

# Exploring the conformal bootstrap at higher points

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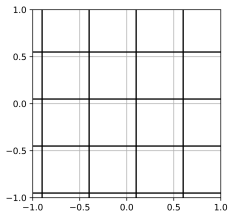
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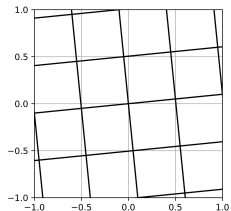
This signals the need for a nonperturbative approach. While general QFTs are still out of the reach of current methods, a subset known as Conformal Field Theories are simple enough that this is possible.

# Conformal symmetry

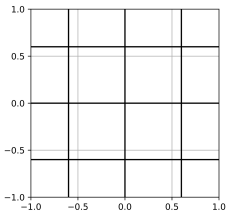
## Translation



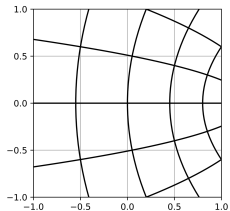
## Rotation



## Dilation



## Special Conformal Transformation



Theories invariant under these transformations are called **conformal field theories (CFTs)**. In a CFT, the two-point function of scalar operators is fixed to be

$$\langle O_1 O_2 \rangle = \frac{\delta_{\Delta_1 \Delta_2}}{(x_{12}^2)^{\frac{\Delta_1 + \Delta_2}{2}}} \quad (1)$$

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while the three-point function is

$$\langle O_1 O_2 O_3 \rangle = \frac{c_{123}}{(x_{12}^2)^{\frac{\Delta_1 + \Delta_2 - \Delta_3}{2}} (x_{13}^2)^{\frac{\Delta_1 + \Delta_3 - \Delta_2}{2}} (x_{23}^2)^{\frac{\Delta_2 + \Delta_3 - \Delta_1}{2}}} \quad (2)$$

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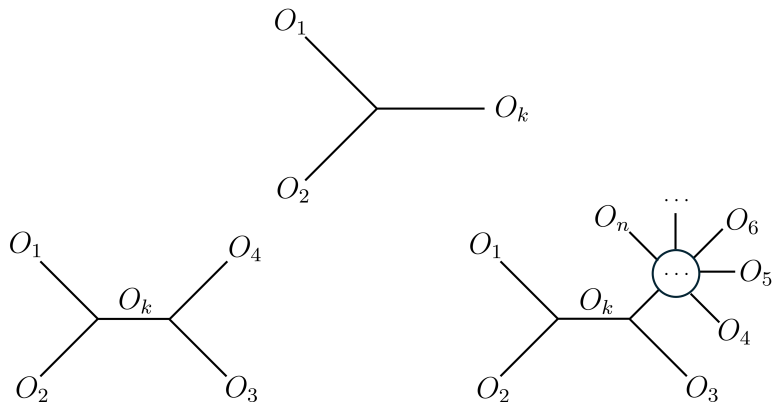
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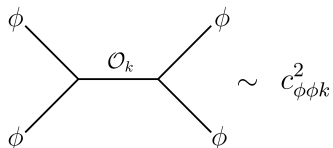
## Remark

The CFT data  $\{\Delta_i, c_{ijk}\}$  encodes all the information about a CFT.

# Conformal symmetry



In general it is easier to consider correlators of identical scalars



A four-point tree-level Feynman diagram. It consists of a central horizontal line representing an internal propagator labeled  $\mathcal{O}_k$ . From the left end of this line, two lines branch out to the left, each ending in a vertex labeled  $\phi$ . Similarly, from the right end of the central line, two lines branch out to the right, each ending in a vertex labeled  $\phi$ . To the right of the diagram is a tilde symbol  $\sim$  followed by the coefficient  $c_{\phi\phi k}^2$ .

# Bootstrap

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$$\begin{array}{c} \phi \\ \diagdown \\ \text{---} \\ \diagup \\ \phi \end{array} \text{---} \text{---} \text{---} \begin{array}{c} \phi \\ \diagup \\ \text{---} \\ \diagdown \\ \phi \end{array} \sim c_{\phi\phi k}^2$$

These functions satisfy the **bootstrap equations**

$$\begin{array}{c} \phi \\ \diagdown \\ \text{---} \\ \diagup \\ \phi \end{array} \text{---} \text{---} \text{---} \begin{array}{c} \phi \\ \diagup \\ \text{---} \\ \diagdown \\ \phi \end{array} = \begin{array}{c} \phi \\ \diagdown \\ \text{---} \\ \diagup \\ \phi \end{array} \begin{array}{c} \phi \\ \diagup \\ \text{---} \\ \diagdown \\ \phi \end{array} \text{---} \text{---} \text{---} \begin{array}{c} \phi \\ \diagup \\ \text{---} \\ \diagdown \\ \phi \end{array}$$

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Five-point functions also contain information about an infinite number of four-point functions through the exchanged operators.



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**Thank you for your time!**