1}p_2^{\mu_2} - p_2^2 p_1^{\mu_1}p_1^{\mu_2} - p_1^2 p_2^{\mu_1}p_2^{\mu_2}\right) + (\mu_1\leftrightarrow\nu_1)\right] + (\mu_2\leftrightarrow\nu_2)\right\}$$



ATT sum rule

Corianò, Lionetti, DM 2502.0318

$$s_\pm = (\sqrt{s_1} \pm \sqrt{s_2})^2$$

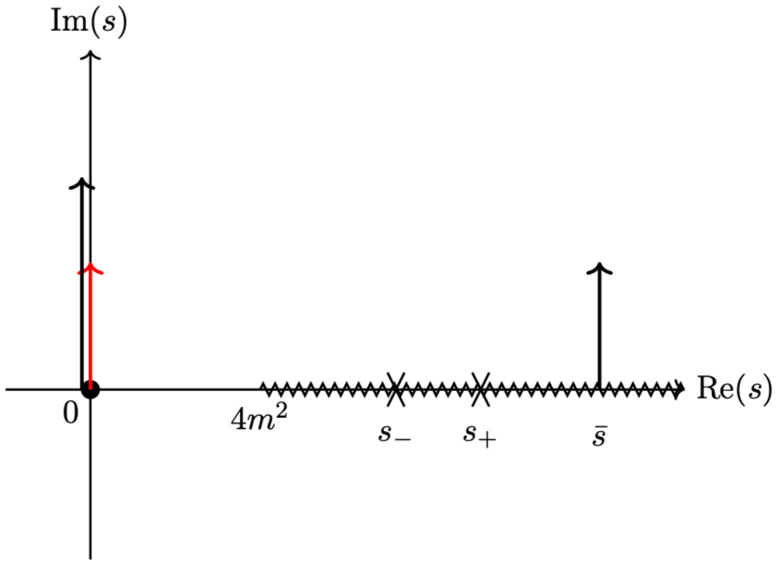
$$\text{disc}(\phi_{ATT})_{pole} = \text{disc}(\phi_{ATT})_0 + \text{disc}(\phi_{ATT})_{s_1+s_2} + \text{disc}(\phi_{ATT})_{s_-} + \text{disc}(\phi_{ATT})_{s_+}.$$

$$\text{disc}(\phi_{ATT})_0 = 2\pi i \frac{\delta(s)}{12\pi^2} - \frac{2\pi i\delta(s)}{12\pi^2(s_1-s_2)^3(s_1+s_2)}\Psi(s,s_1,s_2,m^2)$$

$$\begin{aligned}\Psi(s,s_1,s_2,m^2) \equiv \Big(& 12m^2\big(-\left((s_1-s_2)\left(s_1^2+s_2^2\right)\mathrm{B}_0\left(s,m^2\right)\right) \\ & + s_1\left(s_1^2+2s_1s_2+3s_2^2\right)\mathrm{B}_0\left(s_1,m^2\right) - s_2\left(3s_1^2+2s_1s_2+s_2^2\right)\mathrm{B}_0\left(s_2,m^2\right) \\ & + (s_1-s_2)\left(2m^2\left(s_1^2+s_2^2\right)+s_1s_2(s_1+s_2)\right)\mathrm{C}_0\left(s,s_1,s_2,m^2\right)\big) \\ & + 12m^2\left(s_1^2+s_2^2\right)(s_1-s_2)\Big).\end{aligned}$$

$$\begin{aligned}\text{disc}(\phi_{ATT})_{s_1+s_2} = & \frac{2\pi i\, m^2\delta(s-s_1-s_2)}{4\pi^2(s_1+s_2)}\Big(-2\mathrm{B}_0\left(s_1+s_2,m^2\right)+\mathrm{B}_0\left(s_1,m^2\right) \\ & + \mathrm{B}_0\left(s_2,m^2\right) + \left(4m^2-s_1-s_2\right)\mathrm{C}_0\left(s_1,s_2,s_1+s_2,m^2\right)+2\Big).\end{aligned}$$

$$\frac{1}{\pi}\int_0^\infty \Delta\phi_{ATT}\, ds = \frac{1}{12\pi^2}$$



TJJ diagram

Quark sector

$$
\begin{aligned}
\langle T_{\mu_1\nu_1}(\boldsymbol{p}_1) J^{\mu_2 a_2}(\boldsymbol{p}_2) J^{\mu_3 a_3}(\boldsymbol{p}_3)\rangle_q &= \langle t_{\mu_1\nu_1}(\boldsymbol{p}_1) j^{\mu_2 a_2}(\boldsymbol{p}_2) j^{\mu_3 a_3}(\boldsymbol{p}_3)\rangle_q \\
&+ 2\mathscr{T}_{\mu_1\nu_1}{}^{\alpha}(\boldsymbol{p}_1)\left[\delta^{\mu_3}_{[\alpha} p_{3\beta]}\langle J^{\mu_2 a_2}(\boldsymbol{p}_2) J^{\beta a_3}(-\boldsymbol{p}_2)\rangle_q + \delta^{\mu_2}_{[\alpha} p_{2\beta]}\langle J^{\mu_3 a_3}(\boldsymbol{p}_3) J^{\beta a_2}(-\boldsymbol{p}_3)\rangle_q\right] \\
&+ \frac{1}{d-1}\,\pi_{\mu_1\nu_1}(\boldsymbol{p}_1)\mathcal{A}_q^{\mu_2\mu_3 a_2 a_3},
\end{aligned}
$$

$$
\mathscr{T}_{\mu\nu\alpha}(\boldsymbol{p}) = \frac{1}{p^2}\left[2p_{(\mu}\delta_{\nu)\alpha} - \frac{p_\alpha}{d-1}\left(\delta_{\mu\nu} + (d-2)\frac{p_\mu p_\nu}{p^2}\right)\right]
$$

WIs

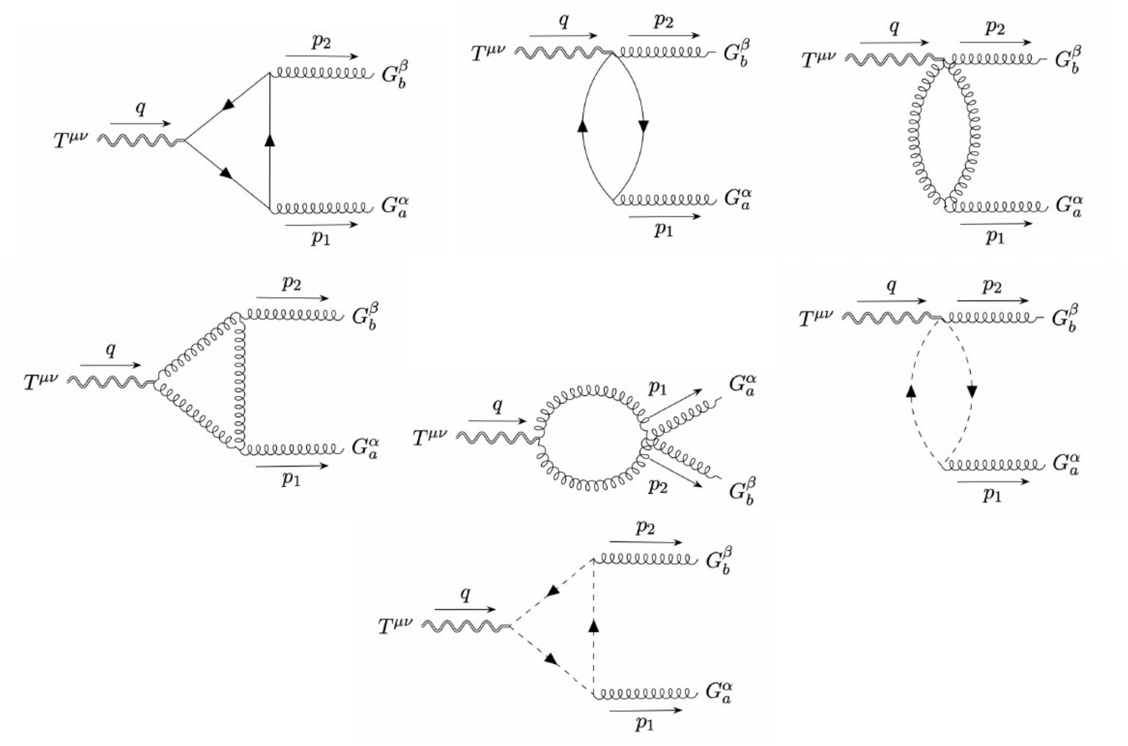
$$
p_1^{\nu_1}\langle T_{\mu_1\nu_1}(p_1) J^{\mu_2 a_2}(p_2) J^{\mu_3 a_3}(p_3)\rangle_q = 2\delta^{\mu_3}_{[\mu_1} p_{3\alpha]}\langle J^{\mu_2 a_2}(p_2) J^{\alpha a_3}(-p_2)\rangle_q + 2\delta^{\mu_2}_{[\mu_1} p_{2\alpha]}\langle J^{\alpha a_2}(p_3) J^{\mu_3 a_3}(-p_3)\rangle_q,
$$

$$
p_{2\mu_2}\langle T_{\mu_1\nu_1}(p_1) J^{\mu_2 a_2}(p_2) J^{\mu_3 a_3}(p_3)\rangle_q = 0,
$$

$$
\langle T(p_1) J^{\mu_2 a_2}(p_2) J^{\mu_3 a_3}(p_3)\rangle_q = \mathcal{A}_q^{\mu_2\mu_3 a_2 a_3},
$$

Gluon sector

$$
\begin{aligned}
\langle T^{\mu_1\nu_1}(p_1) J^{a_2\mu_2}(p_2) J^{a_3\mu_3}(p_3)\rangle_g &= \langle t^{\mu_1\nu_1}(p_1) j^{a_2\mu_2}(p_2) j^{a_3\mu_3}(p_3)\rangle_g + \langle t^{\mu_1\nu_1}(p_1) j^{a\mu_2}_{loc}(p_2) j^{a_3\mu_3}(p_3)\rangle_g + \langle t^{\mu_1\nu_1}(p_1) j^{a_2\mu_2}(p_2) j^{a_3\mu_3}_{loc}(p_3)\rangle_g \\
&+ 2\mathcal{I}^{\mu_1\nu_1\rho}(p_1)\left[\delta^{\mu_3}_{[\rho} p_{3\sigma]}\langle J^{a_2\mu_2}(p_2) J^{a_3\sigma}(-p_2)\rangle_g + \delta^{\mu_2}_{[\rho} p_{2\sigma]}\langle J^{a_3\mu_3}(p_3) J^{a_2\sigma}(-p_3)\rangle_g\right] + \frac{1}{d-1}\pi^{\mu_1\nu_1}(p_1)\left[\mathcal{A}_g^{\mu_2\mu_3 a_2 a_3} + \mathcal{B}_g^{\mu_2\mu_3 a_2 a_3}\right]
\end{aligned}
$$



$$
g_{\mu\nu}\langle T^{\mu\nu}\rangle = b_1\,E_4 + b_2\,C^{\mu\nu\rho\sigma}C_{\mu\nu\rho\sigma} + b_3\,\nabla^2 R + b_4\,F^{\mu\nu}F_{\mu\nu},
$$

$$
S_{pole} \;=\; \beta(g)\int d^4x\, d^4y\, R^{(1)}(x)\,\square^{-1}(x,y)\,F^a_{\alpha\beta}F^{a\alpha\beta}
$$

$$
\beta(g) = \frac{dg}{d\ln(\mu^2)} = -\beta_0\frac{g^3}{16\pi^2}, \qquad \beta_0 = \frac{11}{3}C_A - \frac{2}{3}n_f
$$

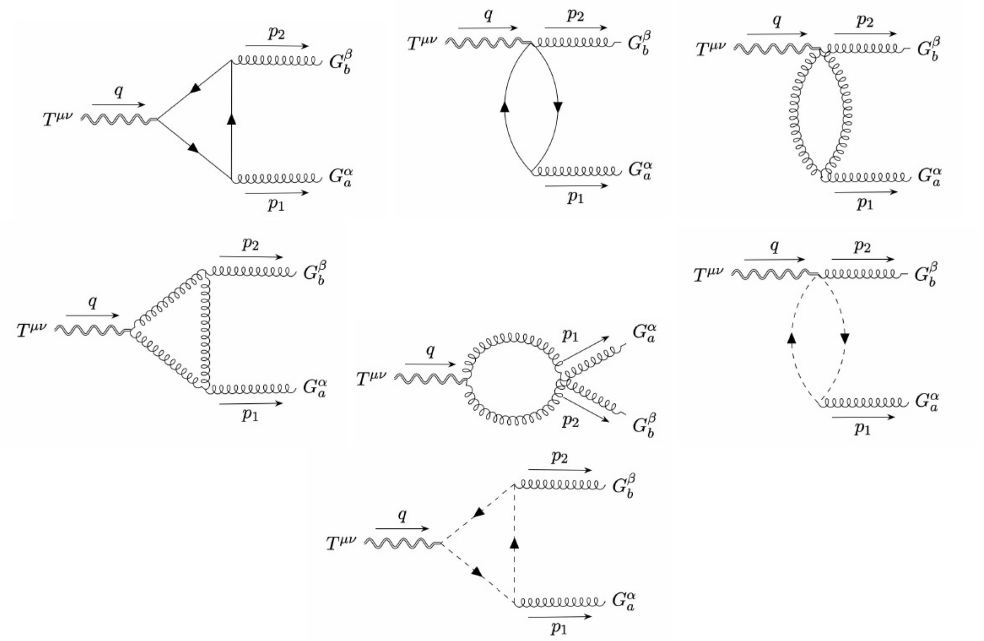
Corianò, Lionetti, DM, Tommasi 2409.05609

TJJ trace of the quark sector

$$g_{\mu\nu}\left\langle T^{\mu\nu}J^{\alpha a}J^{\beta b}\right\rangle_q = \left(\phi_1^{\alpha\beta}(p_1,p_2,q,m)+\phi_2^{\alpha\beta}(p_1,p_2,q,m)\right)\delta^{ab}$$

$$\phi_1^{\alpha\beta ab} = \left(-\frac{2}{3}n_f\frac{g_s^2}{16\pi^2}+\chi_0(p_1,p_2,q,m)\right)\delta^{ab}u^{\alpha\beta}(p_1,p_2)$$

$$\begin{aligned}\chi_0 =& g_s^2m^4H_1C_0\left(p_1^2,p_2^2,q^2,m^2\right)+g_s^2m^2\Big(H_2C_0\left(p_1^2,p_2^2,q^2,m^2\right)\\ &+H_3\bar{B}_0(q^2,m^2)+H_4\bar{B}_0(p_1^2,m^2)+H_5\bar{B}_0(p_2^2,m^2)+H_6\Big)\end{aligned}$$

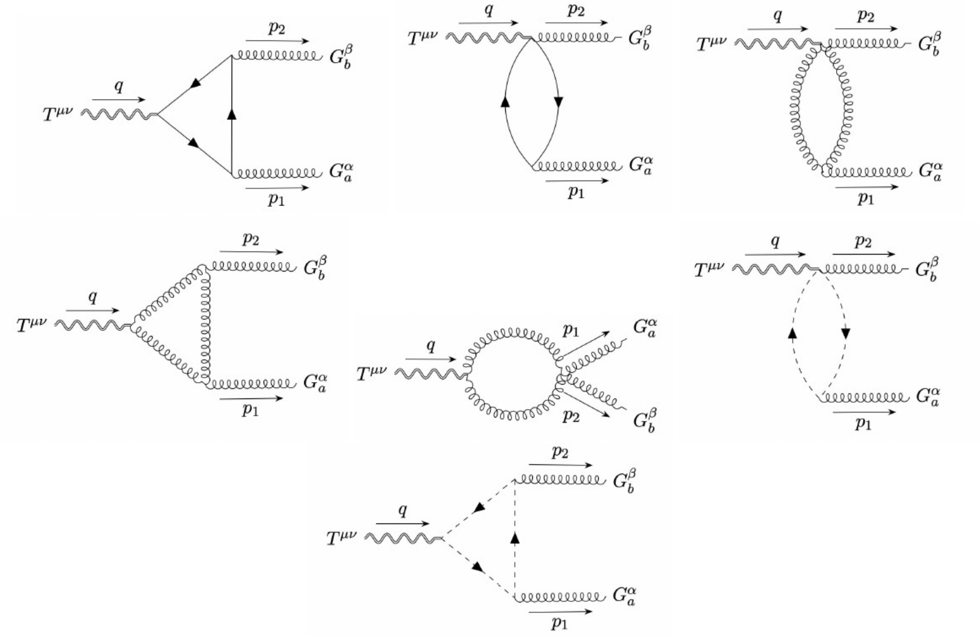


$$\phi_2^{\alpha\beta}(p_1,p_2,q,m)=\chi_1(p_1,p_2,q,m)v^{\alpha\beta}+\left(B_2p_1^\alpha p_1^\beta+B_2(p_1\leftrightarrow p_2)p_2^\alpha p_2^\beta+B_3p_1^\alpha p_2^\beta\right)$$

These are proportional to the equation of motion

TJJ trace of the gluon sector

$$g_{\mu\nu} \langle T^{\mu\nu} J^{\alpha a} J^{\beta b} \rangle_g = \left(\frac{g_s^2}{16\pi^2} \ \ \frac{11}{3} C_A \right) u^{\alpha\beta} \delta^{ab} + \mathcal{B}_g^{\alpha\beta} \delta^{ab}$$



$$\mathcal{B}_g^{\alpha\beta ab} = \chi_g(p_1, p_2, q) u^{\alpha\beta} \delta^{ab} + \phi_g^{\alpha\beta ab}$$

All the for factors are proportional to the equation of motion

$$\phi_g^{\alpha\beta ab} = \chi_g'(p_1, p_2, q) v^{\alpha\beta} \delta^{ab} + C_1\, p_1^\alpha\, p_1^\beta \delta^{ab} + C_2\, p_1^\alpha\, p_2^\beta \delta^{ab} + C_1(p_1 \leftrightarrow p_2)\, p_2^\alpha\, p_2^\beta \delta^{ab}$$

TJJ sum rule

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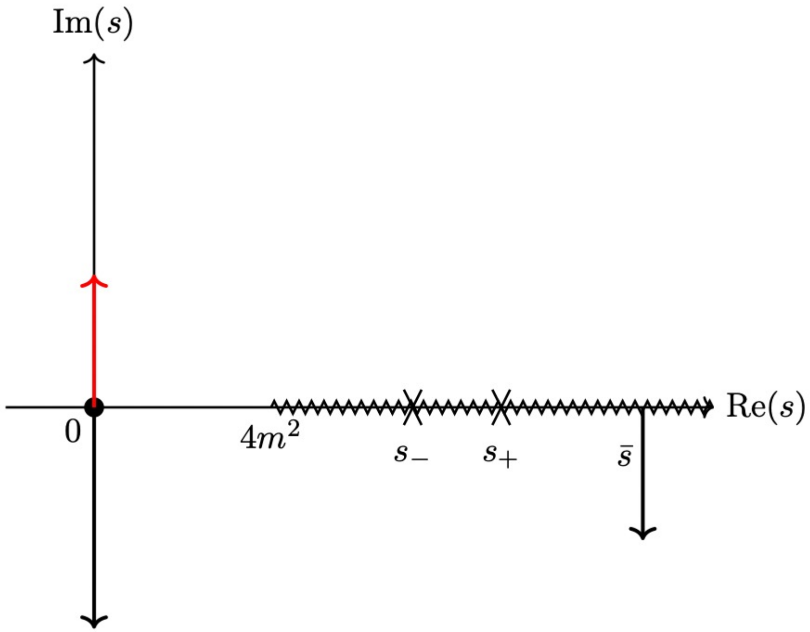
$$\langle T^{\mu\nu}(q)J^{a\alpha}(p_1)J^{b\beta}(p_2)\rangle_{tr} = \frac{1}{3\,q^2}\hat{\pi}^{\mu\nu}(q)\Bigg[\ \left(\frac{g_s^2}{16\pi^2}\left(\frac{11}{3}C_A - \frac{2}{3}n_f\right) + \chi_0(p_1,p_2,q,m) + \chi_g(p_1,p_2,q)\right)\delta^{ab}u^{\alpha\beta}(p_1,p_2)$$
$$+\left(\phi_2^{\alpha\beta ab} + \phi_g^{\alpha\beta ab}\right)\Bigg]$$

$$\Phi_{TJJ}(p_1^2,p_2^2,q^2,m^2) \equiv \frac{1}{3q^2}\mathcal{A} = \frac{1}{3q^2}\left(\frac{g_s^2}{48\pi^2}(11C_A - 2n_f) + \chi_0(p_1,p_2,q,m) + \chi_g(p_1,p_2,q)\right)$$

$$\text{disc}(\Phi_{TJJ})_{pole} = \text{disc}(\Phi_{TJJ})_0 + \text{disc}(\Phi_{TJJ})_{s_1+s_2} + \text{disc}(\Phi_{TJJ})_{s_-} + \text{disc}(\Phi_{TJJ})_{s_+}$$

$$\text{disc}(\Phi_{TJJ})_{cont} = \frac{n_F\, g_s^2\, m^2 K_1(s,s_1,s_2)\,\text{disc}\,B_0\left(s,m^2\right)}{12\pi^2 s(s-s_1-s_2)\lambda^2} + \frac{n_F\, g_s^2\, m^2 K_2(s,s_1,s_2)\,\text{disc}\,C_0\left(s,s_1,s_2,m^2\right)}{24\pi^2 s(s-s_1-s_2)\lambda^2}$$
$$\frac{C_A g_s^2(s_1+s_2)\,\text{disc}\,B_0(s)}{48\pi^2\lambda} + \frac{C_A g_s^2\, K_3\,\text{disc}\,C_0\left(s,s_1,s_2\right)}{96\pi^2 s(s-s_1-s_2)\lambda}$$

$$\int_{4m^2}^{\infty}\Delta\Phi_{TJJ}(s,p_1^2,p_2^2,m^2)\,ds = \frac{g_s^2}{144\pi^2}\left(11C_A - 2n_f\right)$$



TJJ and the dilaton pole

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Off-Shell

$$\lim_{s\to 0} q^2 \, \langle T^{\mu\nu}(q) J^{\alpha a}(p_1) J^{\beta b}(p_2)\rangle = 0.$$

On-Shell

$$\lim_{s\to 0} q^2 \, \langle T^{\mu\nu}(q) J^{\alpha a}(p_1) J^{\beta b}(p_2)\rangle = -\frac{g_s^2}{48\pi^2}\left(\frac{2}{3}n_f - \frac{11}{3}C_A\right)\tilde{\phi}_1^{\mu\nu\alpha\beta} - \frac{g_s^2}{288\pi^2}\,(n_f - C_A)\,\tilde{\phi}_2^{\mu\nu\alpha\beta}.$$

$$\tilde{\phi}_1^{\mu\nu\alpha\beta}(p_1,p_2) \quad = \quad (s\, g^{\mu\nu} - q^\mu q^\nu)\, u^{\alpha\beta}(p_1,p_2)$$

$$\tilde{\phi}_2^{\mu\nu\alpha\beta}(p_1,p_2) \quad = \quad -2\, u^{\alpha\beta}(p_1,p_2)\,[s\, g^{\mu\nu} + 2(p_1^\mu\, p_1^\nu + p_2^\mu\, p_2^\nu) - 4\,(p_1^\mu\, p_2^\nu + p_2^\mu\, p_1^\nu)]$$

Future developments

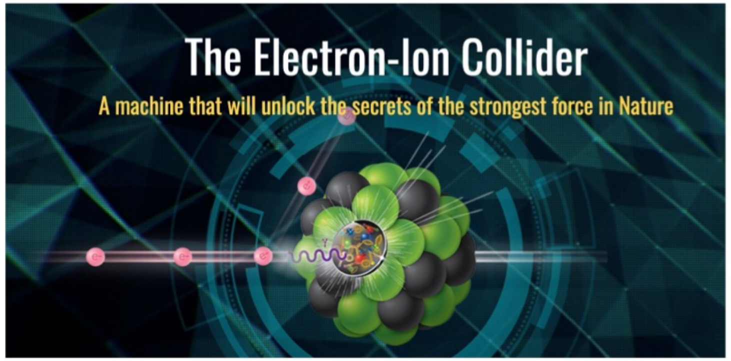
- Anomaly role in the pion and proton GFFs
- Clarify the Anomaly interplay in the Ji's sum rule of the proton
- Full calculation of the NLO for the pion and proton GFF

Thank you for your attention

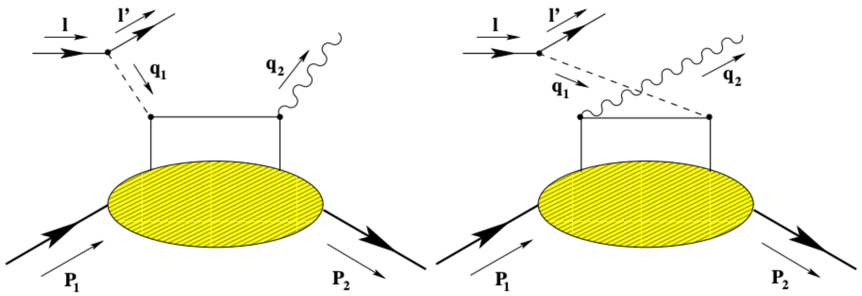
Brookhaven National Laboratory

Jefferson Lab

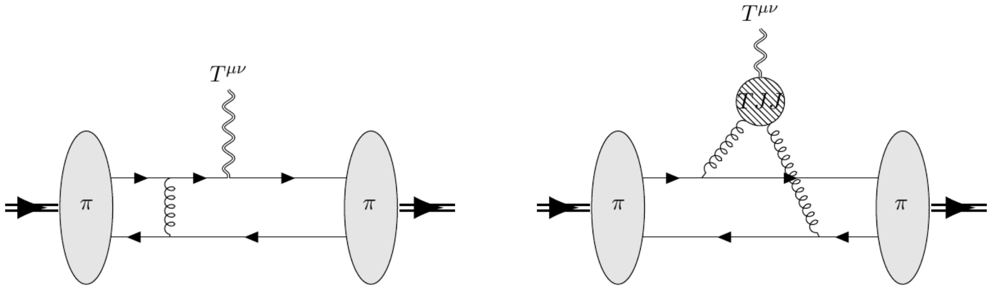
The Electron-Ion Collider (EIC) at Brookhaven National Laboratory is designed to have a highly flexible energy range, with the capability to collide electrons with protons and nuclei at center-of-mass energies ranging from approximately **20 GeV** to **140 GeV (e-p and e-Ions)**



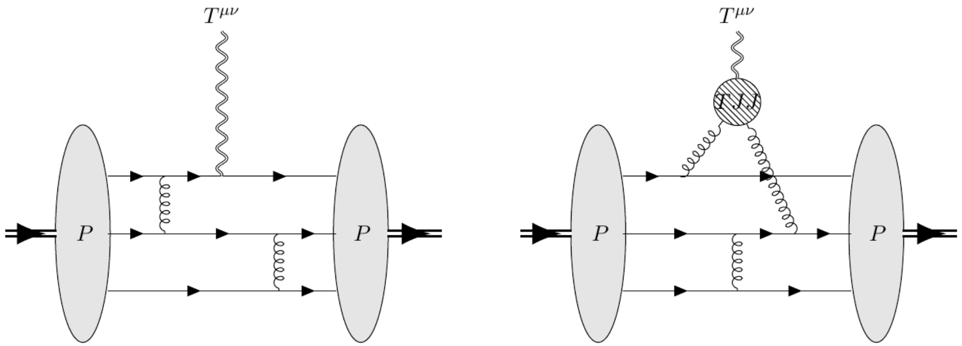
Invariant amplitudes of DVCS
Related to form factors of
GFF of the proton



$$\langle \pi' | \hat{T}_{\mu\nu}(0) | \pi \rangle = \left[2 P_\mu P_\nu \, A(t) + \frac{1}{2} (\Delta_\mu \Delta_\nu - g_{\mu\nu} \Delta^2) \, D(t) + 2 \, m^2 \, \bar{c}(t) \, g_{\mu\nu} \right]$$



$$\langle p', s' | T_{\mu\nu}(0) | p, s \rangle = \bar{u}' \left[A(t) \, \frac{\gamma_{\{\mu} P_{\nu\}}}{2} + B(t) \, \frac{i \, P_{\{\mu} \sigma_{\nu\}\rho} \Delta^\rho}{4M} + D(t) \, \frac{\Delta_\mu \Delta_\nu - g_{\mu\nu} \Delta^2}{4M} + M \sum_{\hat{a}} \bar{c}^{\hat{a}}(t) \, g_{\mu\nu} \right] u$$



Ji's Sum rule

$$J_q + J_g = \frac{1}{2} [A_q(0) + B_q(0)] + \frac{1}{2} [A_g(0) + B_g(0)] = \frac{1}{2}$$

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