



ATLAS Physics Prospects for the High Luminosity LHC



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QUADRO DE REFERÊNCIA ESTRATÉGICO NACIONAL





Highlights from the first LHC physics run

Will show mainly ATLAS results CMS has similar results in most cases

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Data taking during LHC Run I











➢ Measured in the H→_{YY} and H→ZZ→41 decay channels

Better mass precision





Higgs mass meassurement

➢ Measured in the H→_{YY} and H→ZZ→4l decay channels

Better mass precision





CMS: $M_{H} = 125.7 \pm 0.3_{stat} \pm 0.3_{sys}$ GeV



➢ Measured in the H→_{γγ} and H→ZZ→41 decay channels

Better mass precision





Higgs signal strength







Higgs signal strength





Higgs signal strength





Three channels combined:

 $H \rightarrow \gamma \gamma$, $H \rightarrow ZZ$, $H \rightarrow WW$

- Combine J^p sensitive variables
- Combined in a likelihood fit



CLs of alternative J^p hypothesis assuming 0⁺





2013 NOBEL PRIZE IN PHYSICS François Englert Peter W. Higgs

② ② The Nobel Foundation, Photo: Lovisa Englide



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Fundación Príncipe de Asturias



SUSY searches up to now

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013 Model

ATLAS Preliminary

 $\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$

	Model	e, μ, τ, γ	Jets	$\mathbf{E}_{\mathrm{T}}^{\mathrm{miss}}$	∫£ dt[fb	⁻¹] Mass limit		Reference
Inclusive Searches	$\begin{array}{l} \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \tilde{g}\tilde{a}, \tilde{q} \rightarrow q \tilde{v}_{1}^{0} \\ \tilde{g}\tilde{s}, \tilde{g} \rightarrow q \tilde{v}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{v}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{v}_{1}^{0} \\ q \mathcal{W}^{2} \tilde{\chi}_{1}^{0} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\ell'(\ell'_{V}/v)) \tilde{\chi}_{1}^{0} \\ \text{GMSB} (\tilde{c} \text{ NLSP}) \\ \text{GMSB} (\tilde{c} \text{ NLSP}) \\ \text{GGM (bino NLSP)} \\ \text{GGM (hiog NLSP)} \\ \text{GGM (hiog Sino NLSP)} \\ \text{GGM (higgsino LSP)} \\ \text{GGM (higgsino LSP)} \\ \text{Gravitino LSP} \end{array}$	$\begin{matrix} 0 \\ 1 \ e, \mu \\ 0 \\ 0 \\ 1 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 1 - 2 \ \tau \\ 2 \ \gamma \\ 1 \ e, \mu + \gamma \\ \gamma \\ 2 \ e, \mu (Z) \\ 0 \end{matrix}$	2-6 jets 3-6 jets 7-10 jets 2-6 jets 3-6 jets 0-3 jets 0-2 jets 1 b 0-3 jets mono-jet	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 4.7 20.7 4.8 4.8 4.8 5.8 10.5	Image: Constraint of the second se	$ \begin{array}{llllllllllllllllllllllllllllllllllll$	ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 1308.1841 ATLAS-CONF-2013-047 ATLAS-CONF-2013-047 ATLAS-CONF-2013-089 1208.4688 ATLAS-CONF-2013-026 1209.0753 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152 ATLAS-CONF-2012-147
3 rd gen. <i>ẽ med</i> .	$\begin{array}{l} \tilde{g} \rightarrow b \bar{b} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{0} \\ \tilde{g} \rightarrow t \bar{t} \tilde{\chi}_{1}^{1} \\ \tilde{g} \rightarrow b \bar{t} \tilde{\chi}_{1}^{1} \end{array}$	0 0 0-1 e,μ 0-1 e,μ	3 b 7-10 jets 3 b 3 b	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	Ē 1.2 TeV Ē 1.1 TeV Ē 1.34 Te Ē 1.3 TeV	$\begin{array}{c} m(\tilde{\chi}_{1}^{0}) < 600 \ {\rm GeV} \\ m(\tilde{\chi}_{1}^{0}) < 350 \ {\rm GeV} \\ \\ \hline {\bf V} \qquad m(\tilde{\chi}_{1}^{0}) < 400 \ {\rm GeV} \\ \hline {\bf V} \qquad m(\tilde{\chi}_{1}^{0}) < 00 \ {\rm GeV} \end{array}$	ATLAS-CONF-2013-061 1308.1841 ATLAS-CONF-2013-061 ATLAS-CONF-2013-061
3 rd gen. squarks direct production	$ \begin{array}{l} \tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{x}_1^0 \\ \tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{x}_1^0 \\ \tilde{b}_1 \tilde{b}_1, \tilde{b}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ([\text{igh}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ([\text{igh}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ((\text{medum}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ((\text{meav}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ((\text{meav}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{x}_1 \tilde{t}_1 ((\text{mav}), \tilde{t}_1 \rightarrow b \tilde{x}_1^+ \\ \tilde{t}_1 \tilde{t}_1 ((\text{matural GMSB}) \\ \tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z \end{array} \right. $	$\begin{array}{c} 0 \\ 2 \ e, \mu \ (\text{SS}) \\ 1\text{-}2 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 3 \ e, \mu \ (Z) \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b ono-jet/c-t 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.7 20.7	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{split} m(\tilde{\tau}_1^0) -&90 \mbox{ GeV } \\ m(\tilde{\chi}_1^0) -&= 50 \mbox{ GeV } \\ m(\tilde{\chi}_1^0) -&= 55 \mbox{ GeV } \\ m(\tilde{\chi}_1^0) -&= 55 \mbox{ GeV } \\ m(\tilde{\chi}_1^0) -&= 0 \mbox{ GeV } \\ m(\tilde{\chi}_1^0) -&= 150 \mbox{ GeV } \\ m(\tilde$	1308.2631 ATLAS-CONF-2013-007 1208.4305, 1209.2102 ATLAS-CONF-2013-048 ATLAS-CONF-2013-048 ATLAS-CONF-2013-037 ATLAS-CONF-2013-027 ATLAS-CONF-2013-028 ATLAS-CONF-2013-028 ATLAS-CONF-2013-025
EW direct	$ \begin{split} \tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \to \ell\tilde{\chi}_{1}^{0} \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \to \tilde{\ell}\nu(\ell\tilde{\nu}) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{+} \to \tilde{\tau}\nu(\tilde{\nu}) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{0}^{0} \to \tilde{\ell}_{L}\nu\tilde{\ell}_{L}(\ell\tilde{\nu}\nu), \ell\tilde{\nu}\tilde{\ell}_{L}\ell(\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{0}^{0} \to W\tilde{\chi}_{0}^{0}\tilde{\ell}_{L}\tilde{\chi}_{L}(\ell\tilde{\nu}\nu) \\ \tilde{\chi}_{1}^{+}\tilde{\chi}_{0}^{0} \to W\tilde{\chi}_{0}^{0}\tilde{\ell}_{L}\tilde{\chi}_{0}^{0} \end{split} $	2 e,μ 2 e,μ 2 τ 3 e,μ 3 e,μ 1 e,μ	0 0 - 0 2 b	Yes Yes Yes Yes Yes Yes	20.3 20.3 20.7 20.7 20.7 20.7 20.3	$ \begin{array}{c c c c c c c c c c c c c c c c c c c $	$\begin{array}{l} m(\tilde{k}_1^0){=}0 \; \text{GeV} \\ m(\tilde{k}_1^0){=}0 \; \text{GeV}, m(\tilde{\ell}, \tilde{\nu}){=}0.5(m(\tilde{k}_1^+){+}m(\tilde{k}_1^0)) \\ m(\tilde{k}_1^0){=}0 \; \text{GeV}, m(\tilde{\ell}, \tilde{\nu}){=}0.5(m(\tilde{k}_1^+){+}m(\tilde{k}_1^0)) \\ m(\tilde{k}_1^0){=}m(\tilde{k}_2^0), m(\tilde{k}_1^0){=}0, m(\tilde{\ell}, \tilde{\nu}){=}0.5(m(\tilde{k}_1^0){+}m(\tilde{k}_1^0)) \\ m(\tilde{k}_1^0){=}m(\tilde{k}_2^0), m(\tilde{k}_1^0){=}0, sleptons\; decoupled \\ m(\tilde{k}_1^0){=}m(\tilde{k}_2^0), m(\tilde{k}_1^0){=}0, sleptons\; decoupled \end{array}$	ATLAS-CONF-2013-049 ATLAS-CONF-2013-049 ATLAS-CONF-2013-028 ATLAS-CONF-2013-035 ATLAS-CONF-2013-035 ATLAS-CONF-2013-093
Long-lived particles	Direct $\tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-}$ prod., long-lived $\tilde{\chi}_{1}^{\pm}$ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_{1}^{0} \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu})_{+}\tau(e$ GMSB, $\tilde{\chi}_{1}^{0} \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_{1}^{0}$ $\tilde{q}\tilde{q}, \tilde{\chi}_{1}^{-} \rightarrow qq\mu$ (RPV)	Disapp. trk 0 e, μ) 1-2 μ 2 γ 1 μ, displ. vtx	1 jet 1-5 jets - - -	Yes Yes - Yes -	20.3 22.9 15.9 4.7 20.3	x̂1 270 GeV 832 GeV ĝ 832 GeV 832 GeV k ⁰ / ₂ 475 GeV 832 GeV ĝ 230 GeV 475 GeV ĝ 1.0 TeV	$\begin{array}{l} m(\tilde{\chi}_1^+) \! = \!\! 160 \; {\rm MeV}, \; \tau(\tilde{\chi}_1^+) \! = \!\! 0.2 \; {\rm ns} \\ m(\tilde{\chi}_1^0) \! = \!\! 100 \; {\rm GeV}, \; 10 \; \mu {\rm s} \! < \! \tau(\tilde{g}) \! < \!\! 1000 \; {\rm s} \\ 10 \! < \!\! 10 \! < \!\! 10g \! \! 10g \!$	ATLAS-CONF-2013-069 ATLAS-CONF-2013-057 ATLAS-CONF-2013-058 1304.6310 ATLAS-CONF-2013-092
RPV	$\begin{array}{l} LFV pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e + \mu \\ LFV pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e(\mu) + \tau \\ Bilinear RPV CMSSM \\ \tilde{x}_{1}^{+} \tilde{x}_{1}^{-}, \tilde{x}_{1}^{+} \rightarrow W \tilde{x}_{1}^{0}, \tilde{x}_{1}^{0} \rightarrow e e \tilde{v}_{\mu}, e \mu \tilde{v} \\ \tilde{x}_{1}^{+} \tilde{x}_{1}^{-}, \tilde{x}_{1}^{+} \rightarrow W \tilde{x}_{1}^{0}, \tilde{x}_{1}^{1} \rightarrow t \tau \tilde{v}_{e}, e \tau \tilde{v} \\ \tilde{g}^{-} q q q \\ \tilde{g} \rightarrow \tilde{t}_{1}, \tilde{t}_{1} \rightarrow b s \end{array}$	$2 e, \mu 1 e, \mu + \tau 1 e, \mu e 4 e, \mu 3 e, \mu + \tau 0 2 e, \mu (SS)$	- 7 jets - 6-7 jets 0-3 b	- Yes Yes - Yes	4.6 4.7 20.7 20.7 20.3 20.7	Fr. 1.6 Fr. 1.1 TeV G. Z. 1.1 TeV K_1^2 760 GeV K_1^2 350 GeV G. Z. 916 GeV G. Z. 880 GeV	$ \begin{array}{llllllllllllllllllllllllllllllllllll$	1212.1272 1212.1272 ATLAS-CONF-2012-140 ATLAS-CONF-2013-036 ATLAS-CONF-2013-036 ATLAS-CONF-2013-091 ATLAS-CONF-2013-007
Other	Scalar gluon pair, sgluon $\rightarrow q\bar{q}$ Scalar gluon pair, sgluon $\rightarrow t\bar{t}$ WIMP interaction (D5, Dirac χ)	$2 e, \mu (SS)$	4 jets 1 <i>b</i> mono-jet	- Yes Yes	4.6 14.3 10.5	sgluon 100-287 GeV sgluon 800 GeV M* scale 704 GeV	incl. limit from 1110.2693 $m(\chi)$ 80 GeV, limit of <687 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
	$\sqrt{s} = 7$ leV	vs≐8 leV artial data	√s = full	8 TeV data		10 ⁻¹ 1	Mass scale [TeV]	

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

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Exotics



ATLAS Exotics Searches* - 95% CL Lower Limits (Status: May 2013) Large ED (ADD) : monojet + ET.miss 4.37 TeV M_D (δ=2) lb⁻¹, 7 TeV [1210,4491] Large ED (ADD) : monophoton + ET miss 1.93 TeV M_D (δ=2) 7 TeV [1209.4625 ATLAS Large ED (ADD) : diphoton & dilepton, myy /II Extra dimensions 7 TeV [1211.1150] 4.18 TeV M_S (HLZ δ=3, NLO) Preliminary UED : diphoton + $E_{T,miss}$ 1.40 TeV Compact. scale R L=4.8 fb⁻¹, 7 TeV [1209.0753] 4.71 TeV M_{KK} ~ R⁻¹ S¹/Z₂ ED : dilepton, m_i .=5.0 fb⁻¹, 7 TeV [1209.2535] RS1 : dilepton. m. 2.47 TeV Graviton mass (k/Mpl = 0.1) 20 fb⁻¹, 8 TeV [ATLAS-CONF-2013-017 RS1 : WW resonance, m_{T.bob} L=4.7 fb⁻¹, 7 TeV [1208.2880] **1.23 TeV** Graviton mass $(k/M_{Pl} = 0.1)$ $Ldt = (1 - 20) \text{ fb}^{-1}$ Bulk RS : ZZ resonance, m Graviton mass $(k/M_{Pl} = 1.0)$ =7.2 fb⁻¹, 8 TeV [ATLAS-CONF-2012-15 RS $g_{\mu\nu} \rightarrow t\bar{t}$ (BR=0.925) : $t\bar{t} \rightarrow l+jets, m_{\mu}$ 2.07 TeV g_{KK} mass s = 7.8 TeV ADD BH (M_{TH} /M_D=3) : SS dimuon, N_{ch. part} 1.25 TeV M_D (δ=6) ADD BH $(M_{TH}^{\prime}/M_{D}=3)$: leptons + jets, Σp 1.5 TeV M_D (δ=6) L=1.0 fb⁻¹, 7 TeV [1204,4646 Quantum black hole : dijet, F (m) 4.11 TeV M_D (δ=6) 7 TeV [1210.1718 gggg contact interaction : $\chi(m)$ 7.6 TeV A TeV [1210.1718 0 qqll CI : ee & μμ, m 13.9 TeV A (constructive int.) uutt CI : SS dilepton + jets + E7,miss 3.3 TeV Λ (C=1) Z' (SSM) : m_{ee/uu} 2.86 TeV Z' mass Z' (SSM) : m₇₇ 1.4 TeV Z' mass 4.7 fb⁻¹, 7 TeV [1210.6604 Z' (leptophobic topcolor) : $t\bar{t} \rightarrow l+jets, m$ 1.8 TeV Z' mass 14.3 fb⁻¹, 8 TeV [ATLAS-CO Ś W' (SSM) : m_{T.e/u} 2.55 TeV W' mass =4.7 fb⁻¹, 7 TeV [1209.4446] W' $(\rightarrow tq, g_{p}=1)$: m_{tr} =4.7 fb⁻¹, 7 TeV [1209.6593] 430 GeV W' mass W'_{R} (\rightarrow tb, LRSM) : m_{i} 1.84 TeV W' mass 8 TeV [ATLAS-C Scalar LQ pair (β =1) : kin. vars. in eeii. evii 1^{eff} gen. LQ mass =1.0 fb⁻¹, 7 TeV [1112.4828] 660 GeV 2 Scalar LQ pair (β =1) : kin. vars. in µµjj, µvjj 685 Gev 2nd gen, LQ mass =1.0 fb⁻¹, 7 TeV [1203.3172] =4.7 fb⁻¹, 7 TeV [1303.0526] Scalar LQ pair (β=1) : kin. vars. in ττjj, τvjj 534 GeV 3rd gen. LQ mass 656 Gev ť mass 4th generation : t't' \rightarrow WbWb 4th generation : b'b' \rightarrow SS dilepton + jets + $E_{T_{reles}}$ 1.7 fb⁻¹, 7 TeV [1210.5468 New quarks b' mass 720 GeV Vector-like guark : TT→ Ht+X T mass (isospin doublet) Vector-like quark : CC, ming 1.12 TeV VLQ mass (charge -1/3, coupling κ_{aQ} = v/m_o) =4.6 fb⁻¹, 7 TeV [ATLAS-CONF-2012-137] Excited quarks : y-jet resonance, m 7 TeV [1112.3580 2.46 TeV q* mass Excit. ferm. Excited quarks : dijet resonance, m .84 TeV q* mass Excited b quark : W-t resonance, m 870 Gev b* mass (left-handed coupling) =4.7 fb⁻¹, 7 TeV [1301.1583] Excited leptons : I-y resonance, m 2.2 TeV I* mass (Λ = m(I*)) Techni-hadrons (LSTC) : dilepton, mee/uu **850 GeV** ρ_{τ}/ω_{T} mass $(m(\rho_{\tau}/\omega_{T}) - m(\pi_{T}) = M_{u})$ Techni-hadrons (LSTC) : WZ resonance (IvII), m ρ_{T} mass $(m(\rho_{T}) = m(\pi_{T}) + m_{W}, m(a_{T}) = 1.1m(\rho_{T}))$ Major. neutr. (LRSM, no mixing) : 2-lep + jets 1.5 TeV N mass (m(W_) = 2 TeV) Other N^{\pm} mass (|V_{|}| = 0.055, |V_{|}| = 0.063, |V_{|}| = 0) Heavy lepton N[±] (type III seesaw) : Z-I resonance, m_{21} H_{\perp}^{\pm} (DY prod., BR($H_{\perp}^{\pm} \rightarrow \parallel)=1$) : SS ee ($\mu\mu$), m_{\perp}^{\pm} H^{±±} mass (limit at 398 GeV for μμ) Color octet scalar : dijet resonance, m 1.86 TeV Scalar resonance mass Multi-charged particles (DY prod.) : highly ionizing tracks 490 Gev mass (|q| = 4e) Magnetic monopoles (DY prod.) : highly ionizing tracks 862 GeV mass 111 1 1 1 1 1 10^{-1} 10^{2}

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Mass scale [TeV]

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Heavy Ion Physics highlights in Run I

Jet suppression in central Pb-Pb collisions: \succ





Heavy Ion Physics highlights in Run I









Physics at $\sqrt{s}=14$ TeV



Moving from 8 TeV to 14 TeV the cross sections of the following processes increase by

Top quark production ~ a factor 4

Low mass Higgs boson ~ factor 3

Exotic heavy particles, like Z', up to a factor 10

- Background processes, like W/Z boson production or jet cross sections increase more slowly
- The sensitivity improves considerably!



LHC Upgrade Schedule



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Why do we need so much luminosity?



The SM is incomplete







The stability of the Higgs mass

$$M_{H}^{2} = M_{bare}^{2} + \begin{pmatrix} H \\ H \end{pmatrix} + \begin{pmatrix} t \\ H \end{pmatrix} + \begin{pmatrix} t \\ H \end{pmatrix} + \begin{pmatrix} H \\ H \end{pmatrix} \end{pmatrix} + \begin{pmatrix} H \\ H \end{pmatrix} +$$

Given the observed mass of the Higgs boson, the stop mass should not be much larger than 1-1.5 TeV Accessible at the LHC?



SUSY at high luminosity





Vector-like quarks

 Vector-like quarks appear in certain models to cancel Higgs mass divergencies

Little Higgs, extra-dimensions

Left-handed and right-handed components transform in the same way under SU(2)



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In the absence of direct evidence for new physics, the Higgs will be

- Fundamental to test the validity of the SM!
- Probe of new physics

With 10 times more data:

- Precise measurements of Higgs production/decay rates, couplings and mass.
- Search for rare/ new/ invisible decay modes.
- Measurements of tensor structure of Higgs couplings and possible CP violating contributions.
- Search for additional Higgs bosons beyond the Standard Model.With 100 times more data
- Measurement of Higgs self-coupling becomes possible



Higgs physics in the next 10 years

In the absence of direct evidence for new physics, the Higgs will be

- Fundamental to test the validity of the SM!
- Probe of new physics
 - With 10 times more data:



- Precise measurements of Higgs production/decay rates, couplings and mass.
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Higgs @ 14 TeV





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Higgs signal strength projections

ATLAS Simulation Preliminary

 $\sqrt{s} = 14 \text{ TeV}: \int Ldt = 300 \text{ fb}^{-1}; \int Ldt = 3000 \text{ fb}^{-1}$



ATL-PHYS-PUB-2013-014

- Observe different production modes in each decay channel!
 - ≻ Even for $H \rightarrow \mu \mu$!
- Many channels will have uncertainties at the 10% level with 3000 fb⁻¹



Higgs couplings projections



	Snowmass 2013							
Model	Δκν	Δκ _γ	$\Delta \kappa_{b}$					
Singlet mixing	6%	6%	6%					
2HDMs	1%	1%	10%					
Composite Higgs	-3%	-9%	- (3-9)%					
Vector top partner	–2%	1%	-2%					
Decoupling MSSM	-0.0013%	<1.5%	1.6%					



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$H \rightarrow ZZ$ vertex tensor structure

Generic Higgs coupling vertex to a vector boson CP-Even tree level

$$\begin{aligned} A(X \to VV) &= v^{-1} \left(g_1^{(0)} m_V^2 \epsilon_1^* \epsilon_2^* \right) & \qquad \text{component} \\ &+ g_2^{(0)} f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + g_3^{(0)} f^{*(1),\mu\nu} f_{\mu\alpha}^{*(2)} \frac{q_\nu q^\alpha}{\Lambda^2} \\ &+ g_4^{(0)} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu} \right) & \qquad \text{CP-Even} \\ &\text{loop induced} \end{aligned}$$

- CP-conserving tree level SM: g1=1, g2=g3 = g4 = 0
- CP violation in the Higgs sector: g4 ≠ 0, either g1 or g2, g3 ≠ 0
- Sensitivity to new vertex contributions through angular distributions of the Higgs decay products





- > 8D fit to distributions of kinematic variables $(m_{4l}, m_{Z_{1,2}}, \theta_{1,2}, \phi, \phi_1, \theta^*)$
- Free parameters: real and imaginary part of g2, g4 relative to g1



Sensitive test of the tensor structure at the HL-LHC!

Luminosity	$ g_4 /g_1$	$\mathfrak{R}(g_4)/g_1$	$\Im(g_4)/g_1$	$ g_2 /g_1$	$\Re(g_2)/g_1$	$\Im(g_2)/g_1$
300 fb ⁻¹	1.20	(-0.88, 0.91)	(-1.02, 1.05)	1.02	(-0.84,0.44)	(-1.19, 1.18)
3000 fb ⁻¹	0.60	(-0.30, 0.33)	(-0.39, 0.42)	0.60	(-0.30,0.11)	(-0.71, 0.68)



- Many extensions of the SM predict more than one Higgs
- > 2 Higgs Doublet Model

Add one more complex doublet

Predicts 5 Higgs bosons:

2 CP-even (h, H)

- 1 CP-odd (A),
- 2 charged (H $^{\scriptscriptstyle\pm}$)
- Parameterized in terms of tan $β ≡ v_2/v_1$ and α = mixing angle CP even bosons
- Modified couplings to h (SM-like Higgs boson):

	Coupling strength	Type I	Type II	Type III	Type IV
	KV	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha)$	$\sin(\beta - \alpha)$
	K _u	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$	$\cos(\alpha)/\sin(\beta)$
	Kd	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$
	ĸı	$\cos(\alpha)/\sin(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$-\sin(\alpha)/\cos(\beta)$	$\cos(\alpha)/\sin(\beta)$
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> 95% CL upper limits on $\sigma \times BR$:

Between 5-0.07 fb in the mass range between 200-900 GeV with 3000 fb⁻¹ of integrated luminosity

➢ 3-4 worse limits with only 300 fb⁻¹ of accumulated luminosity



- WIMP: assumed to interact very weakly with visible matter, except for the Higgs boson
- Search for invisible Higgs decays

Assume couplings to other known particles as in SM

Interpreted in terms of dark matter particles coupling only to Higgs sector

on WIMP spin

WIMP mass $< 2m_{_{\rm H}}$







Non-abelian structure of the Electroweak interaction in the SM implies the existence of triple and quartic gauge boson couplings

Weak force

 $L_{\text{WEAK}} = -1/4F^{(a)}_{\mu\nu}F^{(a)\mu\nu} + i\overline{\Psi}_{\text{L}}(\gamma^{\mu}(D_{\mu}) - m)\Psi_{\text{L}}$ where: $\mathbf{D}_{\mu} = \partial_{\mu} + igT^{a}W^{a}_{\mu}$ a is sum over 3 gauge bosons **T**^a = 0.5*Pauli spin matrices (PDG!) W^a_u is weak gauge boson field $\mathbf{F}^{(a)}_{\mu\nu} = \partial_{\mu} \mathbf{W}^{a}_{\nu} - \partial_{\nu} \mathbf{W}^{a}_{\mu} - \mathbf{g}_{s} \varepsilon_{abc} \mathbf{W}^{b}_{\mu} \mathbf{W}^{c}_{\nu}$ $\varepsilon_{abc} = SU(2)$ structure constants F...,F^{µv} contains triple. quartic couplings



- In the SM, triple and quartic gauge boson couplings constrained by the symmetries of the Lagrangian
- Very precise predictions:

WWWW, yyWW, WWZZ exist

ZZZZ not present as quartic gauge boson coupling, but mediated by the Higgs at tree level

- γγZZ: only at loop level
- ≻ WWWW:

Related to the EW Symmetry Breaking mechanism:

Longitudinal modes are Goldstone bosons

Beyond the SM physics can modify the quartic gauge boson couplings

Extra-dimensions



AFP



- Tag and measure protons at ±210 m
- Radiation hard edgeless 3D Silicon detectors
- 10 ps timing detectors

Allow running in high pile up conditions Associate protons with correct primary vertex

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pW p ppp

Forward proton tagging

converts the LHC in a $\gamma\gamma$

collider!



Search for yyWW anomalous couplings

- Reach on anomalous couplings studied using a full simulation of the ATLAS detector, including all pile-up effects; only leptonic decays of *W*s are considered
- Signal appears at high lepton p_T and dilepton mass (central ATLAS) and high diffractive mass (reconstructed using forward detectors)
- Cut on the number of tracks fitted to the primary vertex: very efficient to remove remaining pile-up after requesting a high mass object to be produced (for signal, we have two leptons coming from the W decays and nothing else)





Couplings	OPAL limits	Sensitivity ($\mathcal{Q} \ \mathcal{L} = 30$ (200) fb ⁻¹
	$[GeV^{-2}]$	5σ	95% CL
a_0^W/Λ^2	[-0.020, 0.020]	5.4 10 ⁻⁶	$2.6 \ 10^{-6}$
		(2.7 10 ⁻⁶)	$(1.4 \ 10^{-6})$
a_C^W/Λ^2	[-0.052, 0.037]	$2.0 \ 10^{-5}$	9.4 10^{-6}
		(9.6 10 ⁻⁶)	$(5.2 \ 10^{-6})$
a_0^Z/Λ^2	[-0.007, 0.023]	$1.4 10^{-5}$	6.4 10 ⁻⁶
		$(5.5 \ 10^{-6})$	$(2.5 \ 10^{-6})$
a_C^Z/Λ^2	[-0.029, 0.029]	$5.2 \ 10^{-5}$	2.4 10 ⁻⁵
		(2.0 10 ⁻⁵)	$(9.2 \ 10^{-6})$

- Improve LEP sensitivity by more than 4 orders of magnitude with 30/200 fb⁻¹ at LHC, and of D0/CDF results by ~2 orders of magnitude
- AFP improves the results obtained with central detector only by 2 orders of magnitude
- Reaches the sensitivity needed for extra-dimensions models!!



Vector boson scattering



Divergence in WWWW scattering was the original motivation to introduce the Higgs mechanism

Measure this process: closure test of the SM

Sensitive also to new physics: additional Higgs bosons, other scalars, new heavy bosons, ...





Experimental signature

Two leptons (same sign), large missing $E_{_{T}}$

Two forward/backward jets with m_{ii} > 1 TeV

Sensitive to dimension 8 operators





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LIP Seminar 6 March 2014



FCNC in top quark decays



BR($t \rightarrow$ FCNC) in several models:

	SM	QS	2HDM	FC 2HDM	MSSM	R SUSY	TC2	RS
$t \rightarrow q\gamma$	$\sim 10^{-14}$	$\sim 10^{-9}$	$\sim 10^{-6}$	$\sim 10^{-9}$	$\sim 10^{-6}$	$\sim 10^{-6}$	$\sim 10^{-6}$	$\sim 10^{-9}$
$t \rightarrow qZ$	$\sim 10^{-14}$	$\sim 10^{-4}$	$\sim 10^{-7}$	$\sim 10^{-10}$	$\sim 10^{-6}$	$\sim 10^{-5}$	$\sim 10^{-4}$	$\sim 10^{-5}$
$t \rightarrow qg$	$\sim 10^{-12}$	$\sim 10^{-7}$	$\sim 10^{-4}$	$\sim 10^{-8}$	$\sim 10^{-5}$	$\sim 10^{-4}$	$\sim 10^{-4}$	$\sim 10^{-9}$

Acta Phys. Polon. B35 (2004) 2695–2710; Phys. Rev. D68 (2003) 015002; Phys. Rev. D 75 (2007) 015002

present experimental limits:

	LEP	HERA	Tevatron	LHC
$BR(t \rightarrow q\gamma)$	2.4%	0.47 %	3.2 %	—
$BR(t \rightarrow qZ)$	7.8%	30%	3.2 %	0.05%
			2.0 × 10 ⁻⁴ (tug)	5.7 × 10 ^{–5} (tug)
$BR(t \rightarrow qg)$	17%	13%	3.9 × 10 ⁻³ (tcg)	$2.7 \times 10^{-4} (tcg)$

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FCNC in top quark decays

- With 300 fb⁻¹ ATLAS will improve the sensitivity to FCNC in top quark decays by ~ two orders of magnitude
- With 3000 fb⁻¹ limits in the BR can reach ~10⁻⁵-5×10⁻⁵





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Not only the pp physics will benefit from the Upgrade! Run I:

year	system	√s _{NN} (TeV)	L _{int}
2010	Pb-Pb	2.76	~ 10 μb ⁻¹
2011	рр	2.76	~ 250 nb⁻¹
2011	Pb-Pb	2.76	~ 150 μb ⁻¹
2013	p-Pb	5.02	~ 30 nb ⁻¹
2013	рр	2.76	~ 5 pb ⁻¹

> Run II: 5-10 times higher stats!
 Pb-Pb ~1nb⁻¹ @ √s_{NN} ~ 5.1 TeV

p-Pb (at increased luminosity?)

pp reference at Pb-Pb energy (5.1 TeV)

Extend current results to larger energies. Explore new observables

Run III, IV: experiments request
 >10 nb⁻¹ Pb-Pb
 p-Pb high luminosity, pp reference 5.5 TeV, possibly light ions (e.g. Ar-Ar)



Focus on rare probes, study their coupling with QGP medium and their (medium-modified) hadronization process

- Jets: characterization of energy loss testing ground for the multi-particle aspects of QCD
 - probe of the medium density
- Heavy flavour:
 - mass dependence of energy loss
- Quarkonium dissociation pattern and regeneration

probe of deconfinement and medium temperature





Summary and conclusions



- Very successful LHC Run I
 Higgs discovery
 But no new physics observed
- We know there is something more The SM is not the end of the story
- We will use the LHC to try to find the answer to many of the unknown questions
 - Precision studies of the Higgs sector and SM predictions
 - Searches for deviations from predictions
 - Direct searches for SUSY or other new physics



Backup



Use Higgs rate measurements in different production and decay modes to set limits in the tan β, cos (β-α) plane

Limits depends on the 2HDM type

Considerable improvements in the limits after 3000 fb⁻¹ of integrated luminosity



ATL-PHYS-PUB-2013-015



Cuts	Тор	Dibosons	Drell-Yan	W/Z+jet	Diffr.	$a_0^W/\Lambda^2 = 5 \cdot 10^{-6} \text{ GeV}^{-2}$
$\begin{array}{l} \mbox{timing} < 10 \mbox{ ps} \\ p_T^{lep1} > 150 \mbox{ GeV} \\ p_T^{lep2} > 20 \mbox{ GeV} \end{array}$	5198	601	20093	1820	190	282
M(11)>300 GeV	1650	176	2512	7.7	176	248
nTracks ≤ 3	2.8	2.1	78	0	51	71
$\Delta \phi < 3.1$	2.5	1.7	29	0	2.5	56
$m_X > 800 \text{ GeV}$	0.6	0.4	7.3	0	1.1	50
$p_T^{lep1} > 300 \text{ GeV}$	0	0.2	0	0	0.2	35

Table 9.5. Number of expected signal and background events for 300 fb^{-1} at pile-up $\mu = 46$. A time resolution of 10 ps has been assumed for background rejection. The diffractive background comprises production of QED diboson, QED dilepton, diffractive WW, double pomeron exchange WW.



Diffractive physics with AFP

Central Exclusive production



Higgs

Good mass resolution from p energy loss (~2-3%) Tests J^{CP} & improve b-Higgs coupling measurement Di-jets

Test QCD calculations

Triple and quartic anomalous gauge boson couplings



 Exclusive production of WW via photon exchange (σ_{ww→ww} well kno p



Diffractive mass computed from p energy loss

Large mass tails sensitive to anomalous couplings

Improvement of LEP sensitivity by more than 4 orders of magnitude with 30-200 fb⁻¹





Very good mass resolution (few %)



Vector-like quarks

 Vector-like quarks appear in certain models to cancel Higgs mass divergencies

Little Higgs, extra-dimensions

- Left-handed and right-handed components transform in the same way under SU(2)
- In some models, they interact mainly with third generation quarks

Model independent search:

Scan over BR's in a 2D plane

Assume sum BR = 1







- Use sensitive variables to discriminate against SM backgrounds
- > Example of the T \rightarrow Zt+X channel





Vector-like quark limits

