

Flavour Physics

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Outline

- 1 Flavour in the Standard Model
 - The Flavour Problem
- 2 Flavour Beyond the Standard Model
 - The Flavour Solution (?)
- 3 Conclusion

The Standard Model

Gauge group: $SU(3)_C \times SU(2)_L \times U(1)_Y$

Chiral spin 1/2 fermions (left and right)

Quarks: colour triplets of $SU(3)_C$

Left fermions are doublets of $SU(2)_L$

Spin 0 scalar, doublet of $SU(2)_L$

The Standard Model (1 generation)

Gauge group: $SU(3)_C \times SU(2)_L \times U(1)_Y$

Quarks (Q , u_R , d_R): colour triplets of $SU(3)_C$

LH fields (Q and L): doublets of $SU(2)_L$

e_R just $U(1)_Y$

(ν SM: add ν_R , complete singlet)

Scalar H also doublet of $SU(2)_L$

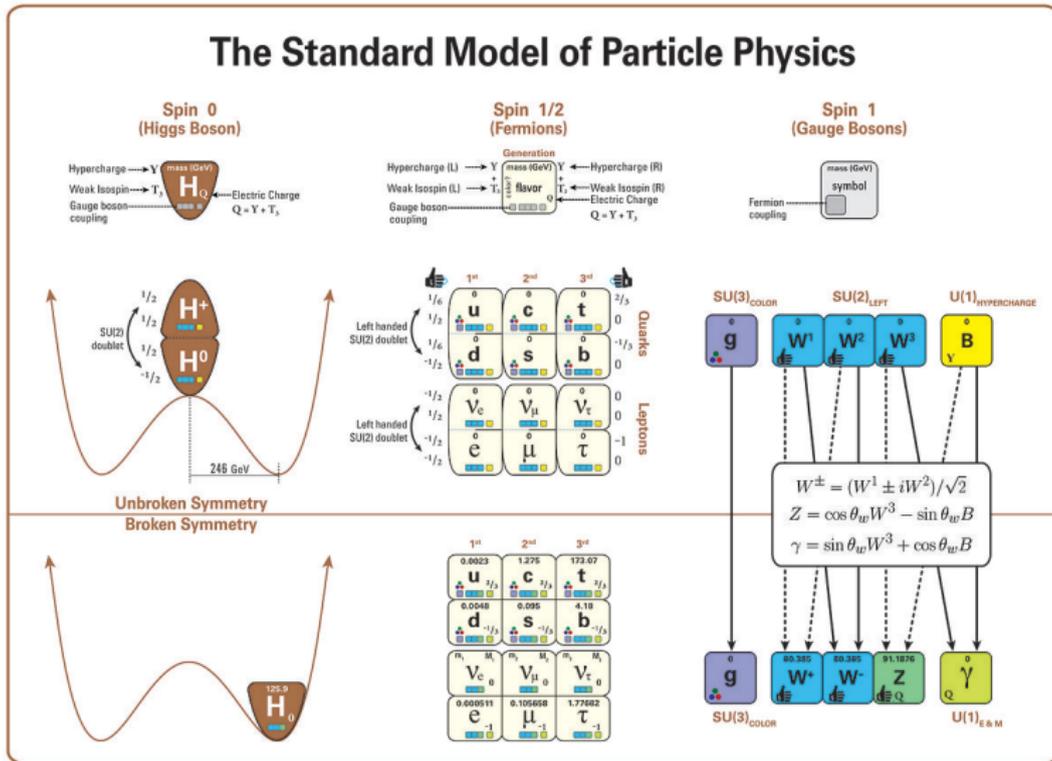
$\langle H \rangle$ breaks $SU(2)_L \times U(1)_Y \rightarrow U(1)_{em}$

Mass terms: $m_f F_\alpha f_R$ not invariant under $SU(2)_L$

But $y_f (\epsilon^{\alpha\beta} H_\alpha F_\beta) f_R$ is...

$y_f \langle H \rangle F f_R \rightarrow m_f F f_R$ with $m_f = y_f \langle H \rangle$

The Standard Model summary



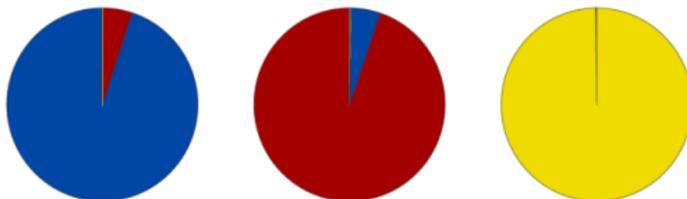
The Standard Model is very successful but...

- Neutrinos have masses (ν SM)
- Dark matter (no viable explanation)
- Matter / antimatter asymmetry (no viable explanation)
- Hierarchy problem (fine-tuning between parameters)
- Strong CP problem (fine-tuning between parameters)
- Gauge couplings (additional free parameters) - GUT?
- **Flavour problem (many additional free parameters) - FS?**

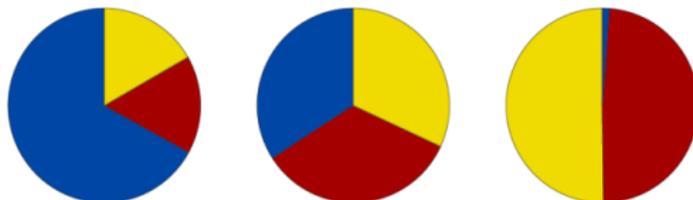
BSM solutions involve additional fields and symmetries

The Standard Model flavour problem: mixing

3 generations of quarks, small mixing



3 generations of leptons, large and peculiar mixing



(mixing between weak and mass eigenstates)

Summary of data: quark mixing

Wolfenstein parametrisation

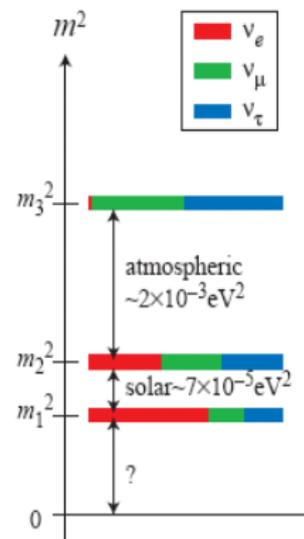
$$V_{CKM} \simeq \begin{pmatrix} 1 & \lambda & \lambda^3 \\ -\lambda & 1 & \lambda^2 \\ \lambda^3 & -\lambda^2 & 1 \end{pmatrix}$$

$\lambda \simeq 0.23$ (Sine of the Cabibbo angle)

Summary of data: lepton mixing

Tri-bi-maximal (TBM) mixing

$$V_{PMNS} \simeq \begin{pmatrix} -\sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & -\sqrt{\frac{1}{2}} \end{pmatrix}$$



Beyond the Standard Model with Family Symmetry

Without $y_f H F f_R$, $\mathcal{L}_{\nu SM}$ has accidental symmetry $SU(3)^6$

FS: upgrade subgroup of $SU(3)^6$ to actual symmetry of \mathcal{L}

- 1 Generations charged differently under FS
- 2 Yukawa couplings no longer invariant
- 3 FS must be broken somehow...

Abelian example: $U(1)$ FS + single familon

$U(1)$ assignments

Field	$U(1)$
H	0
ϕ	-1
Q_3	0
d_{R3}	0
Q_2	1
d_{R2}	1
Q_1	2
d_{R1}	2

$$\mathcal{L}_d = H(y_{33}Q_3d_{R3} + y_{23}(\phi/M_X)Q_2d_{R3} + y_{32}(\phi/M_X)Q_3d_{R2} + y_{22}(\phi/M_X)^2Q_2d_{R2} + \dots)$$

Respective mass matrix

$$M_d \sim m_b \begin{pmatrix} \epsilon^4 & \epsilon^3 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & \epsilon \\ \epsilon^2 & \epsilon & 1 \end{pmatrix}$$

$$\frac{\langle \phi \rangle}{M_X} = \epsilon$$

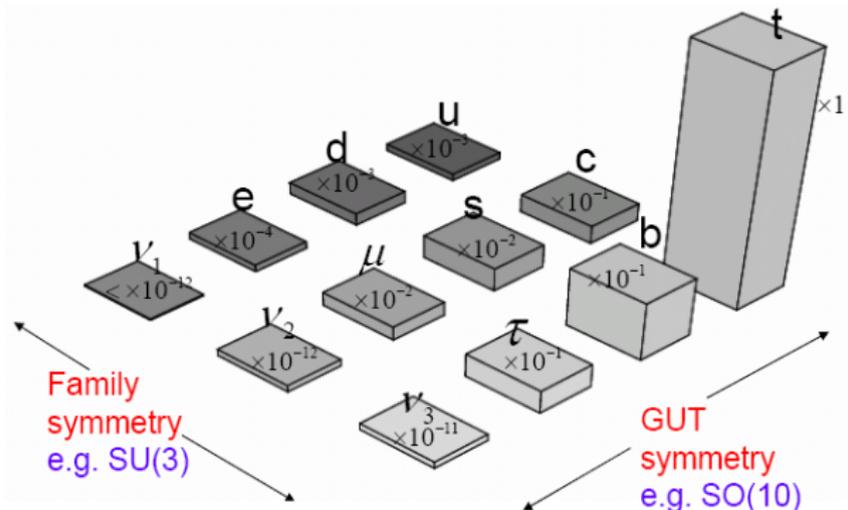
Each entry has a y_{ij} parameter!

Non-Abelian?

3 reasons

- 3 generations explained naturally
- ν SM: $FS \subset SU(3)^6$; $SO(10)$ GUT: $FS \subset SU(3)$
- Lepton mixing strongly suggests non-Abelian FS

$SO(10) \times SU(3)?$



Discrete?

2 generation example

- $V = -m^2(\varphi^i \varphi_i^\dagger) + \lambda(\varphi^i \varphi_i^\dagger)(\varphi^j \varphi_j^\dagger)$ continuous vacua
 $\pm d(\varphi^i \varphi_i^\dagger \varphi^i \varphi_i^\dagger)$ **discrete vacua, special directions**
- Extrema of $|\varphi_1|^4 + |\varphi_2|^4$ for fixed magnitude v :
Positive: $\propto (1, 1)/\sqrt{2} \rightarrow V \sim +2v^4/4$
Negative: $\propto (0, 1) \rightarrow V \sim -v^4$

Summary FS

Interlude

- Discrete symmetries have some interesting advantages
- Magnitudes: Abelian; predictions: non-Abelian
- Very natural extension beyond the SM