

Searching Rare Higgs Decay - Proton-Proton collisions in boosted topology

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E FÍSICA EXPERIMENTAL DE PARTÍCULAS
partículas e tecnologia

[**HIGGS Physics** $H \rightarrow c\bar{c}$ *case study*]

04 July 2025

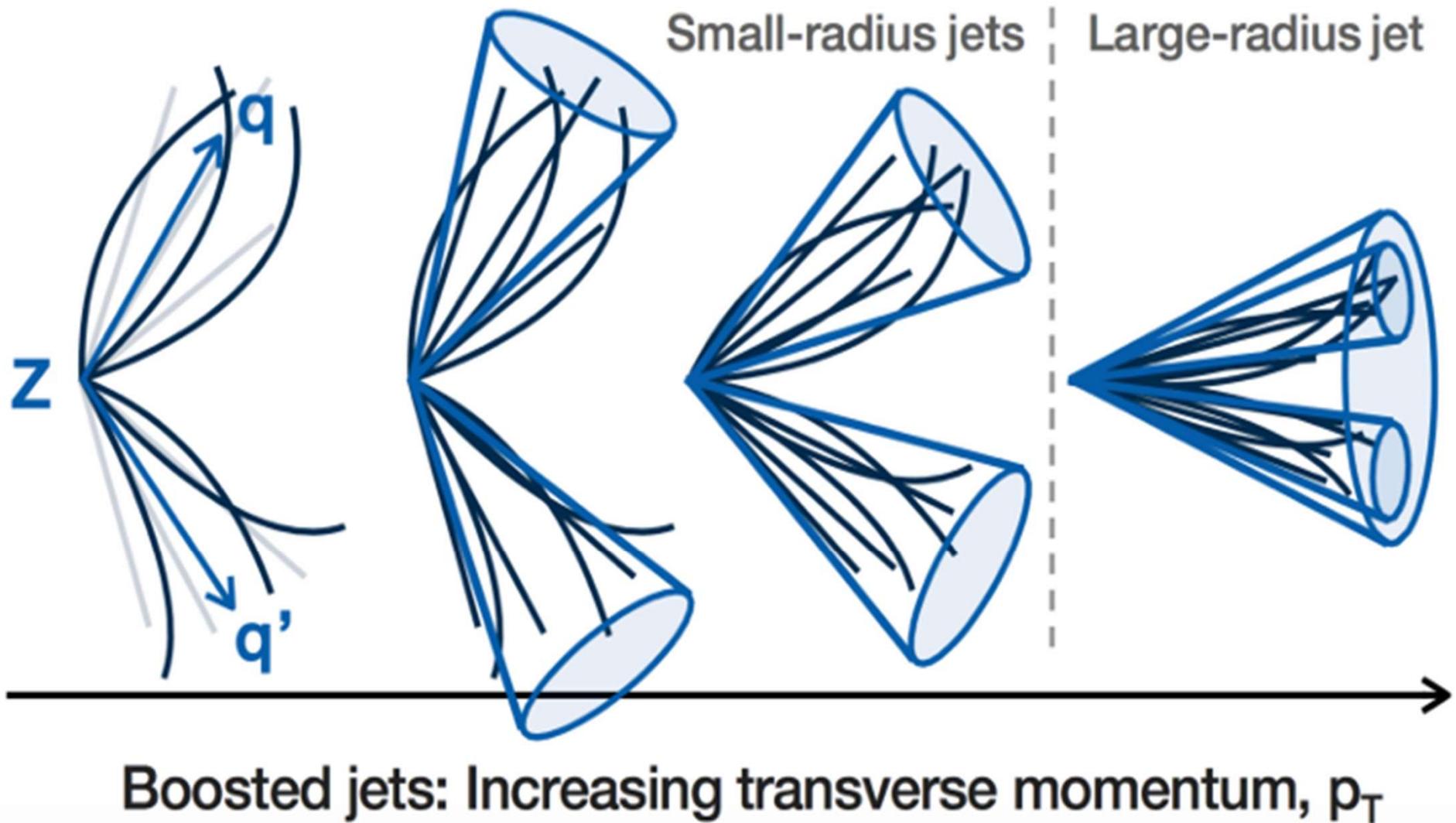
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Program: 14th Course on Physics @ LHC - 2025

Introduction

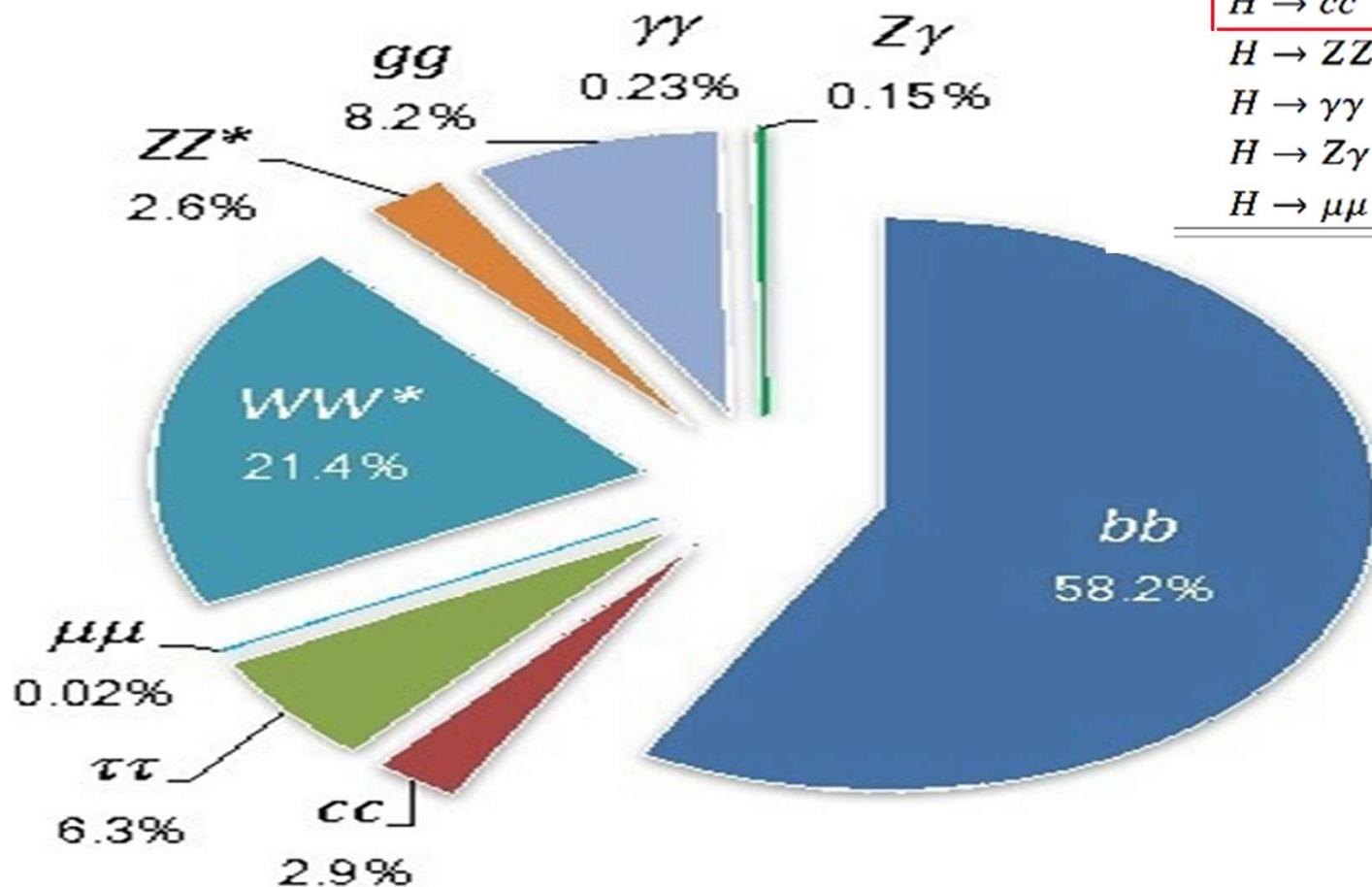
- This study presents SM Higgs Boson produced at $pT > 450$ GeV and decaying into charm-anticharm quark pair.
- P-P collisions at $\sqrt{s} = 13$ TeV were conducted by the CMS experiment at the LHC, with luminosity of 138 fb^{-1} .
- Boosted $H \rightarrow c\bar{c}$ decay products were reconstructed as a large single radius jet and identified using a deep neural network charm tagging technique.
- For validation $Z \rightarrow c\bar{c}$ decay process was measured with jets at high P_t and signal strength of $1.00_{-0.14}^{+0.17} (\text{syst}) \pm 0.08 (\text{theo}) \pm 0.06 (\text{stat})$
- The observed upper limit on cross section $(H)^* \text{ BR}(H \rightarrow c\bar{c})$ is set at 47(39) times the SM prediction, with 95% confidence.

- Quarks/Gluons Hadronize to produce a collimated spray of particles called Jets.
- This study focusses on reconstructing signal from Large-radius jets.



Background – Decay Channels for Higgs

- Couplings and kinematics determine these Branching ratios.



Decay channel	Branching ratio [%]
$H \rightarrow b\bar{b}$	57.5 ± 1.9
$H \rightarrow WW$	21.6 ± 0.9
$H \rightarrow gg$	8.56 ± 0.86
$H \rightarrow \tau\tau$	6.30 ± 0.36
$H \rightarrow c\bar{c}$	2.90 ± 0.35
$H \rightarrow ZZ$	2.67 ± 0.11
$H \rightarrow \gamma\gamma$	0.228 ± 0.011
$H \rightarrow Z\gamma$	0.155 ± 0.014
$H \rightarrow \mu\mu$	0.022 ± 0.001

Background – Decay Channels for Higgs

- $H \rightarrow \text{Muon pairs}$ is the most experimentally accessible decays and has been explored by ATLAS and CMS collaborations.
- CMS recently found Higgs couplings to muons.
- While, $H \rightarrow c\bar{c}$ is more challenging to find due to:
 - a.) Such decays are extremely rare.
 - b.) Extremely large multijet Background especially $H \rightarrow b\bar{b}$.
- What has changed/Improved ?
 - a.) Better flavour tagging techniques.
 - b.) Advanced Jet substructure.
- Prior searches focussed on H production with a Vector boson (VH , where V could be a W or Z).

Background – Decay Channels for Higgs

- The g-g fusion (ggF) and vector boson fusion (VBF) modes benefits from a larger cross section.
- This study focuses on the events from the ggF production.
- The search strategy is same as that of $H \rightarrow b\bar{b}$ decay channel.
- What's different ?
 - Use of new mass decorrelated discriminators to define a charm enriched signal region.
 - On the existing ATLAS and CMS measurements an additional constraint on the decay process is added.

The Search Mechanism



Peter Higgs (1929 – 2024)

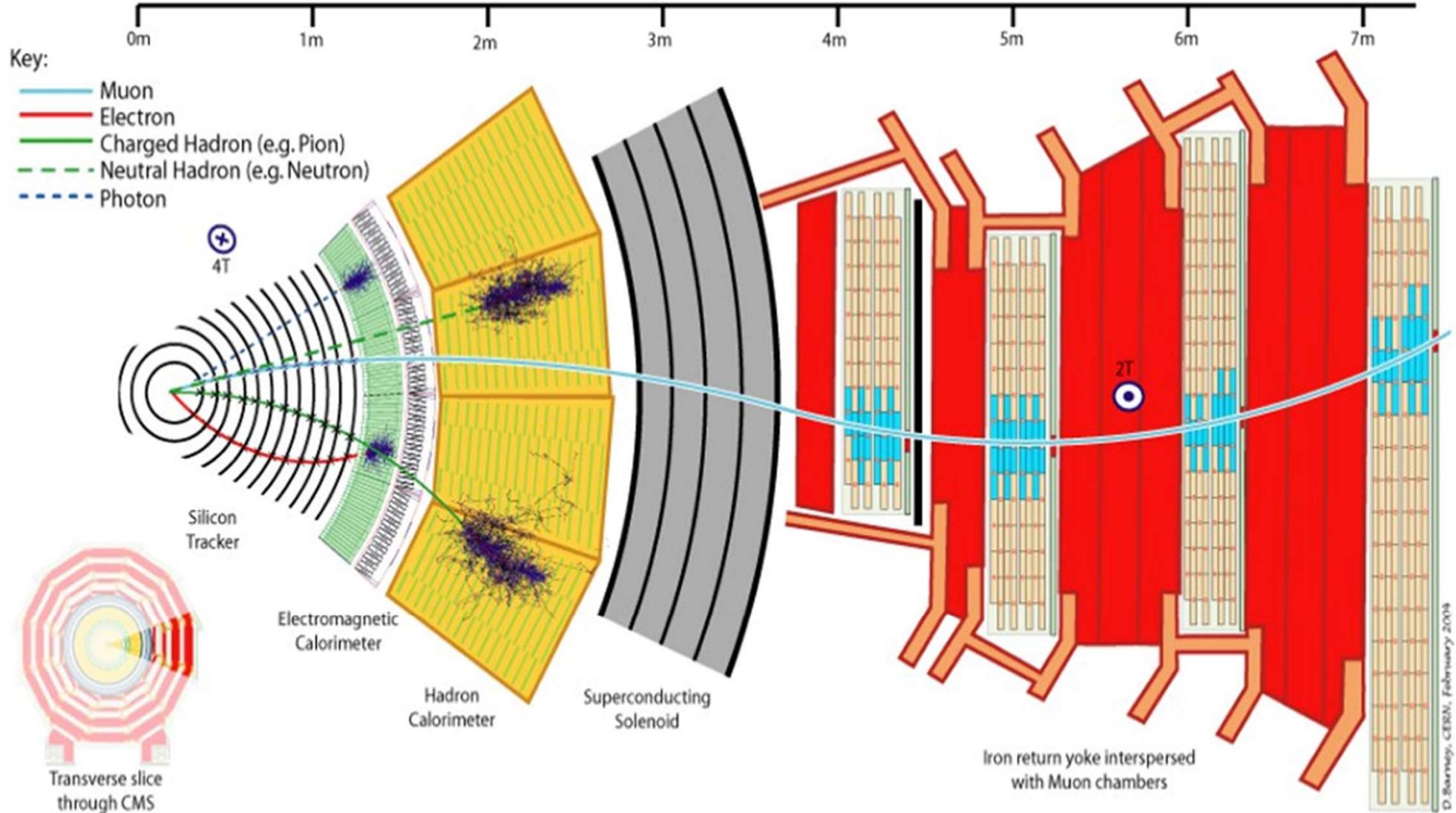
Identification

- The search used a dataset of p-p collisions at $\sqrt{s} = 13\text{TeV}$, collected with CMS detector.
- Candidate events are selected using High pT , large radius jet with substructure observables expected from an $H \rightarrow c\bar{c}$ decay.
- Deep neural Network (DNN) discriminators separates the H signal events from background events, specifically QCD induced multijet events (background events).
- A parallel search for $Z \rightarrow c\bar{c}$ decay is performed to validate the process (as those decays are more abundant).

Identification

- A simultaneous fit of a model of jet mass distributions is then performed across $H \rightarrow c\bar{c}$, $Z \rightarrow c\bar{c}$ signals, QCD multijet background events and other background processes across:
- This allows to measure the signal production cross sections.
- The compact muon solenoid (CMS) is a multipurpose detector, nearly hermetic, and can identify electrons, photons and charged and neutral hadrons, apart from muons of course.
- A “particle-flow” algorithm reconstructs all individual particles in an event, combining information from the silicon tracker, the crystal electromagnet and the brass scintillator hadron calorimeters, operating inside a 3.8 T superconducting solenoid.

Identification



https://indico.cern.ch/event/611242/contributions/2464737/attachments/1442943/2222206/SecondHiggsLecture_Conde.pdf

Reconstruction

- The reconstructed particles are used to build leptons, jets, and missing transverse momentum P_t .
- Simulated signal and background events are produced at matrix element level using various Monte Carlo Event Generators (like PYTHIA).
- The QCD multijet and Z+/W+ processes are modelled at QCD LO accuracy using MADGRAPH5_aMC@NLO v2.4.2 generator.
- The tt(bar) and single top quark processes are modelled at NLO using POWHEG 2.0. Diboson processes are modelled at LO using PYTHIA v8.226.
- The ggF Higgs production is simulated using the HJ-MINLO event generator with $M_h = 125$ GeV, including finite top quark mass effects.
- The POWHEG generator models H production through VBF, VH & tt(bar)H.

Reconstruction

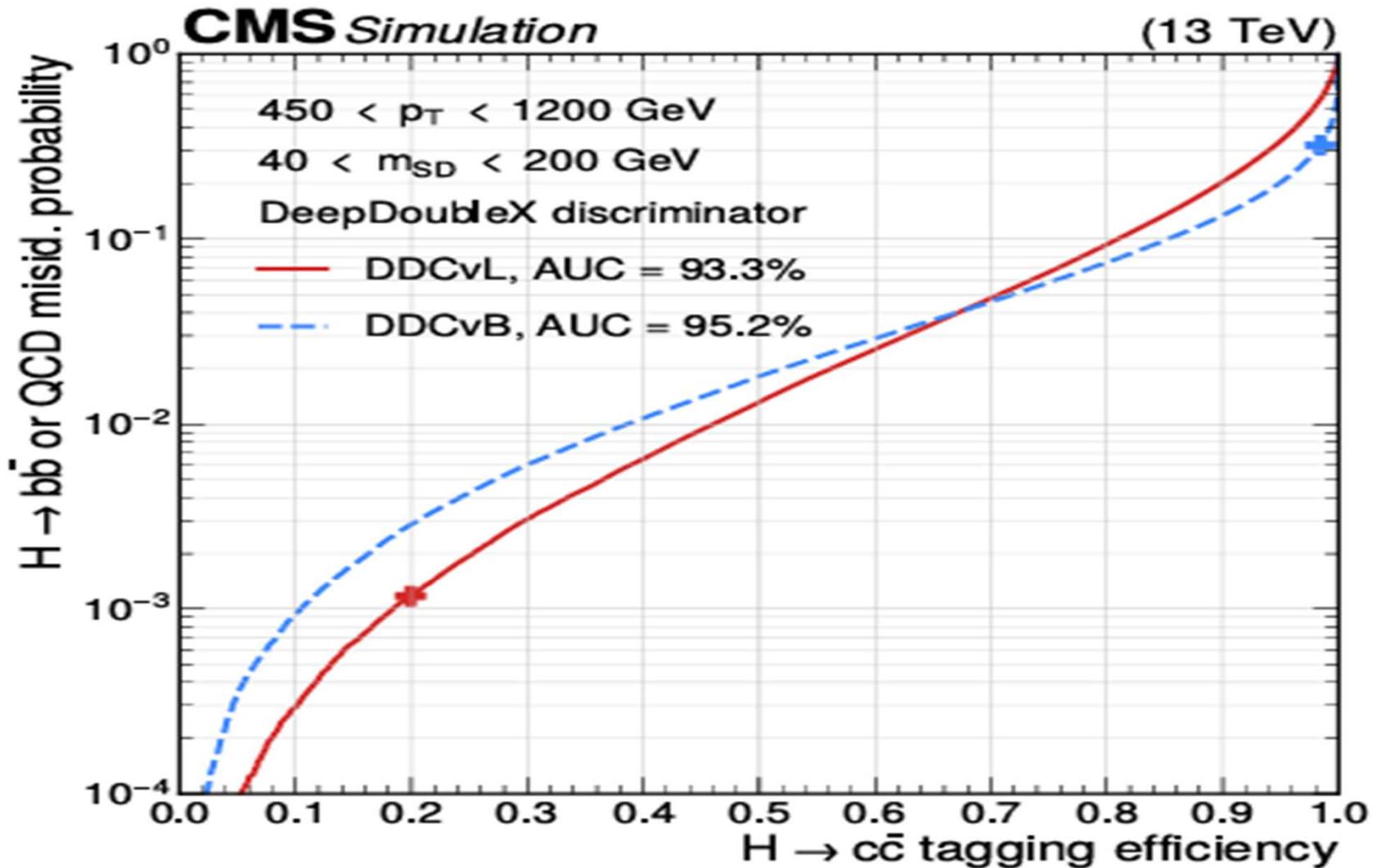
- For Parton showering and hadronization, the POWHEG and MADGRAPH5_aMC@NLO samples are interfaced with PYTHIA v8.205.
- Reconstructed particles are clustered into jets using the anti-Kt algorithm.
- Small radius jets are clustered with distance parameter of 0.4 (AK4 jets) while Large radius jets arising from decays of heavy boosted heavy particles are reconstructed using a distance parameter of 0.8 (AK 8 jets).
- The detector response is modelled with GEANT 4.
- Any additional 'p-p' pile-up is mitigated through charged hadron subtraction and pileup-per-particle identification algorithm.
- Additional jet energy corrections are done as functions of pseudorapidity (η) and pT to account for the detector response.

Reconstruction

- Higgs (H) is constructed as a single AK8 jet with $pT > 450$ GeV.
- The soft drop (SD) algorithm is applied to the jet mass (m_{SD}) to remove soft and wide angle radiation, which reduces the mass of the jets coming from QCD backgrounds events while preserving the mass of heavy boson decays.
- The range of interest is set to $40 < m_{SD} < 201$ GeV.
- To match the tracker acceptance region jets should have $|\eta| < 2.5$.
- 95% of the expected yield is QCD multijet background.
- Other EW processes including Diboson, triboson and $t\bar{t}V$ are estimated from simulations and were found to be negligible.

Reconstruction

- Jet flavor is finally determined by Deep DoubleX DNN algorithm.
- The model is trained to distinguish between two-prolonged H-like signatures of bottom and charm flavors, as well as the QCD background, yielding two per-jet classifiers: charm versus light, DeepDoubleCvL (DDCvL) and charm versus bottom, DeepDoubleCvB (DDCvB).
- The figure (next slide) details out the tagging efficiencies of the two classifiers. $H \rightarrow c\bar{c}$ has an efficiency of 20.6 % while that of $H \rightarrow b\bar{b}$ is 4.8%.
- These efficiencies means the model will be more sensitive to $H \rightarrow c\bar{c}$ detections and less likely to misidentify a b-jet as a c-jet.
- These efficiencies are also applied to $Z \rightarrow c\bar{c}$ and $Z \rightarrow b\bar{b}$ as these classifiers are mass independent.



- The working points are marked with a cross.
- AUC implies to the Area under the curve.
- Both the SR & the CR are divided into 23 evenly spaced bins with jet m_{SD} between 40-200 GeV and 6 p_T bins in the range 475 – 1200 GeV. Data taken from year 2016-2018.

Background Estimations & Corrections

- The $V+$ jet processes and differential $t\bar{t}$ are taken from simulation.
- The $H \rightarrow b\bar{b}$ contribution is taken from simulation and it is fixed to SM expectation.
- The normalizations in the SR and CR are corrected via two scale factors measured in a dedicated $t\bar{t}$ -enriched control region.
- This control region is adapted from the SR, by lowering the H candidate pT threshold, requiring exactly one muon and inverting the selection requirements on missing P_t and b-tagged A_{k4} jets. This gives predominant semi-leptonic $t\bar{t}$ background events in CR.
- The expected SR yield is about 5 times that of $H \rightarrow c\bar{c}$ signal and its impact is negligible w.r.t overall background uncertainty.
- The corrections are displayed in the table next. These corrections are independent of the jet flavor.

Background Estimations & Corrections

Data period	Jet mass correction [GeV]	Jet mass resolution	$N_2^{1,DDT}$, CvB selection	$t\bar{t}$ normalization	CvL selection $t\bar{t}$	CvL selection (W+jets)	CvL selection (signal)
2016	-1.17 ± 0.22	1.021 ± 0.017	0.89 ± 0.02	0.84 ± 0.05	0.93 ± 0.15	0.62 ± 0.09	1.15 ± 0.25
2017	-1.19 ± 0.23	1.019 ± 0.016	0.90 ± 0.02	0.86 ± 0.09	0.93 ± 0.15	0.64 ± 0.09	0.85 ± 0.16
2018	-0.12 ± 0.21	1.090 ± 0.031	0.92 ± 0.02	0.86 ± 0.08	1.00 ± 0.14	0.72 ± 0.08	0.74 ± 0.20

- The efficiency of DDCvL signal selection is estimated using data and simulations samples enriched in $cc(\bar{c})$ pairs and gluon splitting.
- Signal-like events are selected by requiring each of the two SD sub-jets of an AK8 jet to contain a muon, targeting semileptonic decays of b/c hadrons.
- Other systematic uncertainties are assigned to cover mismodelling of H signal, especially ggF & VBF production modes and higher order corrections to W & Z processes.
- The parameter of interest for us is μH i.e. Signal Strength.

Corrections

$$\mu_H = \frac{\sigma_{\text{obs}}(pp \rightarrow H) \cdot \mathcal{B}(H \rightarrow c\bar{c})}{\sigma_{\text{SM}}(pp \rightarrow H) \cdot \mathcal{B}_{\text{SM}}(H \rightarrow c\bar{c})}$$

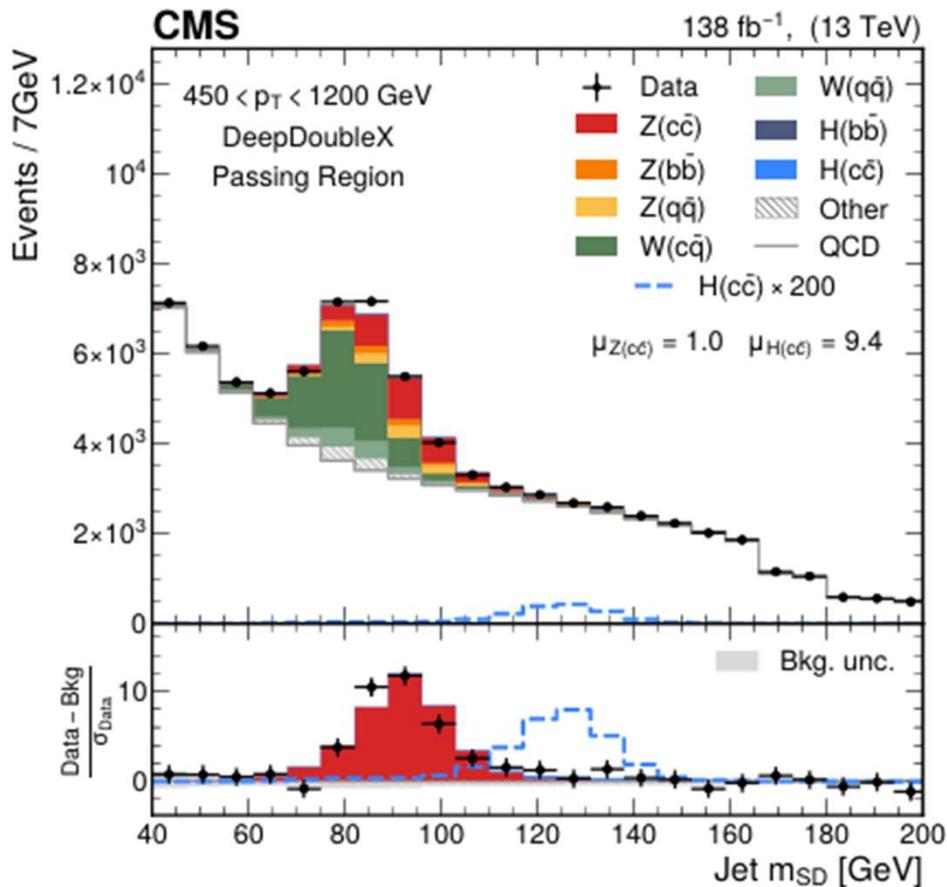
$$\mu = \frac{(\sigma \cdot \text{BR})_{\text{observed}}^{\text{Data}}}{(\sigma \cdot \text{BR})_{\text{expected}}^{\text{Standard Model}}}$$

- μ_H – is dimensionless and is defined as the Observed rate of Higgs production x Branching ratio / SM prediction values.
- If $\mu = 1$, the data matches SM prediction. If $\mu > 1$; excess over SM.
- A maximum likelihood fit is performed to the observed data using binned parameters (m_{SD}, pT), where expected value is sum of the signal contribution + backgrounds accounted by nuisance parameters to account for systematic effects.
- As mentioned earlier to validate the process, μ_Z is also measured using a profile likelihood fit keeping μ_H as nuisance parameter. In this case it was found to be $1.00_{-0.14}^{+0.17}$ (syst) ± 0.08 (theo) ± 0.06 (stat) which corresponds to an excess, both observed and the expected value.

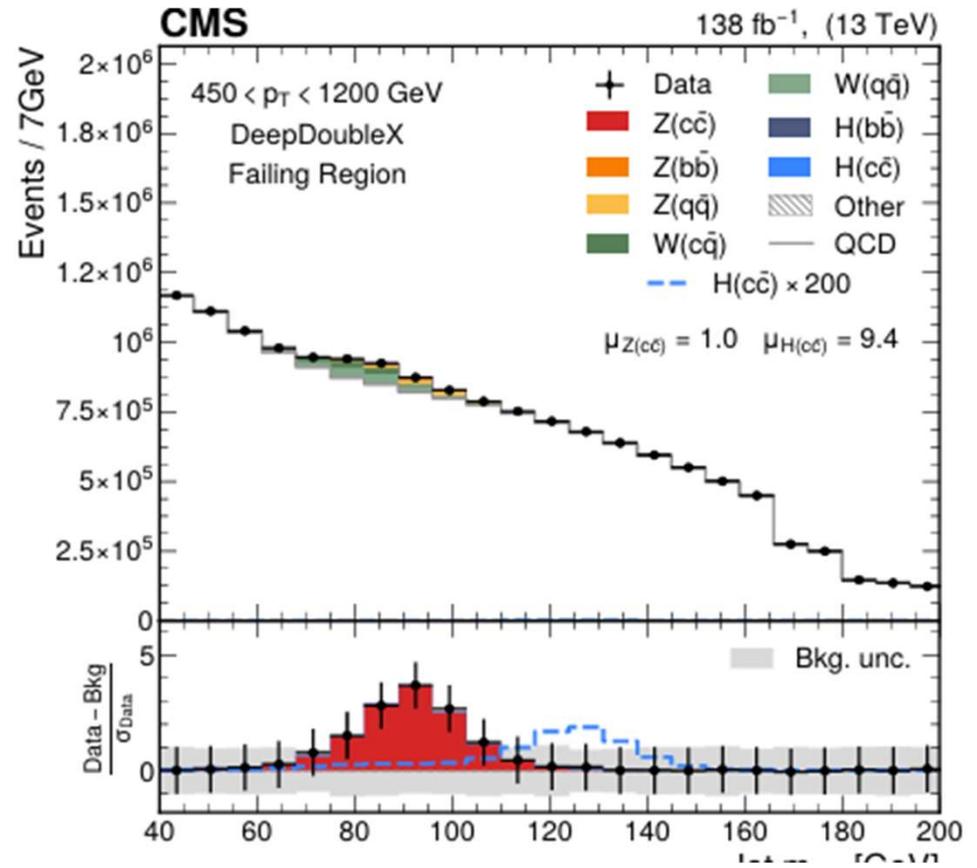
Corrections

- This is primarily due to the systematic uncertainty in the DDCvL signal tagging efficiency.
- To measure μ_H , we fix $\mu_Z \sim 1$, by constraining the expected Z contribution to be within the applicable SM value.
- An observed upper limit is defined for μ_H and found to be 47(39) at 95% confidence level.
- For the best fit value of μ_H (ccbar) = $9.4^{+20.3}_{-19.9}$, the total mSD distributions in the passing and failing regions are shown in the next chart.
- The tt(bar) backgrounds has the highest contributions labelled as “Other”.
- The “blue” dashed line represents $H \rightarrow cc(\text{bar})$ expectation, multiplied by a factor of 200.

Results



Observed Msd – Passing region



Fitted Msd – Failing region

- The data has been collected from 2016 – 2018 combining all the p_T categories.
- The fit is performed using signal-plus-background hypothesis with a single inclusive $\mu_{H(cc)}$.
- Bottom Panel: "Data – Bkg" residuals divided by statistical uncertainty.

To Sum Up...

- We searched a SM (Higgs) (and Z boson for process verification), at High Pt > 450 GeV and later decaying into cc(bar) pairs.
- The integrated Luminosity was kept at 138 fb^{-1} and $\sqrt{s} = 13 \text{ TeV}$ and the data was collected.
- DNN algorithm identifies the jets originating from cc(bar) pairs.
- The $Z \rightarrow \text{cc}(\text{bar})$ process has been observed along with jets at the LHC for the first time with signal strength of $1.00^{+0.19}_{-0.17}$ relative to SM.
- This observation has now set an important reference for any future $X \rightarrow \text{cc}(\text{bar})$ searches.
- The upper limit for $\sigma_H \times \mathcal{B}(H \rightarrow c\bar{c})$ has been set at 47(39) i.e. observed limit = 47 x SM, and expected limit = 39 x SM predictions

Questions ?



Thank you!

