

LABORATÓRIO DE INSTRUMENTAÇÃO E FÍSICA EXPERIMENTAL DE PARTÍCULAS







LOW-x PROJECTILE DENSITY CONTRIBUTIONS IN THE DILUTE-**DENSE CGC FRAMEWORK FOR TWO-PARTICLE CORRELATIONS**

Víctor Vila

[based on arXiv:2303.08711, in collaboration with Cyrille Marquet and Anderson Kendi] LIP — Lisbon Seminar, November 2

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Motivation The ridge







ATLAS Collaboration, "Measurements of long-range azimuthal anisotropies and associated Fourier coefficients for *pp* collisions at $s^{1/2}=5.02$ and 13 TeV and *p*+Pb collisions at $s_{NN}^{1/2}=5.02$ TeV with the ATLAS detector", ATLAS-CONF-2016-026 [arXiv:1609.06213 [nucl-ex]].

• One of the main unveilings of the heavy-ion programs at RHIC and the LHC.

First noticed in relativistic heavy-ion data from RHIC.

An incentive to seek the origin of particle correlations.

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MECHANISMS behind the RIDGE in pp and pA collisions:

PARTICLE CORRELATIONS: ROOTS in the VERY EARLY STAGES of the COLLISION.

GLASMA GRAPH APPROXIMATION — DILUTE-DILUTE LIMIT of the CGC framework:

EXTREMELY EFFECTIVE.

ANISOTROPIC DISTRIBUTION of particles MOTIVATES decomposition in FOURIER HARMONICS with the corresponding COEFFICIENTS v_n.

ALTERNATIVE MECHANISM: HYDRODYNAMICAL EVOLUTION logic consistent with data.

[argument taken from AA collisions]

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Motivation The accidental symmetry breaking of the CGC



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DENSE HADRON OR NUCLEUS: $ho \sim 1/g_s^2$ DILUTE HADRON OR NUCLEUS: $ho \sim 1$

THE DILUTE PROJECTILE LIMIT OF THE DENSE-DENSE RESULT DOES NOT YIELD THE CORRECT DILUTE-DENSE ANSWER, BUT IT HAS BEEN WIDELY USED IN THE CASE OF TWO-GLUON PRODUCTION

DILUTE-DENSE LIMIT OF THE DENSE-DENSE RESULT: $g_s^4 \rho_P^4$ incomplete!

FULL DILUTE-DENSE RESULT: $g_s^4 \rho_P^4 \ g_s^4 \rho_P^3 \ g_s^4 \rho_P^2$

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ODD HARMONICS?

The two-gluon emission beyond the GGA The amplitude

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$$\frac{dN}{d^2pd^2kd\eta d\xi} = \underbrace{\sigma^4 + \sigma^3 + \sigma^2}_{u,z,\bar{u},\bar{z}} p_{,T} \\ \text{A. Kover and M. Lublinsky, "Angular Correlations in Gloon Production at High Energy," phys. Rev. DB3 (2011) 034017 \\ \text{arXiv:1012.3398 (hep-ph]}. \\ \hline \sigma^4 = \int_{u,z,\bar{u},\bar{z}} e^{ik\cdot(z-\bar{z})+ip\cdot(u-\bar{u})} \int_{x_1,x_2,\bar{x}_1,\bar{x}_2} f(\bar{z} - \bar{x}_1) \cdot f(z - x_1) f(\bar{u} - \bar{x}_2) \cdot f(u - x_2) \\ \times \left\{ \rho_{x_1}[S_z^{\dagger} - S_{x_1}^{\dagger}][S_{\bar{z}} - S_{x_2}^{\dagger}]\rho_{\bar{x}_2} \right\} \\ \times \left\{ \rho_{x_1}[S_z^{\dagger} - S_{x_1}^{\dagger}][S_{\bar{z}} - S_{x_2}^{\dagger}]\rho_{\bar{x}_2} \right\} \\ - \frac{1}{2}f(\bar{z} - \bar{x}_1) \cdot f(z - x_1)f(\bar{u} - \bar{x}_1) \cdot f(u - x_2) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_1}[S_z^{\dagger} - S_{x_1}^{\dagger}][S_{\bar{z}} - S_{x_1}]\bar{\rho}_{\bar{x}_1}[S_{\bar{u}}^{\dagger} + S_{x_1}^{\dagger}][S_u - S_{x_2}]\bar{\rho}_{x_2} \right\} \\ + \frac{1}{2}f(\bar{z} - \bar{x}_1) \cdot f(z - x_1)f(\bar{u} - \bar{x}_2) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{x_2}^{\dagger}][S_u + S_{x_1}]\bar{\rho}_{x_1}[S_z^{\dagger} - S_{x_1}^{\dagger}][S_z - S_{x_1}]\bar{\rho}_{\bar{x}_1} \right\} \\ + f(\bar{z} - \bar{u}) \cdot f(z - x_1)f(\bar{u} - \bar{x}_1) \cdot f(u - x_2) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{x_2}^{\dagger}][S_u + S_{x_1}]\bar{\rho}_{x_1}[S_z^{\dagger} - S_{x_1}^{\dagger}][S_z - S_{x_1}]\bar{\rho}_{\bar{x}_1} \right\} \\ + f(\bar{z} - \bar{u}) \cdot f(z - x_1)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{x_1}^{\dagger}][S_z - S_{\bar{u}}]\bar{\rho}_{\bar{x}_1}S_{\bar{u}}^{\dagger} - S_{x_2}]\bar{\rho}_{x_2} \right\} \\ - f(\bar{z} - \bar{x}_1) \cdot f(z - u)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{x_1}^{\dagger}][S_z - S_{\bar{u}}^{\dagger}]\bar{\rho}_{\bar{x}_1}S_{\bar{u}}^{\dagger} - S_{x_1}^{\dagger}]\bar{\rho}_{\bar{x}_1} \right\} \\ - f(\bar{z} - \bar{u}) \cdot f(z - u)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{\bar{u}}^{\dagger}][S_{\bar{u}} - S_{\bar{u}}^{\dagger}]\bar{\rho}_{\bar{x}_1}S_{\bar{u}}^{\dagger} - S_{\bar{u}}^{\dagger}]S_{\bar{u}}^{\dagger}S_{\bar{u}}^{\dagger} \right\} \\ - f(\bar{z} - \bar{u}) \cdot f(z - u)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{\bar{u}}^{\dagger}]S_{\bar{u}}S_{\bar{u}}^{\dagger}S_{\bar{u}}^{\dagger} \right\} \\ - f(\bar{z} - \bar{u}) \cdot f(z - u)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr}\left\{ \bar{\rho}_{x_2}[S_{\bar{u}}^{\dagger} - S_{\bar{u}}^{\dagger}]S_{\bar{u}}S_{\bar{u}}^{\dagger}S_{\bar{u}}^{\dagger}S_{\bar{u}}^{\dagger} \right\} \\ - f(\bar{z} - \bar{u}) \cdot f(z - u)f(\bar{u} - \bar{x}_1) \cdot f(u - x_1) \\ \times \operatorname{Tr$$

The two-gluon emission beyond the GGA **<u>Projectile</u>** and target averaging

Kendi Kohara, A., Marquet, C., Vila, V., arXiv:2303.08711 [hep-ph]

$$\left\langle \hat{\rho}_{x_{2}}^{c} \hat{\rho}_{x_{1}}^{a} \hat{\rho}_{x_{2}}^{b} \hat{\rho}_{x_{1}}^{b} \right\rangle_{P} = \left\langle \hat{\rho}_{x_{2}}^{c} \hat{\rho}_{x_{1}}^{a} \right\rangle_{P} \langle \hat{\rho}_{x_{2}}^{a} \hat{\rho}_{x_{2}}^{b} \rangle_{P} \langle \hat{\rho}_{x_{1}}^{a} \hat{\rho}_{x_{2}}^{b} \rangle_{P} \langle \hat{\rho}_{x_{1}}^{b} \langle \hat{\rho}_{x_{1}}^{b} \rangle_{P} \langle \hat{\rho}_{x_{1$$

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The two-gluon emission beyond the GGA **<u>Projectile</u>** and target averaging



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$$D(\bar{u}, u) = \frac{1}{N_c^2 - 1} \left\langle \operatorname{Tr} S_{\bar{u}}^{\dagger} S_u \right\rangle_T$$
$$Q(\bar{u}, u, z, \bar{z}) = \frac{1}{N_c^2 - 1} \left\langle \operatorname{Tr} S_{\bar{u}}^{\dagger} S_u S_z^{\dagger} S_{\bar{z}} \right\rangle_T$$
$$DD(\bar{u}, u, z, \bar{z}) = \frac{1}{(N_c^2 - 1)^2} \left\langle \operatorname{Tr} S_{\bar{u}}^{\dagger} S_u \operatorname{Tr} S_z^{\dagger} S_{\bar{z}} \right\rangle_T$$

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$$\sigma_4^{(ii)}(k,p) = \langle \tilde{\sigma}(k)\tilde{\sigma}(p)\rangle_T$$
$$\tilde{\sigma}(k) = \mu^2 \int_{z,\bar{z}} e^{ik \cdot (z-\bar{z})} \int_x f(\bar{z}-x) \cdot f(z-x) \operatorname{Tr}[S_z^{\dagger} - S_x^{\dagger}][S_{\bar{z}} - S_x]$$

UNCORRELATED PIECE

TARGET AVERAGES OF ADJOINT WILSON LINES

Kendi Kohara, A., Marquet, C., Vila, V., arXiv:2303.08711 [hep-ph]

$$\begin{split} \sigma_4^{(i)}(k,p) &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z})+ip \cdot (u-\bar{u})} \\ &\times \int_{x_1,\bar{x}_1} f(\bar{z}-\bar{x}_1) \cdot f(z-x_1) \ f(\bar{u}-\bar{x}_1) \cdot f(u-x_1) \\ &\times (N_c^2-1) \mu^4 \Big\{ Q(\bar{z},z,u,\bar{u}) - Q(\bar{z},z,u,\bar{x}_1) - Q(\bar{z},z,x_1,\bar{u}) \\ &+ Q(\bar{z},z,x_1,\bar{x}_1) - Q(\bar{x}_1,z,u,\bar{u}) + D(z,u) + Q(\bar{x}_1,z,x_1,\bar{u}) - D(z,x_1) \\ &- Q(\bar{z},x_1,u,\bar{u}) + Q(\bar{z},x_1,u,\bar{x}_1) + D(\bar{z},\bar{u}) - D(\bar{z},\bar{x}_1) + Q(\bar{x}_1,x_1,u,\bar{u}) \\ &- D(x_1,u) - D(\bar{x}_1,\bar{u}) + 1 \Big\} \end{split}$$

CORRELATED PIECES

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$$(N_c^2 - 1)^2 \left\langle S_{\bar{u}}^{ab} S_u^{cd} S_z^{ef} S_{\bar{z}}^{gh} \right\rangle_T \stackrel{AEM}{\Longrightarrow} \delta_{ac} \delta_{bd} \delta_{eg} \delta_{fh} D(\bar{u}, u) D(z, \bar{z}) + \delta_{ag} \delta_{bh} \delta_{ce} \delta_{df} D(\bar{u}, \bar{z}) D(z, u) + \delta_{ae} \delta_{bf} \delta_{cg} \delta_{dh} D(\bar{u}, z) D(u, \bar{z})$$

Kendi Kohara, A., Marquet, C., Vila, V., arXiv:2303.08711 [hep-ph]

AEM

$$Q(\bar{u}, u, z, \bar{z}) \longrightarrow D(\bar{u}, u) D(z, \bar{z}) + D(\bar{u}, \bar{z}) D(z, u) + \frac{1}{N_c^2 - 1} D(\bar{u}, z) D(u, \bar{z})$$
$$DD(\bar{u}, u, z, \bar{z}) \longrightarrow D(\bar{u}, u) D(z, \bar{z}) + \frac{1}{(N_c^2 - 1)^2} \Big[D(\bar{u}, \bar{z}) D(z, u) + D(\bar{u}, z) D(u, \bar{z}) \Big]$$

$$\left\langle S_u^{b'm} T_{ma}^c S_z^{\dagger ac'} S_{\bar{z}}^{c'b} T_{bn}^c S_{\bar{u}}^{\dagger nb'} \right\rangle_T \longrightarrow N_c (N_c^2 - 1) D(\bar{u}, u) D(z, \bar{z}) - N_c D(\bar{u}, z) D(u, \bar{z})$$

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The two-gluon emission beyond the GGA **Projectile and target averaging**

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AVERAGED SIGMA²

$$\begin{split} \sigma_4^{(i)}(k,p) &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \int_{x,\bar{x}} f(\bar{z}-\bar{x}) \cdot f(z-x) f(\bar{u}-\bar{x}) \cdot f(u-x) \\ &\times (N_c^2 - 1) \mu^4 \Big\{ D(\bar{z},z) D(u,\bar{u}) + D(\bar{z},\bar{u}) D(z,u) - D(\bar{z},z) D(u,\bar{x}) \\ &- D(\bar{z},\bar{x}) D(z,u) - D(\bar{z},z) D(x,\bar{u}) - D(\bar{z},\bar{u}) D(z,x) \\ &+ D(\bar{z},z) D(x,\bar{x}) + D(\bar{z},\bar{x}) D(z,x) - D(\bar{x},z) D(u,\bar{u}) \\ &- D(\bar{x},\bar{u}) D(z,u) + D(z,u) + D(\bar{x},z) D(x,\bar{u}) + D(\bar{x},\bar{u}) D(z,x) \\ &- D(z,x) - D(\bar{z},x) D(u,\bar{u}) - D(\bar{z},\bar{u}) D(x,u) + D(\bar{z},x) D(u,\bar{x}) \\ &+ D(\bar{z},\bar{x}) D(x,u) + D(\bar{z},\bar{u}) - D(\bar{z},\bar{x}) + D(\bar{x},x) D(u,\bar{u}) \\ &+ D(\bar{x},\bar{u}) D(x,u) - D(x,u) - D(\bar{x},\bar{u}) + 1 \Big\} \\ &= \sigma_4^{(iii)}(k,-p) \end{split}$$

$$\begin{split} \sigma_4^{(ii)}(k,p) = &\bar{\sigma}(k)\bar{\sigma}(p), \\ \bar{\sigma}(k) = \mu^2(N_c^2 - 1)\int_{z,\bar{z}} e^{ik\cdot(z-\bar{z})}\int_x f(\bar{z}-x)\cdot f(z-x) \\ &\times \left[D(z,\bar{z}) - D(z,x) - D(x,\bar{z}) + 1\right] \end{split}$$

Kendi Kohara, A., Marquet, C., Vila, V., arXiv:2303.08711 [hep-ph]

$$\begin{split} \sigma_{2}^{(i)}(k,p) &= \int_{u,z,\bar{u},\bar{z}} e^{ik\cdot(z-\bar{z})+ip\cdot(u-\bar{u})} \int_{x} f(\bar{z}-x) \cdot f(z-x) \ f(\bar{u}-x) \cdot f(u-x) \\ &\times \frac{1}{4} N_{c}^{3} \mu^{2} \Big\{ D(u,\bar{u}) D(z,\bar{z}) + D(u,x) D(z,\bar{z}) - D(u,\bar{u}) D(z,x) \\ &- D(u,x) D(z,x) + D(x,\bar{u}) D(z,\bar{z}) + D(z,\bar{z}) - D(x,\bar{u}) D(z,x) - D(z,x) \\ &- D(u,\bar{u}) D(x,\bar{z}) - D(u,x) D(x,\bar{z}) + D(u,\bar{u}) + D(u,x) - D(x,\bar{u}) D(x,\bar{z}) \\ &- D(x,\bar{z}) + D(x,\bar{u}) + 1 \Big\} \end{split}$$

$$\begin{split} \sigma_2^{(ii)}(k,p) &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \times \int_x f(\bar{z}-\bar{u}) \cdot f(z-u) \ f(\bar{u}-x) \cdot f(u-x) \\ &\times N_c^3 \mu^2 D(u,\bar{u}) \Big\{ D(z,\bar{z}) - D(z,\bar{u}) - D(u,\bar{z}) + D(u,\bar{u}) \Big\} \end{split}$$

$$\begin{split} \sigma_{2}^{(iii)}(k,p) &= -\int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \times \int_{x} f(\bar{z}-x) \cdot f(z-u) \ f(\bar{u}-x) \cdot f(u-x) \\ & \times \frac{1}{2} N_{c}^{3} \mu^{2} \Big\{ D(u,\bar{u}) D(z,\bar{z}) - D(u,\bar{u}) D(z,x) - D(u,\bar{u}) D(u,\bar{z}) \\ & + D(u,\bar{u}) D(u,x) + D(u,x) D(z,\bar{z}) - D(u,x) D(z,x) - D(u,x) D(u,\bar{z}) \\ & + D^{2}(u,x) \Big\} \\ \begin{pmatrix} (iv) \\ 2 \end{pmatrix} (k,p) &= -\int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \times \int_{u} f(\bar{z}-\bar{u}) \cdot f(z-x) \ f(\bar{u}-x) \cdot f(u-x) \end{split}$$

$$\begin{array}{l} & J_{u,z,\bar{u},\bar{z}} & J_{x} \\ & \times \frac{1}{2} N_{c}^{3} \mu^{2} \Big\{ D(u,\bar{u}) D(z,\bar{z}) - D(u,\bar{u}) D(z,\bar{u}) - D(u,\bar{u}) D(x,\bar{z}) \\ & + D(u,\bar{u}) D(x,\bar{u}) + D(x,\bar{u}) D(z,\bar{z}) - D(x,\bar{u}) D(z,\bar{u}) - D(x,\bar{u}) D(x,\bar{z}) \\ & + D^{2}(x,\bar{u}) \Big\} \end{array}$$

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The two-gluon emission beyond the GGA Final expressions



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The two-gluon emission beyond the GGA Final expressions

$$\begin{aligned} \sigma_2^{(i)}(k,p) &= \frac{\sigma_0^2}{2N_2} \frac{Q_s^4}{k^6 p^6} e^{-\frac{k^2 + p^2}{2Q_s^2}} & N_2 = \frac{(N_c^2 - 1)^2}{N_c^3} S_\perp \mu^2 \\ &\times \left[k^4 + e^{\frac{k^2}{2Q_s^2}} k^2 Q_s^2 + 4e^{\frac{k^2}{2Q_s^2}} Q_s^4 \right] \left[\left(2e^{\frac{p^2}{2Q_s^2}} - 1 \right) p^4 + e^{\frac{p^2}{2Q_s^2}} p^2 Q_s^2 + 4e^{\frac{p^2}{2Q_s^2}} Q_s^4 \right] \\ &+ (k \leftrightarrow p) \end{aligned}$$

Kendi Kohara, A., Marquet, C., Vila, V., arXiv:2303.08711 [hep-ph]

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Vila, V., arXiv:2303.087

Marquet,

Kendi Kohara,

$$\begin{split} & \tau_4^{(i)} = \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \int_{x,\bar{x}} f(\bar{z}-\bar{x}) \cdot f(z-x) \ f(\bar{u}-\bar{x}) \cdot f(u-x) \\ & \times (N_c^2 - 1) \mu^4 \Big\{ D(\bar{z},z) D(u,\bar{u})_{[1]} + D(\bar{z},\bar{u}) D(z,u)_{[2]} - D(\bar{z},z) D(u,\bar{x})_{[3]} \\ & - D(\bar{z},\bar{x}) D(z,u)_{[4]} - D(\bar{z},z) D(x,\bar{u})_{[5]} - D(\bar{z},\bar{u}) D(z,x)_{[6]} + D(\bar{z},z) D(x,\bar{x})_{[7]} \\ & + D(\bar{z},\bar{x}) D(z,x)_{[8]} - D(\bar{x},z) D(u,\bar{u})_{[9]} - D(\bar{x},\bar{u}) D(z,u)_{[10]} + D(z,u)_{[11]} \\ & + D(\bar{x},z) D(x,\bar{u})_{[12]} + D(\bar{x},\bar{u}) D(z,x)_{[13]} - D(z,x)_{[14]} - D(\bar{z},x) D(u,\bar{u})_{[15]} \\ & - D(\bar{z},\bar{u}) D(x,u)_{[16]} + D(\bar{z},x) D(u,\bar{x})_{[17]} + D(\bar{z},\bar{x}) D(x,u)_{[18]} + D(\bar{z},\bar{u})_{[19]} \\ & - D(\bar{z},\bar{x})_{[20]} + D(\bar{x},x) D(u,\bar{u})_{[21]} + D(\bar{x},\bar{u}) D(x,u)_{[22]} - D(x,u)_{[23]} \\ & - D(\bar{x},\bar{u})_{[24]} + 1_{[25]} \Big\} \end{split}$$

$$\begin{split} \sigma_4^{(i.[1])} &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \int_{x,\bar{x}} f(\bar{z}-\bar{x}) \cdot f(z-x) \ f(\bar{u}-\bar{x}) \cdot f(u-x) \\ &\times (N_c^2 - 1) \mu^4 D(\bar{z},z) D(u,\bar{u}) \end{split}$$

$$\begin{split} \sigma_4^{(i.[1])} &= g_s^4 (N_c^2 - 1) \mu^4 \frac{1}{(2\pi)^4} \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \int_{x,\bar{x}} \\ &\times \frac{d^2 q_1}{2\pi} (-i) e^{iq_1 \cdot (\bar{z}-\bar{x})} \frac{q_1^i}{q_1^2} \frac{d^2 q_2}{2\pi} (-i) e^{iq_2 \cdot (z-x)} \frac{q_2^i}{q_2^2} \frac{d^2 q_3}{2\pi} (-i) e^{iq_3 \cdot (\bar{u}-\bar{x})} \frac{q_3^j}{q_3^2} \\ &\times \frac{d^2 q_4}{2\pi} (-i) e^{iq_4 \cdot (u-x)} \frac{q_4^j}{q_4^2} \frac{d^2 t_1}{2\pi} e^{-it_1 \cdot (\bar{z}-z)} \frac{1}{Q_s^2} e^{-\frac{t_1^2}{2Q_s^2}} \frac{d^2 t_2}{2\pi} e^{-it_2 \cdot (u-\bar{u})} \frac{1}{Q_s^2} e^{-\frac{t_2^2}{2Q_s^2}} \end{split}$$

$$\int d^2 x e^{i(k-p)\cdot x} = (2\pi)^2 \delta^{(2)}(k-p) \qquad S_{\perp} \equiv \int_{\text{proton}} d^2 x$$

$$\begin{split} \sigma_4^{(i.[1])} &= g_s^4 (N_c^2 - 1) \mu^4 \frac{1}{Q_s^4} \int \frac{d^2 t_1}{2\pi} \frac{(t_1 + k)^i}{(t_1 + k)^2} \frac{(t_1 + k)^i}{(t_1 + k)^2} e^{-\frac{t_1^2}{2Q_s^2}} \\ &\times \frac{d^2 t_2}{2\pi} \frac{(t_2 - p)^j}{(t_2 - p)^2} \frac{(t_2 - p)^j}{(t_2 - p)^2} e^{-\frac{t_2^2}{2Q_s^2}} \int_{x,\bar{x}} e^{-i[(t_1 + k) - (t_2 - p)] \cdot (x - \bar{x})} \end{split}$$

$$\begin{split} \sigma_4^{(i.[1])} &= g_s^4 (N_c^2 - 1) \mu^4 \frac{1}{Q_s^4} e^{-\frac{k^2 + p^2}{2Q_s^2}} \int \frac{d^2 s_1}{2\pi} e^{-\frac{s_1^2 - 2k \cdot s_1}{2Q_s^2}} \frac{1}{s_1^2} \frac{d^2 s_2}{2\pi} e^{-\frac{s_2^2 + 2p \cdot s_2}{2Q_s^2}} \frac{1}{s_2^2} \int_{x,\bar{x}} e^{i(s_1 - s_2) \cdot (x - \bar{x})} \\ &= g_s^4 (N_c^2 - 1) \mu^4 \frac{1}{Q_s^4} e^{-\frac{k^2 + p^2}{2Q_s^2}} \int d^2 s e^{-\frac{2s^2 + 2(p - k) \cdot s}{2Q_s^2}} \frac{1}{s^4} \int d^2 x \end{split}$$

$$T_{0,4} = \int d^2 q e^{-\frac{2q^2 + 2(k+p) \cdot q}{Q_s^2}} \frac{1}{q^4}$$

$$\approx \pi Q_s^2 e^{\frac{(k+p)^2}{2Q_s^2}} \frac{2^4}{(k+p)^4} \left[\frac{1}{2} + \frac{2^2 Q_s^2}{(k+p)^2} + \frac{9}{4} \frac{2^4 Q_s^4}{(k+p)^4} \right] + \frac{1}{Q_s^2} \mathcal{O}\left(\frac{Q_s^{10}}{(k+p)^{10}}\right)$$

$$\sigma_4^{(i.[1])} \approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \frac{2^4}{(k-p)^4} \left[\frac{1}{2} + \frac{2^3 Q_s^2}{(k-p)^2} + \frac{9}{4} \frac{2^6 Q_s^4}{(k-p)^4} \right]$$

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Marquet, C., Vila, V., arXiv:2303.08711

Kendi Kohara, A.,

$$\begin{split} I_{1,2} &= \int d^2 q e^{-\frac{q^2 + 2k + q}{Q_s^2}} \frac{q^i}{q^2} = \pi Q_s^2 \frac{k^i}{k^2} \left(1 - e^{\frac{k^2}{Q_s^2}}\right) \\ T_{0,2} &= \int d^2 q e^{-\frac{2q^2 + 2(k + p) \cdot q}{Q_s^2}} \frac{1}{q^2} \approx \pi Q_s^2 e^{\frac{(k + p)^2}{2Q_s^2}} \frac{2^2}{(k + p)^2} \left[\frac{1}{2} + \frac{1}{4} \frac{2^2 Q_s^2}{(k + p)^2} + \frac{1}{4} \frac{2^4 Q_s^4}{(k + p)^4}\right] \\ T_{1,4}^i &= \int d^2 q e^{-\frac{2q^2 + 2(k + p) \cdot q}{Q_s^2}} \frac{q^i}{q^4} \approx -\pi Q_s^2 e^{\frac{(k + p)^2}{2Q_s^2}} \frac{(k + p)^i}{(k + p)^2} \frac{2^2}{(k + p)^2} \left[1 + \frac{2^2 Q_s^2}{(k + p)^2} + \frac{3}{2} \frac{2^4 Q_s^4}{(k + p)^4}\right] \\ T_{2,4}^{ij} &= \int d^2 q e^{-\frac{2q^2 + 2(k + p) \cdot q}{Q_s^2}} \frac{q^i q^j}{q^4} \\ &\approx \pi Q_s^2 e^{\frac{(k + p)^2}{Q_s^2}} \frac{2^2}{(k + p)^2} \left[\frac{1}{2} \frac{(k + p)^i(k + p)^j}{(k + p)^2} + \frac{\delta^{ij}}{8} \frac{2^2 Q_s^2}{(k + p)^2} \left(1 + \frac{2^2 Q_s^2}{(k + p)^2}\right)\right] \\ \sigma_4^{(i,[2])} &\approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k + p)^2}{4Q_s^2}} \frac{p \cdot (k - p)}{p^2(k - p)^2} \frac{2^2}{(k - p)^2} \left[1 + \frac{2^3 Q_s^2}{(k - p)^2} + \frac{3}{2} \frac{2^6 Q_s^4}{(k - p)^4}\right] \\ \sigma_4^{(i,[3])} &= \sigma_4^{(i,[5])} \approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k + p)^2}{4Q_s^2}} \frac{p \cdot (k - p)}{p^2(k - p)^2} \frac{2^2}{(k - p)^2} \left[1 + \frac{2^3 Q_s^2}{(k - p)^2} + \frac{3}{2} \frac{2^6 Q_s^4}{(k - p)^4}\right] \\ \sigma_4^{(i,[4])} &= \sigma_4^{(i,[6])} = \sigma_4^{(i,[10])} = \sigma_4^{(i,[16])} \approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp (2\pi^2) \delta^{(2)}(k + p) \\ &\times \left(e^{-\frac{k^2}{2Q_s^2}} - 1\right) \frac{k^i k^j}{k^4} \left\{\frac{2^2}{2^2} \left[\frac{1}{2} \frac{k^i k^j}{k^2} + \frac{\delta^{ij}}{8} \frac{2^2 Q_s^2}{k^2} \left(1 + \frac{2^2 Q_s^2}{k^2}\right)\right]\right\} \\ \sigma_4^{(i,[7])} &\approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k + p)^2}{4Q_s^2}} \frac{1}{p^2} \frac{2^2}{(k - p)^2} \left[\frac{1}{2} + \frac{1}{4} \frac{2^3 Q_s^2}{(k - p)^2} + \frac{1}{4} \frac{2^6 Q_s^4}{(k - p)^4}\right] \\ \end{array}$$

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$$\begin{split} &\sigma_4^{(i,\mathrm{II})} \approx g_s^4(N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \frac{2^4}{(k-p)^4} \left[\frac{1}{2} + \frac{2^3 Q_s^2}{(k-p)^2} + \frac{9}{4} \frac{2^6 Q_s^4}{(k-p)^4} \right] \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} 2 \left(\frac{p \cdot (k-p)}{p^2(k-p)^2} + \frac{k \cdot (p-k)}{k^2(k-p)^2} \right) \\ &\times \frac{2^2}{(k-p)^2} \left[1 + \frac{2^3 Q_s^2}{(k-p)^2} + \frac{3}{2} \frac{2^6 Q_s^4}{(k-p)^4} \right] \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \left(\frac{1}{k^2} + \frac{1}{p^2} \right) \frac{2^2}{(k-p)^2} \left[\frac{1}{2} + \frac{1}{4} \frac{2^3 Q_s^2}{(k-p)^2} + \frac{1}{4} \frac{2^6 Q_s^4}{(k-p)^4} \right] \\ &- g_s^4(N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \\ &\times \frac{2^2}{(k-p)^2} \left[\frac{k \cdot (k-p) p \cdot (k-p)}{k^2 p^2(k-p)^2} + \frac{2k \cdot p}{k^2 p^2} \frac{1}{4} \frac{2^3 Q_s^2}{(k-p)^2} + \frac{2k \cdot p}{k^2 p^2} \frac{1}{4} \frac{2^6 Q_s^4}{(k-p)^4} \right] \\ &\sigma_4^{(i,\mathrm{II})} \approx g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left\{ \left[\frac{1}{2} \frac{k^i k^j}{k^2} + \frac{\delta^{ij}}{8} \frac{2^2 Q_s^2}{k^2} \left(1 + \frac{2^2 Q_s^2}{k^2} \right) \right] \right\}^2 \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^2} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &\times \frac{2k^i k^j}{k^4} \left\{ \left[\frac{1}{2} \frac{k^i k^j}{k^2} + \frac{\delta^{ij}}{8} \frac{2^2 Q_s^2}{k^2} \left(1 + \frac{2^2 Q_s^2}{k^2} \right) \right] \right\} \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^2} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right)^2 \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^2} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right)^2 \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^2} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^2} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left(e^{-\frac{k^2}{2Q_s^2}} - 1 \right) \\ &+ g_s^4(N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \, \delta^{(2)}(k+p) \frac{2^2}{k^4} \left(e^{-\frac{k^2}{2Q_s^2}} - 1$$

$$\begin{split} \sigma_4^{(i.\mathrm{I})} &\approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \\ &\times \left[\frac{2\left(k^4 + p^4 + 2k^2p^2\cos^2\phi\right)}{k^2p^2\left(k^2 + p^2 - 2kp\cos\phi\right)^2} + \frac{8Q_s^2\left(k^2 + p^2 + 2kp\cos\phi\right)^2}{k^2p^2\left(k^2 + p^2 - 2kp\cos\phi\right)^3} \right. \\ &+ \frac{64Q_s^4\left(k^4 + 16k^2p^2 + p^4 + 8kp\cos\phi(k^2 + p^2) + 2k^2p^2\cos(2\phi)\right)}{k^2p^2\left(k^2 + p^2 - 2kp\cos\phi\right)^4} \end{split}$$

$$\begin{split} \sigma_4^{(i.\mathrm{I})} &= g_s^4 (N_c^2 - 1) \mu^4 S_\perp \frac{2\pi}{Q_s^2} e^{-\frac{(k+p)^2}{4Q_s^2}} \\ &\times \left[\frac{2\left(k^4 + p^4 + 2\left(k \cdot p\right)^2\right)}{k^2 p^2 \left(k - p\right)^4} + \frac{8Q_s^2 \left(k + p\right)^4}{k^2 p^2 \left(k - p\right)^6} \right. \\ &+ \frac{64Q_s^4 \left(k^4 + 4(k \cdot p)^2 + p^4 + 8(k \cdot p)(k^2 + p^2) + 14k^2 p^2\right)}{k^2 p^2 \left(k - p\right)^8} \right] \end{split}$$

$$\sigma_4^{(i.\mathrm{II})} \approx g_s^4 (N_c^2 - 1) \mu^4 S_\perp (2\pi)^2 \,\delta^{(2)} (k+p) \frac{2^2}{k^{12}} e^{-\frac{k^2}{Q_s^2}} \left(k^4 + k^2 e^{\frac{k^2}{2Q_s^2}} Q_s^2 + 2^2 e^{\frac{k^2}{2Q_s^2}} Q_s^4\right)^2$$

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$$\begin{split} \bar{\sigma} &= \int_{z,\bar{z}} e^{ik \cdot (z-\bar{z})} \int_{x} f(\bar{z}-x) \cdot f(z-x) \\ &\times (N_{c}^{2}-1) \mu^{2} \Big[D(z,\bar{z})_{[1]} - D(z,x)_{[2]} - D(x,\bar{z})_{[3]} + 1_{[4]} \Big] \\ \sigma_{4}^{(ii)} &= g_{s}^{4} (N_{c}^{2}-1)^{2} \mu^{4} S_{\perp}^{2} \\ &\times e^{-\frac{k^{2}}{2Q_{s}^{2}}} \left\{ \frac{2}{k^{2}} - \frac{1}{k^{2}} e^{\frac{k^{2}}{2Q_{s}^{2}}} + \frac{1}{2Q_{s}^{2}} \left[\operatorname{Ei} \left(\frac{k^{2}}{2Q_{s}^{2}} \right) - \operatorname{Ei} \left(\frac{k^{2}\lambda}{2Q_{s}^{2}} \right) \right] \right\} \\ &\times \{ (k \to p) \} \end{split}$$

$$\begin{split} \sigma_{2}^{(i)} &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z})+ip \cdot (u-\bar{u})} \int_{x} f(\bar{z}-x) \cdot f(z-x) \ f(\bar{u}-x) \cdot f(u-x) \\ &\times \frac{1}{4} N_{c}^{3} \mu^{2} \Big\{ D(u,\bar{u}) D(z,\bar{z})_{[1]} + D(u,x) D(z,\bar{z})_{[2]} - D(u,\bar{u}) D(z,x)_{[3]} \\ &- D(u,x) D(z,x)_{[4]} + D(x,\bar{u}) D(z,\bar{z})_{[5]} + D(z,\bar{z})_{[6]} - D(x,\bar{u}) D(z,x)_{[7]} \\ &- D(z,x)_{[8]} - D(u,\bar{u}) D(x,\bar{z})_{[9]} - D(u,x) D(x,\bar{z})_{[10]} + D(u,\bar{u})_{[11]} \\ &+ D(u,x)_{[12]} - D(x,\bar{u}) D(x,\bar{z})_{[13]} - D(x,\bar{z})_{[14]} + D(x,\bar{u})_{[15]} + 1_{[16]} \Big\} \\ \sigma_{2}^{(i)} &= g_{s}^{4} N_{c}^{3} \mu^{2} S_{\perp} \frac{1}{k^{6} p^{6}} e^{-\frac{k^{2} + p^{2}}{2Q_{s}^{2}}} \\ &\times \left[k^{4} + e^{\frac{k^{2}}{2Q_{s}^{2}}} k^{2} Q_{s}^{2} + 4e^{\frac{k^{2}}{2Q_{s}^{2}}} Q_{s}^{4} \right] \left[\left(2e^{\frac{p^{2}}{2Q_{s}^{2}}} - 1 \right) p^{4} + e^{\frac{p^{2}}{2Q_{s}^{2}}} p^{2} Q_{s}^{2} + 4e^{\frac{p^{2}}{2Q_{s}^{2}}} Q_{s}^{4} \right] \end{split}$$

The two-gluon emission beyond the GGA Some details on the integration procedure

$$\begin{split} \sigma_{2}^{(ii)} &= \int_{u,z,\bar{u},\bar{z}} e^{ik \cdot (z-\bar{z}) + ip \cdot (u-\bar{u})} \times \int_{x} f(\bar{z}-\bar{u}) \cdot f(z-u) \ f(\bar{u}-x) \cdot f(u-x) \\ &\quad \times N_{c}^{3} \mu^{2} D(u,\bar{u}) \Big\{ D(z,\bar{z})_{[1]} - D(z,\bar{u})_{[2]} - D(u,\bar{z})_{[3]} + D(u,\bar{u})_{[4]} \Big\} \\ \sigma_{2}^{(ii)} &= g_{s}^{4} \frac{1}{(2\pi)^{2} Q_{s}^{4}} N_{c}^{3} \mu^{2} S_{\perp} e^{-\frac{k^{2}+p^{2}}{2Q_{s}^{2}}} \int_{s_{1},s_{2}} \frac{1}{s_{1}^{2}} \frac{1}{(s_{1}-s_{2})^{2}} e^{-\frac{s_{1}^{2}+2k\cdot s_{1}}{2Q_{s}^{2}}} e^{-\frac{s_{2}^{2}-2p\cdot s_{2}}{2Q_{s}^{2}}} \\ &\quad + g_{s}^{4} \frac{1}{(2\pi^{2})} Q_{s}^{4}} N_{c}^{3} \mu^{2} S_{\perp} e^{-\frac{k^{2}+p^{2}}{2Q_{s}^{2}}} \frac{k^{i}}{k^{2}} \int_{s_{1},s_{2}} \frac{s_{1}^{i}}{s_{1}^{2}} \frac{1}{(s_{1}-s_{2})^{2}} e^{-\frac{s_{1}^{2}+2k\cdot s_{1}}{2Q_{s}^{2}}} e^{-\frac{s_{2}^{2}-2p\cdot s_{2}}{2Q_{s}^{2}}} \\ &\quad + g_{s}^{4} \frac{1}{4Q_{s}^{2}} N_{c}^{3} \mu^{2} S_{\perp} \frac{1}{k^{2}} e^{-\frac{(k+p)^{2}}{4Q_{s}^{2}}} \left[\operatorname{Ei} \left(\frac{(k+p)^{2}}{4Q_{s}^{2}} \right) - \operatorname{Ei} \left(\frac{(k+p)^{2}\lambda}{4Q_{s}^{2}} \right) \right] \end{split}$$

$$\begin{split} \sigma_2^{(iii)} &= -\int_{u,z,\bar{u},\bar{z}} e^{ik\cdot(z-\bar{z})+ip\cdot(u-\bar{u})} \times \int_x f(\bar{z}-x) \cdot f(z-u) \ f(\bar{u}-x) \cdot f(u-x) \\ &\times \frac{1}{2} N_c^3 \mu^2 \Big\{ D(u,\bar{u}) D(z,\bar{z})_{[1]} - D(u,\bar{u}) D(z,x)_{[2]} - D(u,\bar{u}) D(u,\bar{z})_{[3]} \\ &+ D(u,\bar{u}) D(u,x)_{[4]} + D(u,x) D(z,\bar{z})_{[5]} - D(u,x) D(z,x)_{[6]} \\ &- D(u,x) D(u,\bar{z})_{[7]} + D^2(u,x)_{[8]} \Big\} \end{split}$$

$$\begin{split} \sigma_{2}^{(iv)} &= -\int_{u,z,\bar{u},\bar{z}} e^{ik\cdot(z-\bar{z})+ip\cdot(u-\bar{u})} \times \int_{x} f(\bar{z}-\bar{u}) \cdot f(z-x) \ f(\bar{u}-x) \cdot f(u-x) \\ &\times \frac{1}{2} N_{c}^{3} \mu^{2} \Big\{ D(u,\bar{u}) D(z,\bar{z})_{[1]} - D(u,\bar{u}) D(z,\bar{u})_{[2]} - D(u,\bar{u}) D(x,\bar{z})_{[3]} \\ &+ D(u,\bar{u}) D(x,\bar{u})_{[4]} + D(x,\bar{u}) D(z,\bar{z})_{[5]} - D(x,\bar{u}) D(z,\bar{u})_{[6]} \\ &- D(x,\bar{u}) D(x,\bar{z})_{[7]} + D^{2}(x,\bar{u})_{[8]} \Big\} \end{split}$$

$$\begin{split} \sigma_2^{[(iii)+(iv)]} &= g_s^4 N_c^3 \mu^2 S_\perp \\ &\times \left\{ \frac{1}{(2\pi)^2 Q_s^4} e^{-\frac{k^2 + p^2}{2Q_s^2}} \int_{s_1, s_2} \left[\left(\frac{s_1^i}{s_1^2} + \frac{k^i}{k^2} \right)^2 \left(\frac{s_2^j}{s_2^2} + \frac{p^j}{p^2} \right) - \frac{1}{k^2} \frac{p^j}{p^2} \right] \frac{(s_1 - s_2)^j}{(s_1 - s_2)^2} \\ &\times e^{-\frac{s_1^2 + 2k \cdot s_1}{2Q_s^2}} e^{-\frac{s_2^2 - 2p \cdot s_2}{2Q_s^2}} \\ &+ \frac{1}{k^2} \frac{p \cdot (k + p)}{p^2 (k + p)^2} \left(e^{-\frac{(k + p)^2}{4Q_s^2}} - 1 \right) \right\} \end{split}$$

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The two-gluon emission beyond the GGA Final expressions

$$\begin{split} & r_{4}^{([i)(+(iii)],1)}(k,p) = 2\pi \frac{\sigma_{0}^{2}}{N_{4}} e^{-\frac{(k+2)^{2}}{4c_{2}^{2}}} & N_{4} = (N_{c}^{2}-1)S_{\perp}Q_{s}^{2} \\ & \times \left[\frac{2Q_{s}^{4}\left(k^{4}+p^{4}+2\left(k\cdot p\right)^{2}\right)}{k^{2}p^{2}\left(k-p\right)^{4}} + \frac{8Q_{s}^{6}\left(k+p\right)^{4}}{k^{2}p^{2}\left(k-p\right)^{6}} \\ & + \frac{64Q_{s}^{8}\left(k^{4}+4\left(k\cdot p\right)^{2}+p^{4}+8\left(k\cdot p\right)\left(k^{2}+p^{2}\right)+14k^{2}p^{2}\right)}{k^{2}p^{2}\left(k-p\right)^{8}} \right] \\ & + (p \rightarrow -p) \\ & \tau_{4}^{([(i)+(iii)],1]}(k,p) = \frac{\sigma_{0}^{2}}{N_{4}}Q_{s}^{2}(2\pi)^{2}\left[\delta^{(2)}(k+p) + \delta^{(2)}(k-p)\right] \\ N_{4} = (N_{c}^{2}-1)S_{\perp}Q_{s}^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}}\left(k^{4}+k^{2}e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{2}+4e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{4}\right)^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}}\left(k^{4}+k^{2}e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{2}+4e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{4}\right)^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{(k+p)} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}}\left(k^{4}+k^{2}e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{2}+4e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{4}\right)^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}}\left(k^{4}+k^{2}e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{2}+4e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{4}\right)^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{(k+p)^{2}} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}}Q_{s}^{2} + \frac{k^{2}}{2Q_{s}^{2}}Q_{s}^{2}\right) - \operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right)^{2} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}} + \frac{1}{2}\left[\operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right) - \operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right)\right] \right\} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{Q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}} + \frac{1}{2}\left[\operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right) - \operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right)\right] \right\} \\ & \times \left\{ \frac{2Q_{s}^{2}}{k^{2}} - \frac{Q_{s}^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}} + \frac{1}{2}\left[\operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right) - \operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right)\right] \right\} \\ & \times \left\{ \frac{k^{2}}{k^{2}} + \frac{k^{2}}{k^{2}}e^{\frac{k^{2}}{2Q_{s}^{2}}} + \frac{1}{2}\left[\operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right) - \operatorname{Ei}\left(\frac{k^{2}}{2Q_{s}^{2}}\right)\right] \right\} \\ & \times \left\{ \frac{k^{2}}{k^{2}} + \frac{k^{2}}{k^{2}}e^{\frac{k^{2}}{2}} + \frac{k^{2}}{k^{2}}e^{\frac{k^{2}}{2}} + \frac{k^{2}}{k^{$$

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 \Rightarrow $\langle \sigma_2 \rangle_{P,T}$ CORRELATIONS are BIGGER than the $\langle \sigma_4 \rangle_{P,T}$ ones.









 $\Rightarrow \sigma_2^{(i)}$ does NOT CONTRIBUTE to the FH as it contains NO ANGULAR DEPENDENCE.



The two-gluon emission beyond the GGA The momentum dependent v_2 and v_3

$$v_n^2(k,p) = \frac{\int d\phi_k d\phi_p e^{in(\phi_k - \phi_p)} \frac{d^2 N^{(2)}}{d^2 k d^2 p}}{\int d\phi_k d\phi_p \frac{d^2 N^{(2)}}{d^2 k d^2 p}}$$



$$\begin{split} \sigma_4^{([(i)+(iii)].\mathbf{I})}(k,p) &= 2\pi \frac{\sigma_0^2}{N_4} \; e^{-\frac{(k+p)^2}{4Q_s^2}} \\ &\times \left[\frac{2Q_s^4 \left(k^4 + p^4 + 2 \left(k \cdot p\right)^2\right)}{k^2 p^2 \left(k - p\right)^4} + \frac{8Q_s^6 \left(k + p\right)^4}{k^2 p^2 \left(k - p\right)^6} \right. \\ &+ \frac{64Q_s^8 \left(k^4 + 4 \left(k \cdot p\right)^2 + p^4 + 8 \left(k \cdot p\right) \left(k^2 + p^2\right) + 14k^2 p^2\right)}{k^2 p^2 \left(k - p\right)^8} \right] \\ &+ \left(p \to -p\right) \end{split}$$

A numerical study RESTRICTED TO an INTERMEDIATE WINDOW of TRANSVERSE MOMENTUM VALUES.

Obtained from the TRADITIONAL PIECE.

MOMENTUM VALUES chosen for BOSE-ENHANCEMENT EFECT CONTRIBUTION, leaving the HBT effect ASIDE.

The two-gluon emission beyond the GGA The momentum dependent v_2 and v_3



$$\begin{split} \langle \sigma_2 \rangle_{P,T}^{sub} &= \frac{\sigma_0^2}{N_2} \Biggl\{ \frac{Q_s^2}{8} \left(\frac{1}{k^2} + \frac{1}{p^2} \right) e^{-\frac{(k+p)^2}{4Q_s^2}} \left[\operatorname{Ei} \left(\frac{(k+p)^2}{4Q_s^2} \right) - \operatorname{Ei} \left(\frac{(k+p)^2 \lambda}{4Q_s^2} \right) \right] \\ &+ \frac{Q_s^4}{2k^2 p^2} \left(e^{-\frac{(k+p)^2}{4Q_s^2}} - 1 \right) \Biggr\} \end{split}$$

 $\sigma_2^{(i)}$ does NOT SHOW any ANGULAR DEPENDENCE.

 $\sigma_2^{(ii),([(iii)+(iv)])}$ contain CORRELATIONS associated with the BACK-TO-BACK DI-JET PEAK.

 v_2^2 values are COMPARABLE with the ones arising from $\sigma_4^{([(i)+(iii)].I)}$.

 v_3^2 shows SIMILAR magnitude as v_2^2 but with OPPOSITE SIGN. A work that extends existing studies on TWO-PARTICLE CORRELATIONS in PROTON-NUCLEUS COLLISIONS within the CGC APPROACH.

This study incorporates **RELEVANT CONTRIBUTIONS OVERLOOKED** in the literature when implementing the DILUTE PROJECTILE LIMIT of the LEADING-ORDER DENSE-DENSE FORMULA.

Averaging over the color charge configurations of the projectile via using the MV MODEL and invoking the AREA ENHANCEMENT MODEL to express all the TARGET AVERAGES in terms of the dipole scattering amplitude ALLOW TO WORK WITH TRACTABLE EXPRESSIONS.

The GBW MODEL is employed to arrive at the FINAL EXPRESSIONS.

 $\Rightarrow \langle \sigma_2 \rangle_{P,T} \text{ CONTRIBUTES AS MUCH AS } \langle \sigma_4 \rangle_{P,T} \text{ TO THE SECOND FOURIER COEFFICIENT,} \\ \text{WHILE A NON-ZERO NEGATIVE } v_3^2 \text{ IS OBTAINED FROM } \langle \sigma_2 \rangle_{P,T} \text{ .}$

ORDERS OF MAGNITUDE OF OBTAINED NUMERICAL RESULTS MATCH THOSE RECENTLY MEASURED IN p+Au AND d+Au COLLISIONS BY THE PHENIX COLLABORATION.

THE STUDY OF THE JET-LIKE CONTRIBUTIONS THAT SHOW UP $\langle \sigma_2 \rangle_{P,T}$ DOES NOT FALL WITHIN THE SCOPE OF THIS WORK.

IT WOULD BE WORTHWHILE TO GO A BIT FURTHER IN EXAMINING THE MAGNITUDE OF THE JET CORRELATION PRESENT ON TOP OF THE RIDGE STRUCTURE.

THANK YOU VERY MUCH FOR YOUR ATTENTION!

November 2 | LIP — Lisbon Seminar