

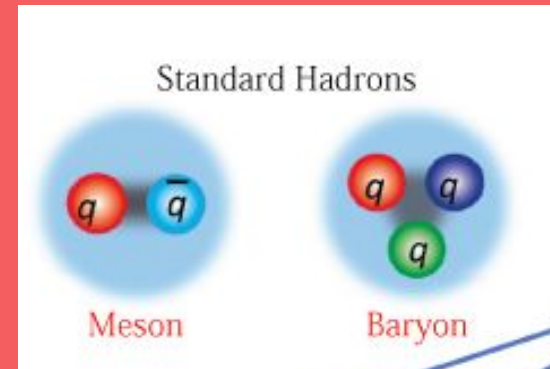
Hyperfine interactions in heavy quarkonia

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Standard Model

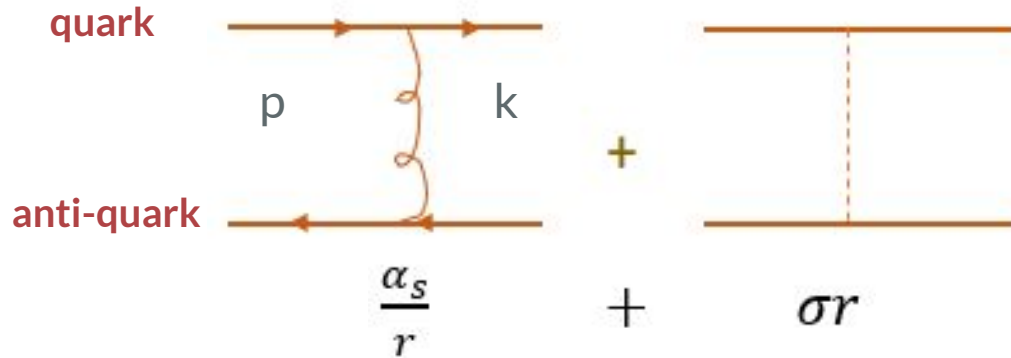
	Fermions			Bosons	
Quarks	u up	c charm	t top	γ photon	Force carriers
	d down	s strange	b bottom	Z Z boson	
Leptons	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
	e electron	μ muon	τ tau	g gluon	
				Higgs boson	

- hadrons: bound states of strongly interacting quarks
- mesons: quark-antiquark
- baryons: 3 quarks



Quark and Anti-quark Potential

- One Gluon Exchange + Linear Confinement: $V(r) = -\frac{\alpha_s}{r} + \sigma r$
- Good description of the bottomonium spectrum
- Main purpose : study effects of hyperfine interaction



First Step

- Write the base of the code to solve the Schrödinger equation in momentum space
- First test case : Calculate the binding energy of deuteron and the hydrogen atom
- Gauss Legendre Quadrature

$$E_b = -2.22452$$

$$\int_0^{\infty} dk f(k) \rightarrow \sum_{i=1}^N w_i f(k_i)$$

n° points = 50	n° points = 100	n° points = 200	n° points = 500	n° points = 1000
-2.225825	-2.224652	-2.224646	-2.224646	-2.224646

Bound state Schrödinger equation in momentum space is a homogeneous integral equation

S-waves ($l=0$):

$$\frac{p^2}{2\mu} \psi_0(p) - \frac{\alpha}{2\pi} \int_0^\infty dk \frac{k}{p} \ln \left(\frac{p+k}{p-k} \right)^2 \psi_0(k) - \frac{4\sigma}{\pi} \int_0^\infty dk k^2 \frac{\psi_0(k) - \psi_0(p)}{(p^2 - k^2)^2} = E \psi_0(p)$$

One-gluon-exchange Linear confinement

Singular at $k = p$

With a subtraction technique, we can make the kernel regular

$$\frac{p^2}{2\mu} \psi_0(p) - \frac{\alpha\pi p}{2} \psi_0(p) - \frac{\alpha}{2\pi} \int_0^\infty dk \ln \left(\frac{p+k}{p-k} \right)^2 \left[\frac{k}{p} \psi_0(k) - \frac{p}{k} \psi_0(p) \right] - \frac{2\sigma}{\pi} \int_0^\infty \frac{dk}{k^2 - p^2} \left[\frac{2k^2}{k^2 - p^2} (\psi_0(k) - \psi_0(p)) - p\psi_0'(p) \right] = E \psi_0(p)$$

Regular

Regular

Coulomb Potential

n	l = 0		l = 1		l = 2	
	Model Values	Exact Values	Model Values	Exact Values	Model Values	Exact Values
1	-13.598320	-13.598321	-3.3995801	-3.3995802	-1.510924	-1.510924
2	-3.3995802	-3.3995802	-1.5109246	-1.510924	-0.849895	-0.849895
3	-1.5109247	-1.5109245	-0.8498951	-0.849895	-0.543933	-0.849895
4	-0.8498953	-0.849895	-0.5439334	-0.543933	-0.377732	-0.377731
5	-0.5439333	-0.849895	-0.377731	-0.377731	-0.277517	-0.277516

→ Apply this to meson bound state problem

Linear Potential

Test case of the Schrodinger equation with linear potential : exact solution is known Airy Functions

n	l = 0		l = 1		l = 2	
	Model Values	Exact Values	Model Values	Exact Values	Model Values	Exact Values
1	2.334960	2.338107	3.359266	3.361258	4.443187	4.248187
2	4.080205	4.087949	4.877074	4.884455	5.652748	5.629114
3	5.506225	5.520560	6.192766	6.207626	6.964232	6.868888
4	6.763920	6.786708	7.381525	7.405667	8.017647	8.009707
5	7.911185	7.944134	8.480134	8.515235	9.126500	9.077007

Linear + Coulomb Potential

$$\begin{aligned}
 & \left[\frac{p^2}{2m_R} + \frac{\sigma\pi}{2p} P'_\ell(1) - \frac{\alpha\pi p}{2} \right] \psi_\ell(p) - \frac{2\sigma}{\pi} \int_0^\infty dk \left\{ \left[\frac{2k^2}{(k^2 - p^2)^2} (P_\ell(y)\psi_\ell(k) - \psi_\ell(p)) - \frac{p\psi'_\ell(p)}{k^2 - p^2} \right] \right. \\
 & \quad \left. - \frac{1}{4p^2} \ln \left(\frac{p+k}{p-k} \right)^2 \left[P'_\ell(y)\psi_\ell(k) - P'_\ell(1) \frac{p}{k} \psi_\ell(p) \right] + \frac{w'_{\ell-1}(y)}{2p^2} \psi_\ell(k) \right\} \\
 & \quad - \frac{\alpha}{\pi} \int_0^\infty dk \left\{ \frac{1}{2} \ln \left(\frac{p+k}{p-k} \right)^2 \left[\frac{k}{p} P_\ell(y)\psi_\ell(k) - \frac{p}{k} \psi_\ell(p) \right] - \frac{k}{p} w_{\ell-1}(y)\psi_\ell(k) \right\} = E\psi_\ell(p).
 \end{aligned}$$

→ Apply to bottomonium and charmonium

$$m_b = 4.7477 \quad \sigma = 0.2140$$

$$m_c = 1.3703 \quad \alpha_s = 0.4174$$

HyperFine Interactions in One Gluon Exchange

Spin-orbit

$$V(\vec{q}^2) = -\frac{4\pi\alpha_s}{q^2} \left\{ 1 - \frac{\vec{k} \cdot \vec{p}}{m^2} + \frac{3i}{4m^2} (\vec{p} \times \vec{k}) \cdot (\vec{\sigma}_1 + \vec{\sigma}_2) \right. \\ \left. - \frac{1}{6m^2} \frac{1}{q^2} (\vec{\sigma}_1 \cdot \vec{\sigma}_2) + \frac{1}{4m^2} \left[\frac{\vec{\sigma}_1 \cdot \vec{q}}{q} \frac{\vec{\sigma}_2 \cdot \vec{q}}{q} - \frac{1}{3} \frac{1}{q^2} (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \right] \right\}$$

spin-spin

Tensor Force

Pauli-Villars subtraction (form factors)

$$\begin{aligned} V(\vec{q}^2) &= -4\alpha_s \pi \left(\frac{1}{q^2} - \frac{\Lambda^2}{q^2 + \Lambda^2} \right) \\ &= -\frac{4\alpha_s \pi}{q^2} \left(\frac{\Lambda^2}{q^2 + \Lambda^2} \right) \end{aligned}$$

$$\Lambda = 2m_b$$

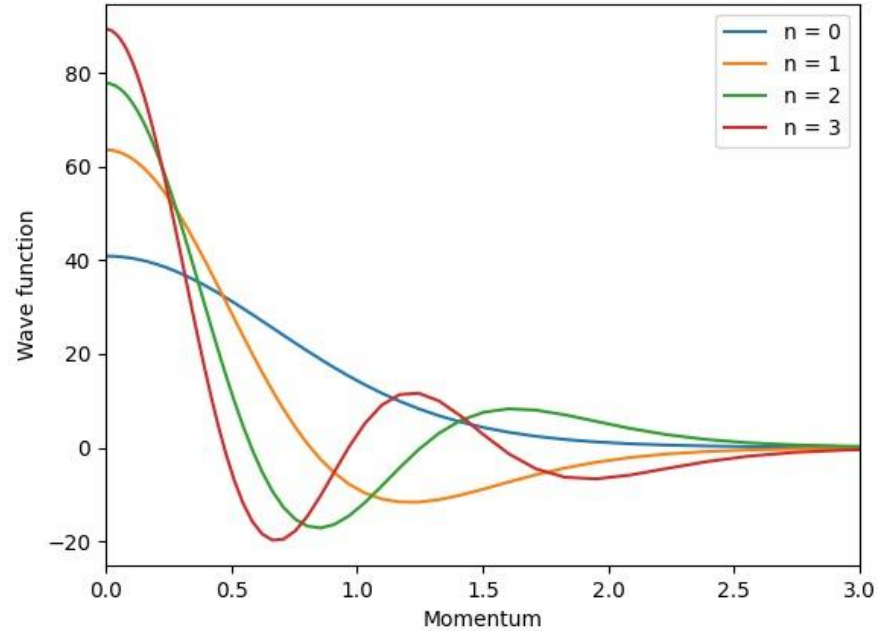
Bottomonium

n	Non-relativistic (without form factor)	Non-relativistic (with form factor)	Spin-spin dependence	Fully Relativistic
1	9.544758	9.543773	9.494522	9.15776
2	10.094968	10.083422	10.059072	9.85597
3	10.457809	10.443936	10.424913	10.20683
4	10.757798	10.742810	10.726324	10.47873
5	11.023094	11.007414	10.992496	10.71291

Charmonium

n	Non-relativistic (without form factor)	Non-relativistic (with form factor)	Spin-spin dependence	Fully Relativistic
1	3.236602	3.183976	3.133244	2.7707
2	3.886227	3.825244	3.791603	3.3697
3	4.379722	4.315851	4.287100	3.8003
4	4.804384	4.738877	4.714281	4.1651
5	5.187036	5.120418	5.098023	4.4921

Bottomonium Pseudoscalar Wave Functions



Summary and Conclusions

- Solved the Schrodinger equation using a linear + Coulomb potential for bottomonium and charmonium
- For pseudoscalar mesons: Included the hyperfine interaction to the potential and we observed substantial effects which makes the masses smaller
- We compared our results with hyperfine interaction with fully relativistic calculations: they lie between a non-relativistic and fully relativistic as expected

What's next?

- Look to vector, scalar, axial, vector, ect mesons
- Include higher partial waves therefore Spin-orbit and tensor force dependence.
- Heavy-light mesons
- Improve the code for linear potential : use better method to calculate the derivative

Questions?