

Diluting Quark Flavour Hierarchies Using Dihedral Symmetry

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Outline

- Motivation
 - The SM and its Freedom
 - Flavour Symmetries
- Quark Flavour Hierarchies
 - Quark Masses
 - The Wolfenstein Parametrization
- D_4 Symmetry
 - For the masses
 - For the Mixings
- Summary

The SM and its Freedom

$$\mathcal{L} = -\frac{1}{4}F^2 + i\bar{\Psi}\not{D}\Psi + \bar{\Psi}\phi\Psi + h.c. + |D\phi|^2 - V(\phi)$$

STANDARD MODEL



The SM and its Freedom

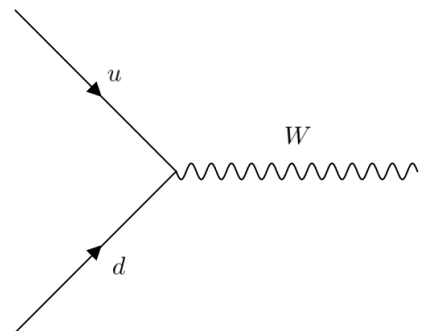
$$\mathcal{L} = -\frac{1}{4}F^2 + i\bar{\Psi}\not{D}\Psi + \bar{\Psi}\phi\Psi + h.c. + |D\phi|^2 - V(\phi)$$

STANDARD MODEL

+ NP?



The SM and its Freedom



Gauge Symmetries

Flavour Symmetries

$$Q_L \equiv (u_L, c_L, t_L) \sim 3$$

mass →	$\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 126 \text{ GeV}/c^2$
charge →	$2/3$	$2/3$	$2/3$	0	0
spin →	$1/2$	$1/2$	$1/2$	1	0
	u up	c charm	t top	g gluon	H Higgs boson
	$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-1/3$	$-1/3$	$-1/3$	0	
	$1/2$	$1/2$	$1/2$	1	
QUARKS	d down	s strange	b bottom	γ photon	
	$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$	$91.2 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$1/2$	$1/2$	$1/2$	1	
	e electron	μ muon	τ tau	Z Z boson	
LEPTONS	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	$80.4 \text{ GeV}/c^2$	
	0	0	0	± 1	
	$1/2$	$1/2$	$1/2$	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
					GAUGE BOSONS

The SM and its Freedom

Free Parameters	Gauge	Flavour
SM	3+1	Masses: 3+3+3 Mixings: 3+1 Scalar: 2
SM+ ν_R	3+1	Masses: 3+3+3+3 Mixings: (3+1)x2 Scalar: 2

The SM and its Freedom

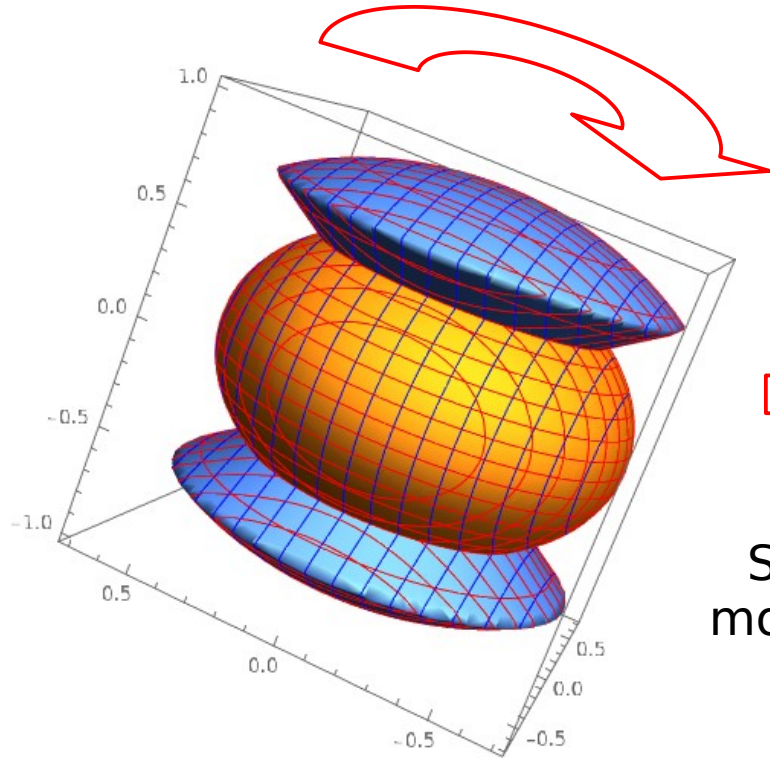
$$\mathcal{L} \supset Y_u \overline{Q}_L \tilde{\phi} u_R + Y_d \overline{Q}_L \phi d_R$$

$$M_{u,d} = v Y_{u,d} \equiv \begin{pmatrix} \times & \times & \times \\ \times & \times & \times \\ \times & \times & \times \end{pmatrix}$$

$$U^\dagger M M^\dagger U = \text{Diag}(m_u, m_c, m_t), \quad V^\dagger M M^\dagger V = \text{Diag}(m_d, m_s, m_b)$$

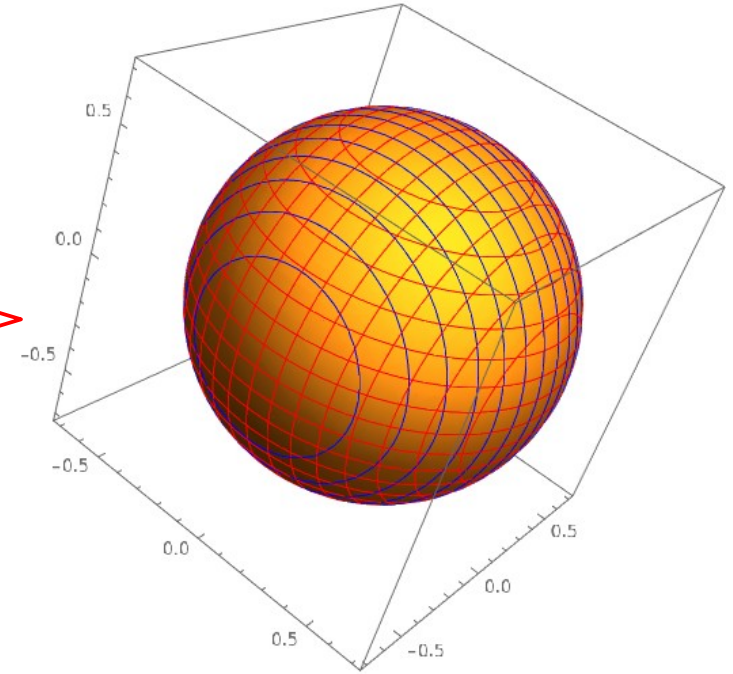
$$V_{\text{CKM}} = U^\dagger V \equiv \begin{pmatrix} \times & \times & \times \\ \times & \times & \times \\ \times & \times & \times \end{pmatrix}$$

Flavour Symmetries



“ $SO(2)$ ”

Enlarging the
Symmetry leads to
more restrictive cases



“ $SO(3)$ ”

Flavour Symmetries

$$\mathcal{L} \supset Y_u \overline{Q}_L \tilde{\phi} u_R + Y_d \overline{Q}_L \phi d_R$$

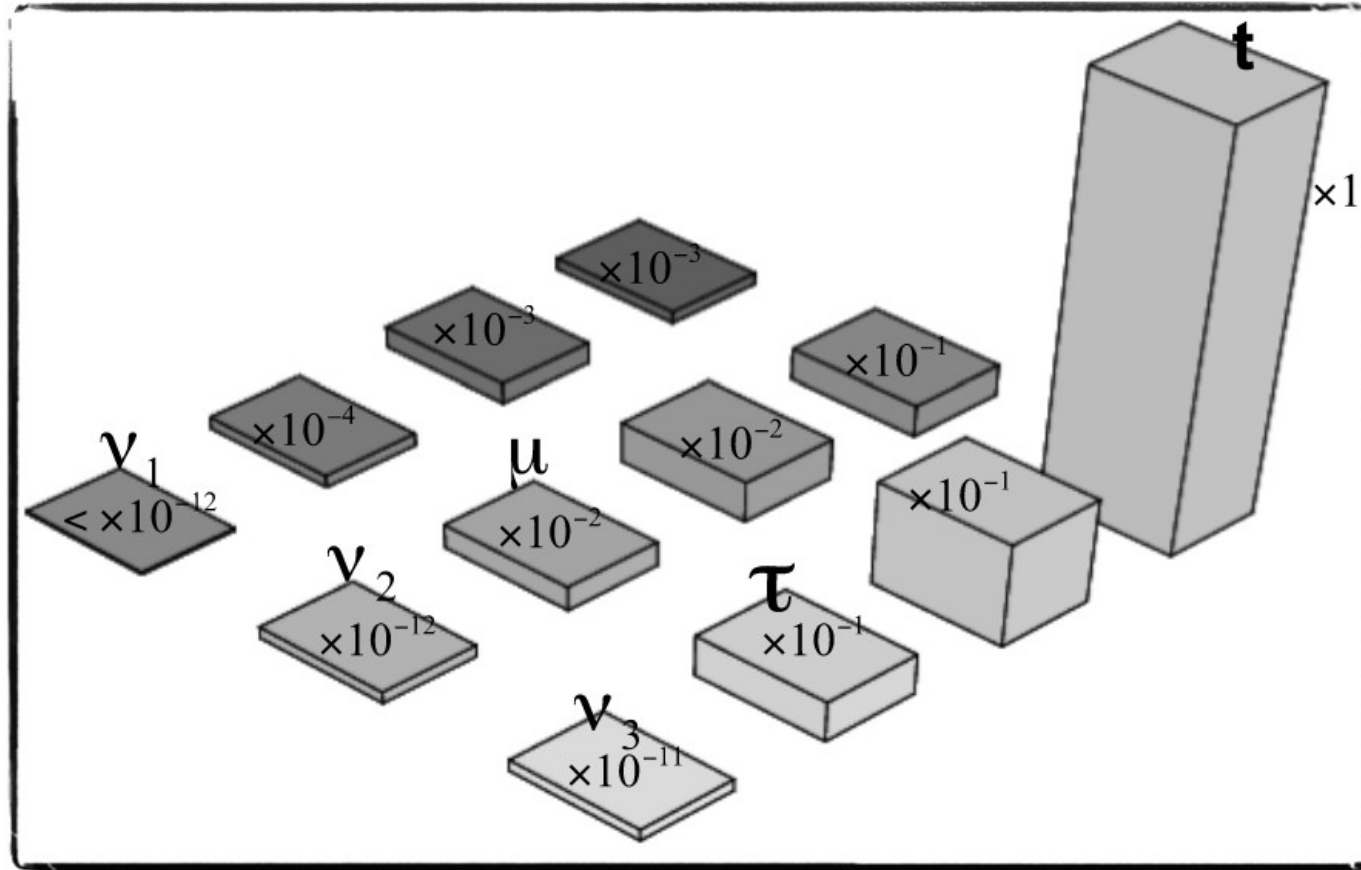
$$M_{u,d} = vY_{u,d} \equiv \begin{pmatrix} m_{u,d} & 0 & 0 \\ 0 & m_{u,d} & 0 \\ 0 & 0 & m_{u,d} \end{pmatrix}$$

$$U^\dagger M M^\dagger U = \text{Diag}(m_u, m_c, m_t), \quad V^\dagger M M^\dagger V = \text{Diag}(m_d, m_s, m_b)$$

$$V_{\text{CKM}} = U^\dagger V \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Quark Flavour Hierarchies

Masses

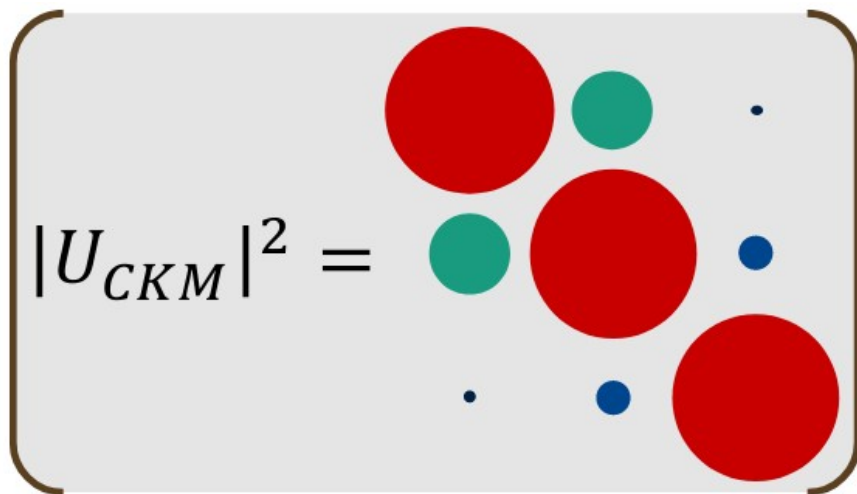


Blatantly Stolen
from Steve King

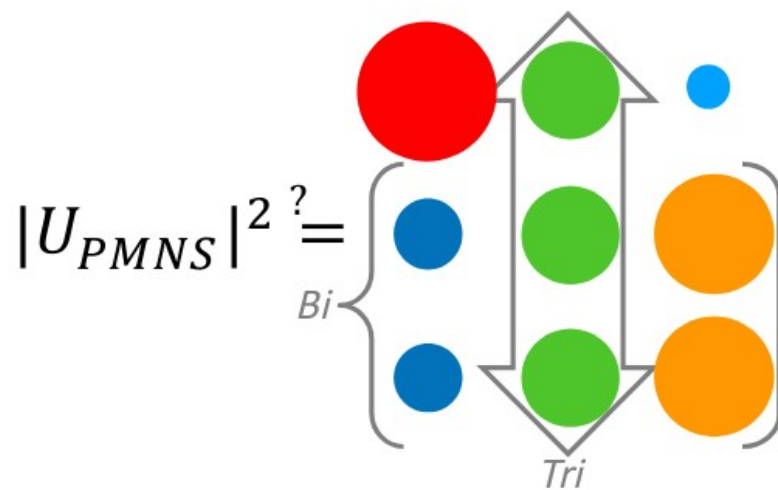
Quark Flavour Hierarchies

Mixing Hierarchies

CKM



PMNS



Blatantly Stolen
from Steve King

Parametrization of the Kobayashi-Maskawa Matrix

Lincoln Wolfenstein

Department of Physics, Carnegie-Mellon University, Pittsburgh, Pennsylvania 15213

(Received 22 August 1983)

The quark mixing matrix (Kobayashi-Maskawa matrix) is expanded in powers of a small parameter λ equal to $\sin\theta_c = 0.22$. The term of order λ^2 is determined from the recently measured B lifetime. Two remaining parameters, including the CP -nonconservation effects, enter only the term of order λ^3 and are poorly constrained. A significant reduction in the limit on ϵ'/ϵ possible in an ongoing experiment would tightly constrain the CP -nonconservation parameter and could rule out the hypothesis that the only source of CP nonconservation is the Kobayashi-Maskawa mechanism.

PACS numbers: 11.30.Er, 12.10.Ck, 13.25.+m

$$V \approx \begin{pmatrix} 1 - \lambda^2/2 & -\lambda & \mathcal{O}(\lambda^3) \\ \lambda & 1 - \lambda^2/2 & \mathcal{O}(\lambda^2) \\ \mathcal{O}(\lambda^3) & \mathcal{O}(\lambda^2) & 1 \end{pmatrix} \quad \lambda \approx 0.22$$

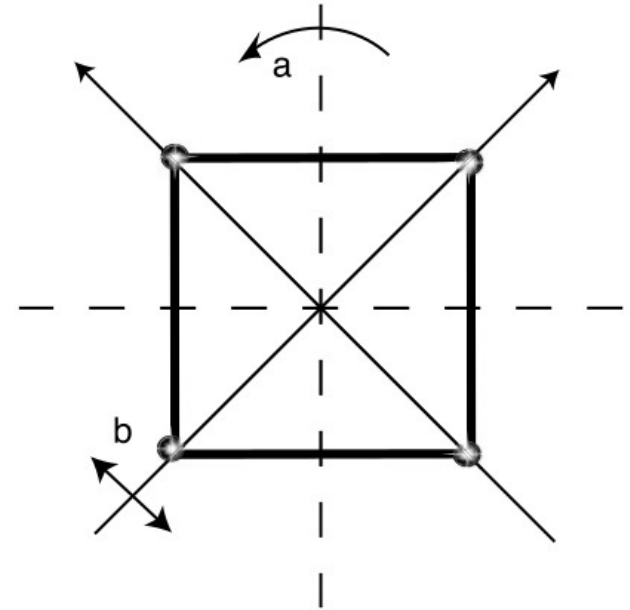
D₄ Symmetry

Group Theory

$$a = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}, \quad b = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix},$$

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix}_2 \otimes \begin{bmatrix} y_1 \\ y_2 \end{bmatrix}_2 = [x_1 y_1 + x_2 y_2]_{1_{++}} \oplus [x_1 y_2 - x_2 y_1]_{1_{--}} \\ \oplus [x_1 y_2 + x_2 y_1]_{1_{-+}} \\ \oplus [x_1 y_1 - x_2 y_2]_{1_{+-}},$$

$$\mathbf{1}_{r,s} \otimes \mathbf{1}_{r',s'} = \mathbf{1}_{r \cdot r', s \cdot s'}.$$



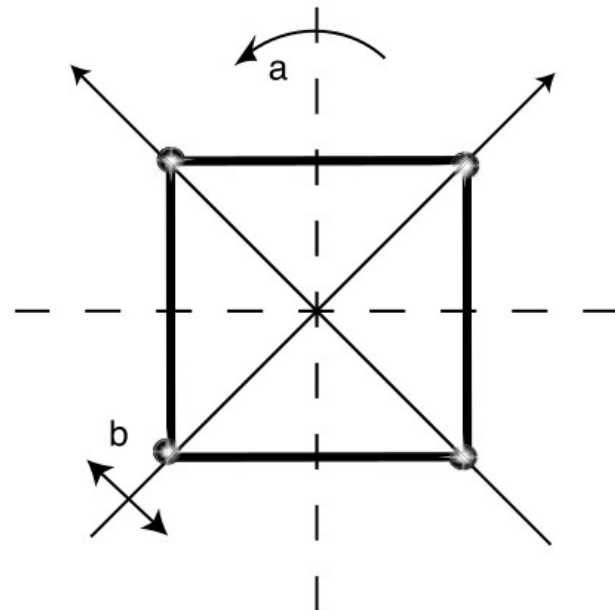
D₄ Symmetry

Group Theory

$$a = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}, \quad b = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix},$$

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix}_2 \otimes \begin{bmatrix} y_1 \\ y_2 \end{bmatrix}_2 = \underbrace{[x_1 y_1 + x_2 y_2]}_{\text{red circle}} \mathbf{1}_{++} \oplus [x_1 y_2 - x_2 y_1] \mathbf{1}_{--} \\ \oplus [x_1 y_2 + x_2 y_1] \mathbf{1}_{-+} \\ \oplus [x_1 y_1 - x_2 y_2] \mathbf{1}_{+-},$$

$$\mathbf{1}_{r,s} \otimes \mathbf{1}_{r',s'} = \mathbf{1}_{r \cdot r', s \cdot s'}.$$



D₄ Symmetry

Model

$$\mathbf{2} : \begin{bmatrix} Q_{1L} \\ Q_{2L} \end{bmatrix}, \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}, \quad (4a)$$

$$\begin{aligned} \mathbf{1}_{++} : n_{1R}, \quad \mathbf{1}_{--} : n_{2R}, n_{3R}, \phi_u, \quad \mathbf{1}_{-+} : p_{2R}, p_{3R}, \phi_d, \\ \mathbf{1}_{+-} : Q_{3L}, p_{1R}. \end{aligned} \quad (4b)$$

D₄ Symmetry

Model

$$2 : \begin{bmatrix} Q_{1L} \\ Q_{2L} \end{bmatrix}, \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}, \quad (4a)$$

$$\begin{aligned} 1_{++} : n_{1R}, \quad 1_{--} : n_{2R}, n_{3R}, \phi_u, \quad 1_{-+} : p_{2R}, p_{3R}, \phi_d, \\ 1_{+-} : Q_{3L}, p_{1R}. \end{aligned} \quad (4b)$$

D₄-Symmetric 4HDM

D₄ Symmetry

Model

$$\mathbf{2} : \begin{bmatrix} Q_{1L} \\ Q_{2L} \end{bmatrix}, \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}, \quad (4a)$$

$$\mathbf{1}_{++} : n_{1R}, \quad \mathbf{1}_{--} : n_{2R}, n_{3R}, \phi_u, \quad \mathbf{1}_{-+} : p_{2R}, p_{3R}, \phi_d, \\ \mathbf{1}_{+-} : Q_{3L}, p_{1R}. \quad (4b)$$

No Extra Energy Scales: All contribute to EWSB

$$\phi_1 = \begin{pmatrix} 0 \\ v_1 \end{pmatrix}, \quad \phi_2 = \begin{pmatrix} 0 \\ v_2 \end{pmatrix}, \quad \phi_u = \begin{pmatrix} 0 \\ v_u \end{pmatrix}, \quad \phi_d = \begin{pmatrix} 0 \\ v_d \end{pmatrix},$$

$$v_1^2 + v_2^2 + v_u^2 + v_d^2 = v_{\text{SM}}^2, \quad v_{12}^2 = v_1^2 + v_2^2, \quad \frac{v_2}{v_1} = \tan \beta$$

Model

$$\mathbf{2} : \begin{bmatrix} Q_{1L} \\ Q_{2L} \end{bmatrix}, \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}, \quad (4a)$$

$$\mathbf{1}_{++} : n_{1R}, \quad \mathbf{1}_{--} : n_{2R}, n_{3R}, \phi_u, \quad \mathbf{1}_{-+} : p_{2R}, p_{3R}, \phi_d, \\ \mathbf{1}_{+-} : Q_{3L}, p_{1R}. \quad (4b)$$

$$\begin{aligned} -\mathcal{L}_u &= A_u(\bar{Q}_{1L}\tilde{\phi}_1 - \bar{Q}_{2L}\tilde{\phi}_2)p_{1R} + B_u(\bar{Q}_{1L}\tilde{\phi}_2 \\ &\quad + \bar{Q}_{2L}\tilde{\phi}_1)p_{2R} + C_u(\bar{Q}_{1L}\tilde{\phi}_2 + \bar{Q}_{2L}\tilde{\phi}_1)p_{3R} \\ &\quad + X_u\bar{Q}_{3L}\phi_u p_{2R} \\ &\quad + Y_u\bar{Q}_{3L}\phi_u p_{3R}, \\ -\mathcal{L}_d &= A_d(\bar{Q}_{1L}\phi_1 + \bar{Q}_{2L}\phi_2)n_{1R} \\ &\quad + B_d(\bar{Q}_{1L}\phi_2 - \bar{Q}_{2L}\phi_1)n_{2R} \\ &\quad + C_d(\bar{Q}_{1L}\phi_2 - \bar{Q}_{2L}\phi_1)n_{3R} \\ &\quad + X_d\bar{Q}_{3L}\phi_d n_{2R} + Y_d\bar{Q}_{3L}\phi_d n_{3R}, \end{aligned}$$

$$M_u = \begin{pmatrix} A_u v_1 & B_u v_2 & C_u v_2 \\ -A_u v_2 & B_u v_1 & C_u v_1 \\ 0 & X_u v_u & Y_u v_u \end{pmatrix},$$

$$M_d = \begin{pmatrix} A_d v_1 & B_d v_2 & C_d v_2 \\ A_d v_2 & -B_d v_1 & -C_d v_1 \\ 0 & X_d v_d & Y_d v_d \end{pmatrix},$$

D₄ Symmetry

Model

$$\mathbf{2} : \begin{bmatrix} Q_{1L} \\ Q_{2L} \end{bmatrix}, \begin{bmatrix} \phi_1 \\ \phi_2 \end{bmatrix}, \quad (4a)$$

$$\mathbf{1}_{++} : n_{1R}, \quad \mathbf{1}_{--} : n_{2R}, n_{3R}, \phi_u, \quad \mathbf{1}_{-+} : p_{2R}, p_{3R}, \phi_d, \\ \mathbf{1}_{+-} : Q_{3L}, p_{1R}. \quad (4b)$$

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β

Block-Diagonalization

$$M_u = \begin{pmatrix} A_u v_1 & B_u v_2 & C_u v_2 \\ -A_u v_2 & B_u v_1 & C_u v_1 \\ 0 & X_u v_u & Y_u v_u \end{pmatrix},$$

$$M_d = \begin{pmatrix} A_d v_1 & B_d v_2 & C_d v_2 \\ A_d v_2 & -B_d v_1 & -C_d v_1 \\ 0 & X_d v_d & Y_d v_d \end{pmatrix},$$

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$$O_\beta = \begin{pmatrix} \cos \beta & -\sin \beta & 0 \\ \sin \beta & \cos \beta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

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$$O_\beta = \begin{pmatrix} \cos \beta & -\sin \beta & 0 \\ \sin \beta & \cos \beta & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$(M_u^2)_{\text{Block}} \equiv O_\beta M_u M_u^\dagger O_\beta^\dagger = \begin{pmatrix} A_u^2 v_{12}^2 & 0 & 0 \\ 0 & (B_u^2 + C_u^2) v_{12}^2 & (C_u Y_u + B_u X_u) v_{12} v_u \\ 0 & (C_u Y_u + B_u X_u) v_{12} v_u & (Y_u^2 + X_u^2) v_u^2 \end{pmatrix},$$

$$(M_d^2)_{\text{Block}} \equiv O_\beta^\dagger M_d M_d^\dagger O_\beta = \begin{pmatrix} A_d^2 v_{12}^2 & 0 & 0 \\ 0 & (B_d^2 + C_d^2) v_{12}^2 & -(C_d Y_d + B_d X_d) v_{12} v_d \\ 0 & -(C_d Y_d + B_d X_d) v_{12} v_d & (Y_d^2 + X_d^2) v_d^2 \end{pmatrix}$$

Full Diagonalization

$$(M_u^2)_{\text{Block}} \equiv O_\beta M_u M_u^\dagger O_\beta^\dagger = \begin{pmatrix} A_u^2 v_{12}^2 & 0 & 0 \\ 0 & (B_u^2 + C_u^2) v_{12}^2 & (C_u Y_u + B_u X_u) v_{12} v_u \\ 0 & (C_u Y_u + B_u X_u) v_{12} v_u & (Y_u^2 + X_u^2) v_u^2 \end{pmatrix},$$
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Full Diagonalization

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$$O_\theta^{u,d} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{u,d} & -\sin \theta_{u,d} \\ 0 & \sin \theta_{u,d} & \cos \theta_{u,d} \end{pmatrix},$$

Full Diagonalization

$$(M_u^2)_{\text{Block}} \equiv O_\beta M_u M_u^\dagger O_\beta^\dagger = \begin{pmatrix} A_u^2 v_{12}^2 & 0 & 0 \\ 0 & (B_u^2 + C_u^2) v_{12}^2 & (C_u Y_u + B_u X_u) v_{12} v_u \\ 0 & (C_u Y_u + B_u X_u) v_{12} v_u & (Y_u^2 + X_u^2) v_u^2 \end{pmatrix},$$

$$(M_d^2)_{\text{Block}} \equiv O_\beta^\dagger M_d M_d^\dagger O_\beta = \begin{pmatrix} A_d^2 v_{12}^2 & 0 & 0 \\ 0 & (B_d^2 + C_d^2) v_{12}^2 & -(C_d Y_d + B_d X_d) v_{12} v_d \\ 0 & -(C_d Y_d + B_d X_d) v_{12} v_d & (Y_d^2 + X_d^2) v_d^2 \end{pmatrix}$$

$$O_\theta^{u,d} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{u,d} & -\sin \theta_{u,d} \\ 0 & \sin \theta_{u,d} & \cos \theta_{u,d} \end{pmatrix},$$

$$\tan 2\theta_u = \frac{2(C_u Y_u + B_u X_u) v_{12} v_u}{(Y_u^2 + X_u^2) v_u^2 - (B_u^2 + C_u^2) v_{12}^2},$$

$$\tan 2\theta_d = -\frac{2(C_d Y_d + B_d X_d) v_{12} v_d}{(Y_d^2 + X_d^2) v_d^2 - (B_d^2 + C_d^2) v_{12}^2}.$$

$$D_u^2 = O_\theta^u O_\beta (M_u M_u^\dagger) O_\beta^\dagger O_\theta^{u\dagger} \equiv \text{diag}(m_u^2, m_c^2, m_t^2),$$

$$D_d^2 = O_\theta^d O_\beta^\dagger (M_d M_d^\dagger) O_\beta O_\theta^{d\dagger} \equiv \text{diag}(m_d^2, m_s^2, m_b^2).$$

D₄ Symmetry

Goal:

$$V_{\text{CKM}} = \begin{pmatrix} \cos 2\beta & -\cos \theta_d \sin 2\beta & -\sin 2\beta \sin \theta_d \\ \cos \theta_u \sin 2\beta & \cos 2\beta \cos \theta_d \cos \theta_u + \sin \theta_d \sin \theta_u & \cos 2\beta \cos \theta_u \sin \theta_d - \cos \theta_d \sin \theta_u \\ \sin 2\beta \sin \theta_u & -\cos \theta_u \sin \theta_d + \cos 2\beta \cos \theta_d \sin \theta_u & \cos \theta_d \cos \theta_u + \cos 2\beta \sin \theta_d \sin \theta_u \end{pmatrix}$$

+

$$\sin 2\beta = \lambda, \quad \sin \theta_{u,d} \approx \mathcal{O}(\lambda^2), \quad \cos \theta_{u,d} \approx \mathcal{O}(1).$$

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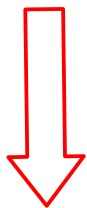
$$V \approx \begin{pmatrix} 1 - \lambda^2/2 & -\lambda & \mathcal{O}(\lambda^3) \\ \lambda & 1 - \lambda^2/2 & \mathcal{O}(\lambda^2) \\ \mathcal{O}(\lambda^3) & \mathcal{O}(\lambda^2) & 1 \end{pmatrix}$$

Can we do it?

$$m_u^2 = A_u^2 v_{12}^2,$$

$$m_c^2 \approx \frac{(B_u Y_u - C_u X_u)^2}{(Y_u^2 + X_u^2)} v_{12}^2.$$

$$m_t^2 \approx (Y_u^2 + X_u^2) v_u^2.$$



$$\frac{m_c}{m_t} \approx \frac{v_{12}}{v_u} \sim \mathcal{O}(\lambda^2)$$

$$\theta_u \approx \frac{(C_u Y_u + B_u X_u)}{(Y_u^2 + X_u^2)} \frac{v_{12}}{v_u} \approx \mathcal{O}\left(\frac{v_{12}}{v_u}\right),$$

$$\theta_d \approx -\frac{(C_d Y_d + B_d X_d)}{(Y_d^2 + X_d^2)} \frac{v_{12}}{v_d} \approx \mathcal{O}\left(\frac{v_{12}}{v_d}\right),$$

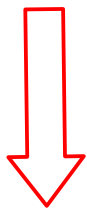
D₄ Symmetry

Can we do it?

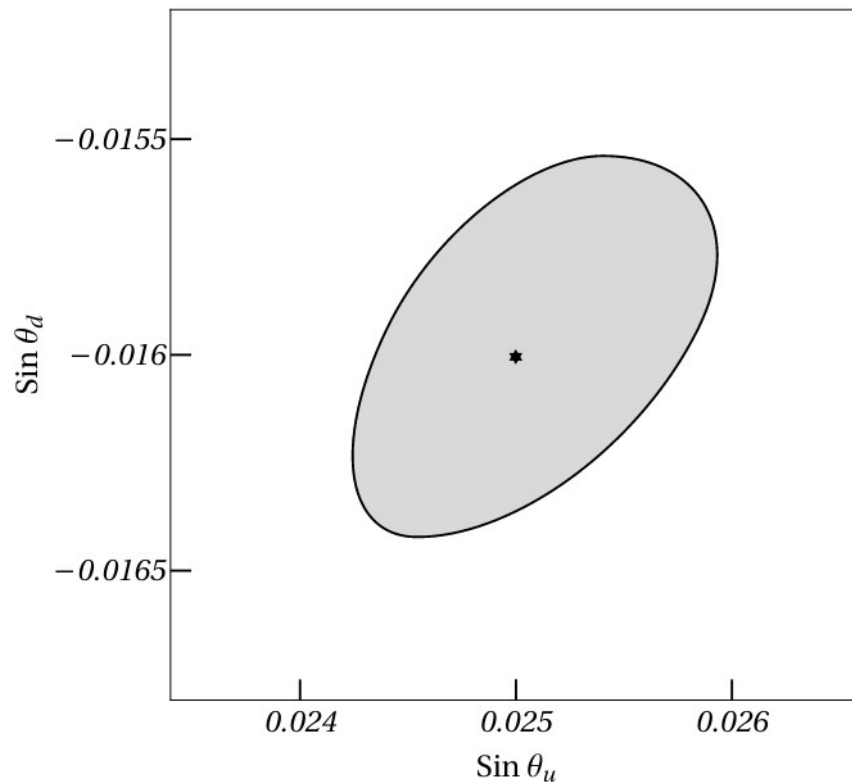
$$m_u^2 = A_u^2 v_{12}^2,$$

$$m_c^2 \approx \frac{(B_u Y_u - C_u X_u)^2}{(Y_u^2 + X_u^2)} v_{12}^2.$$

$$m_t^2 \approx (Y_u^2 + X_u^2) v_u^2.$$



$$\frac{m_c}{m_t} \approx \frac{v_{12}}{v_u} \sim \mathcal{O}(\lambda^2)$$



D₄ Symmetry

Can we do it?

yes

D₄ Symmetry

Can we do it?

Well, yes
But Why?

Why we did it:

Aesthetics:

- Dillution of the Quark Yukawa hierarchies
- Relating Quark mass hierarchies to CKM hierarchies
- Connected Quark Masses and Mixings to Scalar Sector

Observable Consequences:

- FCNCs
- Extra Higgs Phenomenology
- Small Parametric Space

Thank You For
Your Attention!

Back-Up Slides

FCNCs

$$N_d^1 \approx \frac{1}{\sqrt{2}v_{12}} \begin{pmatrix} m_d \cos \beta & -m_s \sin \beta & m_b \theta_d \sin \beta \\ -m_d \sin \beta & -m_s \cos \beta & m_b \theta_d \cos \beta \\ -m_d \theta_d \sin \beta & -m_s \theta_d \cos \beta & m_b \theta_d^2 \cos \beta \end{pmatrix}$$

$$N_d^2 \approx \frac{1}{\sqrt{2}v_{12}} \begin{pmatrix} m_d \sin \beta & m_s \cos \beta & -m_b \theta_d \cos \beta \\ m_d \cos \beta & -m_s \sin \beta & m_b \theta_d \sin \beta \\ m_b \theta_d \cos \beta & -m_s \theta_d \sin \beta & m_b \theta_d^2 \sin \beta \end{pmatrix}$$

$$N_d^d \approx \frac{1}{\sqrt{2}v_d} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -m_b \theta_d \\ 0 & 0 & m_b \end{pmatrix}.$$

Scalar Potential

$$V(\phi) = V_{\text{quadratic}} + V_{\text{quartic}},$$

$$\begin{aligned} V_{\text{quartic}} = & \lambda_1 \left(\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 \right)^2 + \lambda_2 \left(\phi_1^\dagger \phi_2 - \phi_2^\dagger \phi_1 \right)^2 + \lambda_3 \left(\phi_1^\dagger \phi_2 + \phi_2^\dagger \phi_1 \right)^2 + \lambda_4 \left(\phi_1^\dagger \phi_1 - \phi_2^\dagger \phi_2 \right)^2 \\ & + \lambda_5 \left(\phi_u^\dagger \phi_u \right)^2 + \lambda_6 \left(\phi_d^\dagger \phi_d \right)^2 + \lambda_7 \left(\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 \right) \left(\phi_d^\dagger \phi_d \right) + \lambda_8 \left(\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 \right) \left(\phi_u^\dagger \phi_u \right) \\ & + \lambda_9 \left(\phi_u^\dagger \phi_u \right) \left(\phi_d^\dagger \phi_d \right) + \lambda_{10} \left[\left(\phi_u^\dagger \phi_1 \right) \left(\phi_1^\dagger \phi_u \right) + \left(\phi_u^\dagger \phi_2 \right) \left(\phi_2^\dagger \phi_u \right) \right] \\ & + \lambda_{11} \left[\left(\phi_d^\dagger \phi_1 \right) \left(\phi_1^\dagger \phi_d \right) + \left(\phi_d^\dagger \phi_2 \right) \left(\phi_2^\dagger \phi_d \right) \right] \\ & + \lambda_{12} \left[\left(\phi_u^\dagger \phi_1 \right)^2 + \left(\phi_u^\dagger \phi_2 \right)^2 + \left(\phi_1^\dagger \phi_u \right)^2 + \left(\phi_2^\dagger \phi_u \right)^2 \right] \\ & + \lambda_{13} \left[\left(\phi_d^\dagger \phi_1 \right)^2 + \left(\phi_d^\dagger \phi_2 \right)^2 + \left(\phi_1^\dagger \phi_d \right)^2 + \left(\phi_2^\dagger \phi_d \right)^2 \right] \\ & + \lambda_{14} \left[\left(\phi_u^\dagger \phi_1 \right) \left(\phi_d^\dagger \phi_1 \right) + \left(\phi_u^\dagger \phi_2 \right) \left(\phi_d^\dagger \phi_2 \right) + \left(\phi_1^\dagger \phi_u \right) \left(\phi_1^\dagger \phi_d \right) + \left(\phi_2^\dagger \phi_u \right) \left(\phi_2^\dagger \phi_d \right) \right] \\ & + \lambda_{15} \left[\left(\phi_u^\dagger \phi_1 \right) \left(\phi_1^\dagger \phi_d \right) + \left(\phi_u^\dagger \phi_2 \right) \left(\phi_2^\dagger \phi_d \right) + \left(\phi_1^\dagger \phi_u \right) \left(\phi_d^\dagger \phi_1 \right) + \left(\phi_2^\dagger \phi_u \right) \left(\phi_d^\dagger \phi_2 \right) \right] \\ & + \lambda_{16} \left(\phi_1^\dagger \phi_1 - \phi_2^\dagger \phi_2 \right) \left[\left(\phi_u^\dagger \phi_d \right) + \left(\phi_d^\dagger \phi_u \right) \right], \end{aligned} \quad (35)$$

$$\begin{aligned} V_{\text{quadratic}} = & \mu_{11}^2 \left(\phi_1^\dagger \phi_1 \right) + \mu_{22}^2 \left(\phi_2^\dagger \phi_2 \right) + \mu_{uu}^2 \left(\phi_u^\dagger \phi_u \right) + \mu_{dd}^2 \left(\phi_d^\dagger \phi_d \right) \\ & + \mu_{12}^2 \left(\phi_1^\dagger \phi_2 + \text{h.c.} \right) + \mu_{1u}^2 \left(\phi_1^\dagger \phi_u + \text{h.c.} \right) + \mu_{1d}^2 \left(\phi_1^\dagger \phi_d + \text{h.c.} \right) \\ & + \mu_{2u}^2 \left(\phi_2^\dagger \phi_u + \text{h.c.} \right) + \mu_{2d}^2 \left(\phi_2^\dagger \phi_d + \text{h.c.} \right) + \mu_{ud}^2 \left(\phi_u^\dagger \phi_d + \text{h.c.} \right). \end{aligned} \quad (36)$$

Escaping Bounds

$$\begin{aligned} V_{\text{MS}} = & \mu_{11}^2 \phi_1^\dagger \phi_1 + \mu_{22}^2 \phi_2^\dagger \phi_2 + \mu_{uu}^2 \phi_u^\dagger \phi_u + \mu_{dd}^2 \phi_d^\dagger \phi_d \\ & + \left(\mu_{12}^2 \phi_1^\dagger \phi_2 + \mu_{1u}^2 \phi_1^\dagger \phi_u + \mu_{1d}^2 \phi_1^\dagger \phi_d + \mu_{2u}^2 \phi_2^\dagger \phi_u + \mu_{2d}^2 \phi_2^\dagger \phi_d + \mu_{ud}^2 \phi_u^\dagger \phi_d + \text{h.c.} \right) \\ & + \lambda \left(\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_u^\dagger \phi_u + \phi_d^\dagger \phi_d \right)^2, \end{aligned}$$

- The alignment limit emerges automatically (there is a mass eigenstate which is SM-like).
- The nonstandard masses are disentangled from the EW scale (the nonstandard masses are tier-wise degenerate and decoupled).
- There is no restriction on the relative hierarchies of the VEVs.
- Unitarity and boundedness from below constraints are trivially satisfied, because there is only one quartic parameter λ , which is related to the SM-like Higgs mass as $m_h^2 = 4\lambda v^2$.
- The ρ -parameter is trivially satisfied by the tier-wise degeneracy of the masses [31, 32].